Spinorial Tensors

If we substitute in

$$\begin{split} & t_{ab} = \left(-\frac{1}{\sqrt{2}} \, \mathcal{T}_{a}^{A\dot{A}} \, \right) \left(-\frac{1}{\sqrt{2}} \, \mathcal{T}_{b}^{B\dot{B}} \right) \, t_{A\dot{A}B\dot{B}} \\ & = \frac{1}{2} \, \mathcal{T}_{a}^{A\dot{A}} \, \mathcal{T}_{b}^{B\dot{B}} \, \left(\, \mathcal{T}_{AB} \, \mathcal{E}_{\dot{A}\dot{B}} \, + \, \mathcal{T}_{\dot{A}\dot{B}} \, \mathcal{E}_{AB} \right) \\ & = \frac{1}{2} \, \mathcal{T}_{AB} \, \mathcal{T}_{a}^{(A\, \dot{A}\dot{B}\dot{I})} \, \, \mathcal{E}_{\dot{A}\dot{B}} \, + \, \mathcal{T}_{\dot{A}\dot{B}} \, \mathcal{E}_{\dot{A}\dot{B}} \, \mathcal{E}_{\dot{A}\dot{B}} \\ & = \frac{1}{2} \, \mathcal{T}_{AB} \, \mathcal{S}_{ab}^{A\dot{B}} \, + \, \mathcal{T}_{\dot{A}\dot{B}} \, \mathcal{S}_{ab}^{\dot{A}\dot{B}} \\ & = \frac{1}{2} \, \mathcal{T}_{AB} \, \mathcal{S}_{ab}^{\dot{A}\dot{B}} \, + \, \mathcal{T}_{\dot{A}\dot{B}} \, \mathcal{S}_{ab}^{\dot{A}\dot{B}} \\ \mathcal{S}_{ab}^{\dot{A}\dot{B}} = \mathcal{T}_{a}^{(A\, \dot{I}\dot{B}\dot{I})} \, \mathcal{E}_{\dot{B}}^{\dot{A}\dot{B}} \, + \, \mathcal{T}_{\dot{A}\dot{B}} \, \mathcal{S}_{ab}^{\dot{A}\dot{B}} = \mathcal{T}_{a}^{(A\, \dot{I}\dot{B}\dot{I})} \, \mathcal{E}_{\dot{B}}^{\dot{A}\dot{B}} \\ \mathcal{S}_{ab}^{\dot{A}\dot{B}} = \mathcal{T}_{a}^{\dot{A}\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}_{\dot{A}\dot{B}}^{\dot{A}\dot{B}}^{\dot{A}\dot{B}} \, \mathcal{T}$$

Maxwell's equations

Homogeneous: $\nabla \cdot \vec{B} = 0$, $\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$, $\vec{E} \in \mathbb{R}^3$.

Inhomogeneous: $\nabla \cdot \vec{E} = \beta$, $\nabla \times \vec{B} + \frac{\partial \vec{E}}{\partial +} = \vec{J}$

Tensorial form:

Homework: Prove that from the tensorial form we can get the normal form of maxwell equations.

The spinorial equivalent of Fab

with $f_{AB} = \frac{1}{2} + \frac{\dot{R}}{A}_{B\dot{R}}$, $f_{\dot{A}\dot{B}} = \frac{1}{2} + \frac{\dot{R}}{AR\dot{B}}$ symmetrics.

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Homework: Prove that effectively fix = fax.

So the spinorial equivalent to the maxwell equations are $-\partial_{BB}(f^{AB}E^{\dot{A}\dot{B}}+f^{\dot{A}\dot{B}}E^{AB})=J^{A\dot{A}}$ and $-i\partial_{B\dot{B}}(f^{AB}E^{\dot{A}\dot{B}}+f^{\dot{A}\dot{B}}E^{AB}=0$

Combining them

$$\partial_{B}^{\dot{A}} f^{AB} = 2 J^{A\dot{A}}, \quad \partial_{\dot{B}}^{\dot{A}\dot{B}} = 2 J^{A\dot{A}}$$

Due to the signature, the second equation is the complex conjugate of the first.

Using the spinorial tensors

from db*Fab=0, then daFbc+dbFab+dcFab=0, Implies the existence of a vector field Aa, such that

If ABB IS the spinorial equivalent to Aa

Analogously, from dbFab=0, exist a vector field Åa.

*Fab= da Åb-db Åa

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 $f_{bc} = i \partial_{(B} \dot{A}_{c)\dot{B}}, \quad f_{bc} = i \partial^{(B} \dot{A}_{Bic)}$

the complex vector field.

 $\phi_b = \frac{1}{2} (A_a - i \dot{A}_a)$

It follows that

 $\partial_{(\mathcal{B}}\dot{\mathcal{B}} + \partial_{(\mathcal{B})\dot{\mathcal{B}}} = \underbrace{1}_{2} (\partial_{(\mathcal{B}}\dot{\mathcal{B}} + \partial_{(\mathcal{B})\dot{\mathcal{B}}} - \partial_{(\mathcal{B})\dot{\mathcal{B}}}\dot{\mathcal{A}}_{(\mathcal{B})\dot{\mathcal{B}}}) = 0.$

i.e.,

dis b cis= O.

Maxwell equation in spinorial form.