

Second Quantization

Representation of number of occupation.

$\{|v_1\rangle, |v_2\rangle, \dots\}$ Complete base of state

$|n_{v_1}, n_{v_2}, \dots\rangle$ State of N-particles

$$\sum_j n_{v_j} = N$$

$\hat{n}_{v_j} |n_{v_j}\rangle = n_{v_j} |n_{v_j}\rangle$ Number of operator

$$n_{v_j} = \begin{cases} 0, 1 & \text{Fermions} \\ 0, 1, 2, \dots & \text{Bosons} \end{cases}$$

$$\mathcal{F} = \mathcal{F}_0 \oplus \mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \dots \quad \text{Fock space}$$

$$\mathcal{F}_N = \text{span} \{ |n_{v_1}, n_{v_2}, \dots\rangle : \sum_j n_{v_j} = N \}$$

N Fermions

$$0 \quad |0, 0, 0, \dots\rangle$$

$$1 \quad |1, 0, 0, \dots\rangle, |0, 1, 0, \dots\rangle, |0, 0, 1, \dots\rangle, \dots$$

$$2 \quad |1, 1, 0, 0, \dots\rangle, |0, 1, 1, 0, \dots\rangle, |1, 0, 1, 0, \dots\rangle, \dots$$

\vdots

N Bosons

$$0 \quad |0, 0, 0, \dots\rangle$$

$$1 \quad |1, 0, 0, \dots\rangle, |0, 1, 0, \dots\rangle$$

$$2 \quad |2, 0, 0, \dots\rangle, |0, 2, 0, \dots\rangle, |1, 1, 0, \dots\rangle$$

Creation and annihilation operator for bosons

$$\hat{b}_{v_j}^+ | \dots, n_{v_{j-1}}, n_{v_j}, n_{v_{j+1}}, \dots \rangle \quad \text{Creation operator}$$

$$= B_+(n_{\nu_j}) | \dots, n_{\nu_{j-1}}, n_{\nu_j} + 1, n_{\nu_{j+1}}, \dots \rangle$$

$B_+(n_{\nu_j})$ normalization constant.

$$\langle n_{\nu_j} + 1 | \hat{b}_{\nu_j}^+ | n_{\nu_j} \rangle \quad \text{Elements} \neq 0$$

$$\langle n_{\nu_j} + 1 | \hat{b}_{\nu_j}^+ | n_{\nu_j} \rangle^* = \langle n_{\nu_j} | (\hat{b}_{\nu_j}^+)^{\dagger} | n_{\nu_j} + 1 \rangle$$

$$\hat{b}_{\nu_j} := (\hat{b}_{\nu_j}^+)^{\dagger} \rightarrow \text{Annihilation operator}$$

$$\hat{b}_{\nu_j} | \dots, n_{\nu_{j-1}}, n_{\nu_j}, n_{\nu_{j+1}}, \dots \rangle$$

$$= B_-(n_{\nu_j}) | \dots, n_{\nu_{j-1}}, n_{\nu_j} - 1, n_{\nu_{j+1}}, \dots \rangle$$

$$[\hat{A}, \hat{B}] = \hat{A}\hat{B} - \hat{B}\hat{A} \quad \text{commutator.}$$

Is clear that

$$\hat{b}_{\nu_j}^+ \hat{b}_{\nu_k}^+ | \rangle = \hat{b}_{\nu_k}^+ \hat{b}_{\nu_j}^+ | \rangle \rightarrow [\hat{b}_{\nu_j}^+, \hat{b}_{\nu_k}^+] = 0$$

and

$$[\hat{b}_{\nu_j}, \hat{b}_{\nu_k}] = 0 \quad \text{if } j \neq k$$

but, if $k=j$ we have to be careful

$$\hat{b}_{\nu_j} | \dots, 0, \dots \rangle = 0 \rightarrow B_-(0) = 0$$

$$\text{and, } \hat{b}_{\nu_j}^+ | \dots, 0, \dots \rangle = | \dots, 1, \dots \rangle$$

where we have taken $B_+(0) = 1$,

then,

$$\hat{b}_{\nu_j}^+ \hat{b}_{\nu_j}^+ | 0 \rangle = | 0 \rangle \neq \hat{b}_{\nu_j}^+ \hat{b}_{\nu_j} | 0 \rangle = 0$$

Summarizing,

$$[\hat{b}_{\nu_j}^+, \hat{b}_{\nu_k}^+] = 0, [\hat{b}_{\nu_j}, \hat{b}_{\nu_k}] = 0, [\hat{b}_{\nu_j}, \hat{b}_{\nu_k}^+] = \delta_{\nu_j, \nu_k}$$

\hat{b} and \hat{b}^+ are not hermitian.

$\hat{b}_\nu^\dagger \hat{b}_\nu$ Operator:

$$[\hat{b}_\nu^\dagger \hat{b}_\nu, \hat{b}_\nu] = \hat{b}_\nu^\dagger \hat{b}_\nu \hat{b}_\nu - \hat{b}_\nu \hat{b}_\nu^\dagger \hat{b}_\nu = \hat{b}_\nu^\dagger \hat{b}_\nu \hat{b}_\nu - (1 + \hat{b}_\nu^\dagger \hat{b}_\nu) \hat{b}_\nu = -\hat{b}_\nu$$

$$[\hat{b}_\nu^\dagger \hat{b}_\nu, \hat{b}_\nu^\dagger] = \hat{b}_\nu^\dagger \hat{b}_\nu \hat{b}_\nu^\dagger - \hat{b}_\nu^\dagger \hat{b}_\nu^\dagger \hat{b}_\nu = \hat{b}_\nu^\dagger \hat{b}_\nu \hat{b}_\nu^\dagger - \hat{b}_\nu^\dagger (\hat{b}_\nu \hat{b}_\nu^\dagger - 1) = \hat{b}_\nu^\dagger$$

$$\langle \phi | \hat{b}_\nu^\dagger \hat{b}_\nu | \phi \rangle \geq 0 \quad \text{Norm of } \hat{b}_\nu | \phi \rangle$$

let $\hat{b}_\nu^\dagger \hat{b}_\nu | \phi_\lambda \rangle = \lambda | \phi_\lambda \rangle$ with $\lambda > 0$

let us take $\lambda = \lambda_0$ particular

$$\begin{aligned} (\hat{b}_\nu^\dagger \hat{b}_\nu^\dagger) \hat{b}_\nu | \phi_{\lambda_0} \rangle &= (\hat{b}_\nu^\dagger \hat{b}_\nu^\dagger - 1) \hat{b}_\nu | \phi_{\lambda_0} \rangle \\ &= \hat{b}_\nu (\hat{b}_\nu^\dagger \hat{b}_\nu - 1) | \phi_{\lambda_0} \rangle = \hat{b}_\nu (\lambda_0 - 1) | \phi_{\lambda_0} \rangle = (\lambda_0 - 1) \hat{b}_\nu | \phi_{\lambda_0} \rangle \end{aligned}$$

therefore,

$$(\hat{b}_\nu^\dagger \hat{b}_\nu) \hat{b}_\nu | \phi_{\lambda_0} \rangle = (\lambda_0 - 1) \hat{b}_\nu | \phi_{\lambda_0} \rangle$$

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Eigenvalues equation

$$\begin{aligned} (\hat{b}_\nu^\dagger \hat{b}_\nu) \hat{b}_\nu \hat{b}_\nu | \phi_{\lambda_0} \rangle &= (\hat{b}_\nu^\dagger \hat{b}_\nu - 1) \hat{b}_\nu \hat{b}_\nu | \phi_{\lambda_0} \rangle \\ &= \hat{b}_\nu (\hat{b}_\nu^\dagger \hat{b}_\nu - 1) \hat{b}_\nu | \phi_{\lambda_0} \rangle \\ &= \hat{b}_\nu (\lambda_0 - 2) \hat{b}_\nu | \phi_{\lambda_0} \rangle = (\lambda_0 - 2) \hat{b}_\nu \hat{b}_\nu | \phi_{\lambda_0} \rangle \end{aligned}$$

then

$$\begin{aligned} &\hat{b}_\nu \hat{b}_\nu (\lambda_0 - 2) | \phi_{\lambda_0} \rangle \\ \rightarrow &(\hat{b}_\nu^\dagger \hat{b}_\nu) \hat{b}_\nu^2 | \phi_{\lambda_0} \rangle = (\lambda_0 - 2) \hat{b}_\nu^2 | \phi_{\lambda_0} \rangle \\ \rightarrow &\hat{b}_\nu^\dagger \hat{b}_\nu (\hat{b}_\nu^m | \phi_{\lambda_0} \rangle) = (\lambda_0 - m) (\hat{b}_\nu^m | \phi_{\lambda_0} \rangle) \end{aligned}$$

$$\text{if } \lambda_0 \in \mathbb{Z}^+ \rightarrow \exists m = \lambda_0 \text{ such that } \hat{b}_\nu^\dagger \hat{b}_\nu (\hat{b}_\nu^m | \phi_{\lambda_0} \rangle) = 0$$

The process ends.

$$\text{if } \lambda_0 \notin \mathbb{Z}^+ \exists m \text{ such that } \hat{b}_\nu^\dagger \hat{b}_\nu (\hat{b}_\nu^m | \phi_{\lambda_0} \rangle) < 0$$

then $\lambda < 0$! \rightarrow Contradiction

therefore $\lambda := n = 0, 1, 2, \dots$

$$|\phi_\lambda\rangle := |n_\nu\rangle \rightarrow \hat{b}_\nu^\dagger \hat{b}_\nu |n_\nu\rangle = n_\nu |n_\nu\rangle \quad \text{number operator}$$

and $\hat{b}_\nu |n_\nu\rangle \sim |n_\nu - 1\rangle$

Now,

$$\begin{aligned} \hat{b}_\nu^\dagger \hat{b}_\nu \hat{b}_\nu^\dagger |\phi_{\lambda_0}\rangle &= \hat{b}_\nu^\dagger (1 + \hat{b}_\nu^\dagger \hat{b}_\nu) |\phi_{\lambda_0}\rangle \\ &= \hat{b}_\nu^\dagger (1 + \lambda_0) |\phi_{\lambda_0}\rangle \\ &= (\lambda_0 + 1) \hat{b}_\nu^\dagger |\phi_{\lambda_0}\rangle \end{aligned}$$

then $\hat{b}_\nu^\dagger \hat{b}_\nu \hat{b}_\nu^\dagger^n |\phi_{\lambda_0}\rangle = (\lambda_0 + n) \hat{b}_\nu^\dagger^n |\phi_{\lambda_0}\rangle$

thus, $(\hat{b}_\nu^\dagger \hat{b}_\nu) \hat{b}_\nu^\dagger |n_\nu\rangle = (n+1) \hat{b}_\nu^\dagger |n_\nu\rangle$

Finally $\hat{b}_\nu^\dagger |n_\nu\rangle \sim |n_\nu + 1\rangle$

So, we have to

$$\|\hat{b}_\nu |n_\nu\rangle\|^2 = (\hat{b}_\nu |n_\nu\rangle)^\dagger (\hat{b}_\nu |n_\nu\rangle) = \langle n_\nu | \hat{b}_\nu^\dagger \hat{b}_\nu |n_\nu\rangle = n_\nu$$

$$\|\hat{b}_\nu^\dagger |n_\nu\rangle\|^2 = (\hat{b}_\nu^\dagger |n_\nu\rangle)^\dagger (\hat{b}_\nu^\dagger |n_\nu\rangle) = \langle n_\nu | \hat{b}_\nu \hat{b}_\nu^\dagger |n_\nu\rangle = n_\nu + 1$$

$$\hat{b}_\nu^\dagger \hat{b}_\nu = \hat{n}_\nu \quad \hat{b}_\nu^\dagger \hat{b}_\nu |n_\nu\rangle = n_\nu |n_\nu\rangle \quad n_\nu = 0$$

$$\hat{b}_\nu |n_\nu\rangle = \sqrt{n_\nu} |n_\nu - 1\rangle$$

$$\hat{b}_\nu^\dagger |n_\nu\rangle = \sqrt{n_\nu + 1} |n_\nu + 1\rangle$$

$$\begin{aligned} (\hat{b}_\nu^\dagger)^{n_\nu} |0\rangle &= \sqrt{0+1} \sqrt{1+1} \sqrt{2+1} \cdots \sqrt{n_\nu-1+1} |n_\nu\rangle \\ &= \sqrt{1 \cdot 2 \cdot 3 \cdots n_\nu} |n_\nu\rangle = \sqrt{n_\nu!} |n_\nu\rangle \end{aligned}$$

Equivalence between states in first and second quantization.

$$\hat{S}_+ |n_{n_1}\rangle_1 |n_{n_2}\rangle_2 \cdots |n_{n_N}\rangle_N \leftrightarrow \hat{b}_{n_1}^\dagger \hat{b}_{n_2}^\dagger \cdots \hat{b}_{n_N}^\dagger |0\rangle$$