Riemann-Lebesque Lemma.

let f(t) be continuous in (a,b), then

$$I(x) = \int_{a}^{b} f(t) e^{ixt} dt = O\left(\frac{1}{x}\right), \quad as \quad x - \infty$$

Provided that the integral.

converges.

Proof:

$$I(x) = \int_{a}^{a+\pi/x} f(t) e^{ixt} dt + \int_{a+\pi/x}^{b} f(t) e^{ixt} dt.$$

$$I(x) = \int_{a}^{b-\pi/x} f(t) e^{ixt} dt + \int_{b-\pi/x}^{b} f(t) e^{ixt} dt.$$

changing
$$t = t - \frac{\pi}{x}$$

$$\int_{a}^{a+\pi/x} f(t) e^{ixt} dt = -\int_{a}^{b-\pi/x} f(t-\pi) e^{ix(t-\pi)} dt$$

then
$$I(x) = \frac{1}{2} \int_{0}^{a+\overline{I}} f(\xi) e^{ixt} dt + \frac{1}{2} \int_{0}^{a+\overline{I}} f(\xi) e^{ixt} dt$$

$$+ \frac{1}{2} \int_{0}^{b+\overline{I}} [f(\xi) - f(\xi + \overline{I})] e^{ixt} dt, \text{ we want } x \longrightarrow \infty.$$

Mean Value Theorem: f(t) continuous and bounded on [a,b].

$$\int_{a}^{b} f(t)dt = f'(c)(b-a), \text{ for some red-number } ct[a,b]$$

the first two integrals
$$\int_{a}^{a+\frac{\pi}{X}} f(t)e^{ixt}dt = f'(\frac{\pi}{X})(\frac{\pi}{X}) \longrightarrow O(\frac{1}{X})$$

and
$$\int_{P^{-\frac{x}{x}}} f(t) e^{ixt} dt = f'\left(\frac{x}{\pi}\right)\left(\frac{x}{\pi}\right) \longrightarrow \mathcal{O}\left(\frac{1}{x}\right).$$

Finally as
$$f(t)$$
 is continuous $\forall t \in [a,b]$

$$\lim_{X \to \infty} \int_{0}^{b-\pi/x} [f(t) - f(t+\pi)] e^{ixt} dt = 0$$

$$T(x) \sim O\left(\frac{1}{x}\right)$$

We may extend the Riemann-Lebesge Lemma to generalized Fourier lategrals as long as |f(t)| is integrable, $\Psi(t)$ is continuous differentiable and $\Psi'(t) \neq 0$. Take.

$$T(x) = \int_{\alpha}^{b} f(\xi) e^{ix \varphi(\xi)} d\xi$$

$$= \frac{1}{x_{i}} \int_{\alpha}^{b} \frac{f(\xi)}{\varphi'(\xi)} d\xi (e^{ix \varphi(\xi)}) d\xi$$

$$T(x) = \frac{1}{x_{i}} \frac{f(\xi)}{\varphi'(\xi)} e^{ix \varphi(\xi)} \Big|_{\alpha}^{b} - \frac{1}{x_{i}} \int_{\alpha}^{b} \frac{d\xi}{d\xi} \left(\frac{f(\xi)}{\varphi'(\xi)} e^{ix \varphi(\xi)} d\xi \right)$$
This vanishes more expediq than $\frac{1}{x}$ as $x \to \infty$.

$$T(x) \sim \frac{1}{ix} \frac{f(t)}{\psi'(t)} e^{ix\psi(t)} \Big|_{\alpha}^{b} as x \longrightarrow \infty$$

This method does not work for stationary points $\psi'(t)=0$. The method of stationary phase will give the asymptotic behaviour of generalized Fourier integrals with stationary points.

Choose the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and $\Psi'(t) = 0$ for as the integral such that $\Psi'(a) = 0$, and Ψ'

E70 is a small parameter

Note:

$$\int_{a+\epsilon}^{b} f(t) e^{ix\psi(t)} dt - \frac{1}{ix} \frac{f(t)}{\psi'(t)} e^{ix\psi(t)} \Big|_{a+\epsilon}^{b} - \frac{1}{x} \quad as \quad x \to \infty.$$

In the first integral change

$$f(t) \longrightarrow f(a)$$
 $\psi(t) \longmapsto \psi(a) + \psi(a) (t-a)^{p}$

As before (for laplace)

the leading contribution comes from a verghbourhood of the stationary point.

Where $\varphi'(a) = \varphi'(a) = ... = \varphi^{(p-1)}(a) = 0$

$$I_{1}(X) = \int_{a}^{a+\epsilon} \{(a) \exp \left\{ ix \left[\psi(a) + \frac{1}{\ell!} \psi^{p}(a) (t-a)^{\ell} \right] \right\} dt$$

Next, replace & -> ~ (this will introduce error terms &(1/x))

Let s=t-a, ds=dt

$$I_1(x) = f(a) e^{ix\psi(a)} \int e^{ix\psi(a)} \int e^{ix\psi(a)} \int e^{ix\psi(a)} ds$$

Define
$$\frac{1}{15} \frac{15^{1} \times \psi^{(p)}(a)}{(a)} = \frac{1}{15} \frac{15^{1} \times \psi^{(p)}(a)}{(a)} = \frac{1}{15}$$

$$\Rightarrow I_{1}(x) = f(a) e^{i(x\psi(a) \pm \pi/2p)} e^{-u} e^{\pm i\pi/2p} \left(\frac{p!}{x!\psi'''(a)!}\right)^{1/p} \frac{1}{p} u^{\frac{1}{p}-1} du$$

$$= f(a) e^{i(x\psi(a) \pm \pi/2p)} \left(\frac{p!}{x!\psi'''(a)!}\right)^{1/p} \frac{1}{p} \int_{0}^{\infty} e^{-u} u^{\frac{1}{p}-1} du$$

$$T_{i}(x) \sim f(a) e^{i(x\psi(a)\pm\pi/2p)} \left(\frac{p!}{x!\psi'''(a)!}\right)^{ip} \frac{r'(yp)}{p}$$
 as $x \rightarrow \infty$

Example

$$T(x) = \int_{0}^{\pi/2} e^{ix\cos t} dt$$

$$f(t) = 1 \qquad \psi(t) = \cos t$$

$$0 = \psi'(t) = -\sin t \qquad \Rightarrow \quad t = 0 \qquad \text{stationary}$$

$$\psi''(t) = -\cos t \qquad \Rightarrow \quad \psi''(0) = -1 < 0$$

$$= \int T(x) \sim e^{i(x-\pi/4)} \left(\frac{2}{x}\right)^{1/2} \frac{\pi(1/2)}{2}$$

$$= e^{i(x-\pi/4)} \sqrt{\pi} \qquad \text{as} \qquad x \longrightarrow \infty$$

Example (Bessel Functions)

$$\overline{J}_{n}(x) = \prod_{t=0}^{\infty} \int_{0}^{\infty} \cos(x \sin(t) - nt) dt$$

We want the leading behaviour of $J_n(n)$ as $n\to\infty$ = $\int J_n(n) = Re\left(\frac{1}{\pi}\int_{-\pi}^{\pi} e^{in(\sin t \cdot t)} dt\right)$

$$\Psi(\xi) = SIN\xi - \xi$$

$$o = \Psi'(t) = cost - 1$$

$$\psi''(\xi) = -S/N \xi$$

$$J_{n}(n) \sim \frac{1}{\pi} \operatorname{Pe}\left(e^{i(n(0)-\pi/6)}\left(\frac{3!}{N}\right)^{1/3} \frac{\pi(1/3)}{3}\right)$$

$$=\frac{1}{11}\frac{\sqrt{3}}{2}\left(\frac{6}{N}\right)^{1/3}\frac{\Gamma(1/3)}{3}=\frac{\Gamma(1/3)}{11}\frac{3^{1/2}\cdot 3^{1/3}\cdot 3^{1}}{2^{1}2^{-1/3}}\frac{1}{N^{1/3}}$$

$$=\frac{11(1/3)}{11}\frac{1}{3^{1/6}2^{2/3}}\frac{1}{11^{1/3}}$$

descend.