

$$ds^2 = a^2(\tau) \left[-d\tau^2 + \frac{dr^2}{1-Kr^2} + r^2 d\Omega^2 \right]$$

As in newtonian, define Hubble parameter

$$H := \frac{\dot{a}}{a}$$

$$v \sim H d$$

Dynamics: $H \& I \longrightarrow$ RW metric.

$G_{\mu\nu} = 8\pi T_{\mu\nu}$ & RW metric \longrightarrow Friedmann equations

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu + p g_{\mu\nu} \quad \text{Perfect fluid.}$$

$\rho :=$ Energy density.

$p :=$ Isotropic pressure.

Choose $u^\mu = (1, 0, 0, 0)$ comoving in the cosmological rest frame, then

$$T_{\mu\nu} = \begin{pmatrix} \rho & 0 \\ 0 & p g_{ij} \end{pmatrix}$$

Friedmann equation $:= F$

$$H^2 = \left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} \rho - \frac{K}{a^2} \quad \longrightarrow \text{Only first-time derivatives.}$$

\longrightarrow Integral of motion.

(associated to energy through T_{00})

Evolution equation $:= E$

$$\frac{\ddot{a}}{a} + \frac{1}{2} \left(\frac{\dot{a}}{a} \right)^2 = -4\pi p - \frac{K}{2a^2} \quad \longrightarrow \text{2nd-time derivative equation of motion.}$$

Note:

- I. Friedmann equation is a constraint, we are not allowed to specify \dot{a} since it is dependent on ρ and K .
- II. Friedmann equation relates the rate of increase of the scale of factor to the rate of increase of all matter in the universe.

$$\text{III. } F \oplus E \longrightarrow \frac{\ddot{a}}{a} + \frac{1}{2} \left(\frac{8\pi}{3} \rho - \frac{K}{a^2} \right) = -4\pi \rho - \frac{K}{2a^2}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi}{3} (\rho + 3p)$$

IV. If we include a cosmological constant

F:

$$H^2 = \frac{\Lambda}{3} - \frac{K}{a^2} + \frac{8}{3} \pi \rho.$$

E:

$$\frac{\ddot{a}}{a} + \frac{1}{2} \left(\frac{\dot{a}}{a} \right)^2 = -4\pi \rho + \frac{\Lambda}{2} - \frac{K}{2a^2}$$

Differentiating Friedmann equation:

$$\begin{aligned} \left[2 \left(\frac{\dot{a}}{a} \right) \left(\frac{\ddot{a}a - \dot{a}^2}{a^2} \right) = \frac{8\pi}{3} \dot{\rho} + \frac{2K}{a^3} \dot{a} \right] \frac{3}{8\pi} \\ = \frac{3}{4\pi} \left(\frac{\dot{a}}{a} \right) \left(\frac{\ddot{a}a - \dot{a}^2}{a^2} \right) = \dot{\rho} + \frac{3}{4\pi a^2} \left(\frac{\dot{a}}{a} \right) \quad (\text{I}) \end{aligned}$$

Multiply E by $-\frac{3}{4\pi} \left(\frac{\dot{a}}{a} \right)$

$$-\frac{3}{4\pi} \left(\frac{\dot{a}}{a} \right) \left(\frac{\ddot{a}}{a} \right) - \frac{3}{8\pi} \left(\frac{\dot{a}}{a} \right) \left(\frac{\dot{a}}{a} \right)^2 = 3\rho \left(\frac{\dot{a}}{a} \right) + \frac{3K}{8\pi a^2} \left(\frac{\dot{a}}{a} \right) \quad (\text{II})$$

(I) + (II)

$$-\frac{9}{8\pi} \left(\frac{\dot{a}}{a} \right) \left(\frac{\ddot{a}}{a^2} \right) - \frac{9}{8\pi} \frac{K}{a^2} \left(\frac{\dot{a}}{a} \right) = \dot{\rho} + 3\rho \left(\frac{\dot{a}}{a} \right)$$

$$-\frac{9}{8\pi} \left(\left(\frac{\dot{a}}{a} \right)^2 + \frac{K}{a^2} \right) H = \dot{\rho} + 3\rho H$$

$$-\frac{9}{8\pi} \left(\frac{8}{3} \pi \rho \right) H = \dot{\rho} + 3\rho H.$$

$$\dot{\rho} + 3\rho H + 3\rho H = 0$$

$$\dot{\rho} + 3H(\rho + p) = 0$$

consistency condition.

fluid equation (see liddle's modern cosmology).

$$\dot{\rho} a^3 + 3H \rho a^3 (\rho + p) = 0$$

$$\begin{aligned} \frac{d}{dt}(\rho a^3) &= \dot{\rho} a^3 + \rho 3a^2 \dot{a} \\ &= \dot{\rho} a^3 + 3 \frac{\dot{a}}{a} a^3 \rho \\ &= \dot{\rho} a^3 + 3H \rho a^3 \end{aligned}$$

$$\boxed{\frac{d}{dt}(\rho a^3) + p \frac{d}{dt} a^3 = 0}$$

As before,

$$V \propto a^3, \quad \epsilon = \rho V$$

$$d\epsilon + p dV = 0$$

Thermodynamics
1st law

Pressure works on
the expansion.

Conservation of energy

$$\nabla_\mu T^{\mu\nu} = 0$$

Define the critical energy density

$$\rho_c := \frac{3}{8\pi} H^2$$

(Spatial sections
become flat $K=0$).

And the density parameter

$$\Omega_{\text{total}} := \frac{\rho}{\rho_c}$$

relating total energy density
to its local geometry.

$$\Omega_{\text{total}} \begin{cases} > 1 & ; \quad K = 1 \\ = 1 & ; \quad K = 0 \\ < 1 & ; \quad K = -1 \end{cases}$$

Friedmann

$$\rho = \rho_c + \frac{3}{8\pi} \frac{K}{a^2}$$

$$\Omega_K := \frac{\rho_K}{\rho_c} \quad ; \quad \rho_K := \frac{3K}{8\pi a^2}$$

$$\Omega_{\text{total}} = 1 + \Omega_K$$

If we consider a cosmological constant

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi T_{\mu\nu}$$

$$G_{\mu\nu} = 8\pi (T_{\mu\nu} + T'_{\mu\nu})$$

$$T'_{\mu\nu} := -\frac{\Lambda}{8\pi} g_{\mu\nu}$$

then,

$$\rho_\Lambda = \frac{\Lambda}{8\pi} \quad \text{constant}$$

$$p_\Lambda = -\frac{\Lambda}{8\pi} = -\rho_\Lambda$$

$$\Omega_\Lambda = \frac{\rho_\Lambda}{\rho_c}$$

So far, two equations (E & F) and three unknowns (a, ρ, p).

Equations of state

$$p = p(\rho) \quad (\text{in general } p = p(\rho, T))$$

$$p = w\rho; \quad w \text{ parameter (constant)}$$

$$w=0 \longrightarrow \text{Dust (matter dominated)}$$

$$w=\frac{1}{3} \longrightarrow \text{Gas of radiation.}$$

$$w=-1 \longrightarrow \text{Cosmological constant.}$$

$$\begin{aligned} 0 &= \dot{\rho} + 3H(\rho + p) = \dot{\rho} + 3H(\rho + w\rho) \\ &= \dot{\rho} + 3H\rho(1+w) \end{aligned}$$

$$\frac{d\rho}{\rho} = -3(1+w)\frac{da}{a}$$

$$\ln(\rho) = -3(1+w)\ln(aC)$$

$$\rho(a) = \frac{c'}{a(t)^{3(1+w)}}$$

$$\text{If } w=0; \text{ then } \rho a^3 = c' \quad (\text{Energy}).$$