Fermi Gas

$$ln(\Xi(T,V,M))=\sum_{i=1}^{n}ln_{i}^{2}1+exp(-\beta(E_{i}-M))_{i}^{2}$$

$$\langle n_j \rangle = 1$$
 $\exp(\beta(E_j - M)) + 1$
Fermi-Dirac
Distribution

$$\Phi(T,V,M) = -\frac{1}{\beta}\ln\left(\bigoplus(T,V,M)\right) = -pV$$

then

$$p(T, M) = K_BT \lim_{V \to \infty} \frac{1}{V} \ln(\Xi(T, V, M))$$

$$0 = \sum_{j} E_{j} \langle N_{j} \rangle = \sum_{j} \frac{E_{j}}{\langle x \rho(\beta(E_{j} - M)) + 1}$$

$$N = \sum_{j} \langle N_{j} \rangle = \sum_{j} \frac{1}{(N_{j} + N_{j}) + 1}$$

$$E_j = E_{j,n} = \frac{h^2 k^2}{2m}$$
 Free fermions

In the thermodynamic limit

$$\ln(\Xi) = \gamma \frac{\sqrt{2\pi}}{(2\pi)^3} \int d^3\vec{k} \ln \left(1 + \exp(-\beta \left(\frac{\hbar^2 \vec{k}^2}{2m} - \mu \right) \right) \right)$$

$$\langle n_j \rangle = \left(exp \left(\frac{Bt^2 K^2}{2m} - BM \right) + 1 \right)^{-1}$$

$$N = \frac{1}{(2\pi)^3} \int d^3\vec{K} \left(\exp\left(\frac{Bt^2K^2}{2m} - BM\right) + 1\right)^{-1}$$

$$U = \gamma \frac{\sqrt{(2\pi)^3}}{\sqrt{2m}} \int d^3 \vec{K} \frac{\hbar^2 K^2}{2m} \left(\exp \left(\frac{\beta \hbar^2 K^2}{2m} - \beta M \right) + 1 \right)^{-1}$$

$$\ln(\Xi) = YV \int_{0}^{\infty} D(E) \ln \left\{ 1 + \exp(-\beta(E - M)) \right\} dE$$

$$U = YV \int_{0}^{\infty} D(E) f(E) dE$$

$$f(E) := \frac{1}{\exp(\beta(E_{j} - M)) + 1}$$

$$\mathcal{D}(\mathcal{E}) = \frac{1}{4\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \mathcal{E}^{1/2} = C \mathcal{E}^{1/2}$$

Integrating by parts

$$\ln(\Xi) = YVC\int_{0}^{\infty} E^{1/2} \ln \left(1 + \exp(-\beta(E-\mu))\right) dE$$

$$= \frac{2}{3} E^{3/2} \ln \left(1 + \exp(-\beta(E - M)) \right) \left(\int_{0}^{\infty} - \int_{0}^{2/2} \frac{-\beta \exp(-\beta(E - M))}{1 - \exp(-\beta(E - M))} dE \right)$$

therefore,

$$\ln(\Xi) = \frac{2}{3} \beta \gamma V C \int_{0}^{\infty} E^{3/2} f(E) dE$$

$$= \frac{2}{3} \frac{\gamma V}{k_{BT}} \int_{0}^{\infty} E D(E) f(E) dE$$

$$= \frac{2}{3} \frac{V}{K_{BT}} \langle E \rangle = \frac{2}{3} \frac{V}{K_{ET}} u = \frac{2}{3} \frac{U}{K_{BT}}$$

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