

# Rigid Body Dynamics (I)

COMP259 March 28, 2006

Nico Galoppo von Borries

# From Particles to Rigid Bodies

- Particles
  - No rotations
  - Linear velocity  $v$  only
- Rigid bodies
  - Body rotations
  - Linear velocity  $v$
  - Angular velocity  $\omega$

# Outline

- Rigid Body Preliminaries
  - Coordinate system, velocity, acceleration, and inertia
- State and Evolution
- Quaternions
- Collision Detection and Contact Determination
- Colliding Contact Response

# Coordinate Systems

- Body Space (Local Coordinate System)
  - bodies are specified relative to this system
  - center of mass is the origin (for convenience)
    - We will specify body-related physical properties (inertia, ...) in this frame

# Coordinate Systems

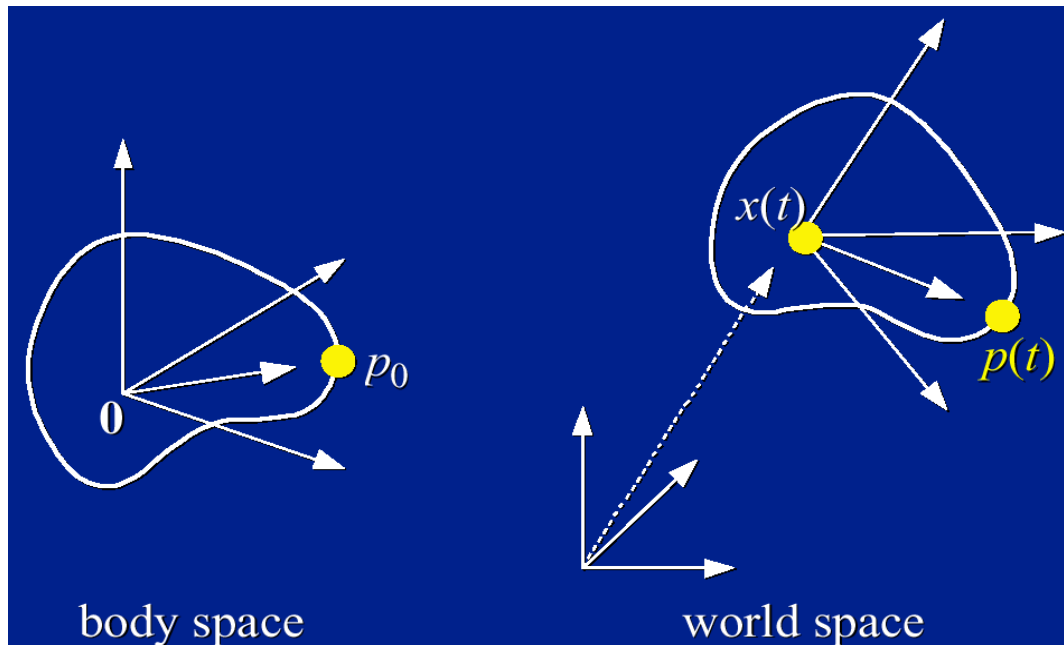
- World Space

- bodies are transformed to this common system

$$p(t) = R(t) p_0 + x(t)$$

- $x(t)$  represents the position of the body center
- $R(t)$  represents the orientation
  - Alternatively, use *quaternion* representation

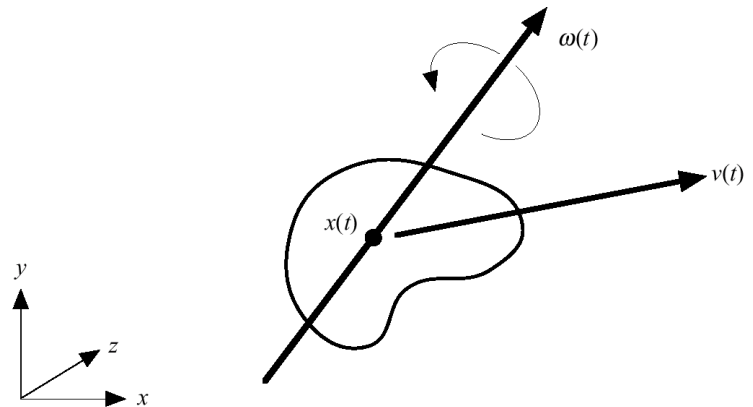
# Coordinate Systems



**Meaning of  $R(t)$ : columns represent the coordinates of the body space base vectors  $(1,0,0)$ ,  $(0,1,0)$ ,  $(0,0,1)$  in world space.**

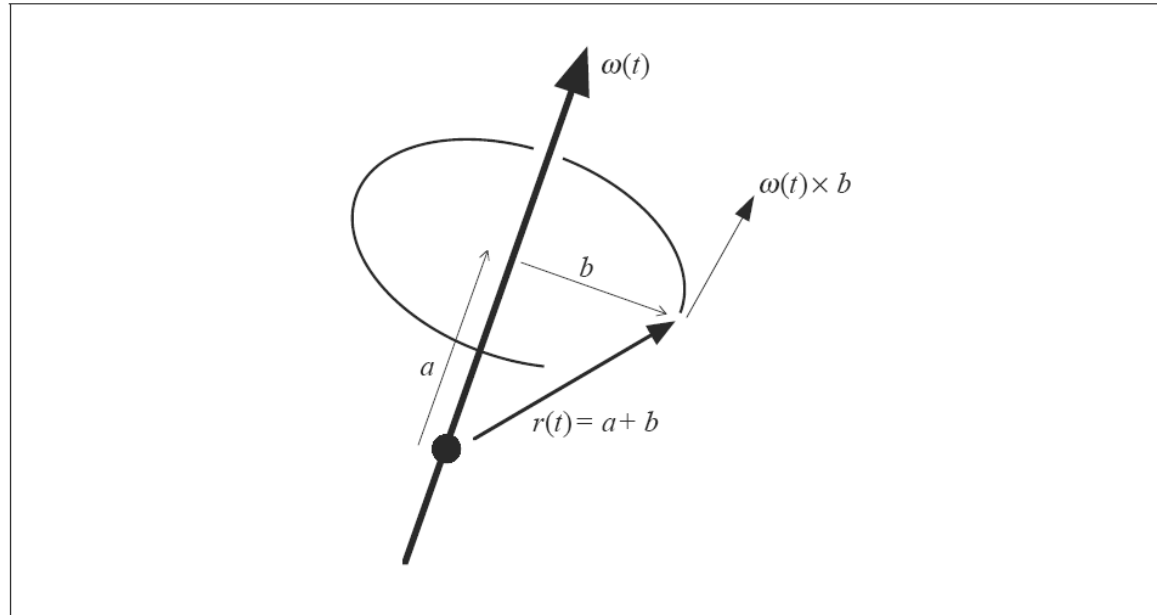
# Kinematics: Velocities

- How do  $x(t)$  and  $R(t)$  change over time?
- Linear velocity  $v(t) = dx(t)/dt$  is the same:
  - Describes the velocity of the center of mass (m/s)
- Angular velocity  $\omega(t)$  is new!
  - Direction is the axis of rotation
  - Magnitude is the angular velocity about the axis (degrees/time)
  - There is a simple relationship between  $R(t)$  and  $\omega(t)$



# Kinematics: Velocities

Then



$$\dot{R} = \left( \omega(t) \times \begin{pmatrix} r_{xx} \\ r_{xy} \\ r_{xz} \end{pmatrix} \quad \omega(t) \times \begin{pmatrix} r_{yx} \\ r_{yy} \\ r_{yz} \end{pmatrix} \quad \omega(t) \times \begin{pmatrix} r_{zx} \\ r_{zy} \\ r_{zz} \end{pmatrix} \right)$$



# Angular Velocities

- $\mathbf{R}(t)$  and  $\boldsymbol{\omega}(t)$  are related by:

$$\frac{d}{dt}\mathbf{R}(t) = \begin{pmatrix} 0 & -\omega_z(t) & \omega_y(t) \\ \omega_z(t) & 0 & -\omega_x(t) \\ -\omega_y(t) & \omega_x(t) & 0 \end{pmatrix} \mathbf{R}(t)$$
$$= \boldsymbol{\omega}(t)^* \mathbf{R}(t)$$

# Dynamics: Accelerations

- How do  $v(t)$  and  $dR(t)/dt$  change over time?
- First we need some more machinery
  - Forces and Torques
  - Momentums
  - Inertia Tensor
- Simplify equations by formulating accelerations terms of *momentum derivatives* instead of velocity derivatives

# Forces and Torques

- External forces  $F_i(t)$  act on particles
  - Total external force  $F = \sum F_i(t)$
- Torques depend on distance from the center of mass:

$$\tau_i(t) = (r_i(t) - x(t)) \times F_i(t)$$

- Total external torque

$$\tau = \sum ((r_i(t) - x(t)) \times F_i(t))$$

- $F(t)$  doesn't convey any information about where the various forces act
- $\tau(t)$  does tell us about the *distribution* of forces

# Linear Momentum

- Linear momentum  $P(t)$  lets us express the effect of total force  $F(t)$  on body (simple, because of conservation of energy):  

$$F(t) = dP(t)/dt$$
- Linear momentum is the product of mass and linear velocity
  - $P(t) = \sum m_i dr_i(t)/dt$   
 $= \sum m_i v(t) + \omega(t) \times \sum m_i (r_i(t) - x(t))$
- But, we work in body space:
  - $P(t) = \sum m_i v(t) = Mv(t)$  (linear relationship)
  - Just as if body were a particle with mass  $M$  and velocity  $v(t)$
  - Derive  $v(t)$  to express acceleration:  

$$dv(t)/dt = M^{-1} dP(t)/dt$$
- Use  $P(t)$  instead of  $v(t)$  in state vectors

# Angular momentum

- Same thing, angular momentum  $L(t)$  allows us to express the effect of total torque  $\tau(t)$  on the body:

$$\dot{L}(t) = \tau(t)$$

- Similarly, there is a linear relationship between momentum and velocity:

$$L(t) = I(t) \omega(t)$$

- $I(t)$  is inertia tensor, plays the role of mass
- Use  $L(t)$  instead of  $\omega(t)$  in state vectors

# Inertia Tensor

- 3x3 matrix describing how the shape and mass distribution of the body affects the relationship between the angular velocity and the angular momentum  $I(t)$
- Analogous to mass – rotational mass
- We actually want the inverse  $I^{-1}(t)$  for computing  $\omega(t)=I^{-1}(t)L(t)$

# Inertia Tensor

$$\begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{yx} & I_{yy} & -I_{yz} \\ -I_{zx} & -I_{zy} & I_{zz} \end{bmatrix}$$

Bunch of volume integrals:

$$I_{xx} = \int_{B_i} (y^2 + z^2) dV \quad I_{yy} = \int_{B_i} (z^2 + x^2) dV \quad I_{zz} = \int_{B_i} (x^2 + y^2) dV$$

$$I_{xy} = I_{yx} = \int_{B_i} xy dV \quad I_{xz} = I_{zx} = \int_{B_i} zx dV \quad I_{yz} = I_{zy} = \int_{B_i} yz dV$$

# Inertia Tensor

- Compute  $I$  in body space  $I_{\text{body}}$  and then transform to world space as required
  - $I(t)$  varies in world space, but  $I_{\text{body}}$  is constant in body space for the entire simulation
- $I(t) = R(t) I_{\text{body}} R^{-1}(t) = R(t) I_{\text{body}} R^T(t)$
- $I^{-1}(t) = R(t) I_{\text{body}}^{-1} R^{-1}(t) = R(t) I_{\text{body}}^{-1} R^T(t)$
- Intuitively: transform  $\omega(t)$  to body space, apply inertia tensor in body space, and transform back to world space

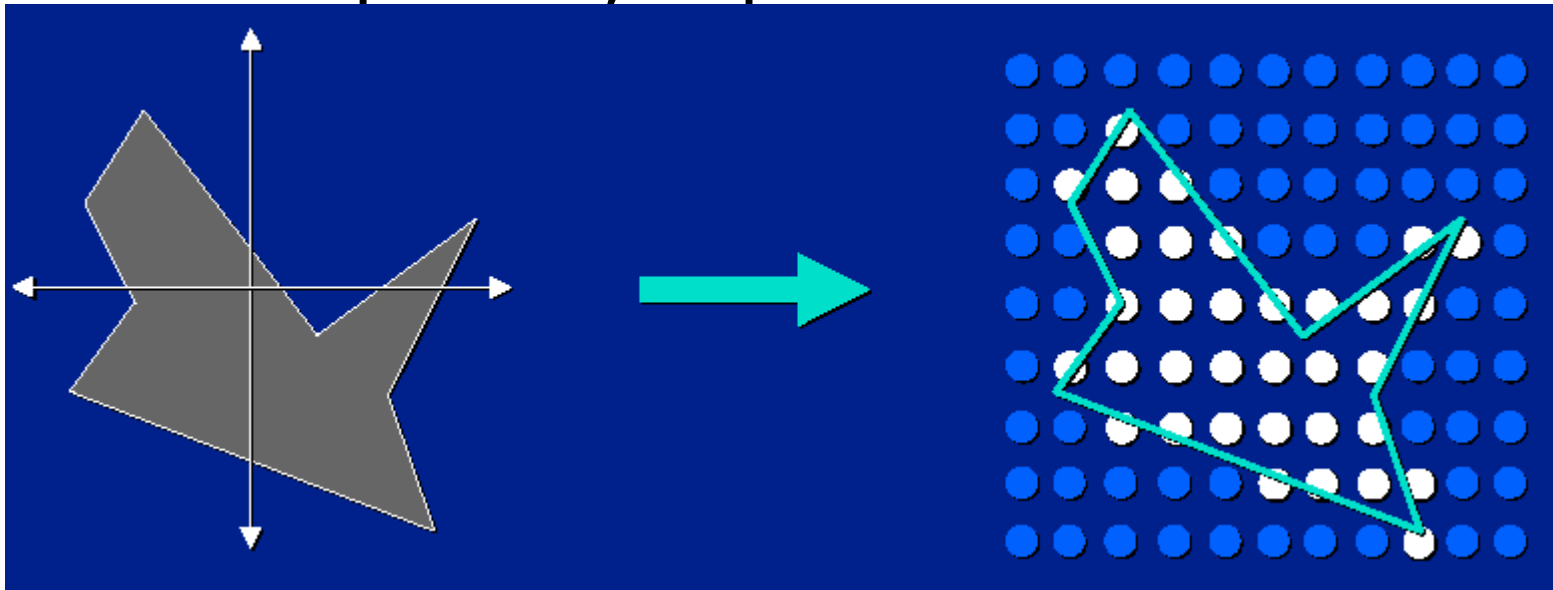


# Computing $I_{\text{body}}^{-1}$

- There exists an orientation in body space which causes  $I_{xy}$ ,  $I_{xz}$ ,  $I_{yz}$  to all vanish
  - Diagonalize tensor matrix, define the eigenvectors to be the local body axes
  - Increases efficiency and trivial inverse
- Point sampling within the bounding box
- Projection and evaluation of Greene's thm.
  - Code implementing this method exists
  - Refer to Mirtich's paper at <http://www.acm.org/jgt/papers/Mirtich96>

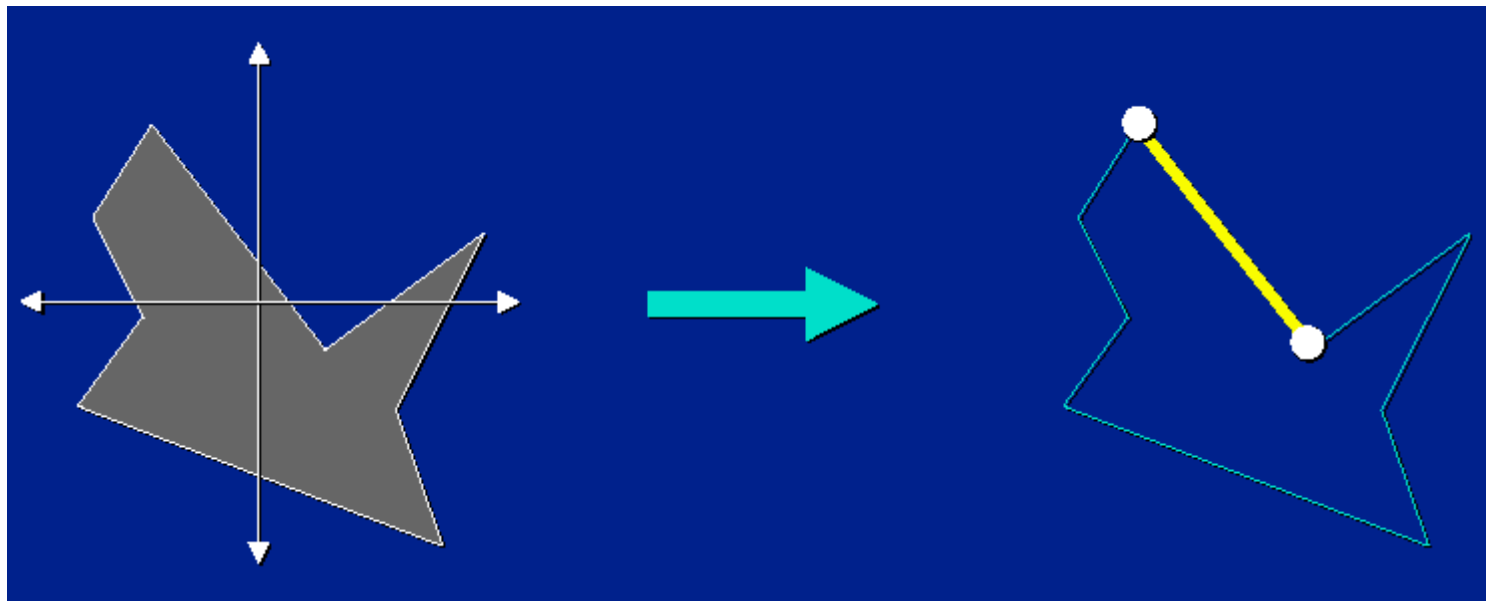
# Approximation w/ Point Sampling

- Pros: Simple, fairly accurate, no B-rep needed.
- Cons: Expensive, requires volume test.



# Use of Green's Theorem

- Pros: Simple, exact, no volumes needed.
- Cons: Requires boundary representation.



# Outline

- Rigid Body Preliminaries
- State and Evolution
  - Variables and derivatives
- Quaternions
- Collision Detection and Contact Determination
- Colliding Contact Response

# New State Space

$$\mathbf{X}(t) = \begin{pmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{pmatrix} \begin{array}{l} \left. \vphantom{\begin{pmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{pmatrix}} \right\} \rightarrow \text{Spatial information} \\ \left. \vphantom{\begin{pmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{pmatrix}} \right\} \rightarrow \text{Velocity information} \end{array}$$

$v(t)$  replaced by linear momentum  $P(t)$

$\omega(t)$  replaced by angular momentum  $L(t)$

Size of the vector:  $(3+9+3+3)N = 18N$

# Taking the Derivative

$$\frac{d}{dt}\mathbf{X}(t) = \frac{d}{dt} \begin{pmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{pmatrix} = \begin{pmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ \tau(t) \end{pmatrix}$$

Conservation of momentum ( $P(t)$ ,  $L(t)$ ) lets us express the accelerations in terms of forces and torques.

# Simulate: next state computation

- From  $X(t)$  certain quantities are computed

$$\begin{cases} I^{-1}(t) = R(t) I_{body}^{-1} R^T(t) \\ v(t) = P(t) / M \\ \omega(t) = I^{-1}(t) L(t) \end{cases}$$

- We cannot compute the state of a body at all times but must be content with a finite number of discrete time points, ***assuming that the acceleration is continuous***
- Use your favorite ODE solver to solve for the new state  $X(t)$ , given previous state  $X(t-\Delta t)$  and applied forces  $F(t)$  and  $\tau(t)$

$$X(t) \leftarrow \text{Solver}::\text{Step}(X(t-\Delta t), F(t), \tau(t))$$

# Simple simulation algorithm

```
X  $\tilde{A}$  InitializeState()  
For t=t0 to tfinal with step  $\Delta t$   
    ClearForces(F(t),  $\tau(t)$ )  
    AddExternalForces(F(t),  $\tau(t)$ )  
    Xnew  $\tilde{A}$  Solver::Step(X, F(t),  $\tau(t)$ )  
    X  $\tilde{A}$  Xnew  
    t  $\tilde{A}$  t +  $\Delta t$   
End for
```



# Outline

- Rigid Body Preliminaries
- State and Evolution
- Quaternions
  - Merits, drift, and re-normalization
- Collision Detection and Contact Determination
- Colliding Contact Response

# Unit Quaternion Merits

- Problem with rotation matrices: numerical drift
- $R(t) = dR(t)/dt * \Delta t R(t_{-1})R(t_{-2})R(t_{-3}) \dots$
- A rotation in 3-space involves 3 DOF
- Unit quaternions can do it with 4
- Rotation matrices  $R(t)$  describe a rotation using 9 parameters
- Drift is easier to fix with quaternions
  - renormalize

# Unit Quaternion Definition

- $q = [s, v]$  --  $s$  is a scalar,  $v$  is vector
- A rotation of  $\theta$  about a unit axis  $u$  can be represented by the unit quaternion:

$$[\cos(\theta/2), \sin(\theta/2) * u]$$

- $\| [s, v] \| = 1$  -- the length is taken to be the Euclidean distance treating  $[s, v]$  as a 4-tuple or a vector in 4-space

# Unit Quaternion Operations

- Multiplication is given by:

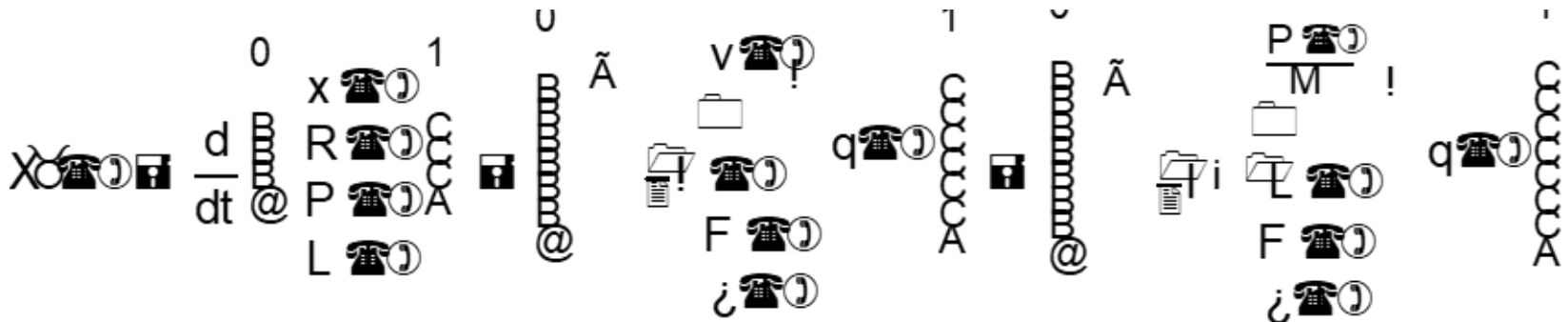
$$[s_1, v_1][s_2, v_2] = [s_1s_2 - v_1 \cdot v_2, s_1v_2 + s_2v_1 + v_1 \times v_2]$$

- $dq(t)/dt = \frac{1}{2} \omega(t) q(t) = [0, \frac{1}{2} \omega(t)] q(t)$

- $$R = \begin{pmatrix} 1 - 2v_y^2 - 2v_z^2 & 2v_xv_y - 2sv_z & 2v_xv_z + 2sv_y \\ 2v_xv_y + 2sv_z & 1 - 2v_x^2 - 2v_z^2 & 2v_yv_z - 2sv_x \\ 2v_xv_z - 2sv_y & 2v_yv_z + 2sv_x & 1 - 2v_x^2 - 2v_y^2 \end{pmatrix}$$

# Unit Quaternion Usage

- To use quaternions instead of rotation matrices, just substitute them into the state as the orientation (save 5 variables)



where

$$\mathbf{R} = \text{QuatToMatrix}(q(t))$$

$$\mathbf{I}^{-1}(t) = \mathbf{R} \mathbf{I}_{\text{body}}^{-1} \mathbf{R}^T$$

# Outline

- Rigid Body Preliminaries
- State and Evolution
- Quaternions
- Collision Detection and Contact Determination
  - Contact classification
  - Intersection testing, bisection, and nearest features
- Colliding Contact Response

# What happens when bodies collide?

- Colliding
  - Bodies *bounce* off each other
  - Elasticity governs 'bounciness'
  - Motion of bodies changes ***discontinuously*** within a discrete time step
  - 'Before' and 'After' states need to be computed
- In contact
  - Resting
  - Sliding
  - Friction

# Detecting collisions and response

- Several choices
  - Collision detection: which algorithm?
  - Response: Backtrack or allow penetration?
- Two primitives to find out if response is necessary:
  - Distance(A,B): cheap, no contact information ! fast intersection query
  - Contact(A,B): expensive, with contact information



# Distance(A,B)

- Returns a value which is the minimum distance between two bodies
- Approximate may be ok
- Negative if the bodies intersect
- Convex polyhedra
  - Lin-Canny and GJK -- 2 classes of algorithms
- Non-convex polyhedra
  - much more useful but hard to get distance fast
  - PQP/RAPID/SWIFT++
- Remark: most of these algorithms give inaccurate information if bodies intersect, except for DEEP

# Contacts(A,B)

- Returns the set of features that are nearest for disjoint bodies or intersecting for penetrating bodies
- Convex polyhedra
  - LC & GJK give the nearest features as a bi-product of their computation – only a single pair. Others that are equally distant may not be returned.
- Non-convex polyhedra
  - much more useful but much harder problem especially contact determination for disjoint bodies
  - Convex decomposition: SWIFT++

# Prereq: Fast intersection test

- First, we want to make sure that bodies will intersect at next discrete time instant
- If not:
  - $X_{\text{new}}$  is a valid, non-penetrating state, proceed to next time step
- If intersection:
  - Classify contact
  - Compute response
  - Recompute new state

# Bodies intersect ! classify contacts

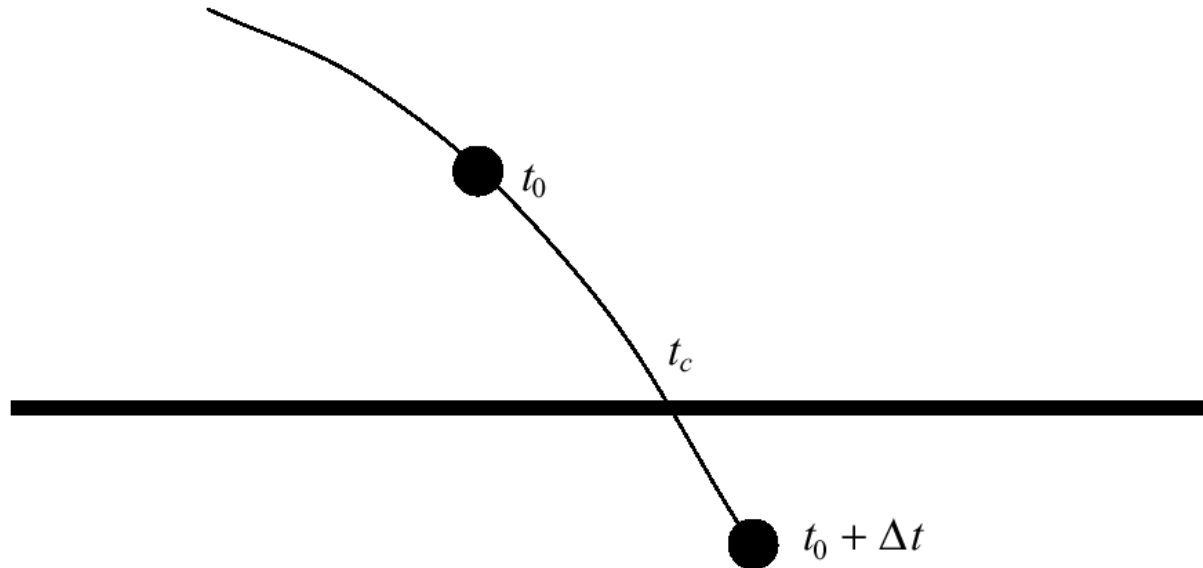
- Colliding contact (Today)
  - $V_{\text{rel}} < -\varepsilon$
  - Instantaneous change in velocity
  - Discontinuity: requires restart of the equation solver
- Resting contact (Thursday)
  - $-\varepsilon < V_{\text{rel}} < \varepsilon$
  - Gradual contact forces avoid interpenetration
  - No discontinuities
- Bodies separating
  - $V_{\text{rel}} > \varepsilon$
  - No response required

# Colliding contacts

- At time  $t_i$ , body A and B intersect and
$$V_{\text{rel}} < -\varepsilon$$
- Discontinuity in velocity: need to stop numerical solver
- Find time of collision  $t_c$
- Compute new velocities  $v^+(t_c) \rightarrow X^+(t)$
- Restart ODE solver at time  $t_c$  with new state  $X^+(t)$

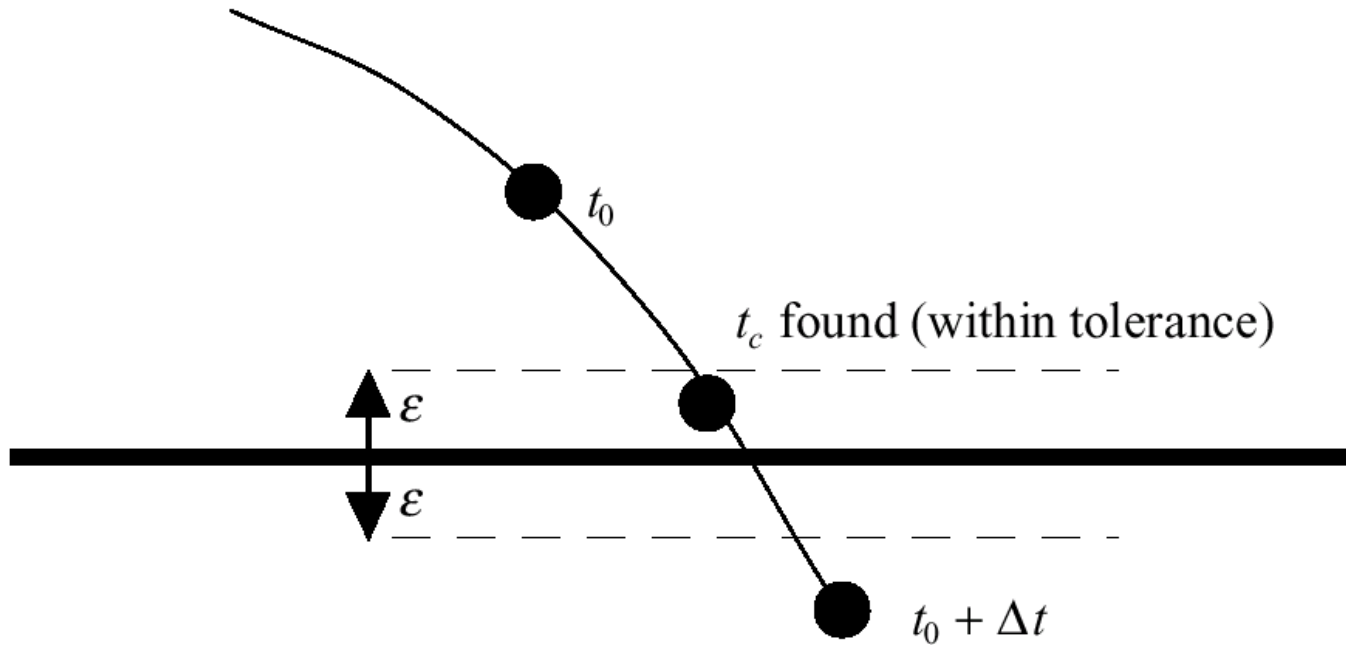
# Time of collision

- We wish to compute when two bodies are “close enough” and then apply contact forces
- Let’s recall a particle colliding with a plane



# Time of collision

- We wish to compute  $t_c$  to some tolerance



# Time of collision

1. A common method is to use **bisection search** until the distance is positive but less than the tolerance
2. Use **continuous collision detection**  
(cf. Dave Knott's lecture)
3.  $t_c$  not always needed  
! **penalty-based methods**



# findCollisionTime(X,t, $\Delta t$ )

```
0 for each pair of bodies (A,B) do
1   Compute_New_Body_States(Scopy, t, H);
2   hs(A,B) = H;    // H is the target timestep
3   if Distance(A,B) < 0 then
4     try_h = H/2;  try_t = t + try_h;
5     while TRUE do
6       Compute_New_Body_States(Scopy, t, try_t - t);
7       if Distance(A,B) < 0 then
8         try_h /= 2;  try_t -= try_h;
9       else if Distance(A,B) <  $\varepsilon$  then
10        break;
11      else
12        try_h /= 2;  try_t += try_h;
13    hs(A,B) = try_t - t;
14  h = min( hs );
```

# Outline

- Rigid Body Preliminaries
- State and Evolution
- Quaternions
- Collision Detection and Contact Determination
- Colliding Contact Response
  - Normal vector, restitution, and force application

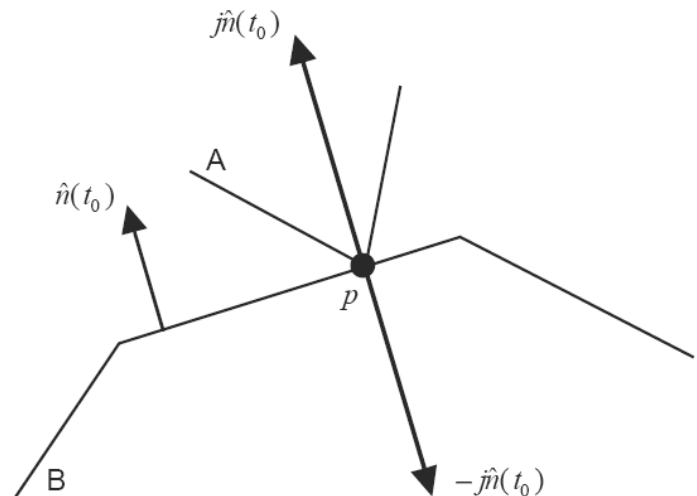
# What happens upon collision

- Impulses provide instantaneous changes to velocity, unlike forces

$$\Delta(P) = J$$

- We apply impulses to the colliding objects, at the point of collision
- For frictionless bodies, the direction will be the same as the normal direction:

$$J = j^T n$$



# Colliding Contact Response

- Assumptions:
  - Convex bodies
  - Non-penetrating
  - Non-degenerate configuration
    - edge-edge or vertex-face
    - appropriate set of rules can handle the others
- Need a contact unit normal vector
  - Face-vertex case: use the normal of the face
  - Edge-edge case: use the cross-product of the direction vectors of the two edges

# Colliding Contact Response

- Point velocities at the nearest points:

$$\dot{p}_a(t_0) = v_a(t_0) + \omega_a(t_0) \times (p_a(t_0) - x_a(t_0))$$

$$\dot{p}_b(t_0) = v_b(t_0) + \omega_b(t_0) \times (p_b(t_0) - x_b(t_0)).$$

- Relative contact normal velocity:

$$v_{rel} = \hat{n}(t_0) \cdot (\dot{p}_a(t_0) - \dot{p}_b(t_0))$$

# Colliding Contact Response

- We will use the empirical law of frictionless collisions:  $v_{rel}^+ = -\epsilon v_{rel}^-$ 
  - Coefficient of restitution  $\epsilon \in [0,1]$ 
    - $\epsilon = 0$  -- bodies stick together
    - $\epsilon = 1$  – loss-less rebound
- After some manipulation of equations...

$$j = \frac{-(1 + \epsilon)v_{rel}^-}{\frac{1}{M_a} + \frac{1}{M_b} + \hat{n}(t_0) \cdot \left( I_a^{-1}(t_0) (r_a \times \hat{n}(t_0)) \right) \times r_a + \hat{n}(t_0) \cdot \left( I_b^{-1}(t_0) (r_b \times \hat{n}(t_0)) \right) \times r_b}$$

# Apply\_BB\_Forces()

- For colliding contact, the computation can be local

```
0  for each pair of bodies (A,B) do
1      if Distance(A,B) <  $\varepsilon$  then
2  Cs = Contacts(A,B);
3  Apply_Impulses(A,B,Cs);
```

# Apply\_Impulses(A,B,Cs)

- The impulse is an instantaneous force – it changes the velocities of the bodies instantaneously:  $\Delta v = J / M$

```
0 for each contact in Cs do
1   Compute n
2   Compute j
3    $J = j^T n$ 
3    $P(A) += J$ 
4    $L(A) += (p - x(t)) \times J$ 
5    $P(B) -= J$ 
6    $L(B) -= (p - x(t)) \times J$ 
```



# Simulation algorithm with Collisions

```
X  $\tilde{\leftarrow}$  InitializeState()
```

```
For  $t=t_0$  to  $t_{\text{final}}$  with step  $\Delta t$ 
```

```
    ClearForces( $F(t)$ ,  $\tau(t)$ )
```

```
    AddExternalForces( $F(t)$ ,  $\tau(t)$ )
```

```
     $X_{\text{new}} \tilde{\leftarrow}$  Solver::Step( $X$ ,  $F(t)$ ,  $\tau(t)$ ,  $t$ ,  $\Delta t$ )
```

```
     $t \tilde{\leftarrow}$  findCollisionTime()
```

```
     $X_{\text{new}} \tilde{\leftarrow}$  Solver::Step( $X$ ,  $F(t)$ ,  $\tau(t)$ ,  $t$ ,  $\Delta t$ )
```

```
     $C \tilde{\leftarrow}$  Contacts( $X_{\text{new}}$ )
```

```
    while (! $C$ .isColliding())
```

```
        applyImpulses( $X_{\text{new}}$ )
```

```
    end if
```

```
     $X \tilde{\leftarrow} X_{\text{new}}$ 
```

```
     $t \tilde{\leftarrow} t + \Delta t$ 
```

```
End for
```

# Penalty Methods

- If we don't look for time of collision  $t_c$  then we have a simulation based on penalty methods: the objects are allowed to intersect.
- **Global** or **local** response
  - **Global**: The penetration depth is used to compute a spring constant which forces them apart (dynamic springs)
  - **Local**: Impulse-based techniques

# Global penalty based response

## Global contact force computation

0 for each pair of bodies (A,B) do

1 if  $\text{Distance}(A,B) < \varepsilon$  then

2     Flag\_Pair(A,B);

**3 Solve For Forces(flagged pairs);**

4 Apply\_Forces(flagged pairs);

# Local penalty based response

## Local contact force computation

0 for each pair of bodies (A,B) do

1 if  $\text{Distance}(A,B) < \varepsilon$  then

2      $C_s = \text{Contacts}(A,B);$

**3      $\text{Apply\_Impulses}(A,B,C_s);$**

# References

- D. Baraff and A. Witkin, "Physically Based Modeling: Principles and Practice," Course Notes, SIGGRAPH 2001.
- B. Mirtich, "Fast and Accurate Computation of Polyhedral Mass Properties," Journal of Graphics Tools, volume 1, number 2, 1996.
- D. Baraff, "Dynamic Simulation of Non-Penetrating Rigid Bodies", Ph.D. thesis, Cornell University, 1992.
- B. Mirtich and J. Canny, "Impulse-based Simulation of Rigid Bodies," in Proceedings of 1995 Symposium on Interactive 3D Graphics, April 1995.
- B. Mirtich, "Impulse-based Dynamic Simulation of Rigid Body Systems," Ph.D. thesis, University of California, Berkeley, December, 1996.
- B. Mirtich, "Hybrid Simulation: Combining Constraints and Impulses," in Proceedings of First Workshop on Simulation and Interaction in Virtual Environments, July 1995.
- COMP259 Rigid Body Simulation Slides, Chris Vanderknyff 2004