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MP232: Applied Mathematics

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1 Prelim: The Exponential Function and Hyperbolic Functions

1.1 Exponential Function

Derivative

$$\frac{d}{dt}(e^{at}) = a e^{at}$$

$$\int e^{at} dt = \frac{1}{a} e^{at} + C$$

1.2 Hyperbolic Functions

Definitions:

$$\sinh(at) = \frac{e^{at} - e^{-at}}{2} \mid \cosh(at) = \frac{e^{at} + e^{-at}}{2} \mid \tanh(at) = \frac{\sinh(at)}{\cosh(at)}.$$

Derivatives

$$\frac{d}{dt}\big(\sinh(at)\big) = a\,\cosh(at), \, \left| \, \frac{d}{dt}\big(\cosh(at)\big) = a\,\sinh(at), \, \left| \, \frac{d}{dt}\big(\tanh(at)\big) = a\,\sinh^2(at).$$

Integrals

$$\int \sinh(at) dt = \frac{1}{a} \cosh(at) + C$$

$$\int \cosh(at) dt = \frac{1}{a} \sinh(at) + C,$$

$$\int \tanh(at) dt = \frac{1}{a} \ln|\cosh(at)| + C.$$

Common Identities

$$\cosh^2 x - \sinh^2 x = 1,$$

$$\sinh(2x) = 2 \sinh x \cosh x,$$

$$\cosh(2x) = \cosh^2 x + \sinh^2 x,$$

$$\tanh(2x) = \frac{2 \tanh x}{1 + \tanh^2 x}.$$

1.3 Partial Fraction Decomposition

Unrepeated Linear Factors: A linear factor is of form (ax + b)

$$\frac{s+1}{s(s-2)(s+3)} = \frac{A}{s} + \frac{B}{s-2} + \frac{C}{s+3}$$

Repeated LinearFactors:

$$\frac{3}{(s+2)^2(s-3)} = \frac{A}{s+2} + \frac{B}{(s+2)^2} + \frac{C}{s-3}$$

Unrepeated Quadratic Factors with complex roots: Where the discriminant $(b^2 - 4ac)$ is negative (complex roots) but the factor is not repeated

$$\frac{3}{(s^2 - s + 1)(s + 2)} = \frac{As + B}{s^2 - s + 1} + \frac{C}{s + 2}$$

Repeated Quadratic Factors with complex roots:

$$\frac{1}{(s^2+1)^2(s-1)} = \frac{As+B}{(s^2+1)^2} + \frac{Cs+D}{s^2+1} + \frac{E}{s-1}$$

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2 Laplace Transforms

2.1 What is a Laplace Transform?

The Laplace Transform, defined for $t \geq 0$, is given by

$$L\{f(t)\}(s) = F(s) = \int_0^\infty e^{-st} dt$$

2.2 Common Laplace Transforms

Example Find the Laplace Transform of f(t) = 1

We have:

$$L\{1\} = \int_0^\infty 1 \cdot e^{-st} dt = \lim_{R \to \infty} \int_0^R e^{-st} dt$$

This integral is equal to:

$$\int_{0}^{R} e^{-st} dt = \frac{e^{-st}}{-s} \bigg|_{t=0}^{t=R} = -\frac{1}{s} [e^{-sR} - 1] = \frac{1 - e^{-sR}}{s}$$

Taking the limit as $R \to \infty$ gives:

$$L\{1\} = \lim_{R \to \infty} \frac{1 - e^{-sR}}{s} = \frac{1}{s}$$

Example Find the Laplace Transform of $f(t) = e^{2t}$

$$\begin{split} L\{e^{2t}\} &= \int_0^\infty e^{2t} e^{-st} \, dt = \int_0^\infty e^{-(s-2)t} \, dt \\ &= \lim_{R \to \infty} \int_0^R e^{-(s-2)t} \, dt \\ &= \lim_{R \to \infty} \left[\frac{e^{-(s-2)t}}{-(s-2)} \right]_{t=0}^{t=R} \\ &= \lim_{R \to \infty} \left(\frac{e^{-(s-2)R} - e^0}{-(s-2)} \right) = \lim_{R \to \infty} \left(\frac{e^{-(s-2)R} - 1}{-(s-2)} \right) \\ &= \frac{1}{s-2} \quad \text{(since } e^{-(s-2)R} \to 0 \text{ as } R \to \infty \text{ provided } s > 2 \text{)} \end{split}$$

Example Find the Laplace Transform of $f(t) = \cosh(at)$

We have:

$$\begin{split} L\{\cosh(at)\} &= L\left\{\frac{e^{at} + e^{-at}}{2}\right\} \quad \text{from the definition of } \cosh(at) \\ &= \frac{1}{2}L\{e^{at}\} + \frac{1}{2}L\{e^{-at}\} \quad \text{by linearity of the Laplace Transform} \\ &= \frac{1}{2}\left(\frac{1}{s-a}\right) + \frac{1}{2}\left(\frac{1}{s+a}\right) \end{split}$$

Hence:

$$L\{\cosh(at)\} = \frac{s}{s^2 - a^2}$$

Noting that $sinh(at) = (e^{at} - e^{-at})/2$, we can find that:

$$L\{\sinh(at)\} = \frac{a}{(s^2 - a^2)}$$

Example Find the Laplace Transform of $\cos(wt)$ and $\sin(wt)$ where w is a constant.

We first compute the Laplace Transform of e^{iwt} using its definition:

$$L\{e^{iwt}\} = \int_0^\infty e^{-st} e^{iwt} dt = \int_0^\infty e^{-(s-iw)t} dt = \frac{1}{s-iw}, \quad \text{for } \Re(s) > 0.$$

To express this in terms of real and imaginary parts, we multiply the numerator and denominator by the complex conjugate of the denominator:

$$\frac{1}{s-iw} = \frac{s+iw}{(s-iw)(s+iw)} = \frac{s+iw}{s^2+w^2}.$$

Since Euler's formula gives:

$$e^{iwt} = \cos(wt) + i\sin(wt),$$

the linearity of the Laplace Transform yields:

$$L\{e^{iwt}\} = L\{\cos(wt)\} + iL\{\sin(wt)\}.$$

Equating the two representations of $L\{e^{iwt}\}$, we have:

$$L\{\cos(wt)\} + iL\{\sin(wt)\} = \frac{s+iw}{s^2 + w^2}.$$

Since the equality must hold for both the real and imaginary parts, we equate them separately:

$$L\{\cos(wt)\} = \frac{s}{s^2 + w^2}$$
 and $L\{\sin(wt)\} = \frac{w}{s^2 + w^2}$

2.3 Linearity of the Laplace Transform

The Laplace Transform is a linear operator, i.e. for any constants a and b:

$$L\{af(t) + bg(t)\} = aL\{f(t)\} + bL\{g(t)\}$$

Proof

$$\begin{split} L\{af(t) + bg(t)\} &= \int_0^\infty e^{-st} (af(t) + bg(t)) \, dt \\ &= a \int_0^\infty e^{-st} f(t) \, dt + b \int_0^\infty e^{-st} g(t) \, dt \\ &= a L\{f(t)\} + b L\{g(t)\} \end{split}$$

2.4 The First Shift Theorem

Theorem First Shift Theorem

If f(t) has a Laplace Transform, F(s), defined for s > k, then e^{at} f(t) has a Laplace Transform, F(s-a) defined for s-a > k and is given by:

$$L\{e^{at}f(t)\} = F(s-a)$$

or, taking the inverse Laplace Transform of both sides:

$$e^{at}f(t) = L^{-1}\{F(s-a)\}\$$

Example Find the Laplace Transform of $e^{at}\cos(wt)$, where a, w are constants.

We know that $L\{\cos(wt)\} = \frac{s}{s^2+w^2}$, so by the First Shift Theorem:

$$L\{e^{at}\cos(wt)\} = \frac{s-a}{(s-a)^2 + w^2}$$
$$= \frac{s-a}{s^2 - 2as + a^2 + w^2}$$

2.5 Existence of the Laplace Transform

Existence of a Laplace transform is not always guaranteed because we're integrating over an infinite integral. For a Laplace Transform to exist for a given s, then the integral must exist:

$$\int_0^\infty e^{-st} f(t) \, dt$$

Theorem Existence Theorem of Laplace Transforms

Suppose f(t) is a piecewise continuous function on $[0, \infty)$. If f(t) satisfies:

$$|f(t)| \le Me^{kt} \ (0 \le t \le \infty)$$

for some constants, M, k, then the Laplace Transform of f(t) exists for s > k. In other words, the Laplace Transform of f(t) exists if f(t) is bounded by an exponential function.

Proof

If s > k, then from the equation above, we have:

$$|F(s)| = \left|\int_0^\infty f(t)e^{-st}\ dt\right| \leq \int_0^\infty |f(t)|e^{-st}\ dt \leq \int_0^\infty Me^{(k-s)t}\ dt = \frac{M}{s-k}$$

2.6 Integration by Parts

Starting with the product rule:

$$\frac{d}{dx}[uv] = u'v + uv',$$

we can express this in differential form as:

$$d(uv) = u \, dv + v \, du.$$

Integrate both sides with respect to x:

$$\int d(uv) = \int_a^b u \, dv + \int_a^b v \, du.$$

The Fundamental Theorem of Calculus tells us that the left-hand side is simply:

$$uv = \int_a^b u \, dv + \int_a^b v \, du.$$

Rearrange to solve for the desired integral:

$$\int_{a}^{b} u \, dv = uv - \int_{a}^{b} v \, du,$$

Example Use integration by parts to find the Laplace of f(t) = t

$$L\{t\} = \int_0^\infty t e^{-st} dt$$

We integrate by parts by setting:

$$u = t$$
, $dv = e^{-st}$, $du = 1$, $v = -\frac{e^{-st}}{s}$

Then intengrating by parts gives:

$$L\{t\} = \left[-\frac{te^{-st}}{s} \right]_0^\infty + \frac{1}{s} \int_0^\infty e^{-st} dt$$
$$= 0 + \frac{1}{s} \left[-\frac{e^{-st}}{s} \right]_0^\infty$$

Hence:

$$L\{t\} = \frac{1}{s^2}$$

Example Use integration by parts to find the Laplace of $f(t) = \cos(t)$

Let:

$$u = e^{-st}$$
, $du = -se^{-st}$, $dv = \cos(t)$, $v = \sin(t)$

Then:

$$\int_{0}^{\infty} e^{-st} \cos(t) \ dt = \left[e^{-st} \sin(t) \right]_{0}^{\infty} + \int_{0}^{\infty} \sin(t) \cdot s e^{-st} dt = 0 + s \int_{0}^{\infty} e^{-st} \sin(t) dt$$

Considering the sin part:

$$u = e^{-st}$$
, $du = -se^{-st}$, $dv = \sin(t)$, $v = -\cos(t)$

$$\int_0^\infty e^{-st} \sin(t) \ dt = 1 - s \int_0^\infty e^{-st} \cos(t) \ dt$$

Substituting this back into the original integral gives:

$$\int_0^\infty e^{-st} \cos(t) \ dt = 1 - s \int_0^\infty e^{-st} \cos(t) \ dt = s - s^2 \int_0^\infty e^{-st} \cos(t) \ dt$$

$$L\{cos(t)\} = \frac{s}{1+s^2}$$

2.7 Table of Laplace Transforms

f(t)	$L\{f(t)\}$
1	$\frac{1}{s}$, $s > 0$
t	$\frac{1}{s^2}, s > 0$
$t^n, n = 0, 1, 2, 3$	$\frac{n!}{s^{n+1}}, s > 0$
e^{at}	$\frac{1}{s-a}$, $s>a$
$\cos(\omega t)$	$\frac{s}{s^2 + \omega^2}$
$\sin(\omega t)$	$\frac{\omega}{s^2 + \omega^2}$
$\cosh(at)$	$\frac{s}{s^2 - a^2}, s > a \ge 0$
$\sinh(at)$	$\frac{a}{s^2 - a^2}, s > a \ge 0$
$e^{at}\cos(\omega t)$	$\frac{s-a}{(s-a)^2+\omega^2}$
$e^{at}\sin(\omega t)$	$\frac{\omega}{(s-a)^2+\omega^2}$
$e^{at}f(t)$	F(s-a)

2.8 Laplace Transforms of Derivatives

Theorem Laplace Transform of Derivatives

Suppose that f(t) and f'(t) are continous and that $|f(t)| \le Me^{kt}$, $\forall t \ge 0$ and for constans M, k. Then the Laplace Transform of f'(t) exists for s > k and is given by:

$$L\left\{\frac{df}{dt}\right\} = sL\left\{f\right\} - f(0) \text{ for } s > k$$

We can easily extend this to higher order derivatives. Assume the Laplace Transform of $f^{(n)}(t)$ exists for s > k and is given by:

$$L\left\{\frac{df^n}{dt^n}\right\} = s^n L\left\{f\right\} - s^{n-1}f(0) - s^{n-2}f'(0) - \dots - f^{(n-1)}(0) \quad \text{for } s > k$$

Example Find $L\{t^2\}$ using the fact $L\{s\} = 1/s$ for s > 0

$$L\{f''\} = s^2 L\{f\} - sf(0) - f'(0)$$

With $f(t) = t^2$. Since f'(t) = 2t, f''(t) = 2, f'(0) = 0, f(0) = 0, gives:

$$L\{2\} = s^2 L\{t^2\} - s \cdot 0 - 0$$

So that:

$$L\{t^2\} = \frac{L\{2\}}{s^2} = \frac{2}{s^3}$$

Example Find $L\{\sin(t)\}$ and $L\{\cos(t)\}$

We again use ther equation:

$$L\{f''\} = s^2 L\{f\} - sf(0) - f'(0)$$

With $f(t) = \sin(t)$, $f'(t) = \cos(t)$, $f''(t) = -\sin(t)$, $\sin(0) = 0$, $\cos(0) = 1$. This gives:

$$L\{-\sin(t)\} = s^2 L\{\sin(t)\} - s \cdot 0 - 1$$

So that:

$$L\{\sin(t)\} = \frac{1}{s^2 + 1}$$

Similarly, we can find:

$$L\{\cos(t)\} = \frac{s}{s^2 + 1}$$

2.9 Solving Initial Value Problems

Consider an example from mechanics: A particle of mass m > 0 lies on rough table, attached to a spring of stiffness k > 0. At any time t > 0, the mass is a distance x(t) from the equillibrium position O, and x(t) is much less than the length of the spring.

The mass is subject to a driving force $F_d(t)$, from Newtons second law, we have:

$$F_d(t) - kx - \gamma \frac{dx}{dt} = m \frac{dx^2}{dt^2}$$

Where $\gamma > 0$ is the **damping constant** and the term $\gamma \frac{dx}{dt}$ models the **fricition due to roughness** of the table, which oppposes direction of motion. The **restoring force** due to the spring is -kx; and always points towards O. The term $m \frac{dx^2}{dt^2}$ is the **acceleration of the mass**. We can rewrite this as:

$$F_d(t) = m\frac{dx^2}{dt^2} + \gamma \frac{dx}{dt} + kx$$

In order to solve this, we also need initial displacement $v_0 = x(0)$ and initial velocity $v_0 = \frac{dx}{dt}(0)$.

Example

$$\frac{dx^2}{dt^2} + 3\frac{dx}{dt} + 2x = 0$$
, $x(0) = 0$, $\frac{dx}{dt}(0) = 1$

1. Take Laplace of governing equation:

$$L\left\{\frac{dx^2}{dt^2}\right\} = s^2 L\{x\} - sx(0) - x'(0) = s^2 L\{x\} - 1$$

$$L\left\{\frac{dx}{dt}\right\} = sL\{x\} - x(0) = sLx$$

Hence:

$$s^{2}L\{x\} - 1 + 3sL\{x\} + 2L\{x\} = 0$$

This is known as the **subsidary equation**. Rearranging:

$$(s^2 + 3s + 2)L\{x\} = 1$$

2. Solve the subsidary equation:

$$L\{x\} = \frac{1}{s^2 + 3s + 2}$$

3. Find the inverse Laplace Transform:

$$x(t) = L^{-1} \left\{ \frac{1}{s^2 + 3s + 2} \right\} = L^{-1} \left\{ \frac{1}{(s+1)(s+2)} \right\}$$

$$\frac{1}{(s+1)(s+2)} = \frac{A}{s+1} + \frac{B}{s+2} = \frac{A(s+2) + B(s+1)}{(s+1)(s+2)}$$

Hence:

$$A(s+2) + B(s+1) = 1 \rightarrow A = 1, B = -1$$

Thus:

$$x = L^{-1} \left\{ \frac{1}{s+1} - \frac{1}{s+2} \right\} = L^{-1} \left\{ \frac{1}{s+1} \right\} - L^{-1} \left\{ \frac{1}{s+2} \right\} = e^{-t} - e^{-2t}$$

2.10 Practice Problems

1. Use the First Shift Theorem $(L\{e^{at}f(t)\} = F(s-a))$ to find the Laplace transform of the following functions:

(a)
$$t^3 e^{-3t}$$
 (b) $e^{-t} \cos(2t)$ (c) $e^{-4t} \cosh(5t)$ (d) $e^{-t} \sin^2(t)$

2. Use the First Shift Theorem $(L^{-1}{F(s-a)}) = e^{at}f(t)$ to find the inverse Laplace transform of the following functions:

(a)
$$\frac{6s-4}{s^2-4s+20}$$
 (b) $\frac{3s+7}{s^2-2s-3}$ (c) $\frac{4s+12}{s^2+8s+16}$

3. Solve the following initial value problems using the method of Laplace transforms:

$$y'' + y' - 6y = 0$$
, $y(0) = 0$, $y'(0) = 1$;
 $y'' - y = t$, $y(0) = 1$, $y'(0) = 1$.

4. Find the inverse Laplace transform of the following functions using the method of partial fractions:

(a)
$$\frac{2s^2-4}{(s+1)(s-2)(s-3)}$$
. (b) $\frac{5s^2-15s-11}{(s+1)(s-2)^3}$. (c) $\frac{3s+1}{(s-1)(s^2+1)}$. (d) $\frac{e^{-5s}}{(s^2+1)(s^2+2)}$.

2.11 The Dirac Delta Function

The **Dirac Delta Function** models extremely brief but intense forces like a hammer hitting a nail. It starts as a function δ_{ε} , that equals $\frac{1}{2\varepsilon}$ over the interval $t \in [-\varepsilon, \varepsilon]$ and 0 elsewhere.

$$\delta_{\epsilon}(t)$$

$$\delta_{\epsilon}(t) = \begin{cases} \frac{1}{2\epsilon} & \text{if } t \in [-\epsilon, \epsilon], \\ 0 & \text{otherwise.} \end{cases}$$

This function creates a rectangular pulse with the following propertites:

Height:
$$\frac{1}{2\varepsilon}$$
 Width: 2ε Area: 1 (always)

As ε approaches 0, the function becomes infinitely tall and thin, but the area remains 1. This limit defines the Dirac Delta Function:

$$\delta(t) = \lim_{\varepsilon \to 0+} \{ \delta_{\varepsilon}(t) \}$$

Properties of the Dirac Delta Function:

$$\delta(t) = 0 \text{ for } t \neq 0$$

$$\int_{-\infty}^{\infty} \delta(t) dt = 1$$

$$\int_{-\infty}^{\infty} \delta(t - t_0) f(t) dt = f(t_0)$$

The Laplace Transform of the Dirac Delta Function is:

$$L\{\delta(t-t_0)\} = \int_0^\infty e^{-st} \delta(t-t_0) dt = \int_{-infty}^\infty e^{-st} \delta(t-t_0) dt = e^{-st_0} \quad \text{for } t_0 > 0$$

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Example : Solve the following initial value problem which governs the behavious of an RLC circuit

$$LQ'' + RQ' + \frac{Q}{C} = V_0 \delta(t - a)$$
$$Q(0) = 0$$
$$'(0) = 0$$

Where a, L, R, C, V_0 are all positive constans and $4l > R^2C$. Note that the applied voltage corresponds to an impulse of stength V_0 at t=a