System Dynamics and Vibrations

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Chapter 1: Elements of analytical dynamics
Part 1

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Introduction

- ➤ Newtonian mechanics (vectorial mechanics)
 - Equations of motion for <u>individual particles</u> expressed in terms of <u>physical coordinates and forces</u> (represented both by vectors)
 - Free-body diagram for each of the masses in the system
 - Involve <u>reaction forces and interacting forces</u> (kinematical constraints) → many equations and unknowns

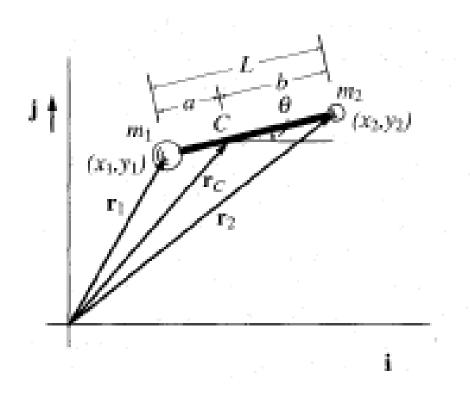
Introduction

- >Analytical mechanics (Lagrange, variational approach)
 - Considers the <u>systems as a whole</u> → reaction and constraint forces are excluded
 - Dynamics problems formulated in terms of: kinetic energy, potential energy, virtual work of non-conservative forces
 - Equations of motion formulated in terms of generalized coordinates and generalized forces → broader and more abstract approach
 - The mathematical formulation is independent of any special system of coordinates

Contents

- Introduction
- Degrees of freedom and generalized coordinates
- The principle of virtual work
- The principle of D'Alembert
- Lagrange's equations

- Equations of motion for a system by the Newtonian approach:
 - Isolate the masses
 - Draw one free-body diagram for each of the masses, including all forces acting upon it:
 - Applied forces
 - Reaction forces
 - Internal forces
 - We obtain normally more equations and unknowns than necessary
 - The motion is described in terms of physical coordinates: may not always be independent → equations of motion + constraint equations



- Movement in x, y plane
- Position vectors:

$$\mathbf{r}_1 = x_1 \mathbf{i} + y_1 \mathbf{j}, \ \mathbf{r}_2 = x_2 \mathbf{i} + y_2 \mathbf{j},$$

• x_1 , x_2 , y_1 , y_2 are not independent

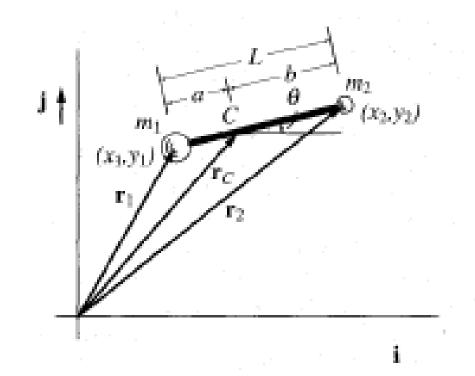
$$(x_2 - x_1)^2 + (y_2 - y_1)^2 = L^2$$

→ constraint equation (+ the equations of motion)

- A better choice of coordinates obviates the difficulties of working with too may coordinates and constraint equations
- For a system of N mass particles with positions defined by the radius vectors \mathbf{r}_i (x_i , y_i , z_i) in a three-dimensional space:

$$n = 3N - c$$

- c number of constraints
- *n* number of degrees of freedom of the system



Much simpler in generalized coordinates:

$$\mathbf{r}_C = \mathbf{r}_C(x_C, y_C)$$

$$\theta$$

• Three independent coordinates:

$$x_{c}, y_{c}, \theta$$

No constraint equation is needed

$$\mathbf{r}_1 = \mathbf{r}_C - a(\cos\theta \mathbf{i} + \sin\theta \mathbf{j}), \ \mathbf{r}_2 = \mathbf{r}_C + b(\cos\theta \mathbf{i} + \sin\theta \mathbf{j}),$$

$$a = \frac{m_2 L}{m_1 + m_2}$$
$$b = \frac{m_1 L}{m_1 + m_2}$$

- Independent coordinates are normally called generalized coordinates: q_1, q_2, \dots, q_n
- Coordinate transformation:

$$x_{1} = x_{1}(q_{1}, q_{2}, ..., q_{n})$$

$$y_{1} = y_{1}(q_{1}, q_{2}, ..., q_{n})$$

$$z_{1} = z_{1}(q_{1}, q_{2}, ..., q_{n})$$

$$x_{2} = x_{2}(q_{1}, q_{2}, ..., q_{n})$$

$$z_{N} = z_{N}(q_{1}, q_{2}, ..., q_{n})$$

- Generalized coordinates q_1, q_2, \dots, q_n are not unique for a given system
- Normally only one or two sets may represent a suitable choice → the equations of motion in terms of these coordinates have the simplest form
- In vibrations, this choice is obvious for most of the cases

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- Basically a statement of the <u>static equilibrium</u> of a mechanical system
- The first variational principle of mechanics
- Tool to transition from Newtonian mechanics to Lagrangian mechanics
- New concepts:
 - Virtual displacements
 - Constraint forces

- System of N particles in a three-dimensional space
- We define <u>virtual displacements</u> as infinitesimal changes in the coordinates x_1 , y_1 , z_1 , x_2 , ..., z_n

$$\delta x_1, \delta y_1, \delta z_1, \delta x_2, ..., \delta z_N$$

- Virtual displacements must be:
 - Consistent with the system constraints, but
 - Arbitrary

- Virtual displacements represent small variations in the coordinates resulting from imagining the system in a slightly displaced position:
 - Virtual displacements take place instantaneously $\rightarrow \delta t = 0$
- Symbol δ emphasizes the virtual character of the instantaneous variations
- Symbol *d* represents actual differentials of position coordinates taking place in the time interval *dt* (forces can change in that time interval)
 - Virtual displacements, being infinitesimal, obey the rules of differential calculus

 We assume that every one of the N particles in the system is acted upon by the resultant force

$$\mathbf{R}_{i} = \mathbf{F}_{i} + \mathbf{f}_{i}, i = 1, 2, ..., N$$

 \mathbf{F}_i is an applied force (gravitational force, aerodynamic lift and drag, magnetic forces, etc.)

 \mathbf{f}_i is a constraint force (for example the force that keeps a particle confined to a given surface)

- For a system in equilibrium every particle must be at rest
- Then:

$$\mathbf{R}_{i} = \mathbf{F}_{i} + \mathbf{f}_{i} = \mathbf{0}, \quad i = 1, 2, ...N$$

$$\overline{\delta W_i} = \mathbf{R}_i \cdot \delta \mathbf{r}_i = 0, \quad i = 1, 2, ...N$$

 δW_i represents the <u>virtual work</u> perfomed by the resultant force vector \mathbf{R}_i over the virtual displacement vector $\delta \mathbf{r}_i$ of particle i

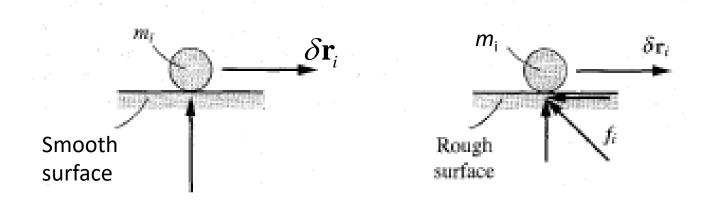
• Summing up over *i*:

$$\overline{\delta W} = \sum_{i=1}^{N} \mathbf{R}_{i} \cdot \delta \mathbf{r}_{i} = 0$$

→ the virtual work for the entire system must vanish

$$\overline{\delta W} = \sum_{i=1}^{N} \mathbf{F}_{i} \cdot \delta \mathbf{r}_{i} + \sum_{i=1}^{N} \mathbf{f}_{i} \cdot \delta \mathbf{r}_{i} = 0$$

 We limit ourselves to systems for which the virtual work performed by the constraint forces is zero:



- We limit ourselves to systems for which the virtual work performed by the constraint forces is zero:
- Particle or body on a smooth surface (no friction)
- Other examples:
 - Articulated bodies
 - Bodies forced to be in contact
 - Bodies rolling and pivoting (without sliding) one on each other
 - → solid mechanics

 We limit ourselves to systems for which the virtual work performed by the constraint forces is zero:

$$\sum_{i=1}^{N} \mathbf{f}_{i} \cdot \delta \mathbf{r}_{i} = 0$$

- Therefore: $\overline{\delta W} = \sum_{i=1}^{N} \mathbf{F}_{i} \cdot \delta \mathbf{r}_{i} = 0$
- → Principle of virtual work: the work performed by the applied forces through infinitesimal virtual displacements compatible with the system constraints is zero

- When the <u>virtual displacements are all independent</u>:
- We can invoke the arbitrariness of the virtual displacements

$$\rightarrow$$
 equation $\overline{\delta W} = \sum_{i=1}^{N} \mathbf{F}_{i} \cdot \delta \mathbf{r}_{i} = 0$

can be satisfied for all possible values of $\delta {f r}_i$

only if
$$\mathbf{F}_i = \mathbf{0}$$
, $i = 1, 2, ...N$ \rightarrow Equilibrium conditions

- If the coordinates \mathbf{r}_i (i = 1, 2, ..., N) are not <u>independent</u> (but related by constraint equations):
- It is more convenient to switch to generalized coordinates q_1 , q_2 , ..., q_n

We can express $\mathbf{r}_i = \mathbf{r}_i(q_1, q_2, ..., q_n)$ i = 1, 2, ...N

generalized coordinates are independent by definition

Virtual displacements are:

$$\delta \mathbf{r}_{i} = \frac{\partial \mathbf{r}_{i}}{\partial q_{1}} \delta q_{1} + \frac{\partial \mathbf{r}_{i}}{\partial q_{2}} \delta q_{2} + \dots + \frac{\partial \mathbf{r}_{i}}{\partial q_{n}} \delta q_{n} = \sum_{k=1}^{n} \frac{\partial \mathbf{r}_{i}}{\partial q_{k}} \delta q_{k}, \quad i = 1, 2, \dots N$$

$$\delta q_k$$
 $(k = 1, 2, ..., n)$

are virtual generalized displacements, all independent

$$\overline{\delta W} = \sum_{i=1}^{N} \mathbf{F}_{i} \cdot \delta \mathbf{r}_{i} = \sum_{i=1}^{N} \mathbf{F}_{i} \cdot \sum_{k=1}^{n} \frac{\partial \mathbf{r}_{i}}{\partial q_{k}} \delta q_{k} = \sum_{k=1}^{n} \left(\sum_{i=1}^{N} \mathbf{F}_{i} \cdot \frac{\partial \mathbf{r}_{i}}{\partial q_{k}} \right) \delta q_{k} = \sum_{k=1}^{n} Q_{k} \delta q_{k} = 0$$

where
$$Q_k = \sum_{i=1}^N \mathbf{F}_i \cdot \frac{\partial \mathbf{r}_i}{\partial q_k}$$
, $(k = 1, 2, ..., n)$

are known as **generalized forces**

- Being δq_k (k=1,2,...,n) all independent, and therefore entirely arbitrary
- Letting $\delta q_1 = 1$, $\delta q_2 = \delta q_3 = ... = \delta q_n = 0$ $\overline{\delta W} = \sum_{k=1}^{n} Q_k \delta q_k = 0 \Rightarrow Q_1 = 0$
- Repeating the argument with k = 2, 3, ..., n in sequence
- → we obtain the equilibrium conditions:

$$Q_k = 0, \quad (k = 1, 2, ..., n)$$