# System Dynamics and Vibrations

Prof. Gustavo Alonso

Chapter 2: Concepts from vibrations
Part 2

School of General Engineering Beihang University (BUAA)

#### Contents

- Introduction
- Modeling of mechanical systems
- System differential equations of motion
- Linearization about equilibrium points
- System and response characteristics. The superposition principle.

$$m\ddot{y} = F(y, \dot{y})$$

Let's consider a solution having the form:

$$y(t) = y_e + x(t)$$

 being x(t) a relatively small displacement from equilibrium

• then: 
$$\dot{y}(t) = \dot{x}(t)$$
$$\ddot{y}(t) = \ddot{x}(t)$$

• Expandig  $F(y,\dot{y})$  in a Taylor series about an equlibrium point  $y_e$ :

$$F(y, \dot{y}) = F(y_e, 0) + \frac{\partial F(y, \dot{y})}{\partial y} \bigg|_{\substack{y=y_e \\ \dot{y}=0}} x + \frac{\partial F(y, \dot{y})}{\partial \dot{y}} \bigg|_{\substack{y=y_e \\ \dot{y}=0}} \dot{x} + O(x^2)$$

$$m\ddot{y} = F(y, \dot{y})$$

$$\frac{1}{m} \frac{\partial F(y, \dot{y})}{\partial y} \Big|_{\substack{y=y_e \\ \dot{y}=0}} = -b$$

$$\frac{1}{m} \frac{\partial F(y, \dot{y})}{\partial \dot{y}} \Big|_{\substack{y=y_e \\ \dot{y}=0}} = -a$$

$$m\ddot{x} + a\dot{x} + bx = 0$$

 We have assumed that displacements from equilibrium are sufficiently small that the nonlinear terms can be ignored

$$m\ddot{y} = F(y, \dot{y})$$
  $\ddot{x} + a\dot{x} + bx = 0$ 

- → linearized equation of motion about equilibrium (small motions assumption)
- The motion characteristics in the neighborhoud of equilibrium depend on parameters a, b

$$\ddot{x} + a\dot{x} + bx = 0$$

• Linear equation with constant coefficients:

$$x(t) = Ae^{st}$$

A: amplitude

s: constant exponent

Combining

$$m\ddot{x} + a\dot{x} + bx = 0$$

$$x(t) = Ae^{st}$$

$$s^{2} + as + b = 0$$

$$s^2 + as + b = 0$$

- → Characteristic equation (algebraic equation)
- The roots are:

$$\frac{s_1}{s_2} = -\frac{a}{2} \pm \sqrt{\left(\frac{a}{2}\right)^2 - b}$$

• So the solution to  $m\ddot{x} + a\dot{x} + bx = 0$  is:

$$x(t) = A_1 e^{s_1 t} + A_2 e^{s_2 t}$$

Asymptotically stable

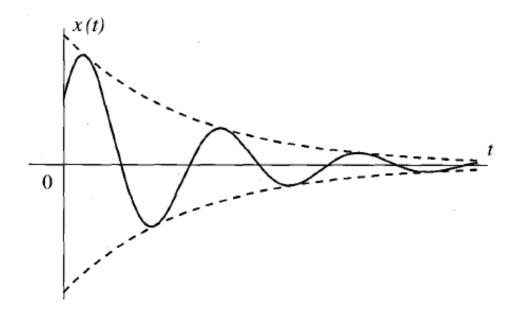
Aperiodic decay -  $\boxed{1}$ Aperiodic decay -  $\boxed{1}$ Aperiodic divergence -  $\boxed{3}$ Aperiodic divergence -  $\boxed{3}$ Aperiodic divergence -  $\boxed{3}$ 

Asymptotically stable solution (a>0, b>0)

#### Aperiodically decay

# 0

#### Decaying oscillation

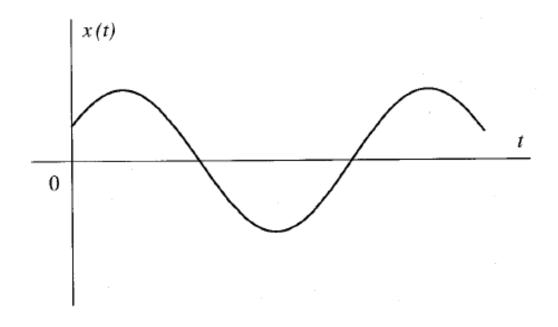


Asymptotically stable

Aperiodic

2. Stable motion (a=0, b>0)

#### Harmonic oscillation



Asymptotically stable

Stable

Aperiodic

divergence - 3Aperiodic

divergence - 3Aperiodic

divergence - 3Aperiodic

divergence - 3Aperiodic

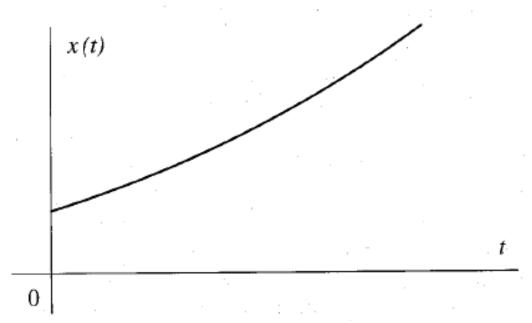
divergence - 3

3. Unstable motion (b<0, b>0 & a<0)

Diverging oscillation

x(t)

#### Aperiodically diverging motion



$$m\ddot{y} = F(y, \dot{y})$$

#### 1. Mass-spring-damper system

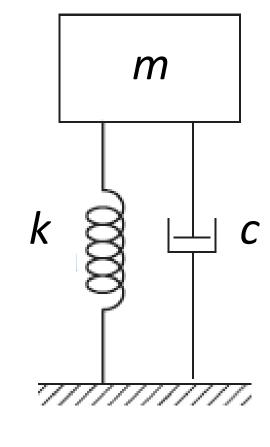
$$m\ddot{y} = -c\dot{y} - ky - mg$$

#### Equilibrium points:

$$F(y, \dot{y}) = --c\dot{y} - ky - mg$$

$$F(y_e, 0) = -ky_e - mg = 0$$

$$y_e = -\frac{mg}{k}$$



$$m\ddot{y} = F(y, \dot{y})$$

#### 1. Mass-spring-damper system

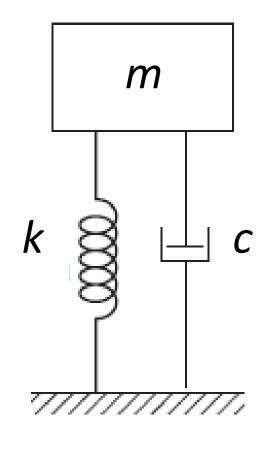
Stability of equilibrium point:

$$y(t) = y_e + x(t) = -\frac{mg}{k} + x(t)$$

$$\ddot{x} + \frac{c}{m}\dot{x} + \frac{k}{m}x = 0$$

$$a = \frac{c}{m}$$

$$b = \frac{k}{m}$$

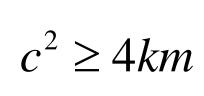


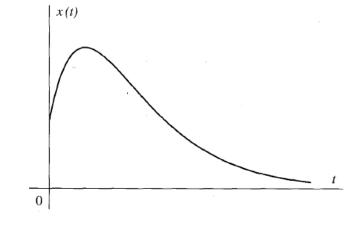
1. Mass-spring-damper system

Stability of equilibrium point:

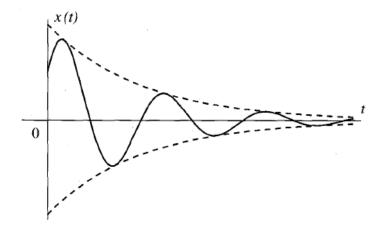
asymptotically stable

$$a = \frac{c}{m} > 0$$
$$b = \frac{k}{m} > 0$$





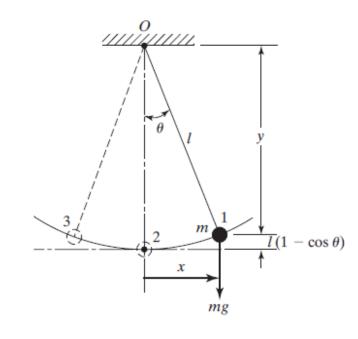
$$c^2 < 4km$$



#### 2. Simple pendulum

$$\ddot{\theta} + \frac{g}{l}\sin\theta = 0$$

#### Equilibrium points:



$$m\ddot{y} = F(y, \dot{y}) \qquad F$$

$$F(y, \dot{y}) = F(\theta, \dot{\theta}) = F(\theta, 0) = -\frac{g}{l}\sin\theta$$

$$m = 1$$

$$y = \theta$$

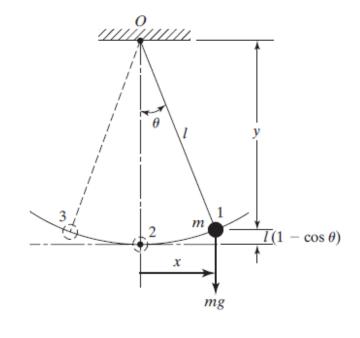
$$F(\theta,0) = -\frac{g}{l}\sin\theta = 0$$

equation of equilibrium

#### 2. Simple pendulum

#### Equilibrium points:

$$F(\theta,0) = -\frac{g}{l}\sin\theta = 0$$
 equation of equilibrium



solutions: 
$$\theta_e = 0, \pm \pi, \pm 2\pi,...$$

$$\theta_{e1} = 0$$
$$\theta_{e2} = \pi$$

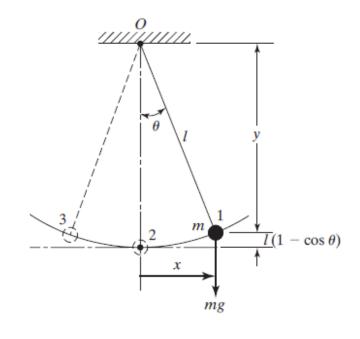
$$\theta_{e2} = \pi$$

#### 2. Simple pendulum

$$\theta_{e1} = 0$$

Stability of equilibrium points:

$$\theta_{e2} = \pi$$



$$\theta(t) = \theta_e + \phi(t)$$

$$\ddot{\phi} + b\phi = 0$$

$$b = -\frac{\partial F(\theta, 0)}{\partial \theta} \bigg|_{\theta = \theta_{e}} = \frac{g}{l} \cos \theta_{e}$$

#### 2. Simple pendulum

Stability of equilibrium points:

$$\theta_{e1} = 0$$

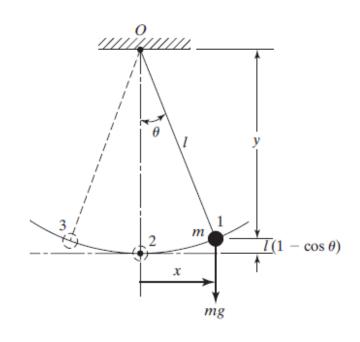
$$b = \frac{g}{l} > 0$$

stable

$$\theta_{e2} = \pi$$

$$b = -\frac{g}{l} < 0$$

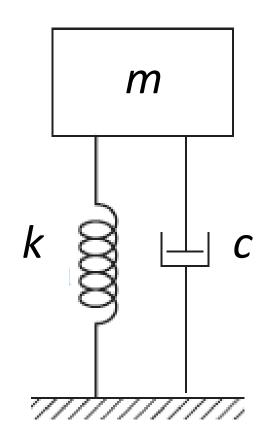
unstable



#### Contents

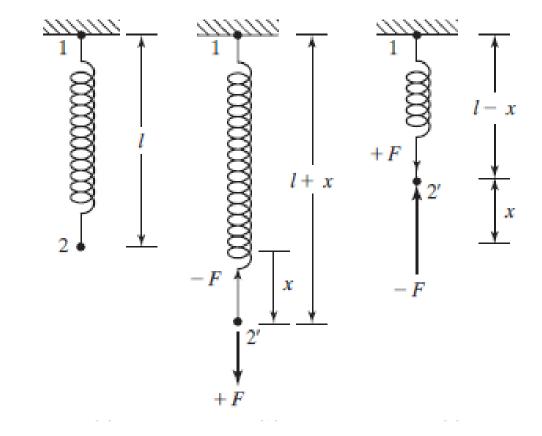
- Introduction
- Modeling of mechanical systems
- System differential equations of motion
- Linearization about equilibrium points
- System and response characteristics. The superposition principle.

## Mass-spring-damper system



## Spring elements

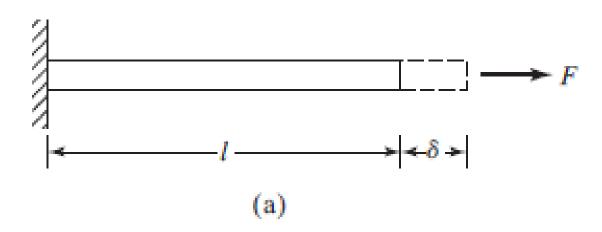
- Type of mechanical link
- Characterized by stiffness, k
- Normally negligible mass and damping
- Any elastic or deformable body or member, such as a cable, bar, beam, shaft or plate, can be considered as a spring.

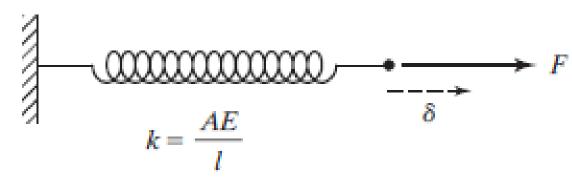


$$F = kx \qquad V = \frac{1}{2}kx^2$$

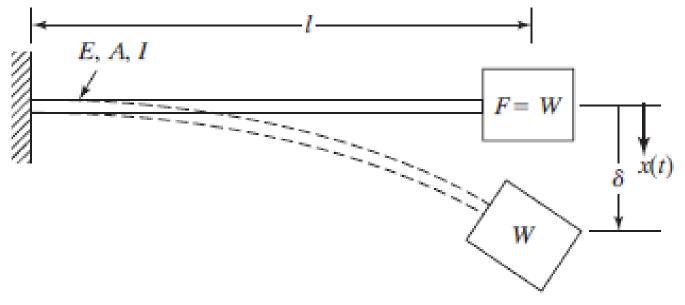
## Spring elements: spring constant of a rod

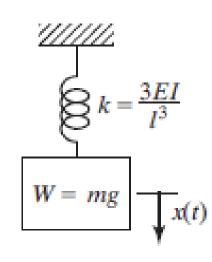
$$\delta = \frac{\delta}{l}l = \varepsilon l = \frac{\sigma}{E}l = \frac{Fl}{AE}$$
$$k = \frac{F}{\delta} = \frac{AE}{l}$$





## Spring elements: spring constant of a cantilever beam





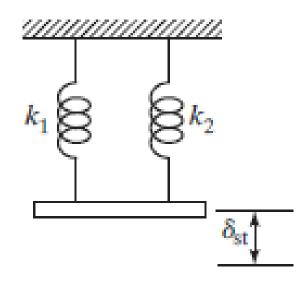
(a) Cantilever with end force

$$\delta = \frac{Wl^3}{3EI} \Longrightarrow k = \frac{W}{\delta} = \frac{3EI}{l^3}$$

(b) Equivalent spring

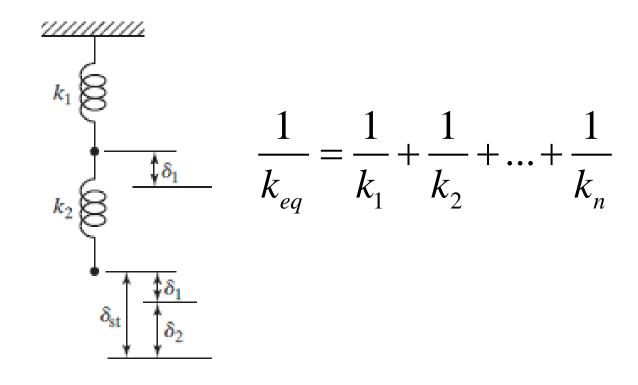
## Spring elements: combination of springs

Springs in parallel



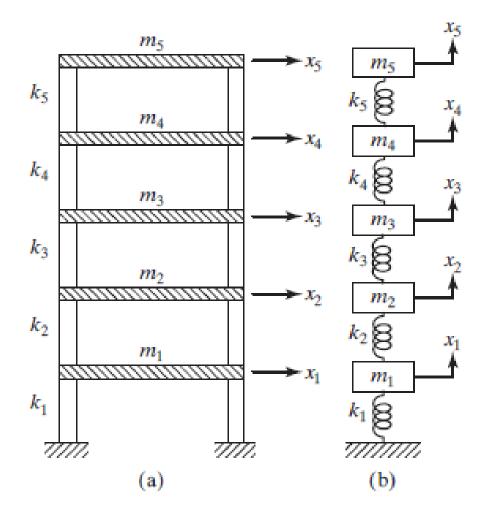
$$k_{eq} = k_1 + k_2 + \dots + k_n$$

• Springs in series



#### Mass or Inertia elements

- Assumed to be a rigid body
- It can gain or lose kinetic energy if the velocity of the body changes
- In many practical cases, several masses appear in combination
- → For a simpe analysis, we can replace these masses by a single equivalent mass



## Damping elements

- Damping is the mechanism by which the vibrational energy is gradually converted into heat or sound
- Dampers are assumed to have neither mass nor elasticity
- Damping forces exists only if there is relative velocity between the two ends of the damper

#### Viscous damping:

- Related to the interaction of the body with the surrounding fluid
- Depends on: the size and shape of the body, the viscosity of the fluid, the frequency of vibration, the velocity of the body, etc.
- The damping force is proportional to the velocity of the vibrating body

## Damping elements

#### Coulomb or dry-friction damping:

- The damping force is constant in magnitude but opposite in direction to that of the motion of the vibrating body
- It is caused by friction between rubbing surfaces that either are dry or have insufficient lubrication

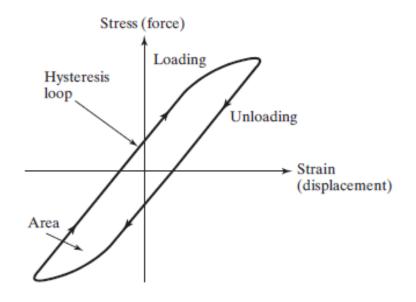
#### Material or Solid or Hysteretic Damping:

- When a material is deformed, energy is absorbed and dissipated by the material
- The effect is due to friction between the internal planes, which slip or slide as the deformations take place.
- When a body having material damping is subjected to vibration, the stress-strain diagram shows a hysteresis loop as indicated
- The area of this loop denotes the energy lost per unit volume of the body per cycle due to damping

## Damping elements

#### Coulomb or dry-friction damping:

- The damping force is constant in magnitude but opposite motion of the vibrating body
- It is caused by friction between rubbing surfaces that eith lubrication



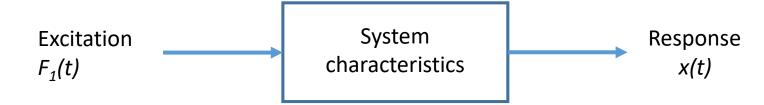
#### Material or Solid or Hysteretic Damping:

- When a material is deformed, energy is absorbed and dissipated by the material
- The effect is due to friction between the internal planes, which slip or slide as the deformations take place.
- When a body having material damping is subjected to vibration, the stress-strain diagram shows a hysteresis loop as indicated
- The area of this loop denotes the energy lost per unit volume of the body per cycle due to damping

#### Contents

- Introduction
- Modeling of mechanical systems
- System differential equations of motion
- Linearization about equilibrium points
- System and response characteristics. The superposition principle.

- How do the system characteristics affect the response of the system?
- A systemis defined as an aggregation of components working together as a single unit.
- System charscteristics depend on:
  - Individual components
  - The manner in which these components are conected to one another



- Systems can be:
  - Linear
  - Nonlinear

• Linear system  $F_1(t)$  Linear system  $x_1(t)$   $x_2(t)$   $x_2(t)$  Linear system  $x_2(t)$  Linear system  $x_2(t)$ 

• Nonlinear system: 
$$x(t) \neq c_1 x_1(t) + c_2 x_2(t)$$

- Principle of superposition (linear system):
  - → If a linear system is acted upon by a linear combination of individual excitations, the individual responses can be first obtained separately and then combined linearly to obtain the total response
- A system is linear if the dependent variable x(t) and all its time derivatives appear in the equation of motion to the first power or zero power only

$$m\ddot{x} + c\dot{x} + kx = F$$
 linear

$$m\ddot{x} + c\dot{x} + k(x + \varepsilon x^3) = F$$
 nonlinear

- Time dependency:
  - Linear time-invariant systems (linear systems with constant coefficients)
    - For example:  $m\ddot{x}(t) + c\dot{x}(t) + kx(t) = F(t)$
    - If the excitation F(t) is delayed by an amount of time  $\tau$ , the resonse x(t) is delayed by the same amount t
  - Time-varying systems (systems with time-dependent coefficients)
    - For example:  $m\ddot{x}(t) + k(1 + a\cos\omega t)x(t) = F(t)$

- Response to <u>initial excitations</u>
- → simplest problem

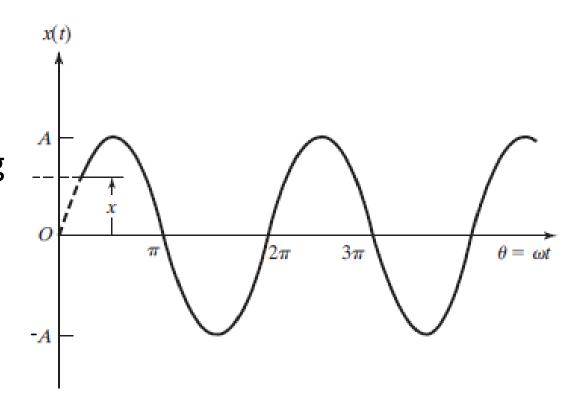
$$m\ddot{x}(t) + c\dot{x}(t) + kx(t) = F(t) = 0$$

- Homogeneous equation 
   solution has exponential form
- Characteristic equation to determine the exponents
- Coefficients of the exponential terms obtained by letting x(t),  $\dot{x}(t)$  evaluated at t=0 match the initial displacement and velocity

- Response to **harmonic excitations**
- → it is also harmonic
- The response has the same frequency as the excitation frequency, but it differs in magnitude and possesses a phase angle relative to the excitation
- Magnitude and phase angle depend on the driving frequency
- The response to harmonic excitations is a steady-state response

- Oscillatory motion may repeat itself regularly (like in the case of a simple pendulum), or
- It may display considerable irregularity (as in the case of ground motion during an earthquake)
- <u>Periodic motion</u>: the motion is repeated after equal intervals of time
- The simplest type of periodic motion is harmonic motion

- Oscillatory motion may repeat itself regularly (like in the case of a simple pendulum), or
- It may display considerable irregularity (as in the case of ground motion during an earthquake)
- <u>Periodic motion</u>: the motion is repeated after equal intervals of time
- The simplest type of periodic motion is harmonic motion

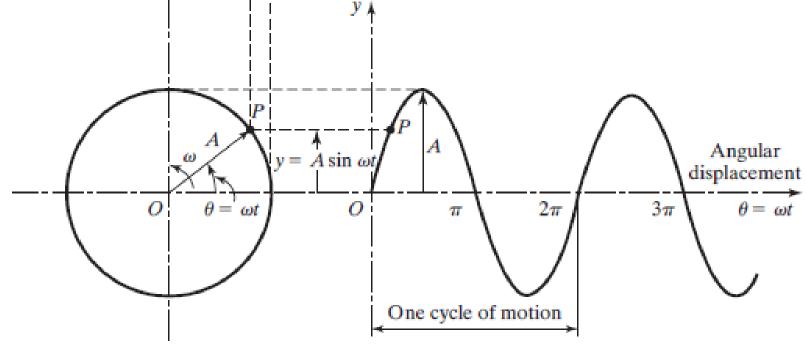


• Harmonic motion can be represented conveniently by means of a vector  $\overrightarrow{OP}$  of magnitude A rotating at a constant angular velocity

$$x(t) = A \sin \theta = A \sin \omega t$$
  

$$\dot{x}(t) = \omega A \cos \omega t$$
  

$$\ddot{x}(t) = -\omega^2 A \sin \omega t = -\omega^2 x(t)$$



• Complex-number representation of harmonic motion

$$\vec{X} = a + ib$$

real and imaginary parts

$$\vec{X} = A\cos\theta + iA\sin\theta = Ae^{i\theta}$$

$$A = (a^2 + b^2)^{\frac{1}{2}}$$

modulus

$$\theta = \tan^{-1} \frac{b}{a}$$

argument

- Response to periodic excitations
- → Periodic excitations can be represented by Fourier series, i.e. series of harmonic functions
- Applying the principle of superposition, the response to periodic excitations can be expressed in the form of a series of harmonic responses → it is a steady-state response

## Nature of excitation: periodic motion

 Any periodic function of time can be represented by Fourier series as an infinite sum of sine and cosine terms

$$x(t) = \frac{a_0}{2} + a_1 \cos \omega t + a_2 \cos 2\omega t + \dots + b_1 \sin \omega t + b_2 \sin 2\omega t + \dots$$

$$= \frac{a_0}{2} + \sum_{n=1}^{\infty} \left( a_n \cos n\omega t + b_n \sin n\omega t \right)$$

$$\omega = \frac{2\pi}{\tau}$$

$$a_0 = \frac{\omega}{\pi} \int_0^{2\pi/\omega} x(t)dt = \frac{2}{\tau} \int_0^{\tau} x(t)dt$$

$$a_n = \frac{\omega}{\pi} \int_0^{2\pi/\omega} x(t)\cos n\omega t dt = \frac{2}{\tau} \int_0^{\tau} x(t)\cos n\omega t dt$$

$$b_n = \frac{\omega}{\pi} \int_0^{2\pi/\omega} x(t)\sin n\omega t dt = \frac{2}{\tau} \int_0^{\tau} x(t)\sin n\omega t dt$$

- Response to an arbitrary excitation:
- → Superposition of impulse forces of different magnitde and applied at different times
- The impulse response is defined by the response to a unit impulse applied at t= 0
- Assuming that the impulse response is known, the response of a linear system with constant coefficients can be expressed as a superposition of impulse responses of different magnitudes and applied at different times.
- This superposition is called the convolution integral, or the superposition integral.

- Response to random excitations:
- → the response is also a random function
- Random functions are characterized by their mean value, mean square value, autocorrelation function, power spectral density function, ...
- Solution obtained by Fourier transfors working in the frequency domain
- Results are defined in terms of probability distributions