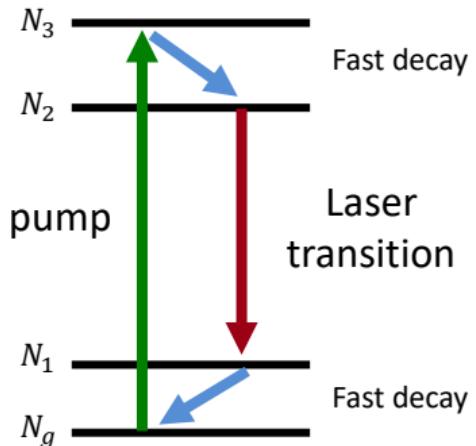


# Self-terminating lasers



Working hypotheses:

- **Four-level laser:**  $N_1 \cong 0$  always

What happens if the lower laser level has a finite lifetime ( $\tau_1 \neq 0$ )?

At equilibrium conditions (steady-state):

$$\frac{N_1}{\tau_1} = \frac{N_2}{\tau_{21}} \Rightarrow \frac{N_2}{N_1} = \frac{\tau_{21}}{\tau_1}$$

To get laser action  $N_2 > N_1 \Rightarrow \tau_{21} > \tau_1$

If  $\tau_{21} < \tau_1 \Rightarrow$  laser action is possible with pulsed pumping only

with  $\Delta t_p < \tau$   
pulse duration      total lifetime of the upper laser level

laser action ends when the accumulation of population in the lower laser level destroys population inversion

→ **Self-terminating lasers**

A Nd:glass laser ( $n = 1.54$ ) oscillating at the fundamental line ( $\lambda = 1054 \text{ nm}$ ) emits in cw an output power  $P_{out} = 320 \text{ mW}$ . The length of the active medium is  $l = 8 \text{ cm}$ . The resonant cavity is a Fabry-Perot cavity of length  $L = 50 \text{ cm}$ , made of a first mirror with reflectivity  $R_1 = 95\%$  and a second mirror (outcoupling mirror) with  $R_2 = 75\%$ .

Assuming that the internal losses for single pass are  $L_i = 16\%$ , determine:

- the photon lifetime
- the number of photons in the cavity
- the critical population inversion
- the saturation intensity

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

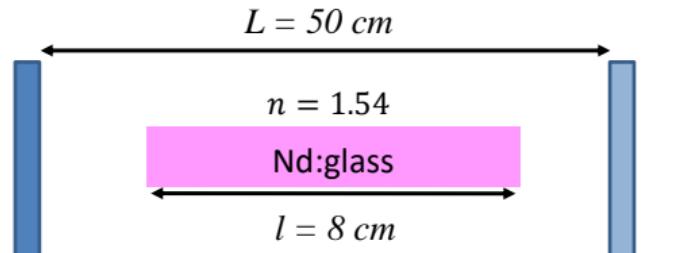
$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 1054 \text{ nm} \quad h\nu = 1.89 \cdot 10^{-19} \text{ J}$$

$$\sigma = 4.0 \cdot 10^{-20} \text{ cm}^2$$

$$\tau = 300 \mu\text{s}$$

$$P_{out} = 320 \text{ mW} \quad L_i = 16\% \quad R_1 = 95\%$$



## 1. the photon lifetime

$$\tau_c = \frac{L_e}{\gamma c} \quad \gamma_1 = -\ln R_1 \cong 0.05 \quad \gamma_2 = -\ln R_2 \cong 0.2877$$

$$\gamma_i = -\ln(1 - L_i) \cong 0.1744 \quad \gamma = \gamma_i + \frac{\gamma_1 + \gamma_2}{2} = 0.3433$$

$$L_e = L + (n - 1)l = 54.32 \text{ cm}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\tau_c = \frac{L_e}{\gamma c} = 7.86 \cdot 10^{-9} \text{ s} = 7.86 \text{ ns}$$

$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 1054 \text{ nm} \quad h\nu = 1.89 \cdot 10^{-19} \text{ J}$$

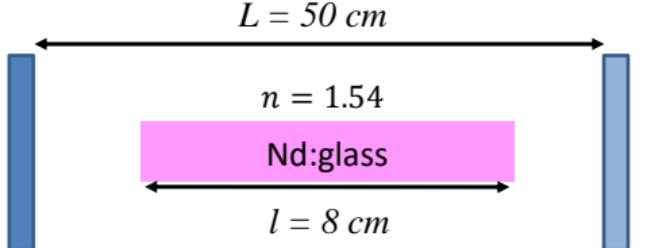
$$\sigma = 4.0 \cdot 10^{-20} \text{ cm}^2$$

$$\tau = 300 \mu\text{s}$$

$$P_{out} = 320 \text{ mW}$$

$$L_i = 16\%$$

$$R_1 = 95\%$$



$$R_2 = 75\%$$

## 2. the number of photons in the cavity

$$P_{out} = \left( \frac{\gamma_2 c}{2L_e} \right) h\nu\phi \quad \phi = \left( \frac{2L_e}{\gamma_2 c} \right) \frac{P_{out}}{h\nu} \cong 2.13 \cdot 10^{10}$$

$$L_e = L + (n - 1)l = 54.32 \text{ cm}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\gamma_2 = -\ln R_2 \cong 0.2877$$

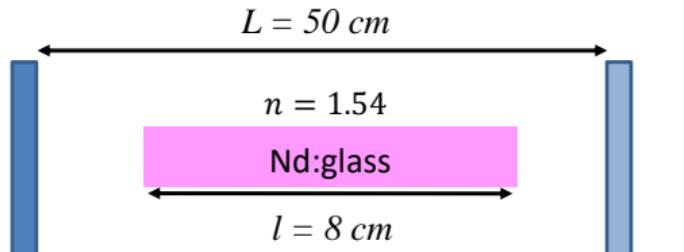
$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 1054 \text{ nm} \quad h\nu = 1.89 \cdot 10^{-19} \text{ J}$$

$$\sigma = 4.0 \cdot 10^{-20} \text{ cm}^2$$

$$\tau = 300 \mu\text{s}$$

$$P_{out} = 320 \text{ mW} \quad L_i = 16\% \quad R_1 = 95\%$$



$$R_2 = 75\%$$

# CW Nd:glass

3. the critical population inversion

$$N_t = 3.2 \cdot 10^{20} \frac{\text{ions}}{\text{cm}^3} \quad \frac{N_c}{N_t} = 0.3\%$$

$$N_c = \frac{\gamma}{\sigma l} = 1.07 \cdot 10^{18} \frac{\text{ions}}{\text{cm}^3}$$

$$\gamma = \gamma_i + \frac{\gamma_1 + \gamma_2}{2} = 0.3433$$

4. the saturation intensity

$$I_s = \frac{hv}{\sigma\tau} = 15.75 \frac{\text{kW}}{\text{cm}^2}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 1054 \text{ nm} \quad hv = 1.89 \cdot 10^{-19} \text{ J}$$

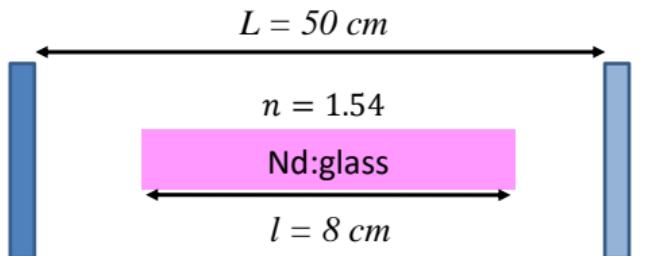
$$\sigma = 4.0 \cdot 10^{-20} \text{ cm}^2$$

$$\tau = 300 \mu\text{s}$$

$$P_{out} = 320 \text{ mW}$$

$$L_i = 16\%$$

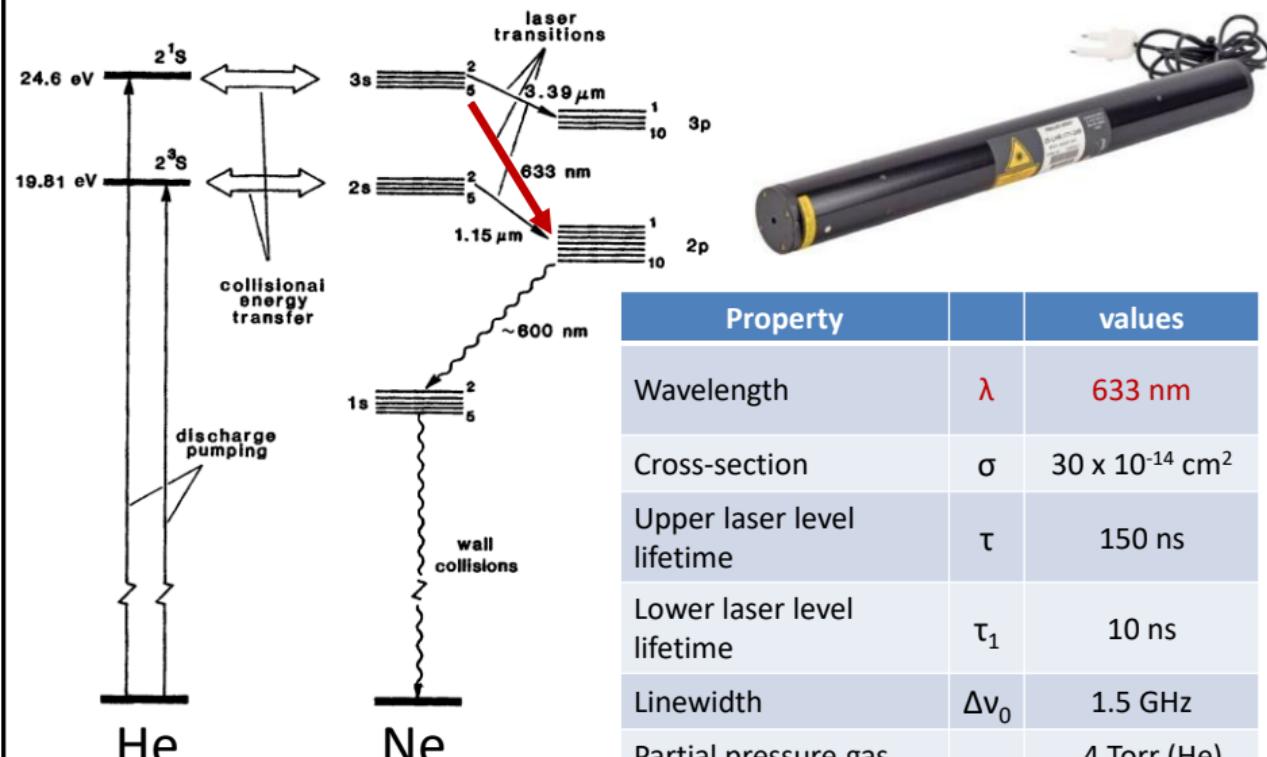
$$R_1 = 95\%$$



$$l = 8 \text{ cm}$$

$$R_2 = 75\%$$

## He-Ne laser



## He-Ne laser

A He-Ne laser ( $n \cong 1$ ) oscillating at the red line ( $\lambda = 633 \text{ nm}$ ) emits in cw an output power  $P_{out} = 30 \text{ mW}$ . The resonant cavity is a Fabry-Perot cavity of length  $L = 60 \text{ cm}$ , totally filled with the gas and made of a first mirror with reflectivity  $R_1 = 98\%$  and a second mirror (outcoupling mirror) with  $R_2 = 80\%$ .

Assuming that the internal losses for single pass are  $L_i = 12\%$ , determine:

1. the photon lifetime
2. the number of photons in the cavity
3. the critical population inversion
4. the saturation intensity

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 633 \text{ nm} \quad h\nu = 3.14 \cdot 10^{-19} \text{ J}$$

$$\sigma = 3.0 \cdot 10^{-13} \text{ cm}^2$$

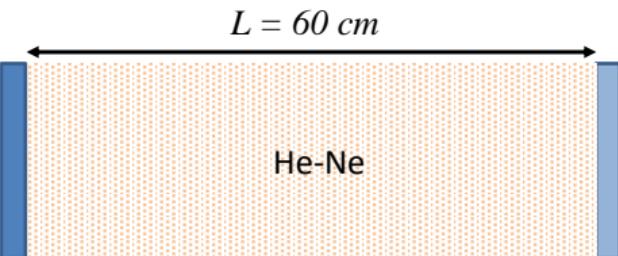
$$\tau = 150 \text{ ns}$$

$$P_{out} = 30 \text{ mW}$$

$$L_i = 12\%$$

$$R_1 = 98\%$$

$$R_2 = 80\%$$



## He-Ne laser

## 1. the photon lifetime

$$\tau_c = \frac{L_e}{\gamma c} \quad \gamma_1 = -\ln R_1 \cong 0.02 \quad \gamma_2 = -\ln R_2 \cong 0.223$$
$$\gamma_i = -\ln(1 - L_i) \cong 0.1278 \quad \gamma = \gamma_i + \frac{\gamma_1 + \gamma_2}{2} = 0.2493$$

$$L_e = nL = 60 \text{ cm}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\tau_c = \frac{L_e}{\gamma c} = 7.86 \cdot 10^{-9} \text{ s} \cong 8 \text{ ns}$$

$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 633 \text{ nm} \quad h\nu = 3.14 \cdot 10^{-19} \text{ J}$$

$$L = 60 \text{ cm}$$

$$\sigma = 3.0 \cdot 10^{-13} \text{ cm}^2$$



$$\tau = 150 \text{ ns}$$

$$P_{out} = 30 \text{ mW}$$

$$L_i = 12\%$$

$$R_1 = 98\%$$

$$R_2 = 80\%$$

## 2. the number of photons in the cavity

$$P_{out} = \left( \frac{\gamma_2 c}{2L_e} \right) h\nu\phi \quad \phi = \left( \frac{2L_e}{\gamma_2 c} \right) \frac{P_{out}}{h\nu} \cong 1.71 \cdot 10^9$$

$$L_e = nL = 60 \text{ cm}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\gamma_2 = -\ln R_2 \cong 0.223$$

$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 633 \text{ nm} \quad h\nu = 3.14 \cdot 10^{-19} \text{ J}$$

$$\sigma = 3.0 \cdot 10^{-13} \text{ cm}^2$$

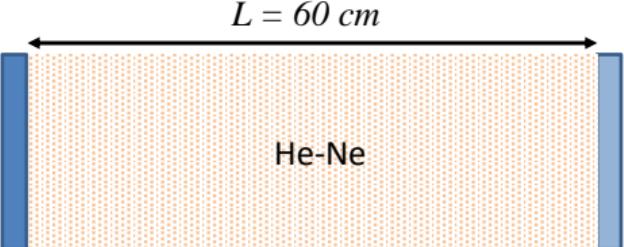
$$\tau = 150 \text{ ns}$$

$$P_{out} = 30 \text{ mW}$$

$$L_i = 12\%$$

$$R_1 = 98\%$$

$$R_2 = 80\%$$



### 3. the critical population inversion

$$N_c = \frac{\gamma}{\sigma l} = 1.385 \cdot 10^{10} \frac{\text{ions}}{\text{cm}^3}$$

$$\gamma = \gamma_i + \frac{\gamma_1 + \gamma_2}{2} = 0.2493$$

$$L = l = 60 \text{ cm}$$

### 4. the saturation intensity

$$I_s = \frac{h\nu}{\sigma\tau} = 6.98 \frac{\text{W}}{\text{cm}^2}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$c = 3 \cdot 10^8 \text{ m/s}$$

$$\lambda = 633 \text{ nm} \quad h\nu = 3.14 \cdot 10^{-19} \text{ J}$$

$$\sigma = 3.0 \cdot 10^{-13} \text{ cm}^2$$

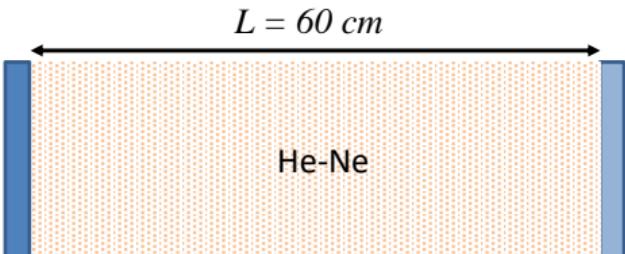
$$\tau = 150 \text{ ns}$$

$$P_{out} = 30 \text{ mW}$$

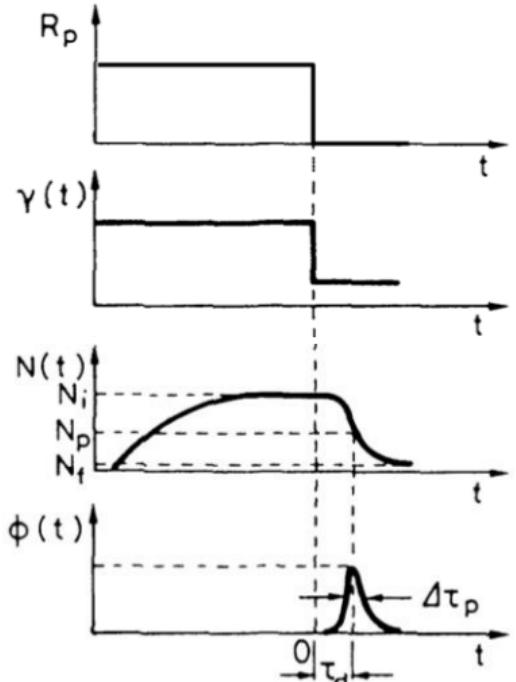
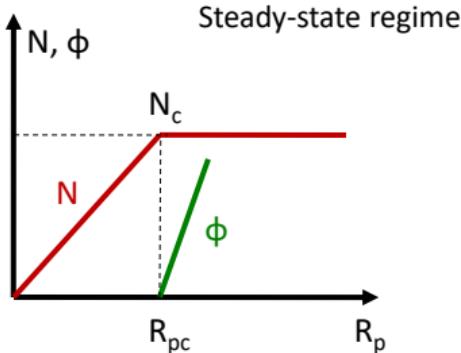
$$L_i = 12\%$$

$$R_1 = 98\%$$

$$R_2 = 80\%$$



## Q-switch

 $\Delta\tau_p = \text{pulse duration}$  $\tau_d = \text{build-up time}$ 

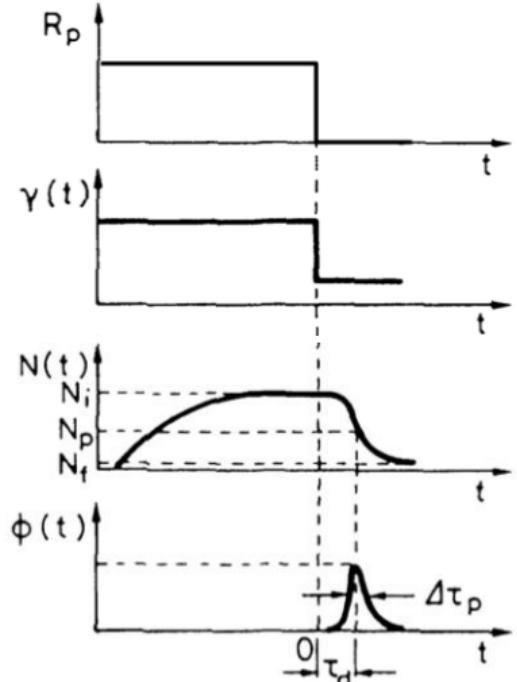
$$N_c = \frac{1}{V_a B \tau_c} = \frac{\gamma}{\sigma l}$$

 $R_p < R_{pc} \Rightarrow \phi = 0 \text{ No laser action!}$ 

$$N(t) = R_p \tau (1 - e^{-t/\tau}) \quad N_\infty = R_p \tau$$

$$Q = 2\pi\nu \tau_c \quad \tau_c = \frac{L_e}{c\gamma} \quad \text{Q-switch}$$

# Q-switch



$\Delta\tau_p$  = pulse duration

$\tau_d$  = build-up time

- **Active Q-switch**
- **Fast Q-switch**
- **Pulsed pumping**

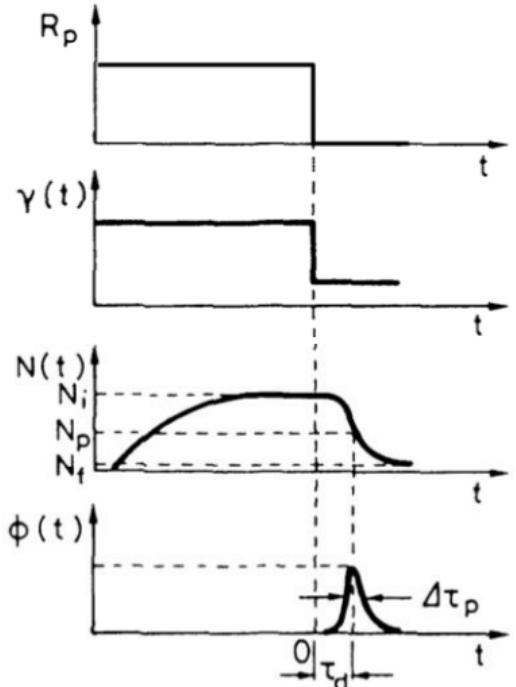
The duration of the **pump pulse** ( $t_p$ ) should be  $t_p \leq \tau$ :

for  $t_p \gg \tau$ ,  $N(t)$  doesn't differ much from its asymptotic value and the pump power would be lost through spontaneous decays instead of being accumulated as population inversion.

$$N(t) = R_p \tau (1 - e^{-t/\tau}) \quad N_\infty = R_p \tau$$

$$Q = 2\pi\nu \tau_c \quad \tau_c = \frac{L_e}{c\gamma} \quad \text{Q-switch}$$

# Q-switch



$\Delta\tau_p$  = pulse duration

$\tau_d$  = build-up time

- **Active Q-switch**
- **Fast Q-switch**
- **Pulsed pumping**

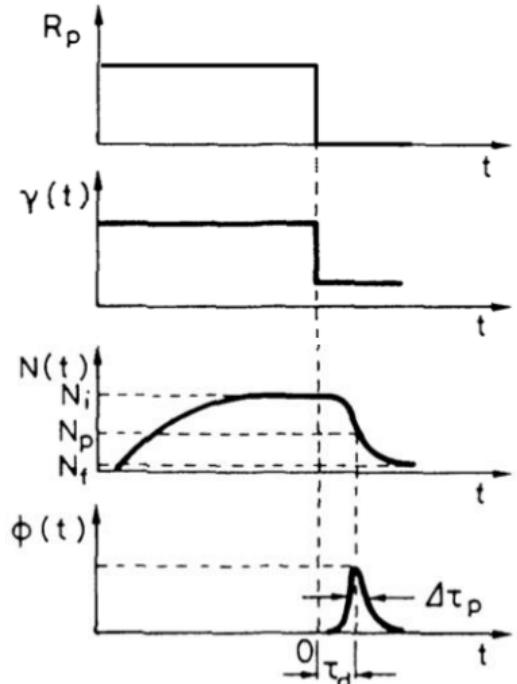
To obtain a sufficiently **large population inversion**, it is necessary to have **long  $\tau$  (ms)**.

This condition is well satisfied by solid-state lasers (Nd, Er, Yb...) and some gas lasers ( $\text{CO}_2$  or  $\text{I}_2$ ). For dye lasers or gas lasers as He-Ne or Ar  $\tau$  is of the order of few ns and the population inversion that can be obtained is too low to make it worth operating in Q-switch mode.

$$N(t) = R_p \tau (1 - e^{-t/\tau}) \quad N_\infty = R_p \tau$$

$$Q = 2\pi\nu \tau_c \quad \tau_c = \frac{L_e}{c\gamma} \quad \text{Q-switch}$$

## Q-switch

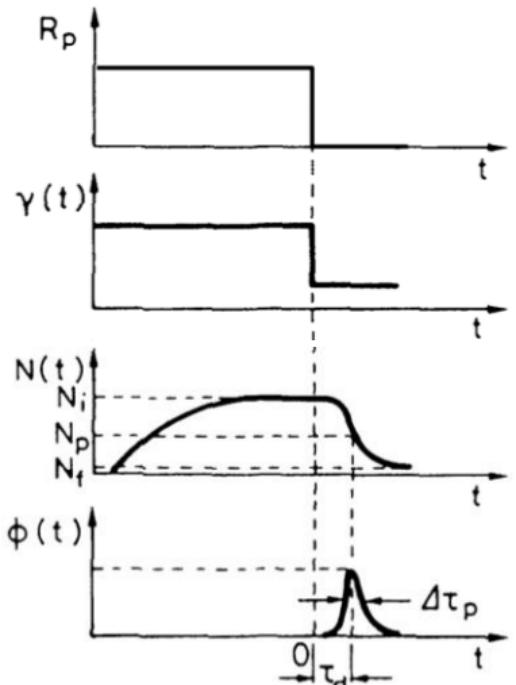
 $\Delta\tau_p$  = pulse duration $\tau_d$  = build-up time

- Active Q-switch
- Fast Q-switch
- Pulsed pumping
- Space-independent rate equations

Working hypotheses:

- Four-level laser:  $N_1 \cong N_3 \cong 0$   $\frac{dN_g}{dt} \cong 0$
- Single mode (longitudinal and transverse)
- Homogeneous broadening
- Uniform energy density of the mode
- Uniform pumping

# Q-switch



$\Delta\tau_p$  = pulse duration

$\tau_d$  = build-up time

$N \equiv N_2 - N_1 \cong N_2$  Population inversion

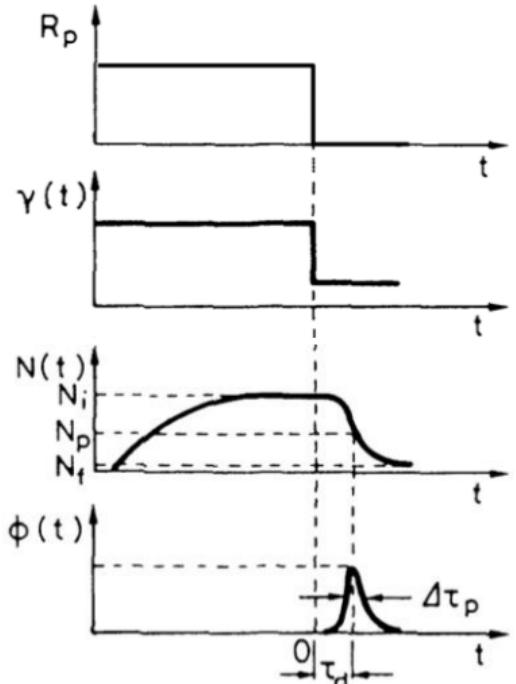
$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right.$$

$$N_i = R_p(0)\tau$$

$E_P$  = pump energy corresponding to a given pumping rate  $R_p$

$$E_P \propto \int R_p dt \implies E_P \propto R_p(0)$$

$$\implies E_P \propto N_i$$



$\Delta\tau_p$  = pulse duration

$\tau_d$  = build-up time

$N \equiv N_2 - N_1 \cong N_2$  Population inversion

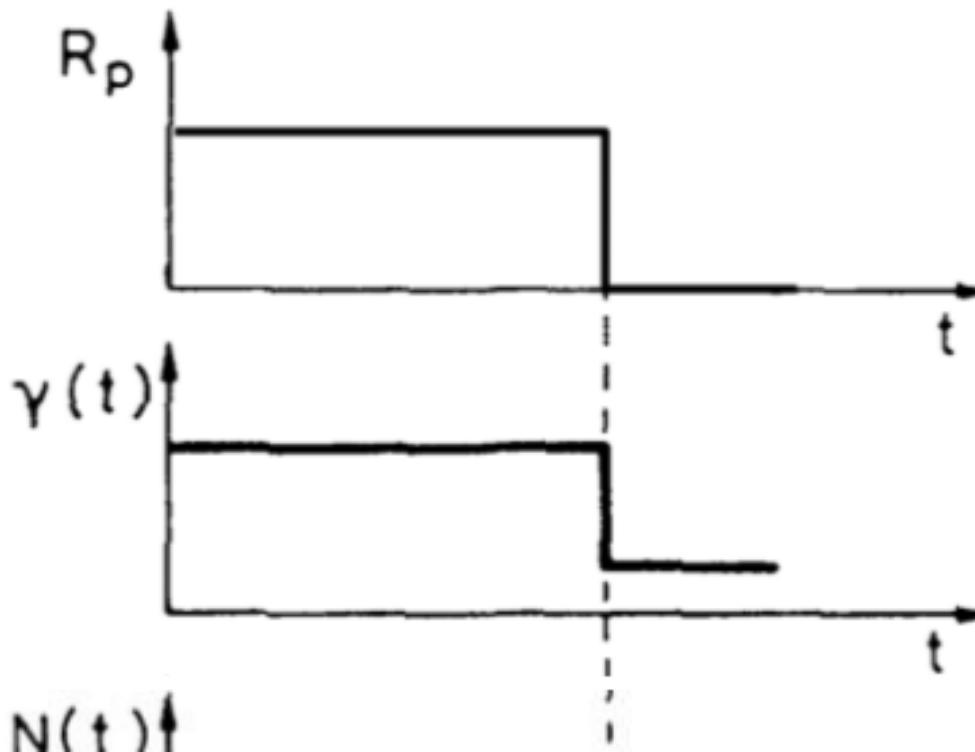
$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right.$$

$$x = \frac{E_p}{E_{PC}} = \frac{N_i}{N_c} \quad \text{over-threshold factor} \quad (*)$$

$$N_c = \frac{\gamma}{\sigma l} \quad \text{critical inversion for normal laser action (at Q-switch element open!)}$$

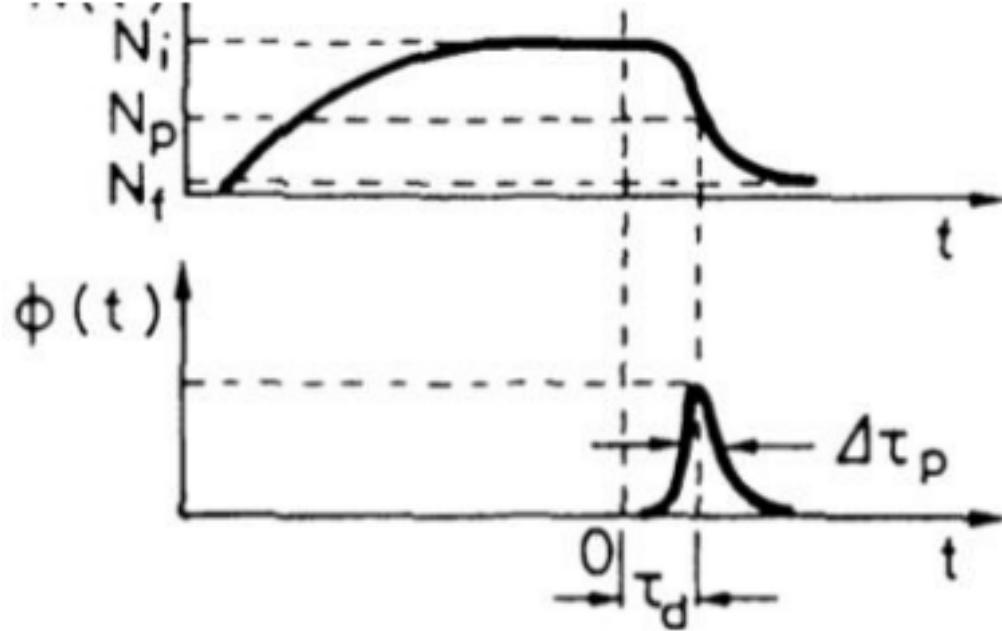
If  $N_c$  and  $x$  are known, equation (\*) allows us to calculate the initial inversion  $N_i$

# Q-switch



$N \equiv N_2 - N_1 \cong N_2$  Population inversion

$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right.$$



$\Delta\tau_p$  = pulse duration

$\tau_d$  = build-up time

$$N(0) = N_i$$

$\phi(0) = \phi_i \approx 1$  for laser action to start

For instantaneous Q-switch, the temporal evolution of  $N(t)$  and  $\phi(t)$  occurs on time scales so short that in the rate equations the terms related to the pumping rate  $R_p$  and the spontaneous emission  $\frac{N}{\tau}$  can be neglected