



Electrons vs photons

Propagation and confinement

$$\left\{ \begin{array}{l} \nabla \times \mathbf{E}(\mathbf{r}, t) = -\frac{\partial \mathbf{B}(\mathbf{r}, t)}{\partial t}, \\ \nabla \times \mathbf{H}(\mathbf{r}, t) = \frac{\partial \mathbf{D}(\mathbf{r}, t)}{\partial t} + \mathbf{j}(\mathbf{r}, t) \\ \nabla \cdot \mathbf{D}(\mathbf{r}, t) = \rho(\mathbf{r}, t), \\ \nabla \cdot \mathbf{B}(\mathbf{r}, t) = 0. \end{array} \right. \quad \begin{array}{l} \nabla \cdot \mathbf{j}(\mathbf{r}, t) + \frac{\partial \rho(\mathbf{r}, t)}{\partial t} = 0 \\ \mathbf{D}(\mathbf{r}, t) = \varepsilon_0 \mathbf{E}(\mathbf{r}, t) + \mathbf{P}(\mathbf{r}, t), \\ \mathbf{H}(\mathbf{r}, t) = \mu_0^{-1} \mathbf{B}(\mathbf{r}, t) - \mathbf{M}(\mathbf{r}, t) \end{array}$$

$$\begin{array}{ll} \mathbf{D} = \varepsilon_0 \varepsilon \mathbf{E} & (\mathbf{P} = \varepsilon_0 \chi_e \mathbf{E}), \\ \mathbf{B} = \mu_0 \mu \mathbf{H} & (\mathbf{M} = \chi_m \mathbf{H}), \\ \mathbf{j}_c = \sigma \mathbf{E}. & \end{array} \quad \begin{array}{l} \varepsilon_1 = n^2 - \kappa^2 \\ \varepsilon_2 = 2n\kappa \\ n^2 = \frac{\varepsilon_1}{2} + \frac{1}{2}\sqrt{\varepsilon_1^2 + \varepsilon_2^2} \\ \kappa = \frac{\varepsilon_2}{2n}. \end{array}$$

$$\mathbf{E} = \mathbf{E}_0 e^{ik \cdot \mathbf{r}} e^{-i\omega t}$$

Plane waves solutions

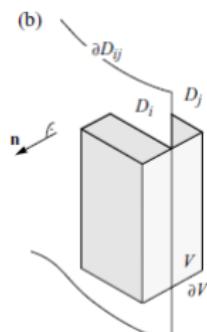
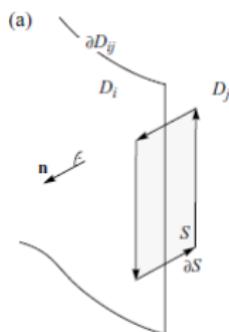
$$\left\{ \begin{array}{l} \nabla \times \mathbf{E}(\mathbf{r}) = i\omega \mathbf{B}(\mathbf{r}), \\ \nabla \times \mathbf{H}(\mathbf{r}) = -i\omega \mathbf{D}(\mathbf{r}) + \mathbf{j}(\mathbf{r}) \\ \nabla \cdot \mathbf{D}(\mathbf{r}) = \rho(\mathbf{r}), \\ \nabla \cdot \mathbf{B}(\mathbf{r}) = 0, \end{array} \right.$$

$$\nabla \times \nabla \times = -\nabla^2 + \nabla \nabla$$

$$\begin{aligned} (\nabla^2 + k_i^2) \mathbf{E}_i &= -i\omega \mu_0 \mu_i \mathbf{j}_i + \frac{\nabla \rho_i}{\varepsilon_0 \varepsilon_i} \\ (\nabla^2 + k_i^2) \mathbf{H}_i &= -\nabla \times \mathbf{j}_i, \end{aligned}$$

Boundary conditions

$$\partial D_{ij}$$



$$\left\{ \begin{array}{l} \mathbf{n} \times (\mathbf{E}_i - \mathbf{E}_j) = 0 \\ \mathbf{n} \times (\mathbf{H}_i - \mathbf{H}_j) = \mathbf{J}_{surf} \quad (\text{surf. curr. density}) \\ \mathbf{n} \cdot (\mathbf{D}_i - \mathbf{D}_j) = \sigma \quad (\text{surf. charge density}) \\ \mathbf{n} \cdot (\mathbf{B}_i - \mathbf{B}_j) = 0 \end{array} \right.$$

$$\left[\begin{array}{l} \nabla \times \mathbf{E}(\mathbf{r}) = i\omega \mathbf{B}(\mathbf{r}), \\ \nabla \times \mathbf{H}(\mathbf{r}) = -i\omega \mathbf{D}(\mathbf{r}) + \mathbf{j}(\mathbf{r}) \\ \nabla \cdot \mathbf{D}(\mathbf{r}) = \rho(\mathbf{r}), \\ \nabla \cdot \mathbf{B}(\mathbf{r}) = 0, \end{array} \right. \quad \left. \begin{array}{l} \nabla \times \nabla \times = -\nabla^2 + \nabla \nabla \\ (\nabla^2 + k_i^2) \mathbf{E}_i = -i\omega \mu_0 \mu_i \mathbf{j}_i + \frac{\nabla \rho_i}{\varepsilon_0 \varepsilon_i} \\ (\nabla^2 + k_i^2) \mathbf{H}_i = -\nabla \times \mathbf{j}_i, \end{array} \right.$$

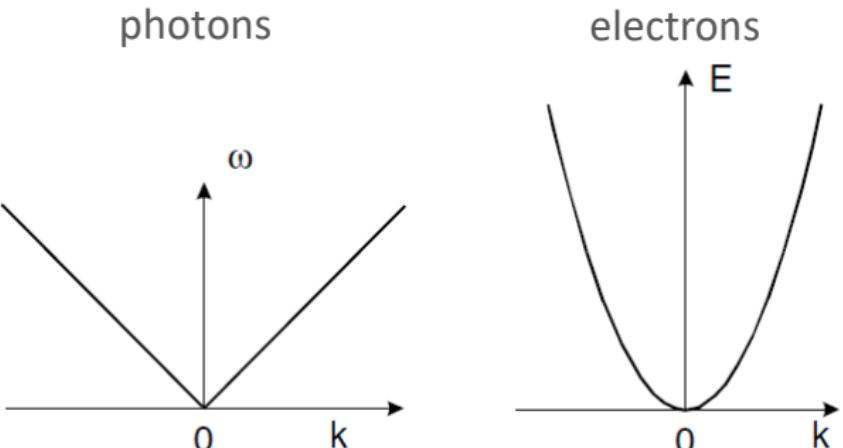
$$\mathbf{K}(\mathbf{K} \cdot \mathbf{E}) - K^2 \mathbf{E} = -\varepsilon(\mathbf{K}, \omega) \frac{\omega^2}{c^2} \mathbf{E},$$

For Transverse waves: $\mathbf{K} \cdot \mathbf{E} = 0$ $K^2 = \varepsilon(\mathbf{K}, \omega) \frac{\omega^2}{c^2}$ Dispersion Relation

For Longitudinal waves: $\mathbf{K} \parallel \mathbf{E}$ $\varepsilon(\mathbf{K}, \omega) = 0$

Photons vs Electrons

	photons	electrons
mass	0	m_0
charge	0	e
statistics	Bosons	Fermions
wavelength	$\lambda = \frac{h}{p} = \frac{c}{f}$	$\lambda = \frac{h}{p} = \frac{h}{mv}$
Eigenvalue Equation	$-\nabla^2 E = k^2 E$ $E = p_1 E$ $-\nabla^2 E = k^2 E$	$\left(-\frac{\hbar^2}{2m}\nabla^2 + V\right)\Psi = \varepsilon \Psi$ $-\nabla^2 \Psi = \left[\frac{2m}{\hbar^2}(\varepsilon - V)\right]\Psi = k^2 \Psi$
Free Space solution	$E = E_0 e^{ik \cdot r} e^{-i\omega t}$	$\Psi = \Psi_0 e^{ik \cdot r} e^{-i\omega t}$
Interaction	$\varepsilon(\omega, r)$	$V(r)$
Classically forbidden zone	Photon tunneling (evanescent wave) $\mathbf{k} \in \mathbb{C}$	Electron tunneling (exp. decay probability) $\mathbf{k} \in \mathbb{C}$
Band Structure	Photonic crystals (periodic modulation of $\varepsilon(\omega, r)$)	Crystals (periodic modulation of $V(r)$)
Cooperative effects	Nonlinear optics G.Mattei	Many-body correlations Cooper pairs


Phase velocity

$$v_p = \frac{\omega}{k}$$

Group velocity

$$|v_g| = |\bar{V}_k \omega| = \frac{1}{\hbar} |\bar{V}_k \varepsilon|$$

Non-dispersive

$$v_p = c$$

G.Mattei

Dispersive

$$v_p = \frac{\hbar}{2m} k$$

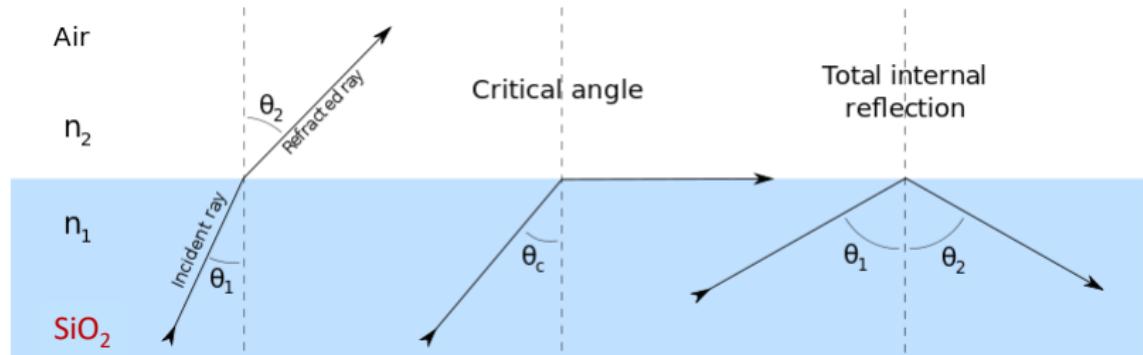
$$v_g = \frac{\hbar}{m} k = 2v_p$$

$$n_1 > n_2$$

Snell

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

$$\sin \theta_2 = \frac{n_1}{n_2} \sin \theta_1$$



$$\theta_2 = 90^\circ$$

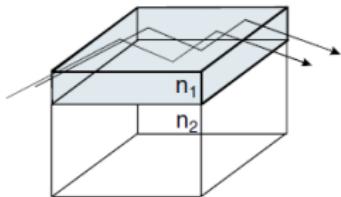
$$\sin \theta_1 = \sin \theta_c = \frac{n_2}{n_1} \sin \theta_2 = \frac{n_2}{n_1}$$

$$\theta_c = \sin^{-1} \left(\frac{n_2}{n_1} \right)$$

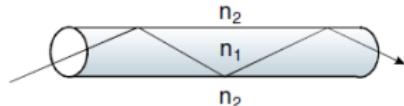
$$\theta_c (SiO_2 - air) = \sin^{-1} \left(\frac{1}{1.46} \right) = 43.23^\circ$$

Photon confinement

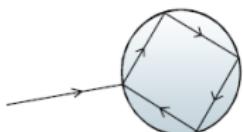
$$n_1 > n_2$$



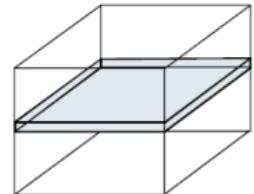
Optical planar waveguide



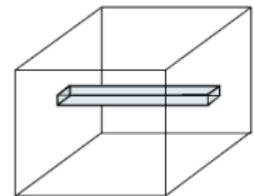
Optical fiber



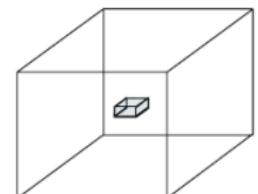
Microsphere optical cavity

1D**Electron confinement**

Quantum well



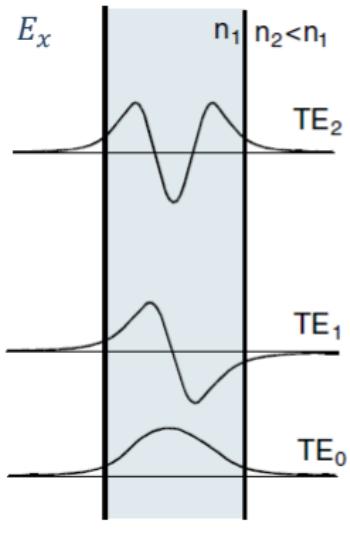
Quantum wire

2D

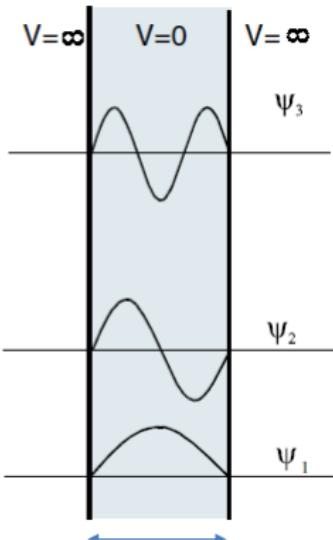
Quantum dot

3D

photons



electrons



Particle-in-a-box

$$\Psi_n(x) = \left(\frac{2}{l}\right)^{1/2} \sin\left(\frac{n\pi x}{l}\right)$$

$$k = n\pi/l \quad n = 1, 2, 3, \dots$$

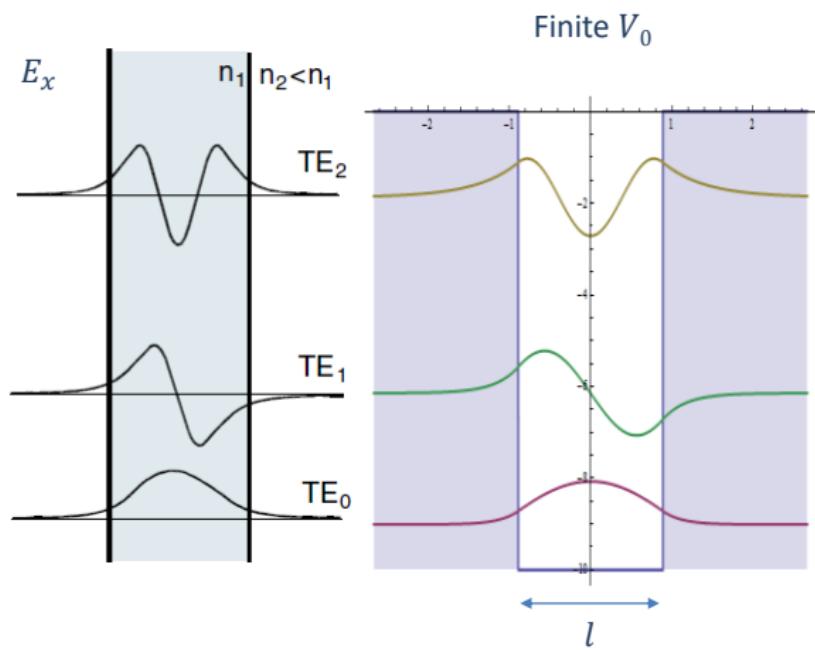
$$l = n\frac{\lambda}{2}$$

$$E_n = \frac{n^2 h^2}{8ml^2}$$

$$\Delta E = (2n + 1) \frac{h^2}{8ml^2}$$

photons

electrons


Particle-in-a-box

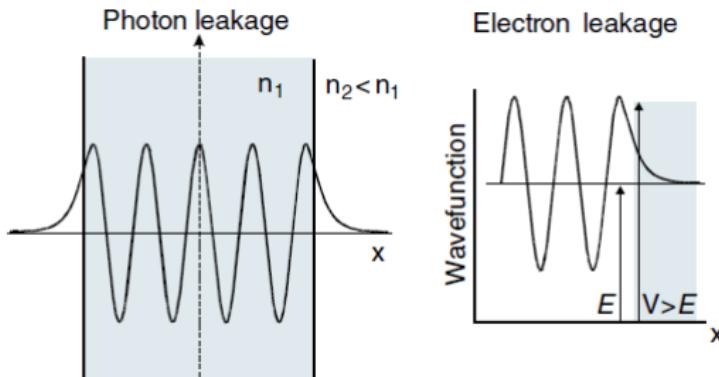
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$$l = n\frac{\lambda}{2}$$

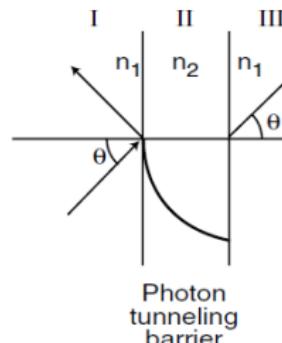
$$E_n = \frac{n^2 h^2}{8ml^2}$$

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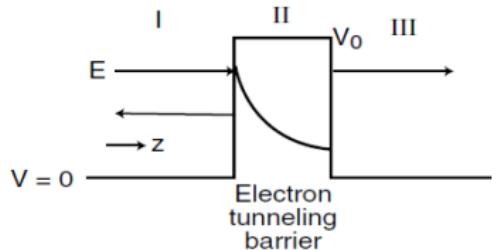


$$E = E_0 e^{-\frac{x}{d}}$$

Frustrated
Total Internal
Reflection

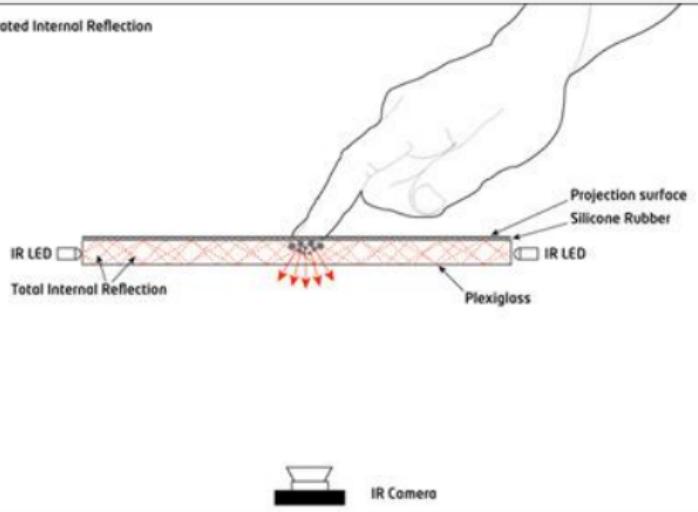


$$\Psi = \Psi_0 e^{-\frac{x}{d}}$$



FTIR and Multitouch Screen

FTIR - Frustrated Internal Reflection



(source: Cornell Univ.,US)



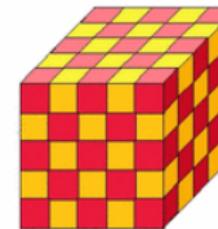
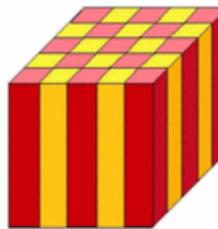
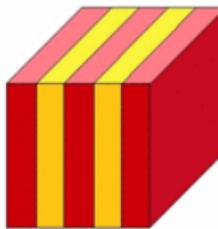
Photonic Crystals

Photonic Crystals

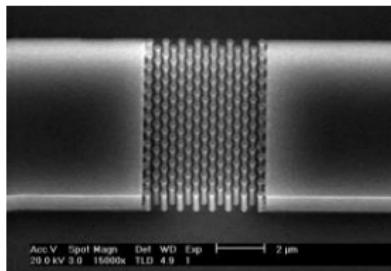
Definition:

A photonic crystal (PC) is a periodic structure in which the dielectric function $\varepsilon(\omega, \mathbf{r})$ varies periodically.

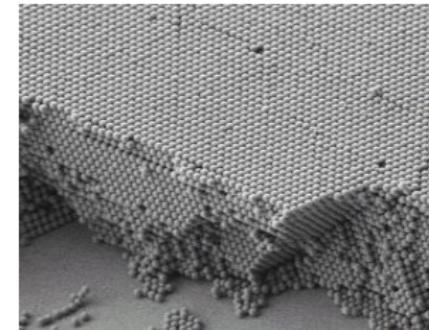
$$\varepsilon(\omega, \mathbf{r}) = \varepsilon(\omega, \mathbf{r} + \mathbf{R})$$



1D: Bragg Reflector



2D: Si pillar crystal



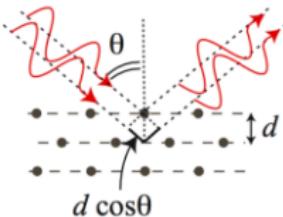
3D: colloidal crystal

1D Photonic Crystals

- Distributed Bragg Reflectors
- Antireflection coatings

Bragg's law:

$$2n_{average}d \cos \theta = m\lambda$$



$$\varepsilon(\omega, \mathbf{r}) = \varepsilon(\omega, \mathbf{r} + \mathbf{R})$$



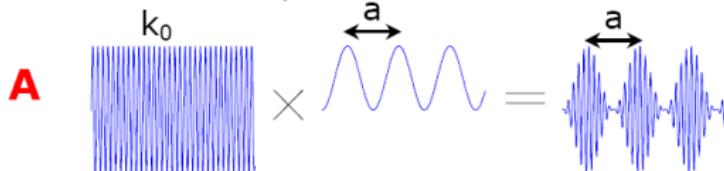
Direct space (Bravais lattice)



Reciprocal space (Bravais lattice)

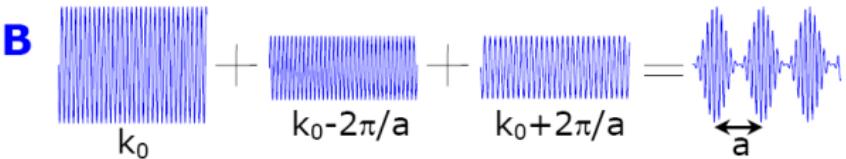
$$\psi_{n\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k} \cdot \mathbf{r}} u_{n\mathbf{k}}(\mathbf{r})$$

A Bloch wave is a plane wave with a modulation:

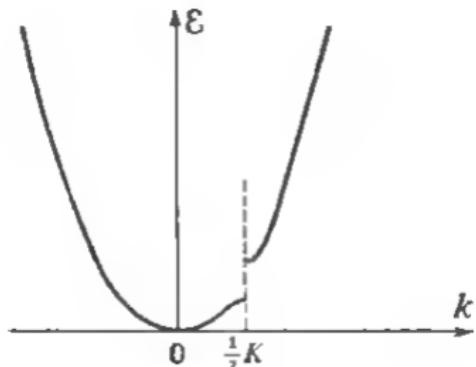
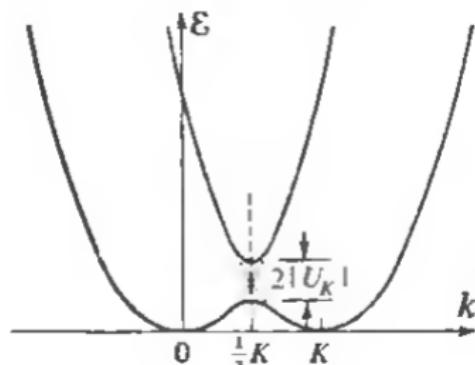
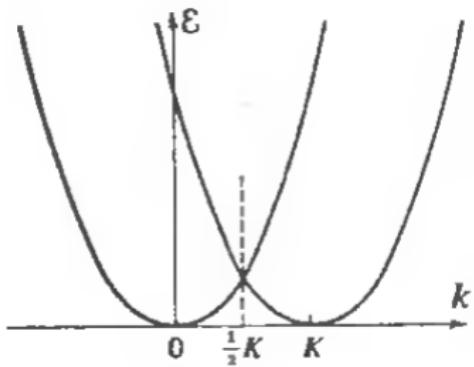
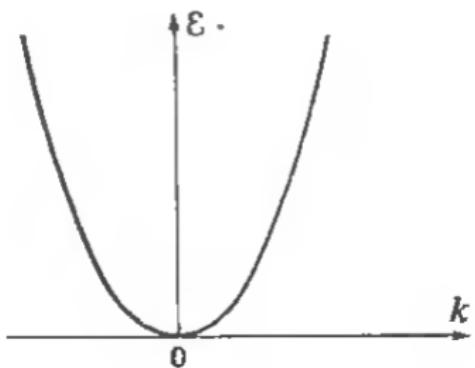


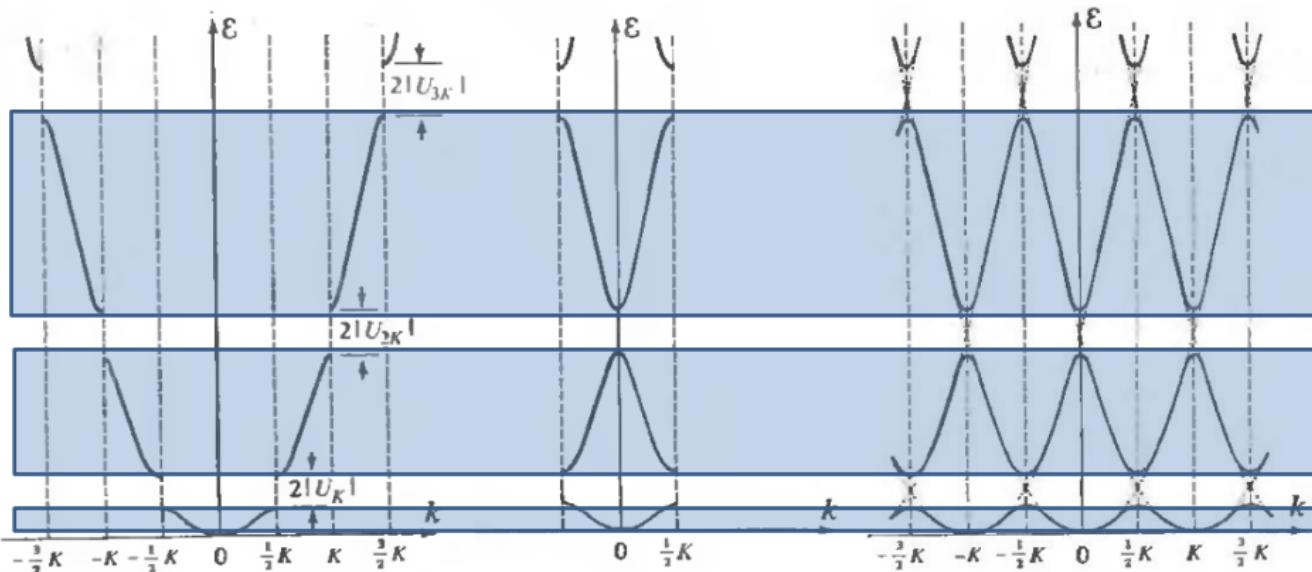
$$\psi(\mathbf{r}) = \sum_m c_m e^{i(k+m2\pi/a)\mathbf{r}}$$

A Bloch wave consists of multiple wavevectors:



Electrons in a weak Periodic Potential



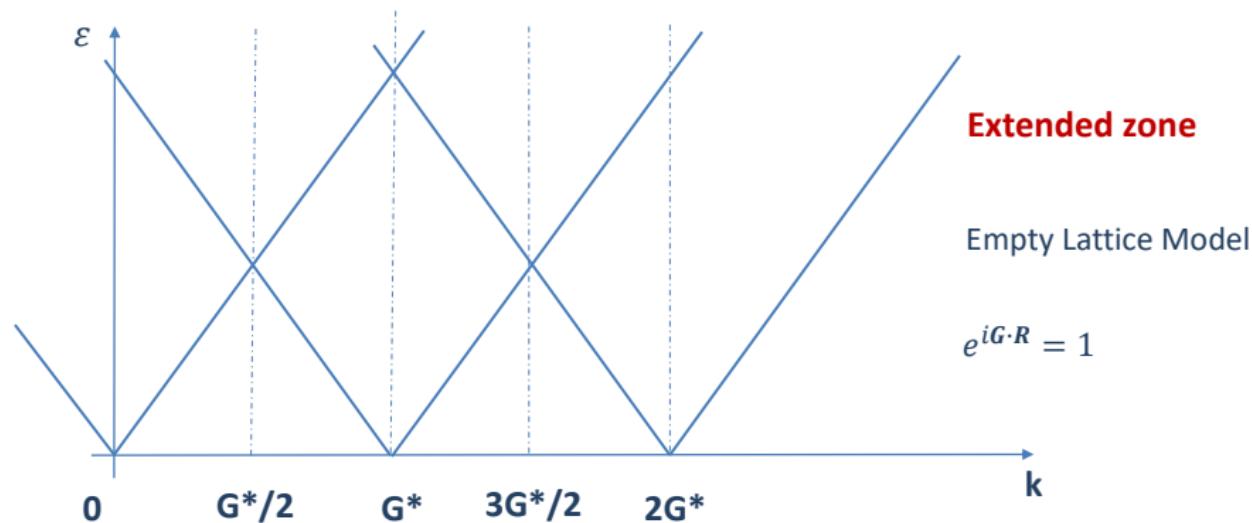
Extended zone
Reduced zone
Repeated zone


$$\epsilon = \frac{1}{2}(\epsilon_q^0 + \epsilon_{q-K}^0) \pm \left[\left(\frac{\epsilon_q^0 - \epsilon_{q-K}^0}{2} \right)^2 + |U_K|^2 \right]^{1/2}$$

$$U(\mathbf{r}) = U(\mathbf{r} + \mathbf{R}) = \sum_K U_K e^{i\mathbf{K} \cdot \mathbf{r}}$$

$$|\psi(\mathbf{r})|^2 \propto (\cos \frac{1}{2}\mathbf{K} \cdot \mathbf{r})^2, \quad \epsilon = \epsilon_q^0 + |U_K|,$$

$$|\psi(\mathbf{r})|^2 \propto (\sin \frac{1}{2}\mathbf{K} \cdot \mathbf{r})^2, \quad \epsilon = \epsilon_q^0 - |U_K|,$$



Remove degeneracy at $\mathbf{G}^*/2$ by modulating the dielectric function

Photonic band structure

$$\varepsilon(\omega, \mathbf{r}) = \varepsilon(\omega, \mathbf{r} + \mathbf{R})$$

$$\varepsilon(\omega, \mathbf{r}) = \sum_{\mathbf{G}} \varepsilon_{\mathbf{G}} e^{i\mathbf{G} \cdot \mathbf{r}}$$

$$\varepsilon_{\mathbf{G}} \equiv \frac{1}{V_c} \int_{V_c} d\mathbf{r} \varepsilon(\omega, \mathbf{r}) e^{-i\mathbf{G} \cdot \mathbf{r}}$$

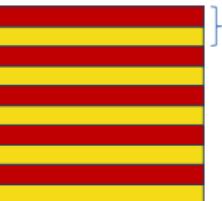


1D Photonic Crystals

$$\nabla^2 \mathbf{E} + \left(\frac{\omega}{c}\right)^2 \varepsilon \mathbf{E} = \mathbf{0}$$

$$e^{i\mathbf{G} \cdot \mathbf{R}} = 1$$

$$\left[\begin{array}{l} \varepsilon(\omega, \mathbf{r}) = \varepsilon(\omega, \mathbf{r} + \mathbf{R}) = \sum_{\mathbf{G}} \varepsilon_{\mathbf{G}} e^{i\mathbf{G} \cdot \mathbf{r}} \quad \varepsilon_{\mathbf{G}} \equiv \frac{1}{V_c} \int_{V_c} d\mathbf{r} \varepsilon(\omega, \mathbf{r}) e^{-i\mathbf{G} \cdot \mathbf{r}} \\ \mathbf{E}(\mathbf{r}) = \sum_{\mathbf{k}} \mathbf{A}(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{r}} \quad \mathbf{E}(\mathbf{r}) \sim \mathbf{A}(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{r}} + \mathbf{A}(\mathbf{k} - \mathbf{G}^*) e^{i(\mathbf{k} - \mathbf{G}^*) \cdot \mathbf{r}} \end{array} \right]$$



$$\left[\begin{array}{l} k^2 \mathbf{A}(\mathbf{k}) - \left(\frac{\omega}{c}\right)^2 \sum_{\mathbf{G}'} \varepsilon_{\mathbf{G}'} \mathbf{A}(\mathbf{k} - \mathbf{G}') = \mathbf{0} \\ (\mathbf{k} - \mathbf{G}^*)^2 \mathbf{A}(\mathbf{k} - \mathbf{G}^*) - \left(\frac{\omega}{c}\right)^2 \sum_{\mathbf{G}'} \varepsilon_{\mathbf{G}' - \mathbf{G}^*} \mathbf{A}(\mathbf{k} - \mathbf{G}') = \mathbf{0} \\ k^2 \mathbf{A}(\mathbf{k}) - \left(\frac{\omega}{c}\right)^2 \varepsilon_0 \mathbf{A}(\mathbf{k}) - \left(\frac{\omega}{c}\right)^2 \varepsilon_{\mathbf{G}} \mathbf{A}(\mathbf{k} - \mathbf{G}^*) = \mathbf{0} \\ (\mathbf{k} - \mathbf{G}^*)^2 \mathbf{A}(\mathbf{k} - \mathbf{G}^*) - \left(\frac{\omega}{c}\right)^2 \varepsilon_{-\mathbf{G}^*} \mathbf{A}(\mathbf{k}) - \left(\frac{\omega}{c}\right)^2 \varepsilon_0 \mathbf{A}(\mathbf{k} - \mathbf{G}^*) = \mathbf{0} \end{array} \right] \quad \mathbf{G}' = \mathbf{0}, \mathbf{G}^*$$

$$\begin{cases} \left(k^2 - \left(\frac{\omega}{c}\right)^2 \varepsilon_0 \right) \mathbf{A}(\mathbf{k}) - \left(\frac{\omega}{c}\right)^2 \varepsilon_{\mathbf{G}} \mathbf{A}(\mathbf{k} - \mathbf{G}^*) = \mathbf{0} \\ - \left(\frac{\omega}{c}\right)^2 \varepsilon_{-\mathbf{G}^*} \mathbf{A}(\mathbf{k}) + \left((\mathbf{k} - \mathbf{G}^*)^2 - \left(\frac{\omega}{c}\right)^2 \varepsilon_0 \right) \mathbf{A}(\mathbf{k} - \mathbf{G}^*) = \mathbf{0} \end{cases}$$

$$\left(\frac{\omega}{c}\right)^4 (\varepsilon_0^2 - \varepsilon_{\mathbf{G}^*} \varepsilon_{-\mathbf{G}^*}) - \left(\frac{\omega}{c}\right)^2 \varepsilon_0 (k^2 + (\mathbf{k} - \mathbf{G}^*)^2) + k^2 (\mathbf{k} - \mathbf{G}^*)^2 = 0$$

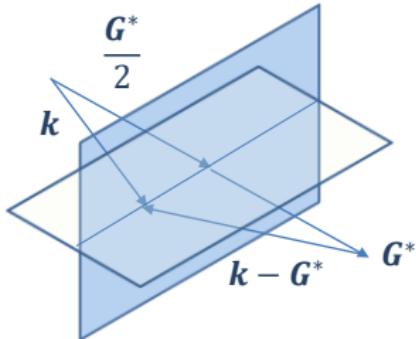
$$\left(\frac{\omega}{c}\right)^2 = \frac{\varepsilon_0 (k^2 + (\mathbf{k} - \mathbf{G}^*)^2)}{2(\varepsilon_0^2 - \varepsilon_{\mathbf{G}^*} \varepsilon_{-\mathbf{G}^*})} \pm \frac{\sqrt{(\varepsilon_0 (k^2 - (\mathbf{k} - \mathbf{G}^*)^2))^2 + 4(\varepsilon_{\mathbf{G}^*} \varepsilon_{-\mathbf{G}^*}) k^2 (\mathbf{k} - \mathbf{G}^*)^2}}{2(\varepsilon_0^2 - \varepsilon_{\mathbf{G}^*} \varepsilon_{-\mathbf{G}^*})}$$

1D Photonic Crystals

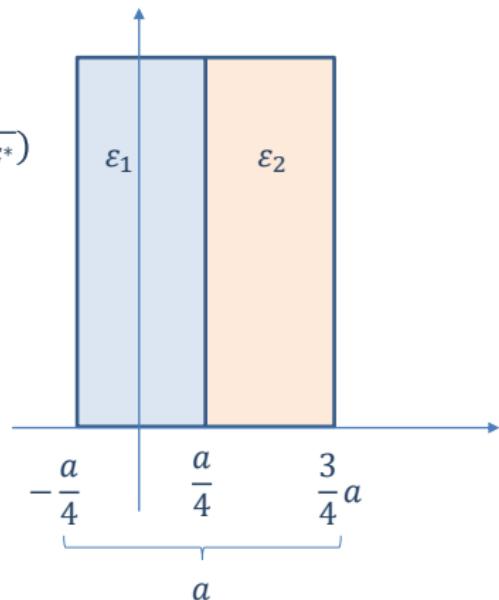
$$\mathbf{k} = \pm \frac{\mathbf{G}^*}{2}$$

$$\left(\frac{\omega}{c}\right)^2 = \frac{(\pi/a)^2}{(\varepsilon_0^2 - \varepsilon_{\mathbf{G}^*} \varepsilon_{-\mathbf{G}^*})} (\varepsilon_0 \pm \sqrt{\varepsilon_{\mathbf{G}^*} \varepsilon_{-\mathbf{G}^*}})$$

$$\omega_{\pm}^2 = \frac{\pi^2 c^2}{a^2 (\varepsilon_0 \mp |\varepsilon_{\mathbf{G}^*}|)}$$



$$\varepsilon(\omega, \mathbf{r}) = \varepsilon(\omega, -\mathbf{r}) \quad \varepsilon_G \in \mathbb{R}$$



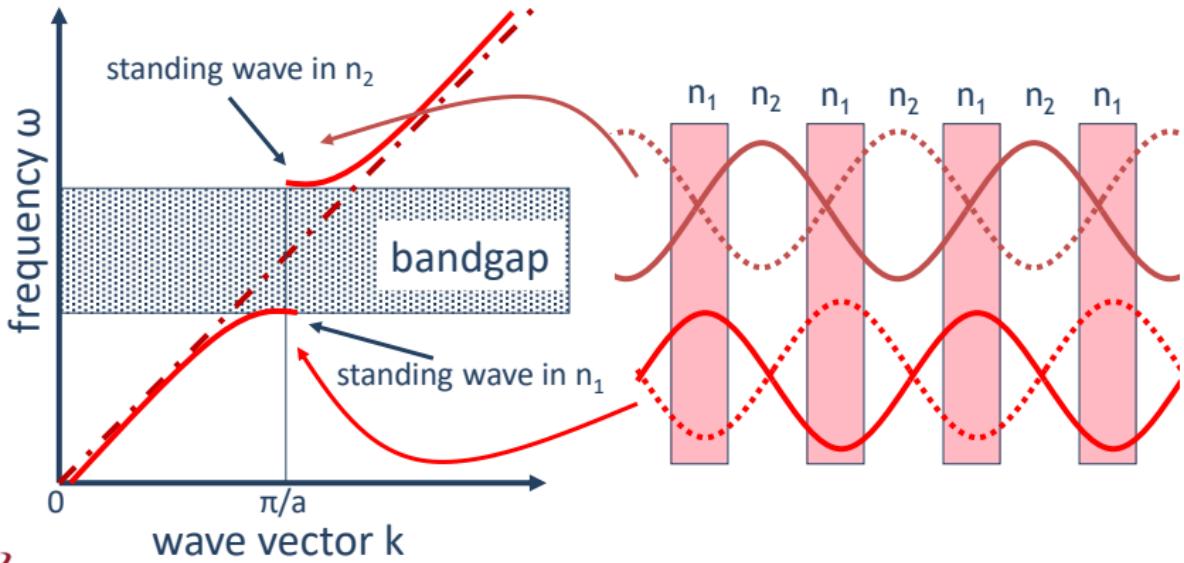
$$\begin{cases} \varepsilon_0 = \frac{\varepsilon_1 + \varepsilon_2}{2} \\ \varepsilon_G = \frac{\varepsilon_1 - \varepsilon_2}{\pi} \end{cases} \quad G = \frac{2\pi}{a}$$

Dispersion relation

$$\omega = \omega(k)$$

n_1 : high index material
 n_2 : low index material

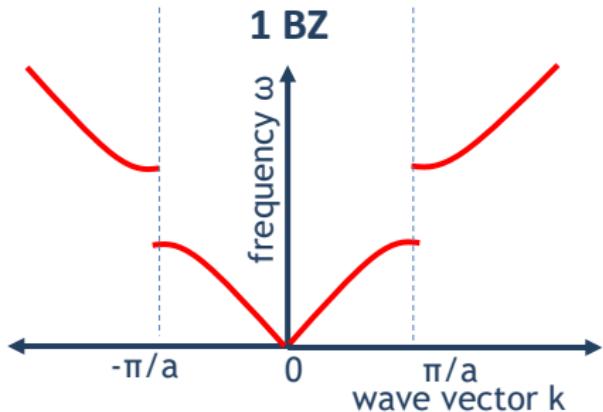
$$\omega = k \frac{c}{n}$$



Dispersion relation

$$\omega = \omega(k)$$

n_1 : high index material
 n_2 : low index material

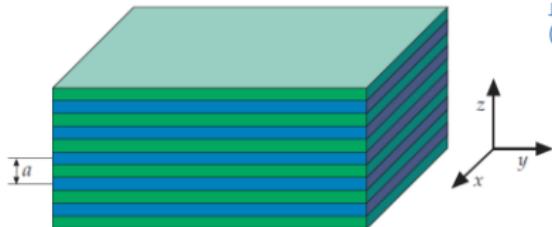


Bloch wave with wave vector \mathbf{k} is equal to Bloch wave with wave vector $\mathbf{k}+m2\pi/a$:

A light pulse with a center frequency ω propagates with *group velocity* v_g , which is altered near the band gap

$$v_g = \frac{d\omega}{dk}$$

1D Photonic Crystals

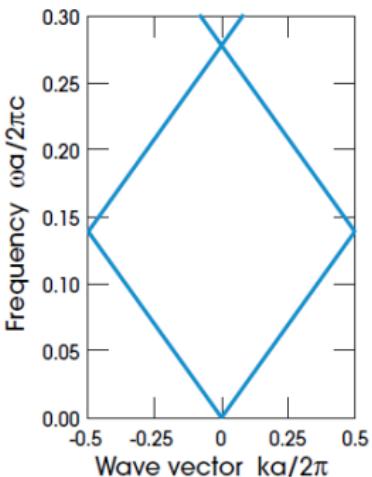


J D Joannopoulos et al. (2008), *Photonic Crystals: Molding the Flow of Light* (2nd ed.), Princeton NJ: Princeton University Press

$$\epsilon(\omega, z) = \epsilon(\omega, z + na)$$

$$d_1 = d_2 = 0.5 a$$

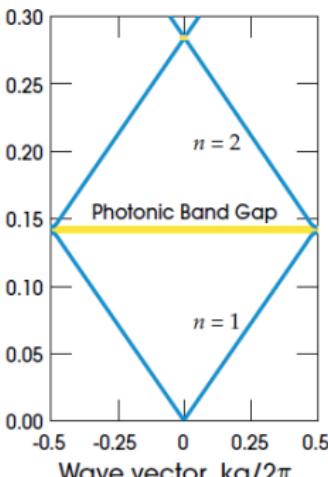
GaAs Bulk



$$\epsilon_1 = 13$$

$$\epsilon_2 = 13$$

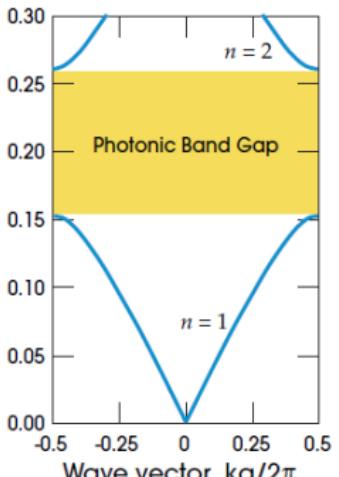
GaAs/GaAlAs Multilayer



$$\epsilon_1 = 13$$

$$\epsilon_2 = 12$$

GaAs/Air Multilayer

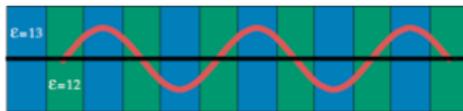
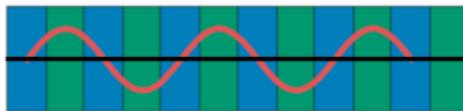


$$\epsilon_1 = 13$$

$$\epsilon_2 = 1$$

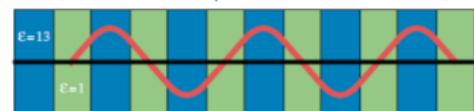
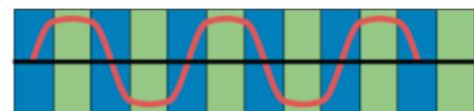
$$d_1 = d_2 = 0.5 a$$

$$\begin{aligned}\varepsilon_1 &= 13 \\ \varepsilon_2 &= 12\end{aligned}$$

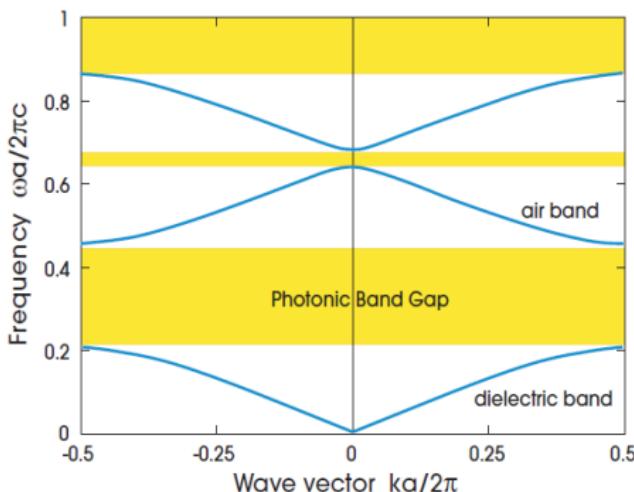
(a) E -field for mode at top of band 1(b) E -field for mode at bottom of band 2(c) Local energy density in E -field, top of band 1(d) Local energy density in E -field, bottom of band 2

J D Joannopoulos et al. (2008), *Photonic Crystals: Molding the Flow of Light* (2nd ed.), Princeton NJ: Princeton University Press

$$\begin{aligned}\varepsilon_1 &= 13 \\ \varepsilon_2 &= 1\end{aligned}$$

(a) E -field for mode at top of band 1(b) E -field for mode at bottom of band 2(c) Local energy density in E -field, top of band 1(d) Local energy density in E -field, bottom of band 2

$$\begin{aligned}\varepsilon_1 &= 13 & d_1 = d &= 0.2 a \\ \varepsilon_2 &= 1 & d_2 = a - d &= 0.8 a\end{aligned}$$



Perturbation Theory:

$$\Delta\omega = -\frac{\omega}{2} \frac{\int d^3\mathbf{r} \Delta\varepsilon(\mathbf{r}) |\mathbf{E}(\mathbf{r})|^2}{\int d^3\mathbf{r} \varepsilon(\mathbf{r}) |\mathbf{E}(\mathbf{r})|^2} + O(\Delta\varepsilon^2).$$

$$\left. \begin{array}{l} \varepsilon_1 = \varepsilon \\ \varepsilon_2 = \varepsilon + \Delta\varepsilon \\ \frac{d}{a} \ll 1 \end{array} \right\}$$

Gap-midgap ratio

$$\frac{\Delta\omega}{\omega_m} \approx \frac{\Delta\varepsilon}{\varepsilon} \cdot \frac{\sin(\pi d/a)}{\pi}$$

More refined theory:

$$\Delta\omega \text{ max if } d_1 n_1 = d_2 n_2.$$

$$\left. \begin{array}{l} \omega_m = \frac{n_1 + n_2}{4n_1 n_2} \cdot \frac{2\pi c}{a} \\ \frac{\Delta\omega}{\omega_m} = \frac{4}{\pi} \sin^{-1} \left(\frac{|n_1 - n_2|}{n_1 + n_2} \right) \end{array} \right\}$$

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