

$$N_1 \cong N_3 \cong 0$$

$$\frac{dN_g}{dt} \cong 0$$

$$N \equiv N_2 - N_1 \cong N_2$$

Population inversion

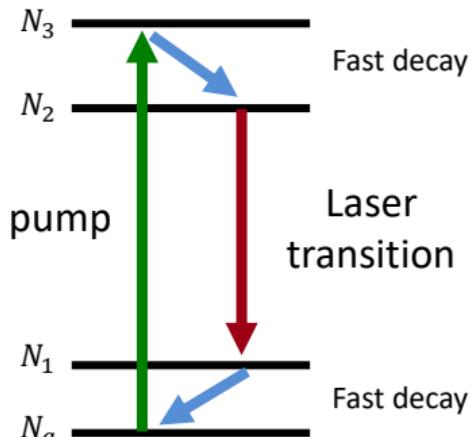
$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right. \quad (*)$$

Let's assume that at $t = 0$ an arbitrarily small number of photons (e.g., $\phi_i = 1$) is present within the cavity due to spontaneous emission.

From equation (**) we get that to have **laser action** (i.e., the amplification of the number of photons in the cavity) it should be:

$$\frac{d\phi}{dt} \geq 0 \Rightarrow V_a BN \geq \frac{1}{\tau_c} \Rightarrow \text{Laser action starts when it is reached a critical population inversion}$$

CW behavior



$$N_1 \cong N_3 \cong 0$$

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Population inversion

$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right. (*)$$

$$\left\{ \begin{array}{l} \frac{d\phi}{dt} = 0 \Rightarrow N_C = \frac{1}{V_a B \tau_c} = \frac{\gamma}{\sigma l} \\ \frac{d\phi}{dt} \geq 0 \Rightarrow V_a B N \geq \frac{1}{\tau_c} \end{array} \right. (**)$$

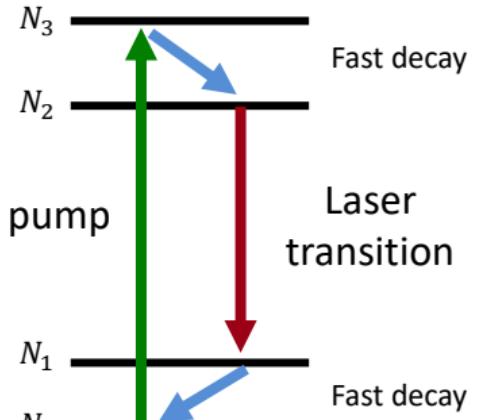
$$B = \frac{\sigma lc}{V_a L_e} \quad \tau_c = \frac{L_e}{\gamma c}$$

$$\frac{d\phi}{dt} = 0 \Rightarrow N_C = \frac{1}{V_a B \tau_c} = \frac{\gamma}{\sigma l}$$

$$\frac{d\phi}{dt} \geq 0 \Rightarrow V_a B N \geq \frac{1}{\tau_c} \Rightarrow$$

Laser action starts when it is reached
a **critical population inversion**

CW behavior



$$N_1 \cong N_3 \cong 0$$

$$\frac{dN_g}{dt} \cong 0$$

$$N \equiv N_2 - N_1 \cong N_2$$

Population inversion

$$\frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \quad (*)$$

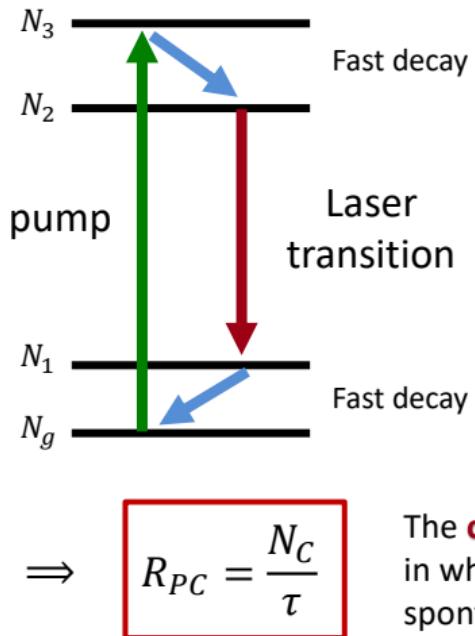
$$\frac{d\phi}{dt} = V_a B \phi N - \frac{\phi}{\tau_c} \quad (**)$$

$$\frac{d\phi}{dt} = 0 \quad \Rightarrow \quad N_C = \frac{1}{V_a B \tau_c} = \frac{\gamma}{\sigma l}$$

$$B = \frac{\sigma lc}{V_a L_e} \quad \tau_c = \frac{L_e}{\gamma c}$$

The corresponding **critical pumping rate** (R_{PC}) is obtained from (*) imposing:

$$\frac{dN}{dt} = 0 \quad (\text{steady-state}) \quad N = N_C \quad \phi = 0$$



$$N_1 \cong N_3 \cong 0 \quad \frac{dN_g}{dt} \cong 0$$

$$N \equiv N_2 - N_1 \cong N_2 \quad \text{Population inversion}$$

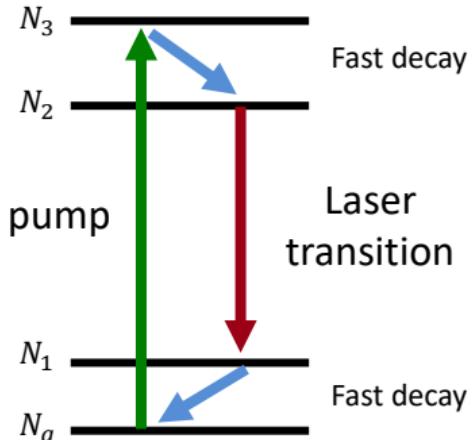
$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right. \quad (*) \quad (**)$$

The **critical pumping rate** corresponds to the situation in which the rate of pump transitions is equal to the spontaneous transitions from level 2.

The corresponding **critical pumping rate** (R_{PC}) is obtained from (*) imposing:

$$\frac{dN}{dt} = 0 \quad (\text{steady-state}) \quad N = N_C \quad \phi = 0$$

CW behavior



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Population inversion

$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right. \quad \begin{array}{l} (*) \\ (**) \end{array}$$

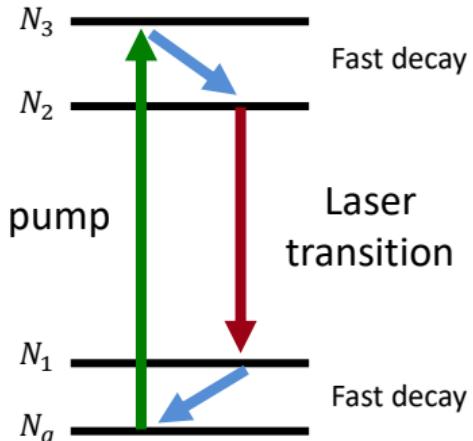
$R_p > R_{PC}$ The number of photons ϕ increases from its initial value determined by spontaneous emission.

If R_p is independent of time, ϕ reaches a steady-state value ϕ_0 which corresponds to a steady-state value of population inversion N_0 .

ϕ_0 and N_0 can be determined imposing:

$$\frac{d\phi}{dt} = 0 = \frac{dN}{dt}$$

CW behavior



$$N_1 \cong N_3 \cong 0 \quad \frac{dN_g}{dt} \cong 0$$

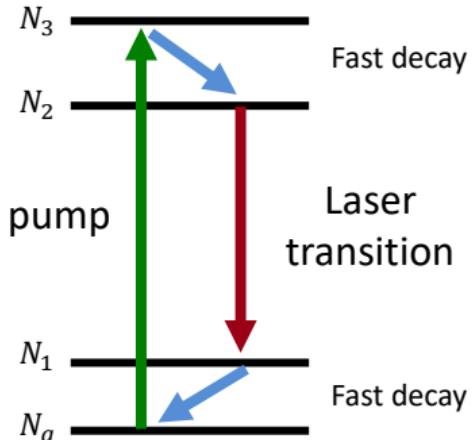
$N \equiv N_2 - N_1 \cong N_2$ **Population inversion**

$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right. \quad (*)$$

$$R_p > R_{PC} \quad \frac{d\phi}{dt} = 0 \quad \Rightarrow \quad N_0 = \frac{1}{V_a B \tau_c} = \frac{\gamma}{\sigma l} = N_C !$$

$$\frac{dN}{dt} = 0 \quad \Rightarrow \quad \phi_0 = V_a \tau_c \left[R_p - \frac{N_0}{\tau} \right]$$

CW behavior



$$N_1 \cong N_3 \cong 0$$

$$\frac{dN_g}{dt} \cong 0$$

$$N \equiv N_2 - N_1 \cong N_2$$

Population inversion

$$\left\{ \begin{array}{l} \frac{dN}{dt} = R_p - B\phi N - \frac{N}{\tau} \\ \frac{d\phi}{dt} = V_a B\phi N - \frac{\phi}{\tau_c} \end{array} \right. \quad \begin{array}{l} (*) \\ (**) \end{array}$$

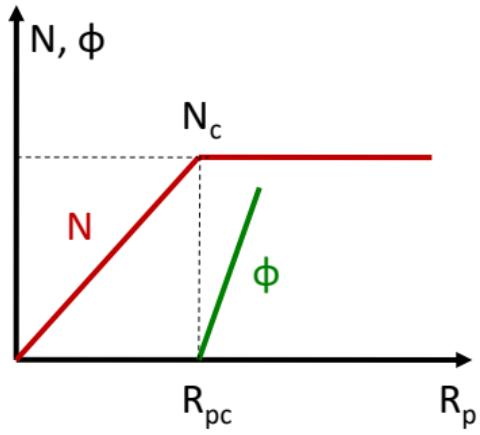
$R_P < R_{PC}$ Below threshold $\Rightarrow \phi = 0$ No laser action!

$$\Rightarrow \frac{dN}{dt} = R_p - \frac{N}{\tau} \Rightarrow \frac{dN}{N - R_p \tau} = -\frac{1}{\tau} dt \Rightarrow \int \frac{dN}{N - R_p \tau} = \int_0^t -\frac{1}{\tau} dt'$$

$$\Rightarrow N(t) = R_p \tau (1 - e^{-t/\tau}) \Rightarrow t \rightarrow \infty \quad N_\infty = R_p \tau \quad \text{Linear function of } R_p$$

CW behavior

Steady-state regime



$$x = \frac{R_p}{R_{PC}} \quad \text{over-threshold factor}$$

$$R_p < R_{PC} \implies \begin{cases} N = R_p \tau \\ \phi = 0 \end{cases}$$

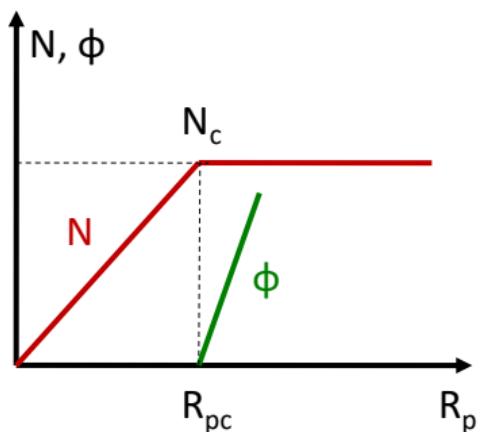
$$R_p = R_{PC} \implies \begin{cases} N = N_c = \frac{\gamma}{\sigma l} \\ \phi = 0 \end{cases}$$

$$R_p > R_{PC} \implies \begin{cases} N = N_0 = N_c \\ \phi_0 = V_a \tau_c \left[R_p - \frac{N_0}{\tau} \right] \end{cases}$$

The pumping rate increases population inversion **below threshold (energy stored in the active medium)**, while it increases the number of photons in the cavity **above threshold (energy stored in the cavity)**.

CW behavior

Steady-state regime



$$\phi_0 = V_a \tau_c \left[R_p - \frac{N_0}{\tau} \right] = V_a N_0 \frac{\tau_c}{\tau} \left[\frac{R_p \tau}{N_0} - 1 \right]$$

$$R_{PC} = \frac{N_C}{\tau} = \frac{N_0}{\tau} \quad N_C = \frac{\gamma}{\sigma l} = N_0$$

$$\phi_0 = V_a N_0 \frac{\tau_c}{\tau} (x - 1) \quad A_b = \frac{V_a}{l}$$

transverse section of the mode (beam area)

$$x = \frac{R_p}{R_{PC}} \quad \text{over-threshold factor}$$

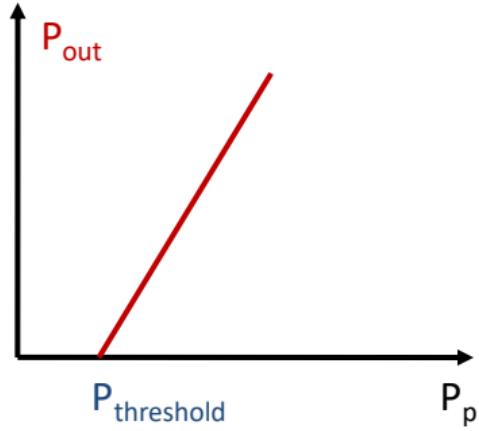
$$x = \frac{P_p}{P_{th}} \quad A_b < A = \frac{V}{l}$$

section of the active medium

$$\phi_0 = \frac{A_b \gamma}{\sigma} \frac{\tau_c}{\tau} \left(\frac{P_p}{P_{th}} - 1 \right) \quad \text{Remembering that}$$

$$P_{out} = \left(\frac{\gamma_2 c}{2 L_e} \right) h \nu \phi$$

Steady-state regime



$$P_{out} = \frac{\gamma_2}{2} \frac{c}{L_e} \frac{A_b \gamma}{\sigma} \frac{\tau_c}{\tau} \left(\frac{P_p}{P_{th}} - 1 \right)$$

$$I_s = \frac{h\nu}{\sigma\tau} \quad \text{Saturation intensity for a 4-level system}$$

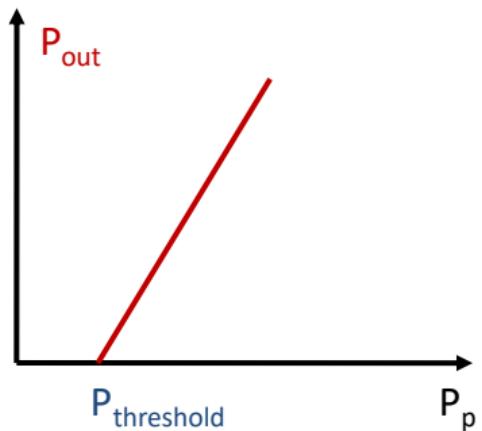
$$\tau_c = \frac{L_e}{\gamma c} \quad \text{Photon lifetime}$$

$$P_{out} = A_b I_s \frac{\gamma_2}{2} \left(\frac{P_p}{P_{th}} - 1 \right)$$

$$\phi_0 = \frac{A_b \gamma \tau_c}{\sigma \tau} \left(\frac{P_p}{P_{th}} - 1 \right) \quad \text{Remembering that} \quad P_{out} = \left(\frac{\gamma_2 c}{2 L_e} \right) h \nu \phi$$

CW behavior

Steady-state regime



$$P_{out} = \frac{\gamma_2}{2} \frac{c}{L_e} \frac{A_b \gamma}{\sigma} \frac{\tau_c}{\tau} \left(\frac{P_p}{P_{th}} - 1 \right)$$

$$I_s = \frac{h\nu}{\sigma\tau} \quad \text{Saturation intensity for a 4-level system}$$

$$\tau_c = \frac{L_e}{\gamma c} \quad \text{Photon lifetime}$$

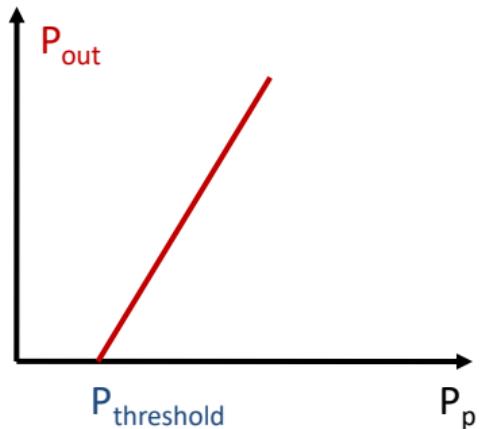
$$P_{out} = A_b I_s \frac{\gamma_2}{2} \left(\frac{P_p}{P_{th}} - 1 \right)$$

Laser slope efficiency

$$\eta_s = \frac{dP_{out}}{dP_p} \quad \Rightarrow \quad \eta_s = \frac{A_b h\nu}{\sigma\tau} \frac{\gamma_2}{2} \frac{1}{P_{th}}$$

Remembering that for optical and electrical pumping: $R_p = \eta_p \left(\frac{P_p}{Alh\nu_{mp}} \right)$

Steady-state regime



$$R_{pc} = \frac{N_C}{\tau} = \frac{\gamma}{\sigma l \tau} = \eta_p \left(\frac{P_{th}}{Alh\nu_{mp}} \right)$$



$$P_{th} = \left(\frac{\gamma}{\eta_p} \right) \left(\frac{h\nu_{mp}}{\tau} \right) \left(\frac{A}{\sigma} \right)$$



$$\eta_s = \eta_p \left(\frac{\gamma_2}{2\gamma} \right) \left(\frac{h\nu}{h\nu_{mp}} \right) \left(\frac{A_b}{A} \right)$$



$$\eta_s = \eta_p \eta_c \eta_q \eta_T$$

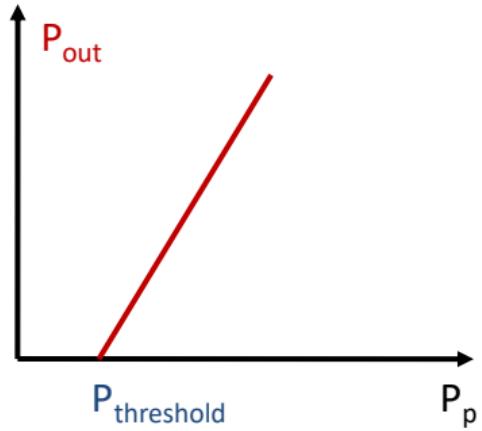
Laser slope efficiency

$$\eta_s = \frac{dP_{out}}{dP_p} \quad \Rightarrow \quad \eta_s = \frac{A_b h \nu}{\sigma \tau} \frac{\gamma_2}{2} \frac{1}{P_{th}}$$

Remembering that for optical and electrical pumping: $R_p = \eta_p \left(\frac{P_p}{Alh\nu_{mp}} \right)$

CW behavior

Steady-state regime



$$\eta_p = \text{pumping efficiency}$$

$$\eta_c = \frac{\gamma_2}{2\gamma} = \text{outcoupling efficiency}$$

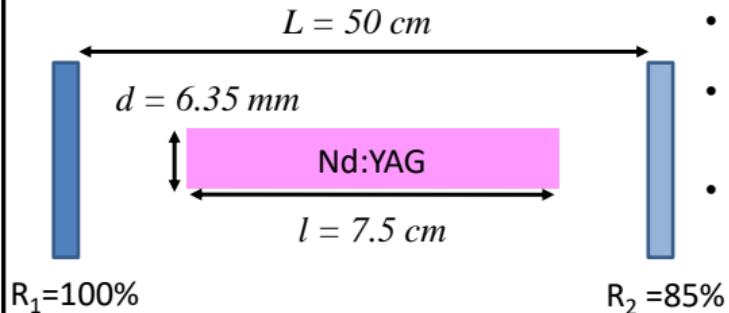
$$\eta_c = \frac{hv}{hv_{mp}} = \text{laser quantum efficiency}$$

$$\eta_T = \frac{A_b}{A} = \text{transverse efficiency}$$

Laser slope efficiency

$$\eta_s = \eta_p \left(\frac{\gamma_2}{2\gamma} \right) \left(\frac{hv}{hv_{mp}} \right) \left(\frac{A_b}{A} \right) = \eta_p \eta_c \eta_q \eta_T$$

CW laser: Nd:YAG

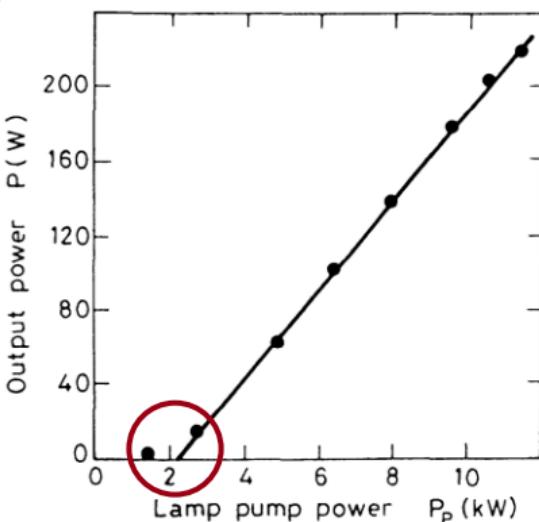


- Optical pumping with elliptical cavity
- High output power in cw:
 $\sim 100 - 200 \text{ W}$
- Multimodal oscillation:
longitudinal and transverse

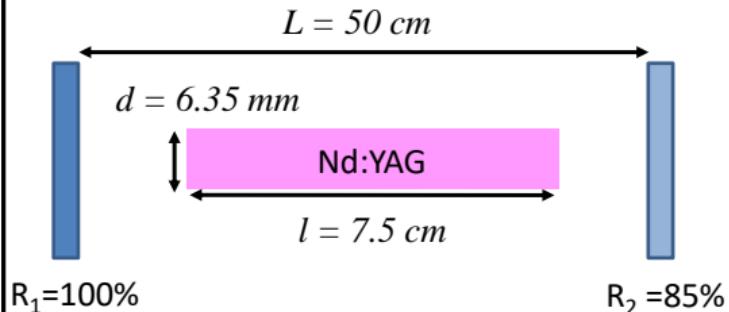
$$P_{out}(W) = 53 \left(\frac{P_p}{P_{th}} - 1 \right) \quad \text{output power}$$

$P_{th} = 2.2 \text{ kW}$ threshold power

$$\eta_s = \frac{dP_{out}}{dP_p} = \frac{53 \text{ W}}{2.2 \text{ kW}} = 2.4\% \quad \text{laser slope efficiency}$$



CW laser: Nd:YAG



$$\lambda = 1.064 \mu m$$

$$h\nu = 1.87 \cdot 10^{-19} J$$

$$\sigma = 2.8 \cdot 10^{-19} cm^2$$

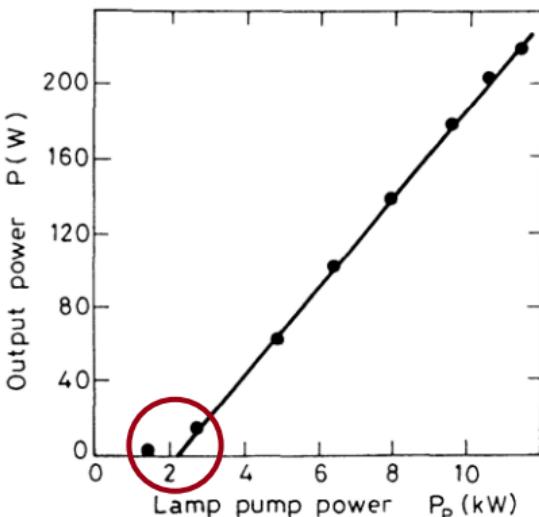
$$\tau = 230 \mu s$$

$$P_{out}(W) = 53 \left(\frac{P_p}{P_{th}} - 1 \right) \quad \text{output power}$$

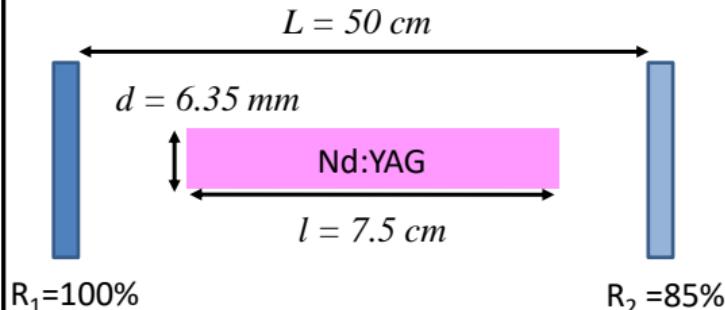
$$P_{out} = A_b I_s \frac{\gamma_2}{2} \left(\frac{P_p}{P_{th}} - 1 \right)$$

$$I_s = \frac{h\nu}{\sigma\tau} = 2.9 \frac{kW}{cm^2} \quad \gamma_2 = -\ln R_2 = 0.162$$

$$A_b \cong 0.23 cm^2 < A = 0.317 cm^2$$



CW laser: Nd:YAG



$$P_{th} = 2.2 \text{ kW}$$

$$\gamma_2 = -\ln R_2 \\ = 0.162$$

$$P_{th} = \left(\frac{\gamma}{\eta_p} \right) \left(\frac{h\nu_{mp}}{\tau} \right) \left(\frac{A}{\sigma} \right)$$

To compare the experimental values of P_{th} and η_s it is necessary to know γ .

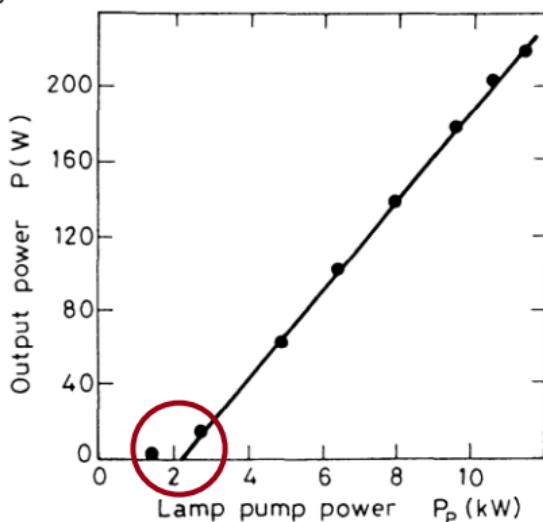
$$\gamma = \gamma_i + \frac{\gamma_1 + \gamma_2}{2} \quad \gamma_1 = -\ln R_1 = 0$$

$$\lambda = 1.064 \mu\text{m}$$

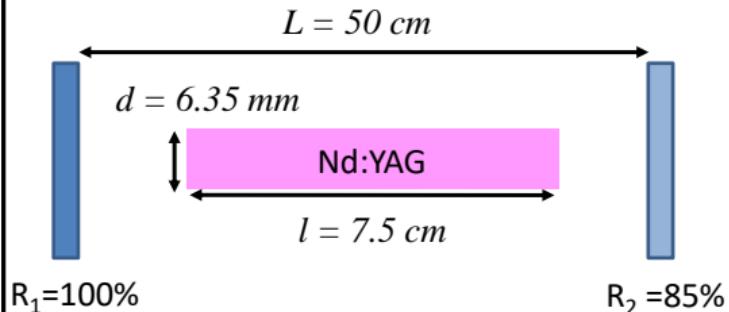
$$h\nu = 1.87 \cdot 10^{-19} \text{ J}$$

$$\sigma = 2.8 \cdot 10^{-19} \text{ cm}^2$$

$$\tau = 230 \mu\text{s}$$



CW laser: Nd:YAG



$$\lambda = 1.064 \mu\text{m}$$

$$hv = 1.87 \cdot 10^{-19} \text{ J}$$

$$\sigma = 2.8 \cdot 10^{-19} \text{ cm}^2$$

$$\tau = 230 \mu\text{s}$$

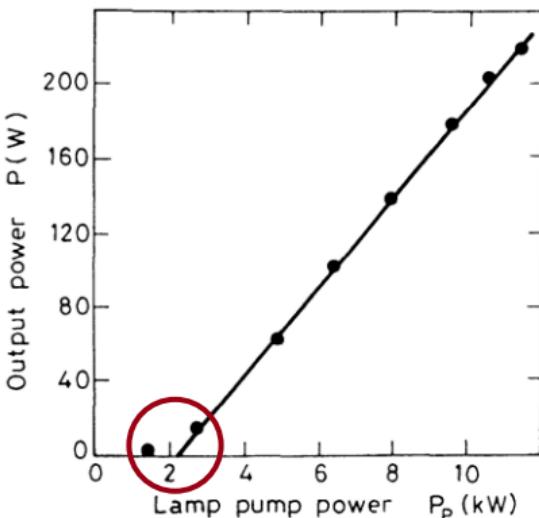
$$P_{th} = 2.2 \text{ kW}$$

$$\gamma_2 = -\ln R_2 \\ = 0.162$$

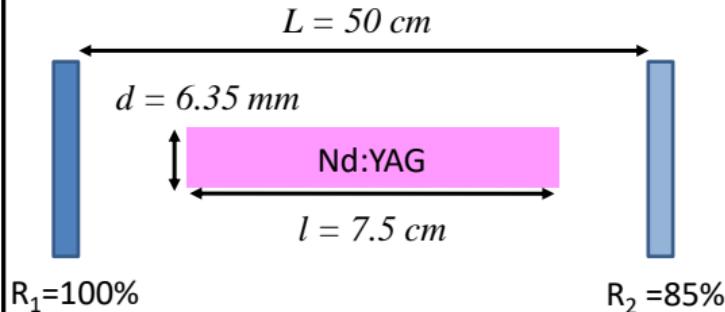
$$P_{th} = \left(\frac{\gamma}{\eta_p} \right) \left(\frac{h\nu_{mp}}{\tau} \right) \left(\frac{A}{\sigma} \right)$$

$$\gamma = \eta_p \left(\frac{\sigma}{A} \right) \left(\frac{\tau}{h\nu_{mp}} \right) P_{th} = \frac{-\ln R_2}{2} + \gamma_i$$

$$-\ln R_2 = -2\gamma_i + 2\eta_p \left(\frac{\sigma}{A} \right) \left(\frac{\tau}{h\nu_{mp}} \right) P_{th}$$



CW laser: Nd:YAG



$$P_{th} = 2.2 \text{ kW}$$

$$\gamma_2 = -\ln R_2 \\ = 0.162$$

$$P_{th} = \left(\frac{\gamma}{\eta_p} \right) \left(\frac{h\nu_{mp}}{\tau} \right) \left(\frac{A}{\sigma} \right)$$

$$\gamma = \eta_p \left(\frac{\sigma}{A} \right) \left(\frac{\tau}{h\nu_{mp}} \right) P_{th} = \frac{-\ln R_2}{2} + \gamma_i$$

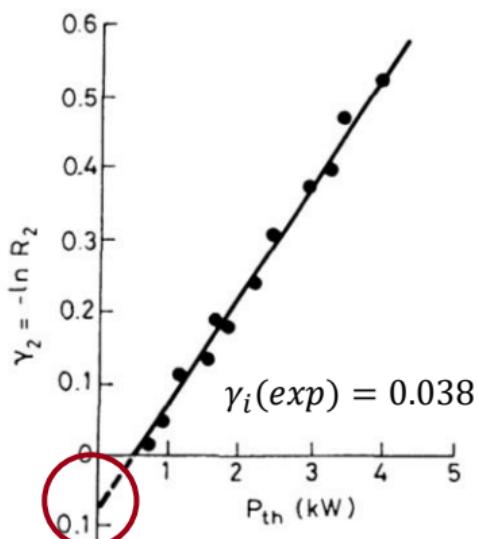
$$-\ln R_2 = \boxed{-2\gamma_i} + 2\eta_p \left(\frac{\sigma}{A} \right) \left(\frac{\tau}{h\nu_{mp}} \right) P_{th}$$

$$\lambda = 1.064 \mu\text{m}$$

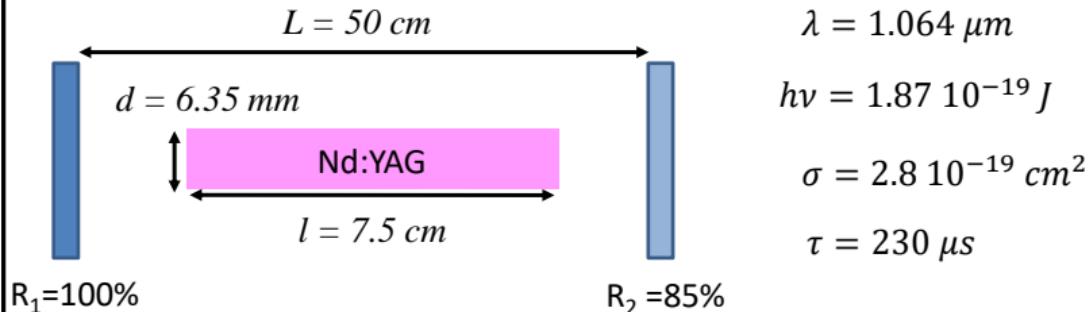
$$h\nu = 1.87 \cdot 10^{-19} \text{ J}$$

$$\sigma = 2.8 \cdot 10^{-19} \text{ cm}^2$$

$$\tau = 230 \mu\text{s}$$



CW laser: Nd:YAG



$$\eta_s(\text{exp}) = 2.4\%$$

$$\eta_s = \eta_p \left(\frac{\gamma_2}{2\gamma} \right) \left(\frac{h\nu}{h\nu_{mp}} \right) \left(\frac{A_b}{A} \right) = \eta_p \eta_c \eta_q \eta_T$$

$$\eta_c = \frac{\gamma_2}{2\gamma} \cong 0.68$$

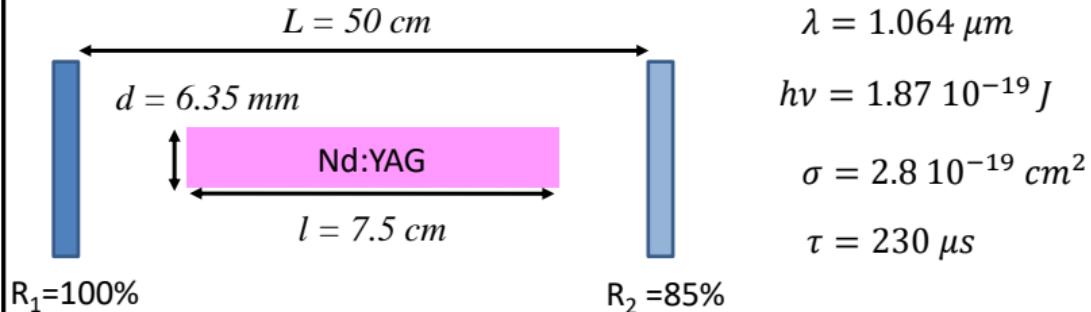
$$\gamma_i(\text{exp}) = 0.038$$

$$\gamma = \gamma_i + \frac{\gamma_1 + \gamma_2}{2} \cong 0.12$$

$$\gamma_2 = 0.162$$

$$\eta_q = \frac{h\nu}{h\nu_{mp}} = \frac{\lambda_{mp}}{\lambda} \cong 0.89 \quad \lambda_{mp} = 0.94 \mu m \quad h\nu_{mp} = 2.11 10^{-19} J$$

$$\eta_T = \frac{A_b}{A} \cong \frac{0.23 \text{ cm}^2}{0.317 \text{ cm}^2} \cong 0.72 \quad \rightarrow \quad \eta_p = \frac{\eta_s}{\eta_c \eta_q \eta_T} \cong 5.5\%$$



$$N_C = \frac{\gamma}{\sigma l} = 5.7 \cdot 10^{16} \frac{\text{ions}}{\text{cm}^3} \quad \gamma \cong 0.12$$

$$N_t = 1.38 \cdot 10^{20} \frac{\text{ions}}{\text{cm}^3} \quad \text{Total population (Nd doping: 1%)}$$

$$\frac{N_C}{N_t} = 4.4 \cdot 10^{-4}$$

The population inversion is a very small fraction of the total population!