



Drude Model

Free electron gas

$$\bar{E} = \bar{E}_0 e^{-i\omega t}$$

$$m\frac{d^2\bar{r}}{dt^2} + m\Gamma\frac{d\bar{r}}{dt} + K\bar{r} = -e\bar{E}_0e^{-i\omega t}$$

$$m\frac{d^2\bar{r}}{dt^2} + m\Gamma\frac{d\bar{r}}{dt} = -e\bar{E}_0e^{-i\omega t}$$
 Free

$$\bar{r}(t) = \frac{e\bar{E}_0}{m} \frac{1}{\omega^2 + i\omega\Gamma} e^{-i\omega t} = \frac{e\bar{E}_0}{m} \frac{\omega^2 + i\omega\Gamma}{\omega^4 + (\omega\Gamma)^2} e^{-i\omega t}$$





Drude Model

Free electron gas

$$\bar{p}(t) \equiv -e\bar{r}(t)$$

$$\bar{P} \equiv n\bar{p} = n\varepsilon_0 \alpha \bar{E} = \varepsilon_0 \chi \bar{E}$$

$$\bar{D} = \varepsilon_0 \varepsilon \bar{E} = \varepsilon_0 \bar{E} + \bar{P} = \varepsilon_0 \bar{E} + \varepsilon_0 \chi \bar{E} = \varepsilon_0 (1 + \chi) \bar{E}$$

$$\varepsilon = 1 + \chi = 1 + n\alpha = 1 - \frac{\omega_P^2}{\omega^2 + i\omega\Gamma} = 1 - \frac{\omega_P^2}{\omega^2 + \Gamma^2} + i\frac{\omega_P^2\Gamma}{\omega(\omega^2 + \Gamma^2)}$$

$$\omega_P^2 \equiv \frac{ne^2}{m\varepsilon_0}$$
 Bulk plasmon frequency $\hbar\omega_P \approx 5 - 20eV$



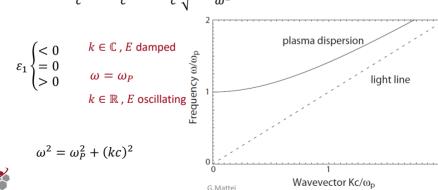
Drude Model

High-frequency limit: $\omega \gg \Gamma$

$$\varepsilon = 1 + \chi \approx 1 - \frac{\omega_P^2}{\omega^2} + i \frac{\omega_P^2 \Gamma}{\omega^3}$$

$$\varepsilon_1 \approx 1 - \frac{\omega_P^2}{\omega^2}$$
 $\varepsilon_2 \approx \frac{\omega_P^2 \Gamma}{\omega^3}$

$$k = k_0 n = \frac{\omega}{c} \sqrt{\varepsilon} \approx \frac{\omega}{c} \sqrt{\varepsilon_1} = \frac{\omega}{c} \sqrt{1 - \frac{\omega_P^2}{\omega^2}}$$





Drude Model

Relationship between ω_P and ω_{SPR} (high-freq. approx: $\omega \gg \Gamma$)

1. To excite ω_P (after Drude)

$$\varepsilon_1(\omega_P)=0$$

2. To excite ω_{SPR} (after Mie)

$$\varepsilon_1(\omega_{SPR}) = -2\varepsilon_m$$

$$\omega_{SPR}^2 = \frac{\omega_P^2}{1 + 2\varepsilon_m}$$

$$\varepsilon_1(\omega_{SPR}) \approx 1 - \frac{\omega_P^2}{\omega_{SPR}^2} = -2\varepsilon_m$$

$$\omega_{SPR} = \frac{\omega_P}{\sqrt{1 + 2\varepsilon_m}}$$

 $\lambda_{SPR} = \lambda_P \sqrt{1 + 2\varepsilon_m} > \lambda_P$

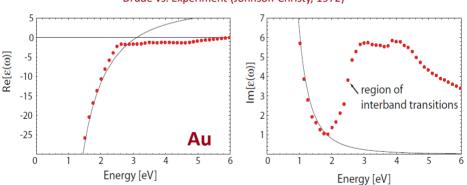


Drude Model

Correction due to higher polarizability of core \boldsymbol{d} electrons

$$\varepsilon = \varepsilon_{\infty} - \frac{\omega_P^2}{\omega^2 + i\omega\Gamma}$$
$$1 \le \varepsilon_{\infty} \le 10$$

Drude vs. Experiment (Johnson-Christy, 1972)





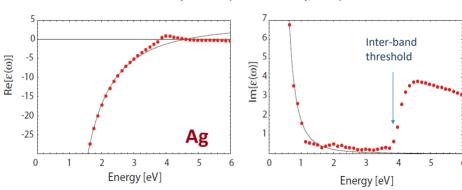


Drude Model

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Drude vs. Experiment (Johnson-Christy, 1972)





P.B. Johnson & R. W. Christy, Phys. Rev. B 6 (1972) 4370

G.Mattei



Drude-Lorentz model

Generalization to bound-electrons

$$\bar{E} = \bar{E}_0 e^{-i\omega t}$$

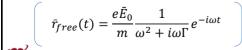
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$$m\frac{d^2\bar{r}}{dt^2} + m\Gamma\frac{d\bar{r}}{dt} + m\omega_0^2\bar{r} = -e\bar{E}_0e^{-i\omega t}$$

$$\bar{r}_{bound}(t) = \frac{e\bar{E}_0}{m} \frac{1}{\omega^2 - \omega_0^2 + i\omega\Gamma} e^{-i\omega t}$$

$$m \omega^2 - \omega_0^2 + i\omega\Gamma^2$$





 $\omega_0^2 \equiv \frac{K}{m}$

$$\varepsilon = \varepsilon(\omega, R)$$

Drude-Lorentz model

Generalization to bound-electrons

$$\varepsilon_{bound} = 1 - \frac{\omega_P^2}{\omega^2 - \omega_0^2 + i\omega\Gamma}$$

$$\omega_{P,bound}^2 \equiv \frac{n_{bound} \ e^2}{m_{bound} \varepsilon_0}$$

Add as many Lorentzian oscillators as the bound states in the system (QM).

$$\varepsilon = \varepsilon_{Drude} + \varepsilon_{bound} = \varepsilon_{s} + \varepsilon_{d} =$$

$$= 1 - \frac{\omega_{P}^{2}}{\omega^{2} + i\omega\Gamma} - \sum_{j} \frac{\omega_{P,j}^{2} f_{j}}{\omega^{2} - \omega_{0,j}^{2} + i\omega\Gamma_{j}}$$





$$\varepsilon = \varepsilon(\omega, R)$$

$$\varepsilon(\omega) = \varepsilon_d(\omega) + \varepsilon_s(\omega)$$
Interband s electrons (Drude)

• $\epsilon_{\rm d}$ (interband) does not depend on R

 $R \sim \lambda_o$

- \bullet Drude correction for ϵ_{bulk}
- Relaxation frequency correction (size-equation) $\Gamma = \Gamma(R)$

$$\Gamma(R) = \Gamma_{bulk} + A \frac{v_F}{R}$$
 $A \sim 1$
 v_F Fermi velocity

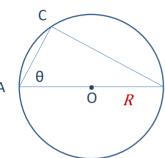


 $\varepsilon(\omega, R) = \varepsilon_d(\omega) + \varepsilon_s(\omega, R)$

Quantum confinement



Semiclassical correction: size equation



$$AC = 2R\cos(\theta)$$

$$\tau_{AC}(\theta) = \frac{AC}{v_E} = \frac{2R\cos(\theta)}{v_E}$$

$$v_F \qquad v_F$$

$$\langle \tau \rangle = \frac{\int_0^{\frac{\pi}{2}} d\theta \ \tau_{AC}(\theta)}{\int_0^{\frac{\pi}{2}} d\theta} = \frac{2R}{v_F \frac{\pi}{2}} = \frac{4R}{v_F \pi}$$

$$\Gamma(R) = \Gamma_{bulk} + \Gamma_{sup} = \Gamma_{bulk} + \frac{1}{\langle \tau \rangle} = \Gamma_{bulk} + \frac{\pi}{4} \frac{v_F}{R}$$



$$\varepsilon(\omega, R) = \varepsilon_d(\omega) + \varepsilon_s(\omega, R)$$

$$\varepsilon(\omega)_{Drude} = 1 - \frac{\omega_P^2}{\omega^2 + \Gamma^2} + i \frac{\omega_P^2 \Gamma}{\omega(\omega^2 + \Gamma^2)}$$

$$\omega^{2} + \Gamma^{2} + \omega(\omega^{2} + \Gamma^{2})$$

$$\varepsilon(\omega, R) = \varepsilon(\omega, \infty) + \omega_{P}^{2} \left(\frac{1}{\omega^{2} + \Gamma_{\infty}^{2}} - \frac{1}{\omega^{2} + \Gamma(R)^{2}} \right) - i \frac{\omega_{P}^{2}}{\omega} \left(\frac{\Gamma_{\infty}}{\omega^{2} + \Gamma_{\infty}^{2}} - \frac{\Gamma(R)}{\omega^{2} + \Gamma(R)^{2}} \right)$$

$$\varepsilon(\omega, R) = \varepsilon(\omega, \infty) + \omega_P^2 \left(\frac{1}{\omega^2 + \Gamma_\infty^2} - \frac{1}{\omega^2 + \Gamma_\infty^2}\right)$$

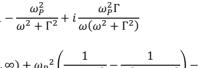
 $\hbar\Gamma(R) = \hbar\Gamma_{bulk} + \hbar A \frac{v_F}{R}$

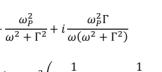
 $E_{\Gamma}(R) = E_{\Gamma}(\infty) \left(1 + \frac{C}{R} \right)$

$$\frac{1}{\omega^2 + \Gamma^2} + i \omega(\omega^2 + \Gamma^2)$$

$$(0, \infty) + \omega_P^2 \left(\frac{1}{\omega^2 + \Gamma_\infty^2} - \frac{1}{\omega^2 + \Gamma_\infty^2} \right)$$

$$-\frac{\omega_p^2}{\omega^2 + \Gamma^2} + i \frac{\omega_p^2 \Gamma}{\omega(\omega^2 + \Gamma^2)}$$











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$$\frac{1}{\Gamma_{\infty}^2} - \frac{1}{\omega^2 + \Gamma(R)^2} - \frac{1}{(1 + \epsilon)^2}$$

2.0-



