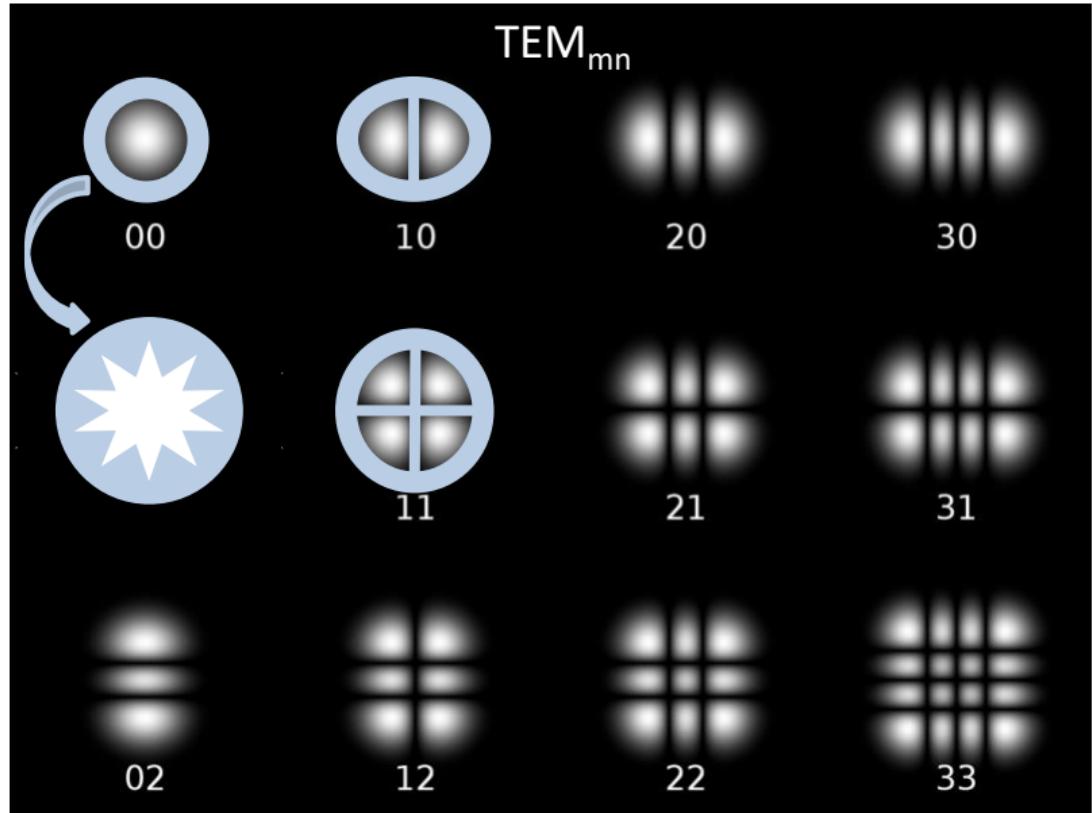
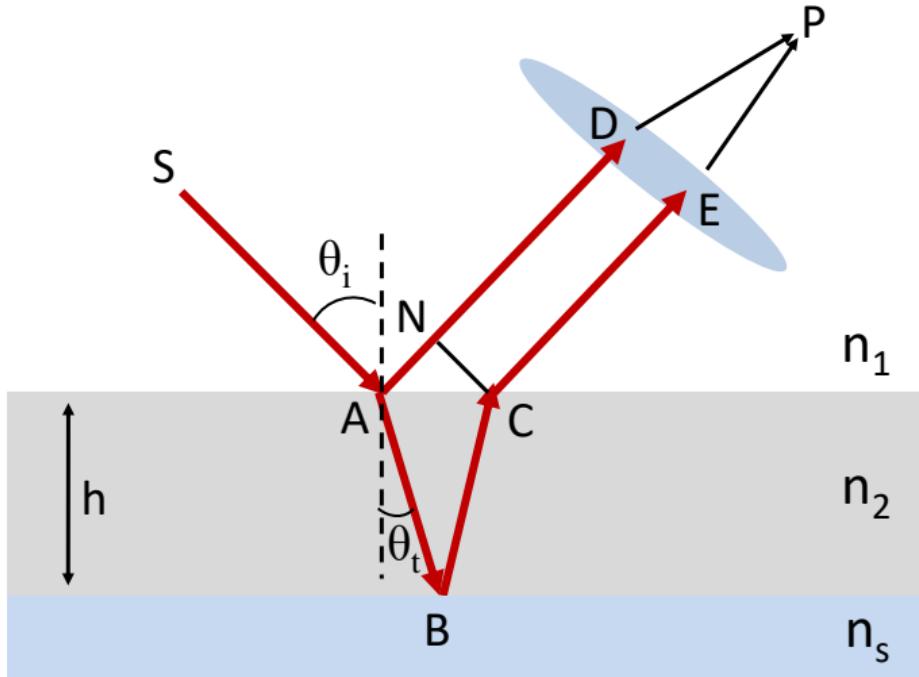


How to get single mode selection?



Thin-film interference



$$\delta = \frac{4\pi}{\lambda_0} n_2 h \cos \theta_t (+\pi)$$

Fringes with equal inclination
(Haidenger's fringes)

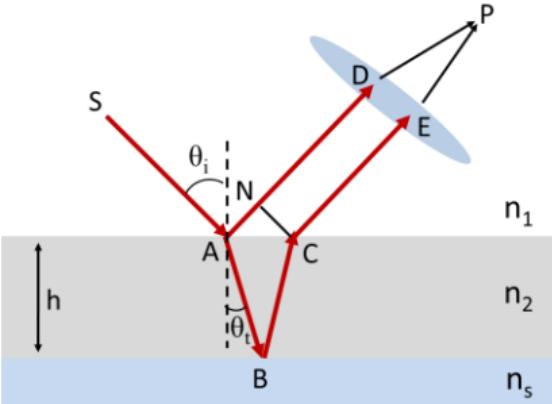
Thin-film interference

The difference of optical path is given by:

$$\Delta L = n_2(\overline{AB} + \overline{BC}) - n_1 \overline{AN}$$

$$\overline{AN} = \overline{AC} \sin \theta_i = 2h \frac{\sin \theta_t}{\cos \theta_t} \sin \theta_i$$

$$\overline{AB} = \overline{BC} = \frac{h}{\cos \theta_t}$$



$$\Delta L = n_2 \frac{2h}{\cos \theta_t} - n_1 2h \frac{\sin \theta_t}{\cos \theta_t} \sin \theta_i = n_2 \frac{2h}{\cos \theta_t} - 2h \frac{\sin \theta_t}{\cos \theta_t} n_2 \sin \theta_t =$$

$$= n_2 \frac{2h}{\cos \theta_t} (1 - \sin^2 \theta_t) = 2h n_2 \cos \theta_t$$

$$\Delta \phi_{path} = k \Delta L = \frac{2\pi}{\lambda_0} \Delta L = \frac{4\pi}{\lambda_0} h n_2 \cos \theta_t \quad \Delta \phi_{refl} = \pi \quad n_2 > n_1$$

$$\Delta \phi_{refl} = 0 \quad n_2 < n_1$$

Exercise:

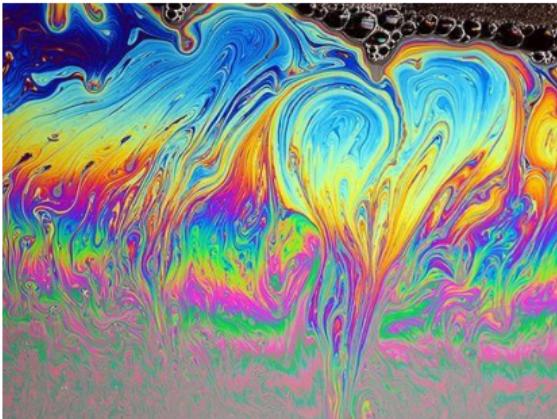
Some oil (with refractive index $n_2 = 1.25$) is on the surface of a water puddle with refractive index $n_s = 1.33$.

1. At normal incidence, what is the minimum oil film thickness reflecting most the red light at $\lambda = 650 \text{ nm}$?
 2. At which angle of incidence the same film reflects most the blue light at $\lambda = 450 \text{ nm}$?
1. $n_1 < n_2 < n_s \implies$ each interface produce a π phase shift in reflection

$$\implies \Delta\phi_{refl} = 0 \quad \delta = \frac{4\pi}{\lambda_0} n_2 h \cos \theta_t = 2\pi \quad \text{for constructive interference}$$

1 at normal incidence

$$h = \frac{\lambda_0}{2n_2} = \frac{650 \text{ nm}}{2 \cdot 1.25} = 260 \text{ nm}$$



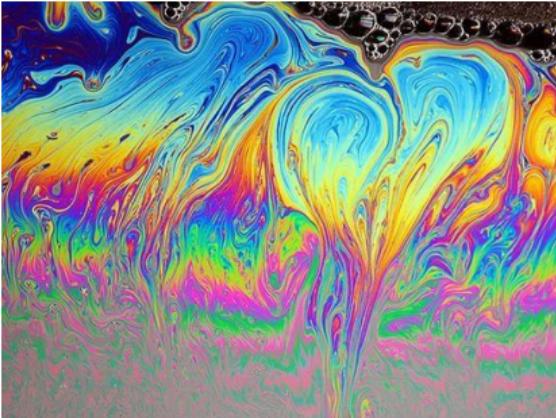
Thin-film interference

Exercise:

Some oil (with refractive index $n_2 = 1.25$) is on the surface of a water puddle with refractive index $n_s = 1.33$.

1. At normal incidence, what is the minimum oil film thickness reflecting most the red light at $\lambda = 650 \text{ nm}$?
2. At which angle of incidence the same film reflects most the blue light at $\lambda = 450 \text{ nm}$?

$$\begin{aligned} 2. \quad 2\pi &= \frac{4\pi}{\lambda_0} n_2 h \cos \theta_t \quad \Rightarrow \quad \cos \theta_t = \frac{\lambda_0}{2n_2 h} = \frac{450 \text{ nm}}{2 \cdot 1.25 \cdot 260 \text{ nm}} = 0.692 \\ &\Rightarrow \theta_t = 42.2^\circ \quad \Rightarrow \quad \sin \theta_i = \frac{n_2}{n_1} \sin \theta_t = 0.902 \\ &\Rightarrow \theta_i = 64.4^\circ \end{aligned}$$



Thin-film interference

A film of soapy water suspended vertically on a metallic ring can get very thin before breaking

$$\Delta\phi_{path} = k\Delta\mathcal{L} = 0$$

$$n_1 = n_s = 1$$

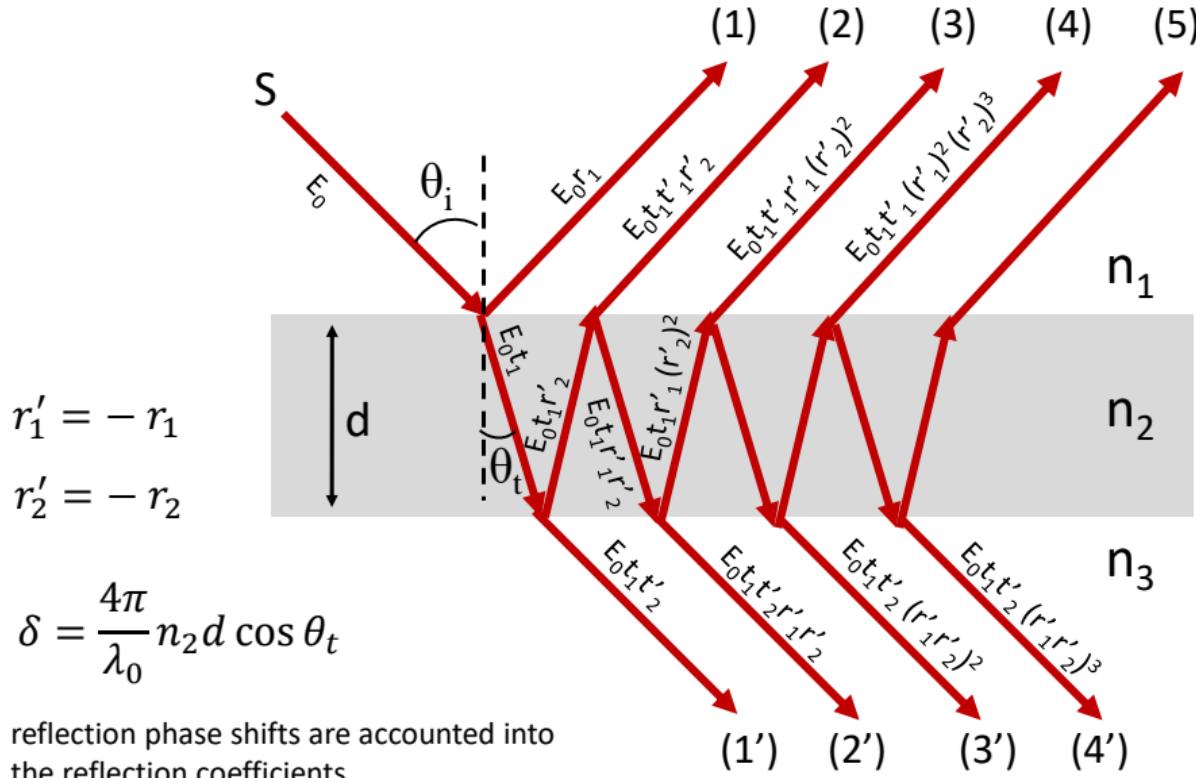
$$\Delta\phi_{refl} = \pi$$

$$\delta = \Delta\phi_{path} + \Delta\phi_{refl} = \pi$$

Destructive interference!



Multiple beams interference



Multiple beams interference

Let's consider the transmitted beam:

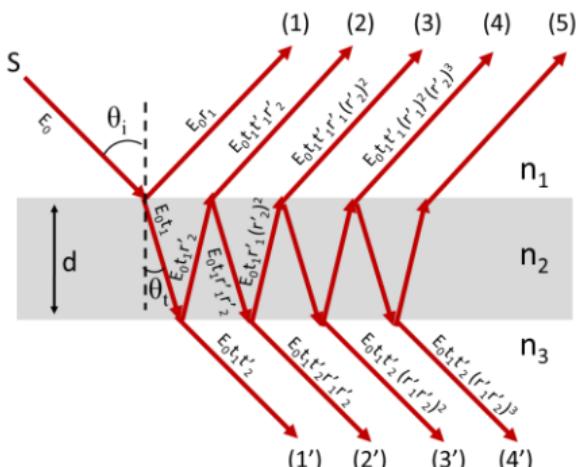
$$E_t = E_0 t_1 t'_2 + E_0 t_1 t'_2 r'_2 r'_1 e^{i\delta} + E_0 t_1 t'_2 (r'_2 r'_1)^2 e^{i2\delta} + \dots + E_0 t_1 t'_2 (r'_2 r'_1)^N e^{iN\delta}$$

$$= E_0 t_1 t'_2 [1 + x + x^2 + \dots x^N] \quad x = r'_2 r'_1 e^{i\delta}$$

$$\sum_{n=1}^N x^n \rightarrow \frac{1}{1-x}$$

$$|x| < 1, \quad N \text{ large}$$

$$E_t = E_0 t_1 t'_2 \frac{1}{1 - r'_2 r'_1 e^{i\delta}}$$



Multiple beams interference

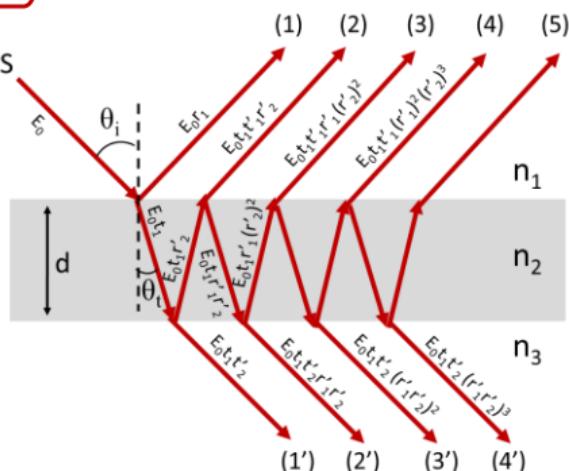
Analogously for the reflected beam:

$$E_r = E_0 r_1 + E_0 t_1 t'_1 r'_2 e^{i\delta} + E_0 t_1 t'_1 r'_1 (r'_2)^2 e^{i2\delta} + \dots + \\ + \dots + E_0 t_1 t'_1 (r'_1)^{N-2} (r'_2)^{N-1} e^{i(N-1)\delta}$$

$$= E_0 r_1 + E_0 t_1 t'_1 r'_2 e^{i\delta} [1 + x + x^2 + \dots x^N] \quad x = r'_2 r'_1 e^{i\delta}$$

$$\frac{1}{1-x}$$

$$E_r = E_0 \left(r_1 + \frac{t_1 t'_1 r'_2 e^{i\delta}}{1 - r'_2 r'_1 e^{i\delta}} \right)$$



Assuming $n_1 = n_3 \Rightarrow r_1 = r_2 = r \quad t_1 = t_2 = t$

$$E_t = E_0 tt' \frac{1}{1 - (r')^2 e^{i\delta}}$$

$$E_r = E_0 \left(r + \frac{tt' r' e^{i\delta}}{1 - (r')^2 e^{i\delta}} \right)$$

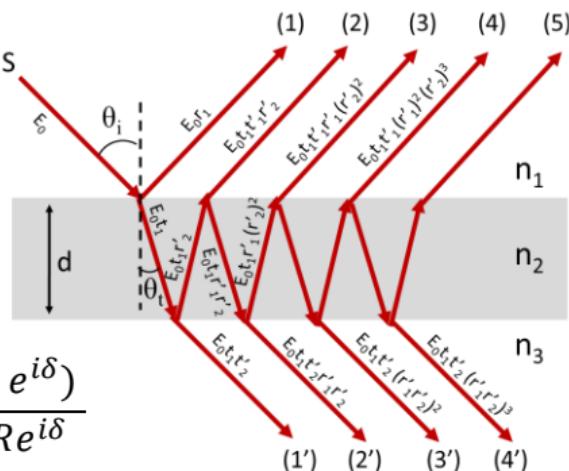
Defining $|r|^2 = |r'|^2 = R \quad tt' = T$

$$\Rightarrow R + \frac{n_3}{n_1} T = 1 \Rightarrow R + T = 1$$

$$E_t = E_0 (1 - R) \frac{1}{1 - Re^{i\delta}}$$

$$r' = -r$$

$$E_r = E_0 \left(r + \frac{(1 - R)r' e^{i\delta}}{1 - Re^{i\delta}} \right) = E_0 \frac{r(1 - e^{i\delta})}{1 - Re^{i\delta}}$$



Multiple beams interference

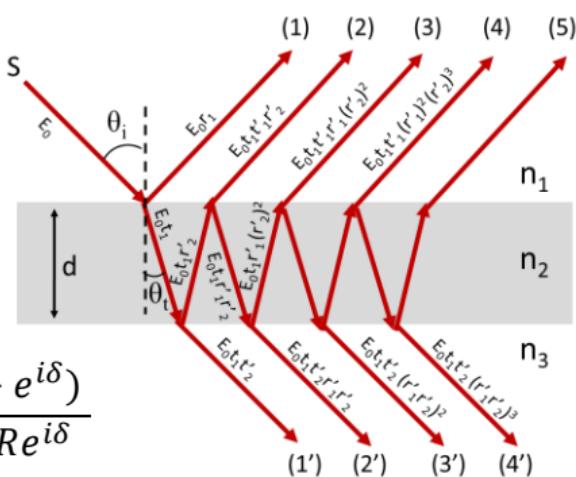
Given $I_t \propto |E_t|^2 \quad I_r \propto |E_r|^2$

$$\begin{aligned} I_t &= I_0(1-R)^2 \frac{1}{1-Re^{i\delta}} \frac{1}{1-Re^{-i\delta}} = I_0(1-R)^2 \frac{1}{1+R^2-2R\cos\delta} \\ &= I_0 \frac{(1-R)^2}{(1-R)^2 + 4R \left(\sin \frac{\delta}{2}\right)^2} \quad \cos\delta = 1 - 2 \left(\sin \frac{\delta}{2}\right)^2 \\ &= I_0 \frac{1}{1 + \frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2} \end{aligned}$$

$$E_t = E_0(1-R) \frac{1}{1-Re^{i\delta}}$$

$$E_r = E_0 \left(r + \frac{(1-R)r'e^{i\delta}}{1-Re^{i\delta}} \right) = E_0 \frac{r(1-e^{i\delta})}{1-Re^{i\delta}}$$

$r' = -r$



Given $I_t \propto |E_t|^2$ $I_r \propto |E_r|^2$

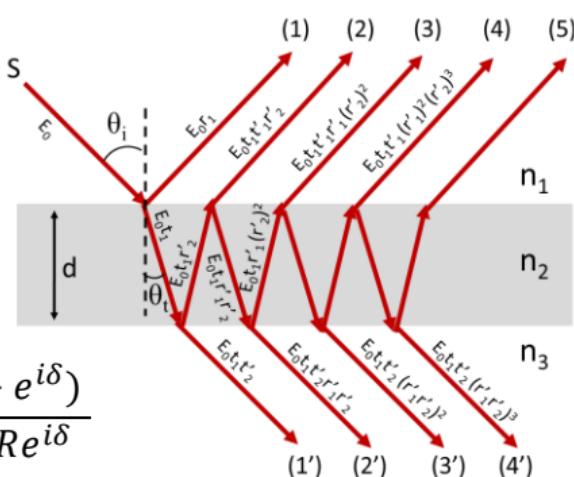
$$I_t = I_0 \frac{1}{1 + \frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2}$$

$$I_r = I_0 \frac{\frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2}{1 + \frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2}$$

$$E_t = E_0 (1-R) \frac{1}{1 - R e^{i\delta}}$$

$$E_r = E_0 \left(r + \frac{(1-R)r' e^{i\delta}}{1 - R e^{i\delta}} \right) = E_0 \frac{r(1 - e^{i\delta})}{1 - R e^{i\delta}}$$

$$\cos \delta = 1 - 2 \left(\sin \frac{\delta}{2}\right)^2$$



Defining $F = \frac{4R}{(1-R)^2}$ **coefficient of finesse** (of the slab)

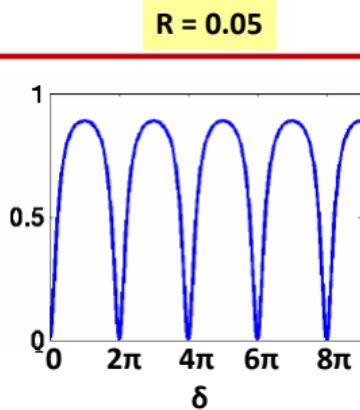
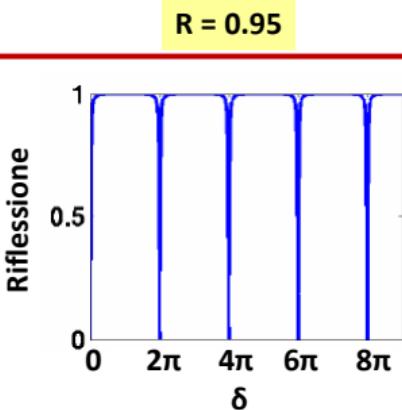
$$I_t = I_0 \frac{1}{1 + \frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2} = I_0 \frac{1}{1 + F \left(\sin \frac{\delta}{2}\right)^2}$$

$$I_r = I_0 \frac{\frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2}{1 + \frac{4R}{(1-R)^2} \left(\sin \frac{\delta}{2}\right)^2} = I_0 \frac{F \left(\sin \frac{\delta}{2}\right)^2}{1 + F \left(\sin \frac{\delta}{2}\right)^2}$$

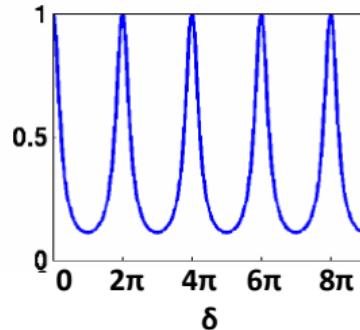
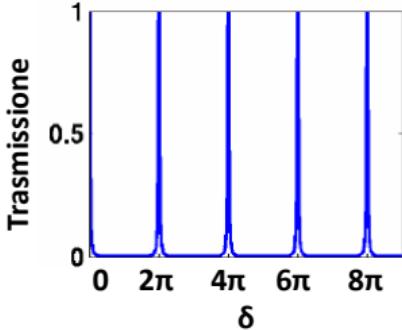
Airy's formulas

Airy's formulas

$$\frac{I_r}{I_0} = \frac{F \sin^2(\delta/2)}{1 + F \sin^2(\delta/2)}$$



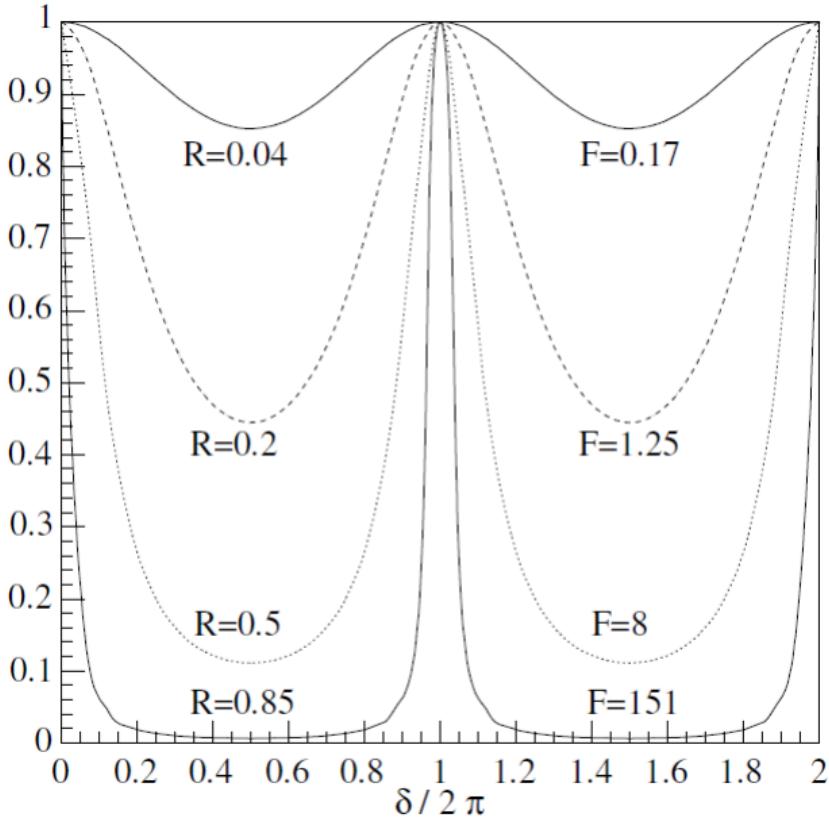
$$\frac{I_t}{I_0} = \frac{1}{1 + F \sin^2(\delta/2)}$$



Airy's formulas

$$\frac{I_t}{I_0} = \frac{1}{1 + F \sin^2(\delta/2)}$$

$$F = \frac{4R}{(1-R)^2}$$



A useful parameter to evaluate the degree of definition of a fringe is its **full width at half maximum** (FWHM, γ)

$$\frac{I_t}{I_0} = \frac{1}{1 + F \left(\sin \frac{\delta_{1/2}}{2} \right)^2} = \frac{1}{2} \quad \Rightarrow \quad \left(\sin \frac{\delta_{1/2}}{2} \right)^2 = \frac{1}{F}$$

$$\Rightarrow \quad \sin \frac{\delta_{1/2}}{2} = \frac{1}{\sqrt{F}} \cong \frac{\delta_{1/2}}{2} \quad F \gg 1$$

$$\boxed{\gamma = 2\delta_{1/2} = \frac{4}{\sqrt{F}}} \quad \text{FWHM}$$

$$\boxed{\mathcal{F} = \frac{2\pi}{\gamma} = \frac{\pi\sqrt{F}}{2} = \frac{\pi\sqrt{R}}{(1-R)}} \quad \text{Finesse (of the fringes)}$$

Reflection

$$\frac{I_r}{I_0} = \frac{F \sin^2(\delta/2)}{1 + F \sin^2(\delta/2)}$$

Transmission

$$\frac{I_t}{I_0} = \frac{1}{1 + F \sin^2(\delta/2)}$$

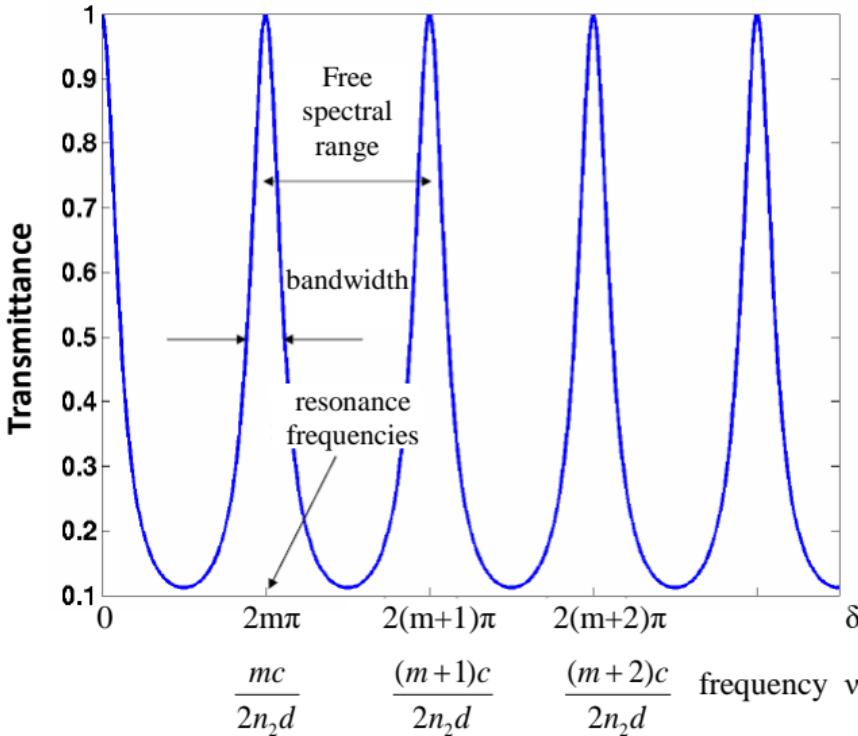
$$F = \frac{4R}{(1-R)^2} \quad \text{Coefficient of finesse (of the slab)}$$

$$\mathcal{F} = \frac{2\pi}{\gamma} = \frac{\pi\sqrt{F}}{2} = \frac{\pi\sqrt{R}}{(1-R)} \quad \text{Finesse (of the fringes)} \quad \gamma = \frac{4}{\sqrt{F}} \quad \text{FWHM}$$

For $R_1 \neq R_2$ replace $R \rightarrow \sqrt{R_1 R_2}$

Fabry-Perot etalon

$$\delta = \frac{4\pi}{\lambda_0} n_2 d \cos \theta_t \quad \Rightarrow \quad \text{at normal incidence} \quad \delta = \frac{4\pi}{\lambda_0} n_2 d = 2m\pi \quad \text{maxima}$$

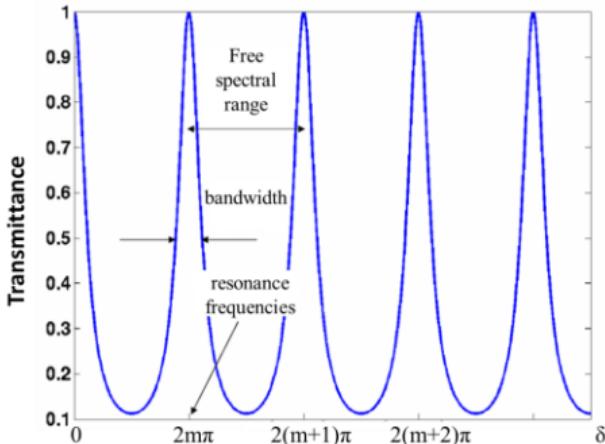


Fabry-Perot etalon

Fabry-Perot bandwidth, $\Delta\nu_{FP}$

$$\delta = \frac{4\pi}{\lambda_0} n_2 d = \frac{4\pi n_2 d}{c} \nu \quad \nu = \frac{c}{\lambda_0}$$

$$\Rightarrow \gamma = \Delta\delta = \frac{4\pi n_2 d}{c} \Delta\nu_{FP} = \frac{2\pi}{\mathcal{F}}$$

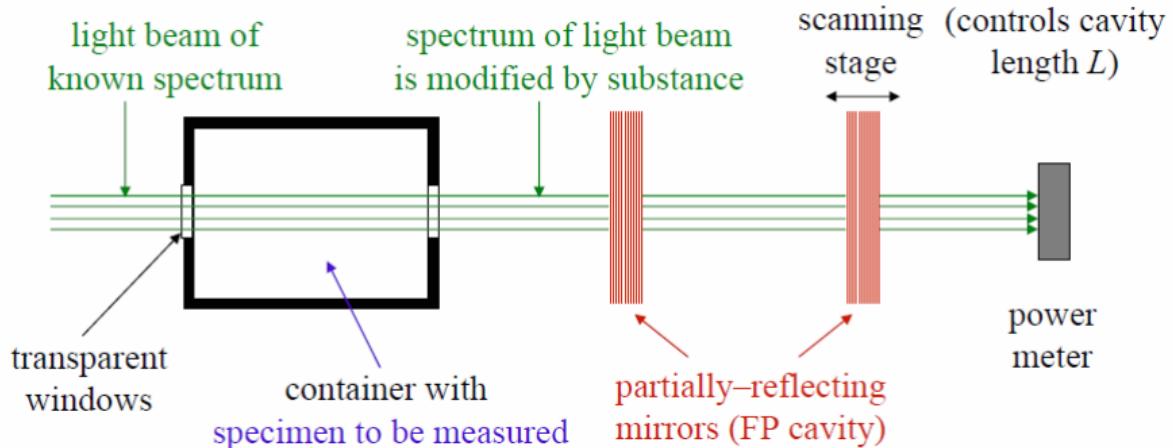


$$\Delta\nu_{FP} = \frac{c}{2n_2 d} \frac{1}{\mathcal{F}}$$

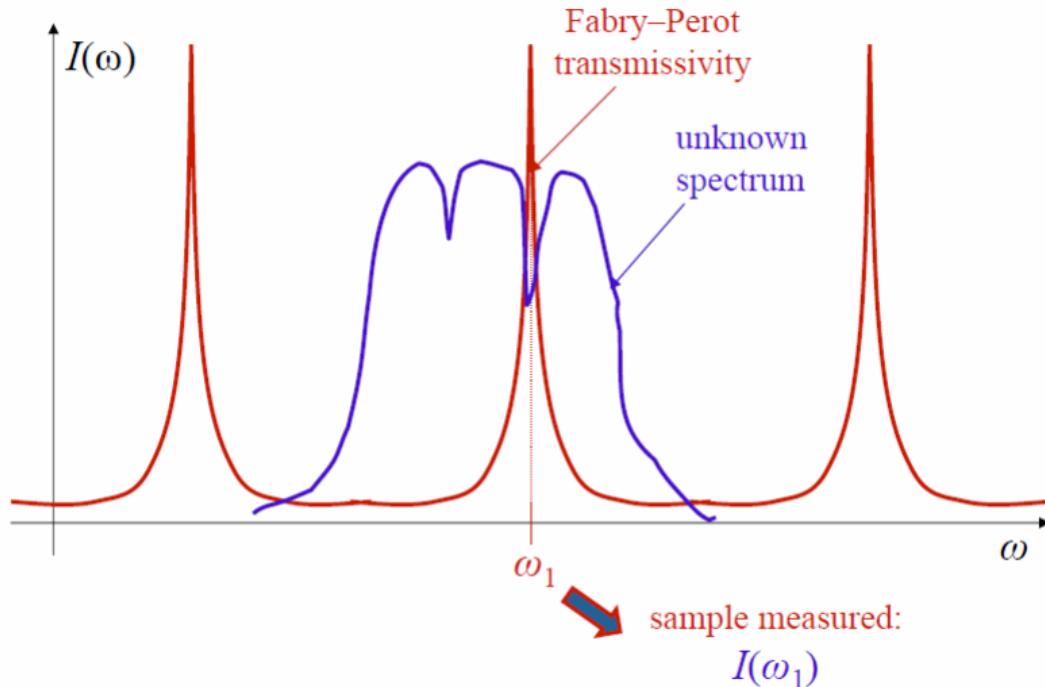
Free spectral range, $\Delta\nu_{FSR}$ $\Rightarrow \Delta\delta = 2\pi$

$$\Delta\nu_{FSR} = \frac{c}{2n_2 d}$$

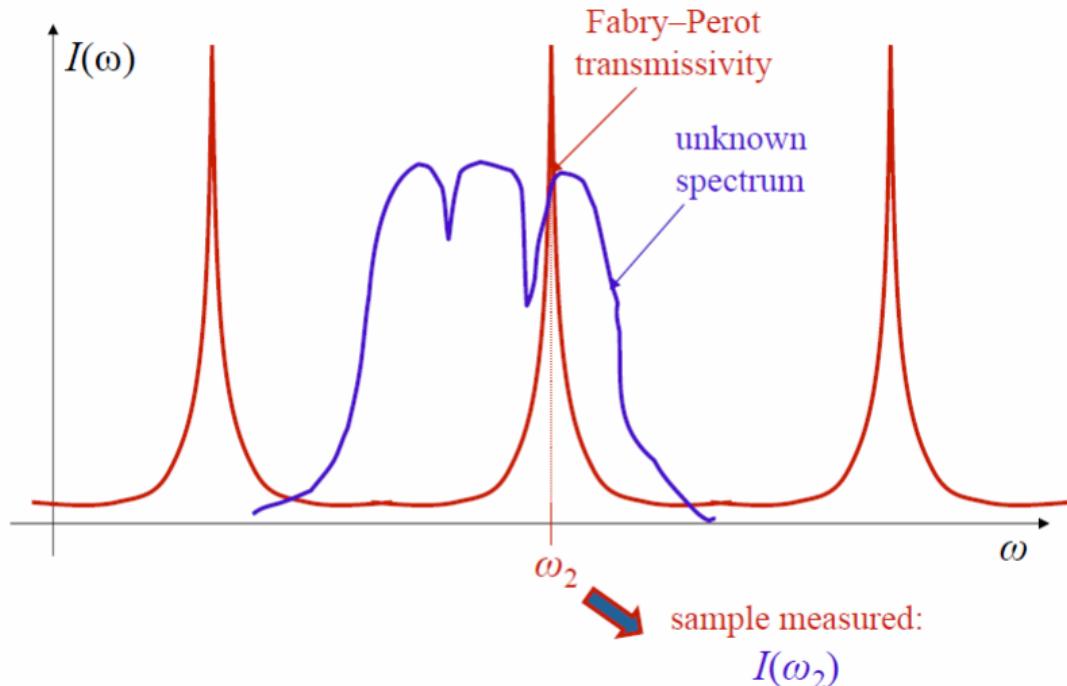
Spectroscopic applications



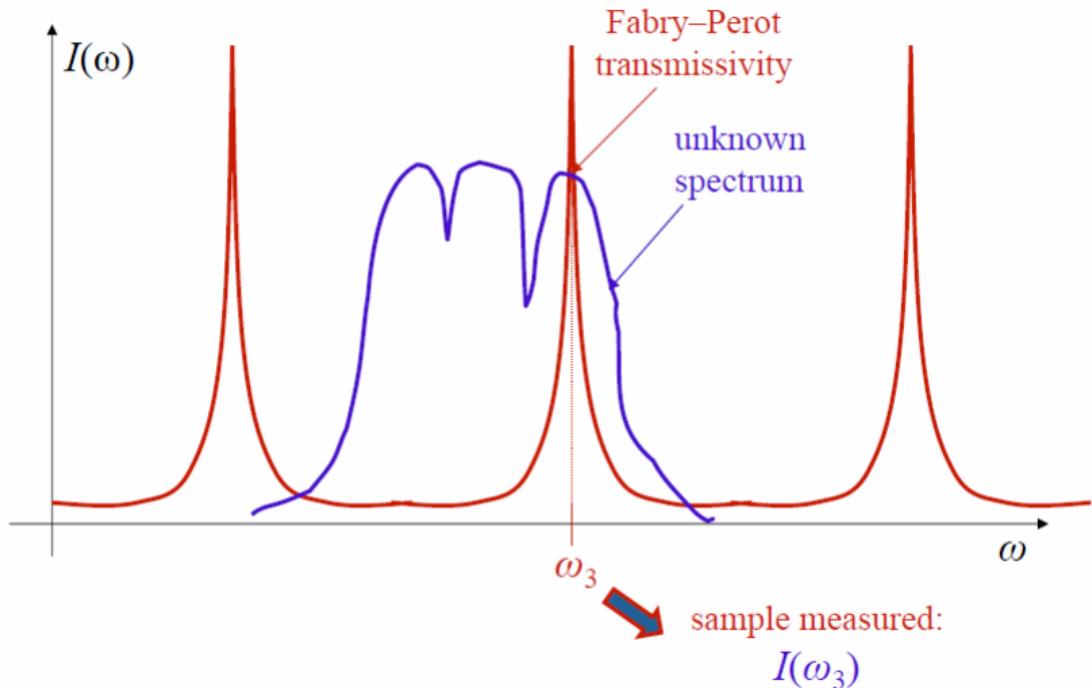
Spectroscopic applications



Spectroscopic applications



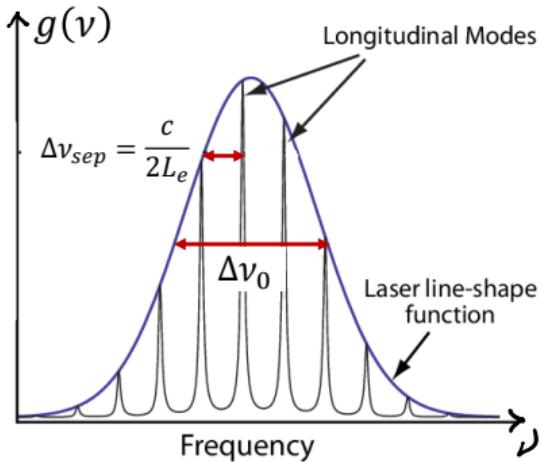
Spectroscopic applications



To have only one longitudinal mode oscillating in the cavity it has to be:

$$\Delta\nu_{sep} = \frac{c}{2L_e} \geq \frac{\Delta\nu_0}{2}$$

$$\Rightarrow L_e \leq \frac{c}{\Delta\nu_0}$$



He-Ne laser @ 633 nm $\Delta\nu_0 \cong 1.5 \text{ GHz}$ $\Rightarrow L_e \leq \frac{c}{\Delta\nu_0} = 20 \text{ cm}$

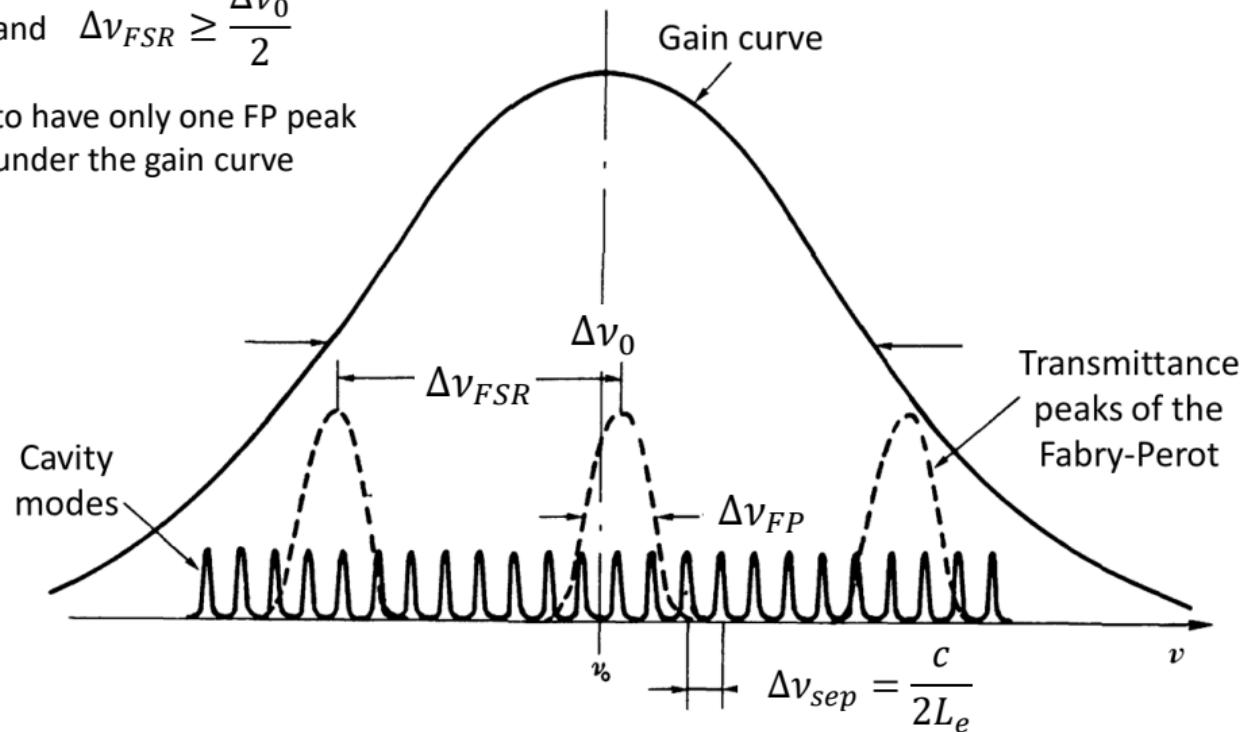
Nd:YAG laser @ 1064 nm $\Delta\nu_0 \cong 120 \text{ GHz}$ $\Rightarrow L_e \leq \frac{c}{\Delta\nu_0} = 2.5 \text{ mm}$!

Longitudinal modes selection

If a Fabry-Perot etalon is inserted in the cavity: $\Delta\nu_{sep} = \frac{c}{2L_e} \geq \frac{\Delta\nu_{FP}}{2} = \frac{\Delta\nu_{FSR}}{2\mathcal{F}}$

and $\Delta\nu_{FSR} \geq \frac{\Delta\nu_0}{2}$

to have only one FP peak
under the gain curve



Longitudinal modes selection

If a Fabry-Perot etalon is inserted in the cavity: $\Delta\nu_{sep} = \frac{c}{2L_e} \geq \frac{\Delta\nu_{FP}}{2} = \frac{\Delta\nu_{FSR}}{2\mathcal{F}}$

$$\text{and } \Delta\nu_{FSR} \geq \frac{\Delta\nu_0}{2}$$

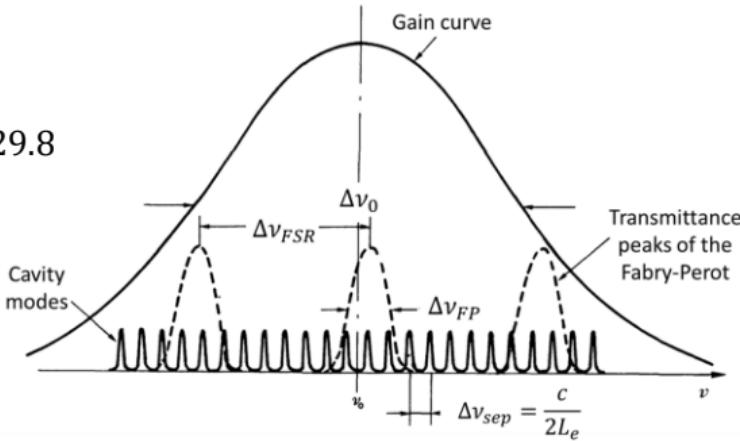
to have only one FP peak under the gain curve $\Rightarrow \frac{\Delta\nu_0}{2} \leq \Delta\nu_{FSR} \leq 2\mathcal{F}\Delta\nu_{sep} = 2\mathcal{F}\frac{c}{2L_e}$

$$\Rightarrow L_e \leq 2\mathcal{F}\frac{c}{\Delta\nu_0}$$

$$R = 90\% \quad \mathcal{F} = \frac{\pi\sqrt{R}}{(1-R)} \cong 29.8$$

$$R = 95\%$$

$$\mathcal{F} = \frac{\pi\sqrt{R}}{(1-R)} \cong 61.2$$



Longitudinal modes selection

Ar laser @ 514.5 nm (green line) $\Delta\nu_0 \cong 3.5 \text{ GHz}$

without the FP etalon: $L_e \leq \frac{c}{\Delta\nu_0} \cong 8.6 \text{ cm}$

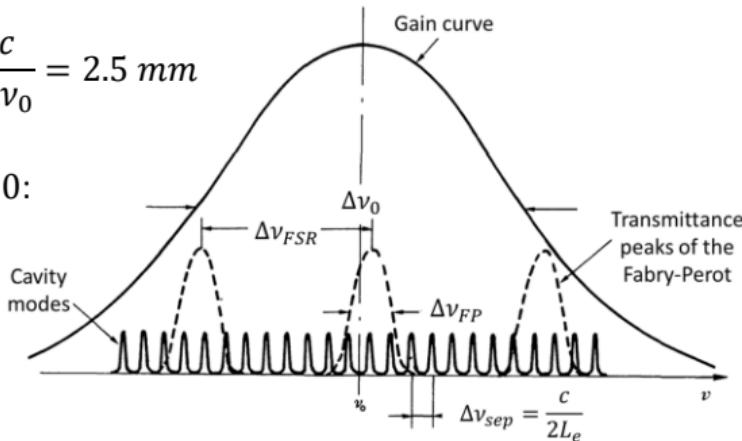
with one FP etalon with $\mathcal{F} = 30$: $L_e \leq 2 \mathcal{F} \frac{c}{\Delta\nu_0} \cong 5.16 \text{ m}$!

Nd:YAG laser @ 1064 nm $\Delta\nu_0 \cong 120 \text{ GHz}$

without the FP etalon: $L_e \leq \frac{c}{\Delta\nu_0} = 2.5 \text{ mm}$

with one FP etalon with $\mathcal{F} = 30$:

$$L_e \leq 2 \mathcal{F} \frac{c}{\Delta\nu_0} \cong 15 \text{ cm}$$



If the condition $L_e \leq 2 \mathcal{F} \frac{c}{\Delta\nu_0}$ is not satisfied

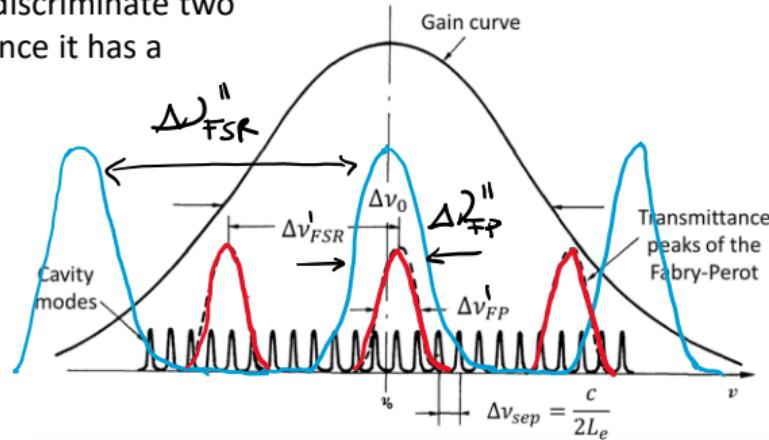
- ⇒ it is not possible to select a single longitudinal mode with one FP only
- ⇒ two or more FPs are needed

Let's consider the case of 2 FP etalons with the same finesse \mathcal{F} and different thickness.

The **thickest etalon** is used to discriminate two adjacent longitudinal modes since it has a smaller bandwidth:

$$\Delta\nu'_{FP} = \frac{c}{2n_2 d'} \frac{1}{\mathcal{F}}$$

The **thinnest etalon** is used to discriminate two transmission peaks of the first etalon



Longitudinal modes selection

$$\Rightarrow \Delta\nu_{sep} = \frac{c}{2L_e} \geq \frac{\Delta\nu'_{FP}}{2} = \frac{\Delta\nu'_{FSR}}{2\mathcal{F}}$$

to discriminate two adjacent longitudinal modes

$$\Delta\nu'_{FSR} \geq \frac{\Delta\nu''_{FP}}{2} = \frac{\Delta\nu''_{FSR}}{2\mathcal{F}}$$

to discriminate two adjacent transmission peaks of the first FP with the second FP

$$\Delta\nu''_{FSR} \geq \frac{\Delta\nu_0}{2}$$

to have only one transmission peak of the second FP under the gain curve

$$\Rightarrow 2\mathcal{F}\Delta\nu_{sep} \geq \Delta\nu'_{FSR} \geq \frac{\Delta\nu''_{FSR}}{2\mathcal{F}}$$

$$(2\mathcal{F})^2\Delta\nu_{sep} \geq \Delta\nu''_{FSR} \geq \frac{\Delta\nu_0}{2}$$

$$\Rightarrow L_e \leq (2\mathcal{F})^2 \frac{c}{\Delta\nu_0}$$

