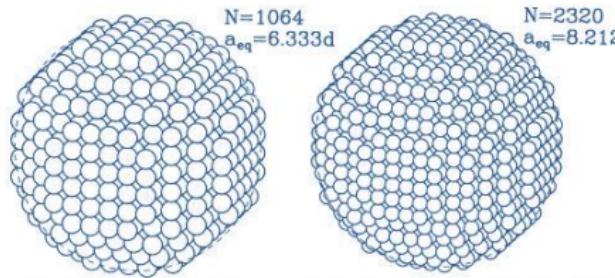
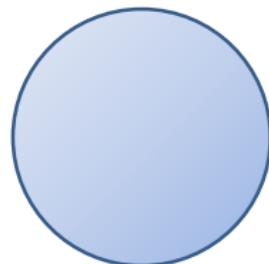




3. Interacting NPs

Discrete Dipole Approximation (DDA)



$$j = 1, \dots, N$$

 \mathbf{r}_j

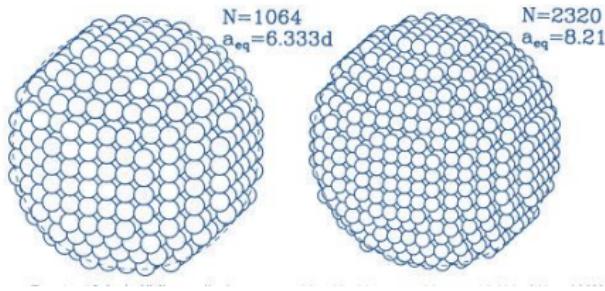
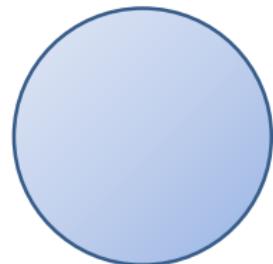
$$\alpha_j^{\text{CM}} = \frac{3d^3}{4\pi} \frac{\epsilon_j - 1}{\epsilon_j + 2}$$

$$\mathbf{E}_{\text{inc},j} = \mathbf{E}_0 \exp(i\mathbf{k} \cdot \mathbf{r}_j - i\omega t) \quad \begin{cases} \mathbf{P}_j = \alpha_j \mathbf{E}_j \\ \mathbf{E}_j = \mathbf{E}_{\text{inc},j} - \sum_{k \neq j} \mathbf{A}_{jk} \mathbf{P}_k \end{cases} \quad \mathbf{P}_j = \alpha_j \left(\mathbf{E}_{\text{inc},j} - \sum_{k \neq j} \mathbf{A}_{jk} \mathbf{P}_k \right)$$

$$\begin{cases} \mathbf{A}_{jk} = \frac{\exp(ikr_{jk})}{r_{jk}} \left[k^2 (\hat{r}_{jk} \hat{r}_{jk} - \mathbf{1}_3) + \frac{ikr_{jk} - 1}{r_{jk}^2} (3\hat{r}_{jk} \hat{r}_{jk} - \mathbf{1}_3) \right] & j \neq k \\ \mathbf{A}_{jj} = \alpha_j^{-1} \quad \text{definition} & j = k \end{cases}$$

$$k \equiv \omega/c, r_{jk} \equiv |\mathbf{r}_j - \mathbf{r}_k|, \hat{r}_{jk} \equiv (\mathbf{r}_j - \mathbf{r}_k)/r_{jk},$$

$$\boxed{\sum_{k=1}^N \mathbf{A}_{jk} \mathbf{P}_k = \mathbf{E}_{\text{inc},j}}$$



$$j = 1, \dots, N$$

$$\mathbf{r}_j$$

$$\alpha_j^{\text{CM}} = \frac{3d^3}{4\pi} \frac{\epsilon_j - 1}{\epsilon_j + 2}$$

$$\mathbf{E}_{\text{inc},j} = \mathbf{E}_0 \exp(i\mathbf{k} \cdot \mathbf{r}_j - i\omega t) \quad \begin{cases} \mathbf{P}_j = \alpha_j \mathbf{E}_j \\ \mathbf{E}_j = \mathbf{E}_{\text{inc},j} - \sum_{k \neq j} \mathbf{A}_{jk} \mathbf{P}_k \end{cases} \quad \mathbf{P}_j = \alpha_j \left(\mathbf{E}_{\text{inc},j} - \sum_{k \neq j} \mathbf{A}_{jk} \mathbf{P}_k \right)$$

solving

$$\sum_{k=1}^N \mathbf{A}_{jk} \mathbf{P}_k = \mathbf{E}_{\text{inc},j} \quad j = 1, \dots, N$$

$$\tilde{\mathbf{A}}\tilde{\mathbf{P}} = \tilde{\mathbf{E}}_{\text{inc}}$$

$$\begin{cases} C_{\text{ext}} = \frac{4\pi k}{|\mathbf{E}_0|^2} \sum_{j=1}^N \text{Im}(\mathbf{E}_{\text{inc},j}^* \cdot \mathbf{P}_j), \\ C_{\text{abs}} = \frac{4\pi k}{|\mathbf{E}_0|^2} \sum_{j=1}^N \left[\text{Im}[\mathbf{P}_j \cdot (\alpha_j^{-1})^* \mathbf{P}_j^*] - \frac{2}{3} k^3 |\mathbf{P}_j|^2 \right] \end{cases}$$

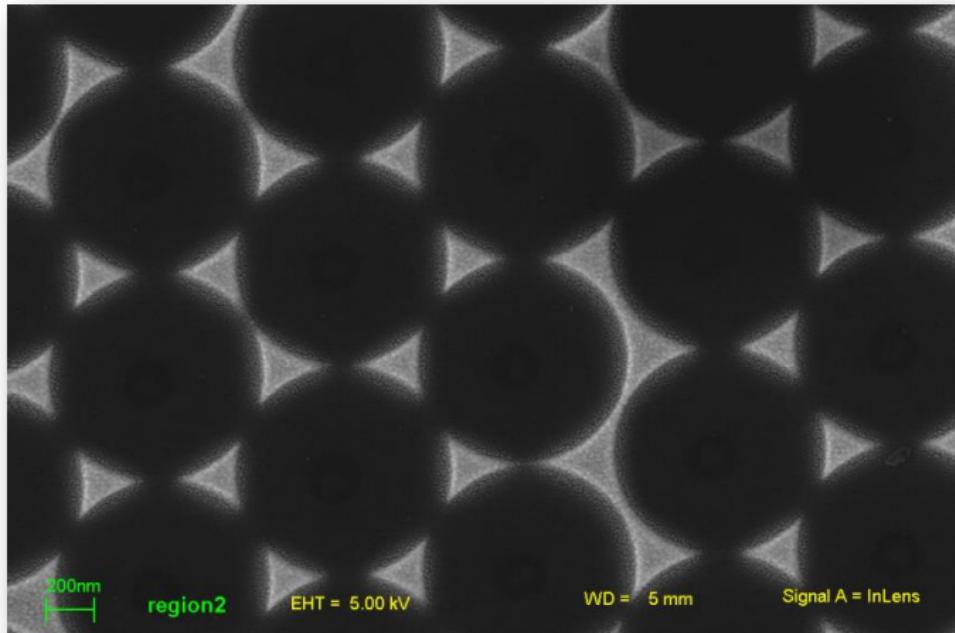
$$\mathbf{E}_{\text{sca}} = \frac{k^2 \exp(ikr)}{r} \sum_{j=1}^N \exp(-ik\hat{r} \cdot \mathbf{r}_j) (\hat{r}\hat{r} - \mathbf{1}_3) \mathbf{P}_j$$

3. Interacting NPs

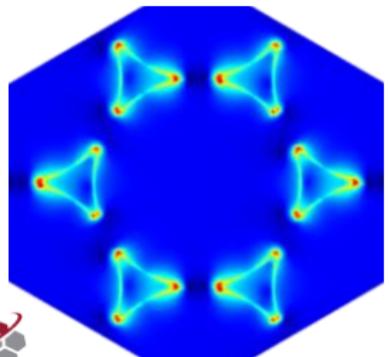
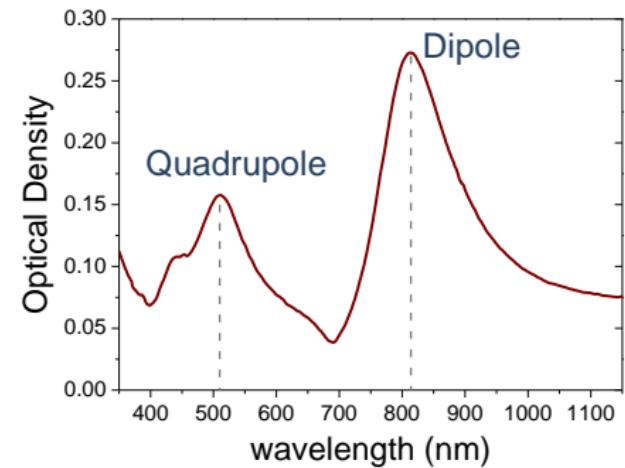
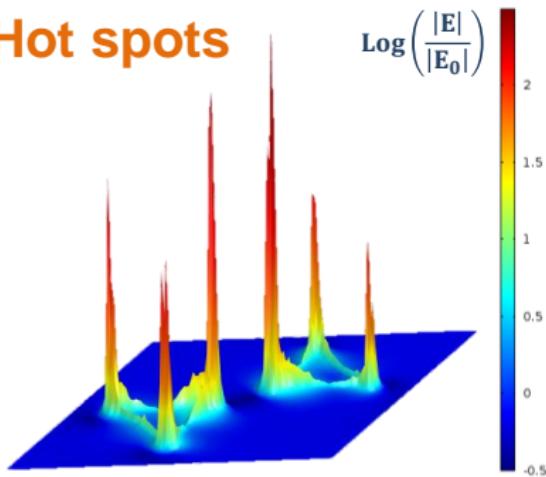
3.3 Array 2D - NSL

NPA

*Nano-Prism
arrays*

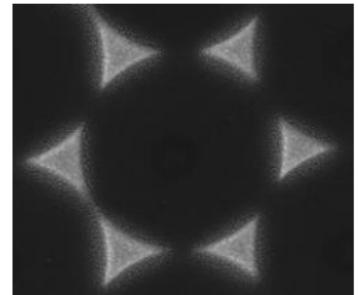


Hot spots

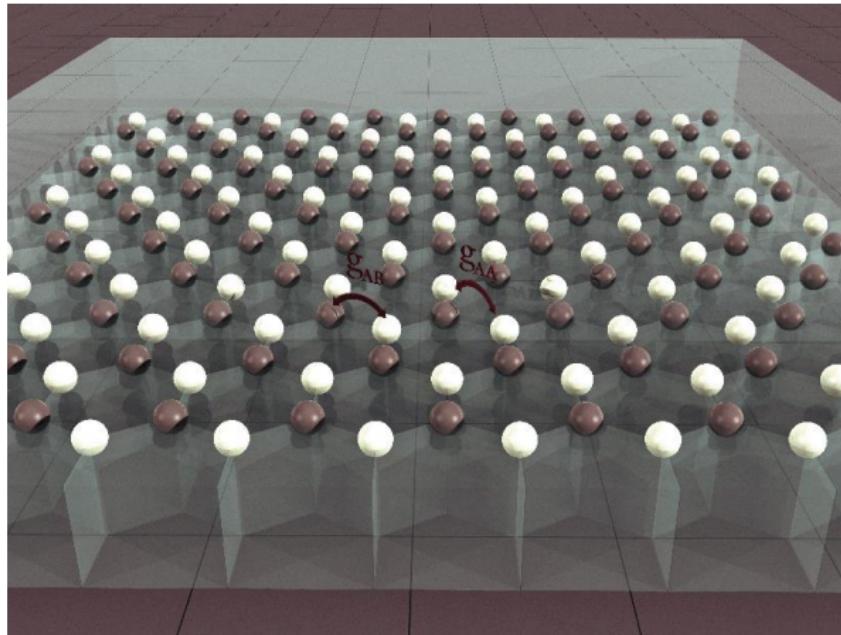


Interaction:

- SPR red-shift
- Local-field enhancement ~ 100 @ SPR
- Hot-spots



Diffractive dipolar coupling in non-Bravais plasmonic lattices



Array 2D - NSL

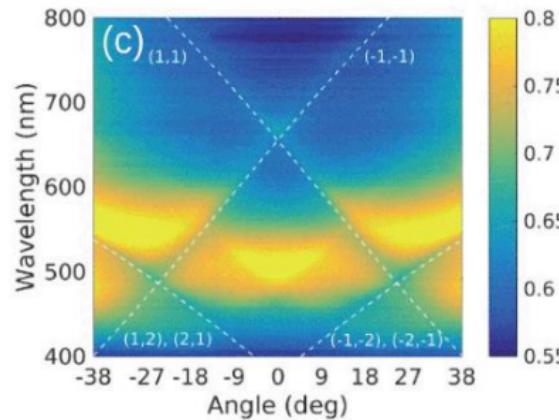
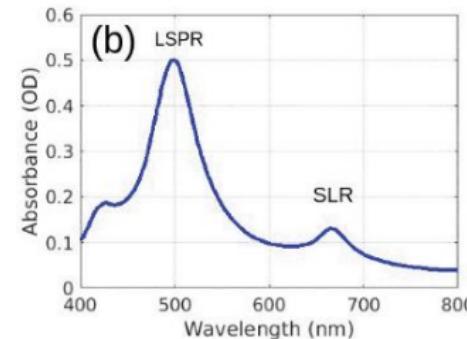
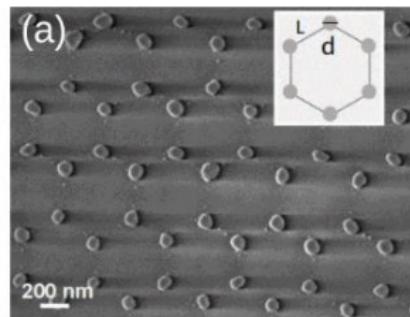
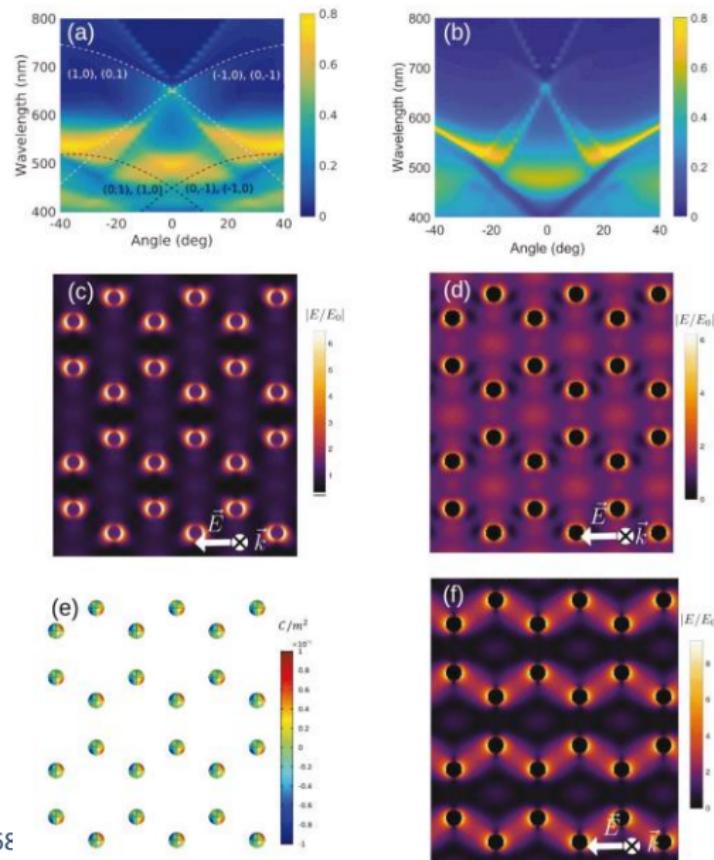


Fig. 1 (a) SEM image of the honeycomb plasmonic lattice. (b) Measured absorbance spectrum of the lattice. (c) Measured s-polarized extinction as a function of the wavelength and angle of incidence along the Γ – M trajectory. Dashed lines indicate RAs calculated with an effective refractive index of 1.47

Fig. 2 (a) Finite element method simulation and (b) spectral representation calculation of the s-polarized extinction map along the Γ – M trajectory. Dashed curves in (a) indicate RAs. (c) Simulated and (d) calculated spatial distribution of the normalized electric field amplitude for the LSPR, around $\lambda = 500$ nm. (e) Simulated surface charge density and (f) calculated spatial distribution of the normalized electric field amplitude for the SLR peak at $\lambda = 660$ nm. (c–f) are plotted in the plane crossing the nanospheres along their diameter, at normal incidence and for s-polarized incident light.



The response of materials to an optical field \mathbf{E} is described by the material polarization \mathbf{P} :

$$\mathbf{P} = \epsilon_0 [\chi^{(1)}\mathbf{E} + \chi^{(2)}\mathbf{E}^2 + \chi^{(3)}\mathbf{E}^3 + \dots + \chi^{(n)}\mathbf{E}^n + \dots]$$

Second-order effects:

- second harmonic generation (SHG)
- optical parametric amplification (OPA, OPO)
- ...

Third-order effects:

- third harmonic generation (THG)
- optical Kerr effect
- four-wave mixing (FWM)
- ...

The **local-field enhancement effects** can be exploited to design plasmonic nanosystems with amplified and tunable nonlinear optical response.



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- four-wave mixing (FWM)
- ...

The **local-field enhancement effects** can be exploited to design plasmonic nanosystems with amplified and tunable nonlinear optical response.



$$\epsilon = \epsilon_0 [\chi^{(1)} + 3\chi^{(3)}\mathbf{E}^2]$$

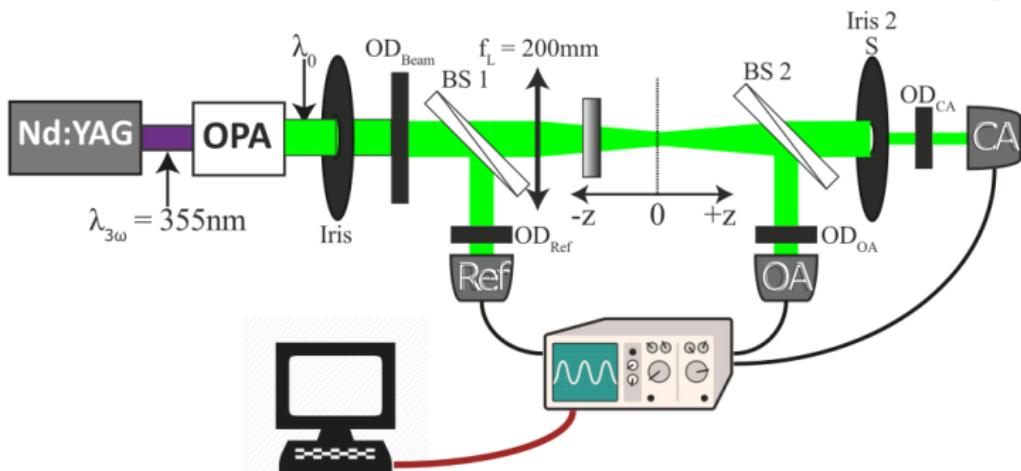
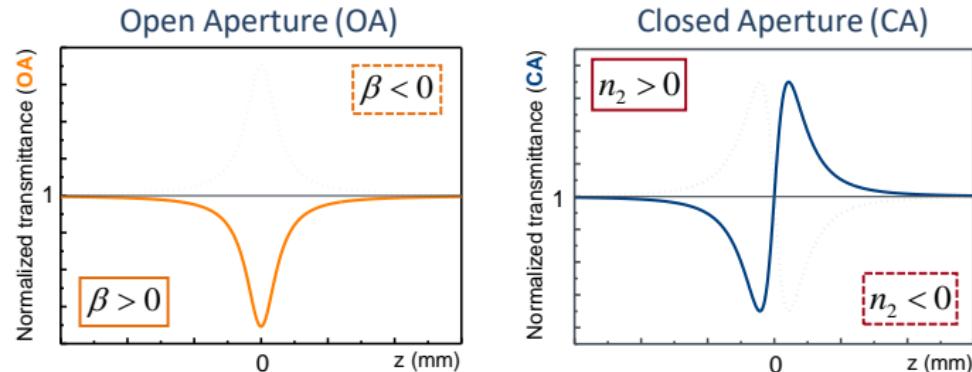
$$\begin{cases} n(I) = n_0 + n_2 I \\ \alpha(I) = \alpha_0 + \beta I \end{cases}$$

optical Kerr effect

The z-scan technique

M. Sheik-Bahae *et al.*
IEEE J. Quantum El. **26** 760 (1990)

$$\begin{cases} n(I) = n_0 + n_2 I \\ \alpha(I) = \alpha_0 + \beta I \end{cases}$$



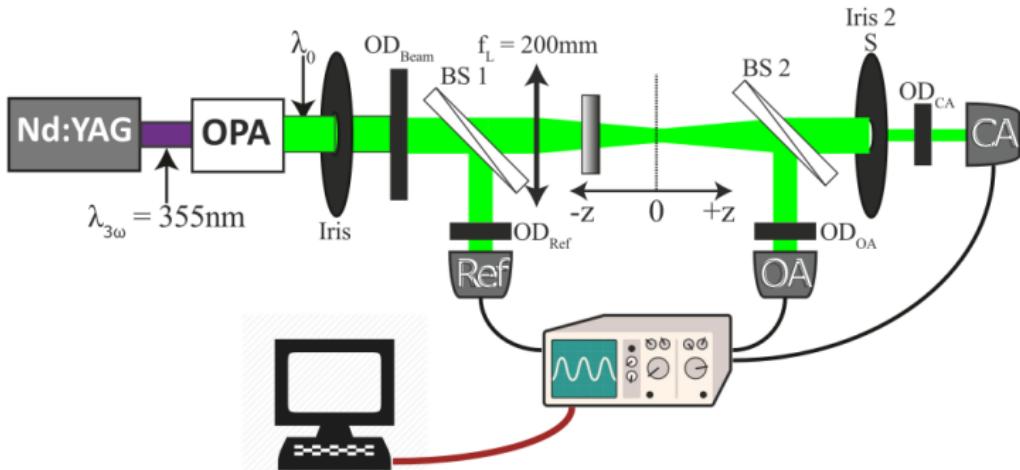
The z-scan set-up

M. Sheik-Bahae *et al.*
IEEE J. Quantum El. **26** 760 (1990)

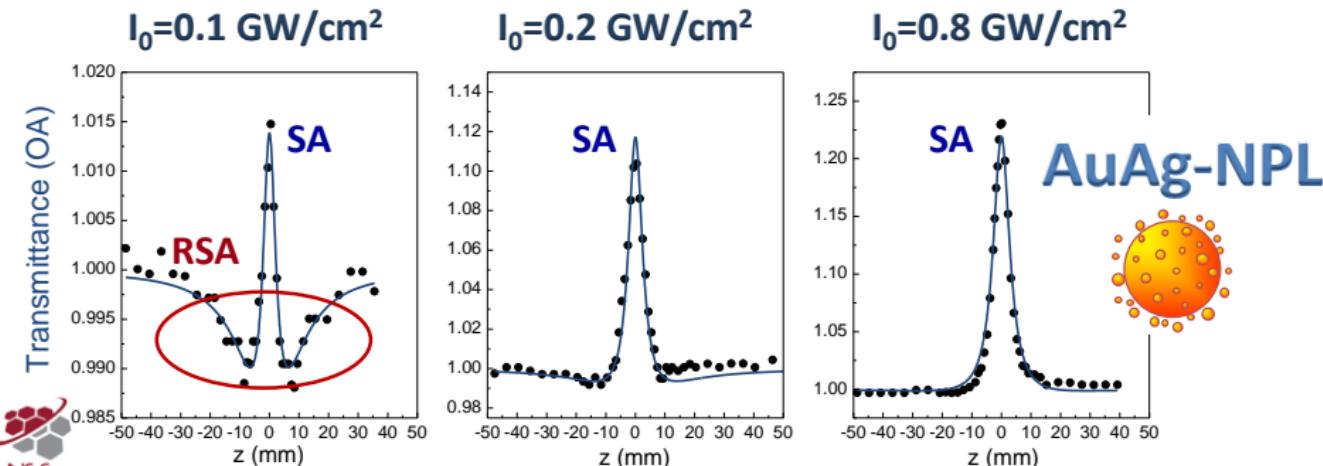
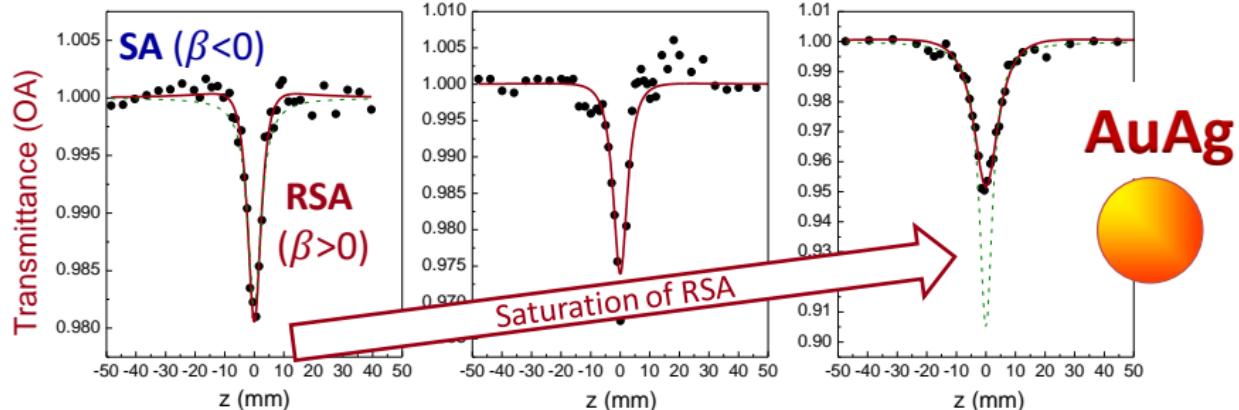
$$\begin{cases} n(I) = n_0 + n_2 I \\ \alpha(I) = \alpha_0 + \beta I \end{cases}$$

mode-locked Nd:YAG + OPA

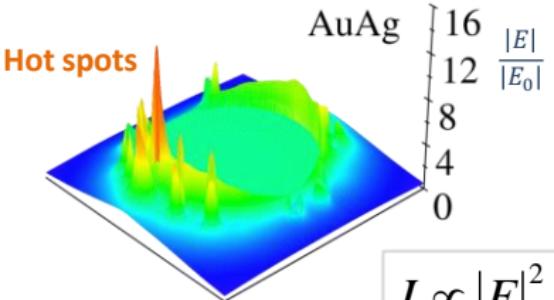
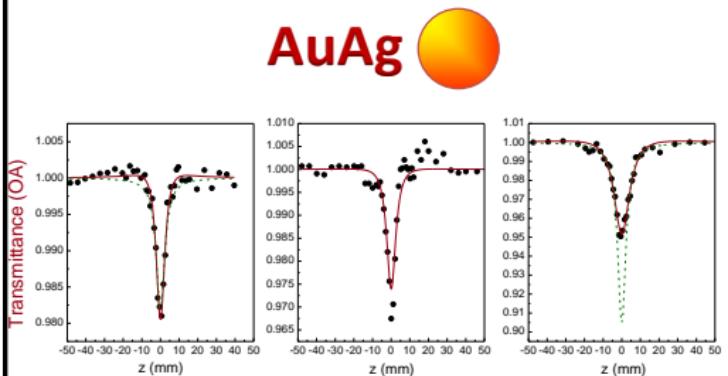
- $E \cong 30 \text{ mJ} (@ 1064 \text{ nm})$
- $\Delta t \cong 15 \text{ ps}$ **pure electronic effects!**
- $RR = 10 \text{ Hz}$
- $\lambda = 420 - 2000 \text{ nm}$ **tunability!**



Open Aperture results

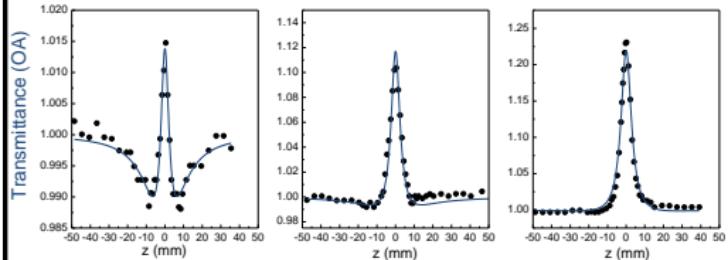


Open Aperture results



$$I \propto |E|^2$$

Effective intensity



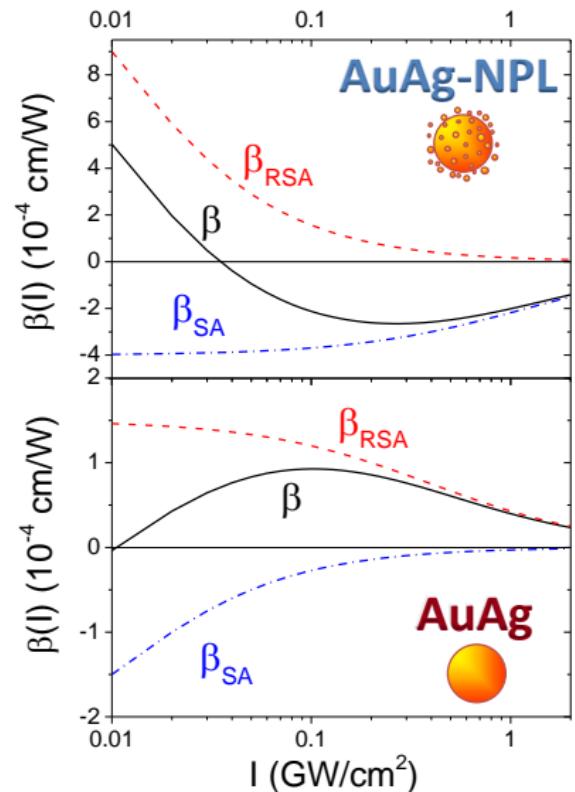
Nonlinear absorption

$$\alpha(I) = \alpha_0 + \beta I$$

Nonlinear absorption coefficient

$$\beta(I) = \beta_{RSA}(I) + \beta_{SA}(I) = \frac{\beta_+}{1 + \frac{I}{I_s^+}} + \frac{\beta_-}{1 + \frac{I}{I_s^-}}$$

	AuAg	AuAg-NPL
$\beta_+ (10^{-3} \text{ cm/W})$	0.15 ± 0.03	1.9 ± 0.7
$I_s^+ (10^8 \text{ W/cm}^2)$	4 ± 2	0.09 ± 0.04
$\beta_- (10^{-4} \text{ cm/W})$	-3 ± 1	-4 ± 1
$I_s^- (10^9 \text{ W/cm}^2)$	0.010 ± 0.001	1.2 ± 0.1



Nonlinear absorption

