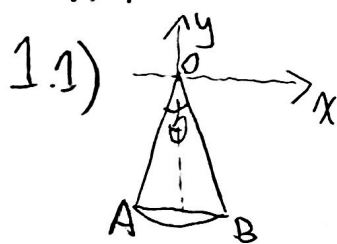


PHY 242 A2



Take O as origin,  $x$  being // to AB,  $y \perp$  to  $x$

No need to consider  $z$  in this question

$$\theta = \angle AOB = \arccos\left(\frac{OA^2 + OB^2 - AB^2}{2 \cdot OA \cdot OB}\right) = \arccos\left(\frac{2R^2 - 2d}{2R^2}\right)$$

$$= \arccos\left(1 - \frac{d}{R^2}\right)$$

Now we notice that due to symmetry of the ring over O, the field will be 0 if the arc  $\widehat{AB}$  is closed

Namely  $\vec{E} + \vec{E}_{\widehat{AB}} = 0 \Rightarrow \vec{E} = -\vec{E}_{\widehat{AB}}$

Consider the cylindrical system, where O is the origin,  $\hat{r}$  is  $\frac{\vec{R}}{R}$ , and  $\hat{\phi}$  is counter-clockwise.  $\psi$  is the angle from OC, where C is



mid point of AB,  $\hat{\phi}$  in  $-\hat{y}$  direction

Then a point  $(r, \psi)$  in cylinder is  $(-\sin\psi \cdot r, -\cos\psi \cdot r)$  in Cartesian system

Again Notice that due to the symmetry of the  $\widehat{AB}$  arc on the  $x$ -axis over O, there will be no  $x$ -component of  $\vec{E}$

Namely,  $\vec{E}_{\widehat{AB}} = E_y \hat{y}$

Given any infinitesimal arc  $d\vec{r}$  on  $\widehat{AB}$ , which has  $(R, \psi)$  as position, and  $d\vec{r} \cdot \hat{\phi}$  as charge,  $E_{y(\psi)} = \frac{1}{4\pi\epsilon_0} \cdot \frac{d\vec{r} \cdot (-\cos\psi \cdot \hat{y})}{R^2}$

Note  $\lambda = \frac{q}{L} = \frac{q}{(2\pi - \theta)R}$ ,  $dr = R \cdot d\psi$

$$\begin{aligned} \therefore \vec{E}_{\widehat{AB}} = \vec{E}_y &= \int_{-\frac{\theta}{2}}^{\frac{\theta}{2}} \frac{1}{4\pi\epsilon_0} \cdot \frac{\cos\psi \cdot d\psi \cdot R \cdot \lambda}{R^2} \hat{y} = \frac{\lambda \hat{y}}{4\pi\epsilon_0 R} \int_{-\frac{\theta}{2}}^{\frac{\theta}{2}} \cos\psi \, d\psi \\ &= \frac{\lambda \hat{y}}{4\pi\epsilon_0 R} \cdot \sin(\psi) \Big|_{-\frac{\theta}{2}}^{\frac{\theta}{2}} \quad \text{Note } \sin\left(\frac{\theta}{2}\right) = \frac{d}{R}, \sin\left(-\frac{\theta}{2}\right) = -\frac{d}{R} \\ &= \frac{\lambda}{4\pi\epsilon_0 R} \cdot \frac{2d}{R} \hat{y} = 14.36 \hat{y} \Rightarrow \vec{E} = -\vec{E}_{\widehat{AB}} = -14.36 \hat{y} \end{aligned}$$

2) 1 cm is very small comparing to 1 m

Again observe the symmetry on x-direction

The arc length of  $\widehat{AB}$  can be approximated by length  $\overline{AB} = 2d$

We can view  $\widehat{AB}$  as a point sitting below O at R distance, carrying

$$\text{charge } \frac{q}{2\pi R} \cdot 2d = \frac{qd}{\pi R}$$

$$\text{Thus } \vec{E}_{AB} = \frac{qd}{\pi R \cdot R} \cdot \frac{1}{4\pi\epsilon_0} = \frac{qd}{4\pi\epsilon_0 R^2} \hat{y} = 14.32 \hat{y}$$

$$\vec{E} = -14.32 \hat{y}$$

2.1) Consider Cartesian System, with origin at center of the shell,  
 $(x, y, z)$  be any three  $\perp$  directions

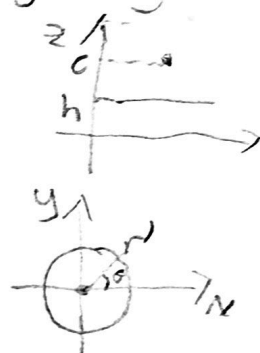
Given any point  $(a, b, c)$  in this coordinate.

Consider any ring of radius  $r'$ , at the plane  $z=h$ , centered on  $z$  axis

The field from any infinitesimal arc  $dr$  on the ring is given by

$$\left\{ \begin{aligned} \vec{E}_x &= \frac{\lambda dr \cdot (a - r' \cos \theta)}{4\pi\epsilon_0 [(a - r' \cos \theta)^2 + (b - r' \sin \theta)^2 + (c - h)^2]^{3/2}} \hat{x} \\ \vec{E}_y &= \frac{\lambda dr}{4\pi\epsilon_0} \cdot \frac{b - r' \sin \theta}{d^{3/2}} \hat{y} \\ \vec{E}_z &= \frac{\lambda dr}{4\pi\epsilon_0} \cdot \frac{c - h}{d^{3/2}} \hat{z} \end{aligned} \right.$$

where  $\lambda$  is charge per length



where  $\theta$  = angle from  $\hat{x}$  on  $x$ - $y$  plane,  $dr = r' d\theta$

$$\therefore \vec{E}_x(h) = \int_0^{2\pi} \vec{E}_x(h, \theta) d\theta = \frac{\lambda r'}{4\pi\epsilon_0} \int_0^{2\pi} \frac{a - r' \cos \theta}{d(\theta)^{3/2}} d\theta$$

Notice that if  $a=b=0$ , namely the point is ON  $z$ -axis, then the  $x, y$ -components are all 0 due to symmetry, and we are left

$$\begin{aligned} \text{with } \vec{E}(h) &= \vec{E}_z(h) = \frac{\lambda}{4\pi\epsilon_0} \cdot \int_0^{2\pi} \frac{c-h}{[(r' \cos \theta)^2 + (r' \sin \theta)^2 + (c-h)^2]^{3/2}} \cdot r' d\theta \cdot \hat{z} \\ &= \frac{\lambda}{4\pi\epsilon_0} \cdot \frac{2\pi r' (c-h)}{(r'^2 + (c-h)^2)^{3/2}} \hat{z} \end{aligned}$$

Again Notice that if we break the shell into infinite rings  $\perp z$ ,  
the radius  $r'$  at  $z=h$  is  $r' = \sqrt{R^2 - h^2}$

$$\therefore \vec{E}(h) = \frac{\lambda}{2\epsilon_0} \cdot \frac{\sqrt{R^2 - h^2} (c-h)}{(R^2 - h^2 + (c-h)^2)^{3/2}} \hat{z}, \text{ for a point } (0,0,c), \text{ where } \lambda = \rho \cdot d$$



$$\text{Then } \vec{E} = \int_{-R}^R \vec{E}(h) = \frac{\rho}{2\epsilon_0} \int_{-R}^R \frac{\sqrt{R^2 - h^2} (c-h)}{(R^2 - h^2 + (c-h)^2)^{3/2}} dh \cdot \hat{z} = \frac{\rho}{2\epsilon_0} \cdot \frac{2R^2}{c^2} \hat{z} = \frac{\rho R^2}{\epsilon_0 c^2} \hat{z}$$

$\rho$  is charge per area

Note:  $\rho = q/4\pi R^2$

$$\therefore \vec{E} = \frac{q}{4\pi R^2} \cdot \frac{R^2}{\epsilon_0 c^2} \hat{z} = \frac{q}{4\pi \epsilon_0 c^2} \hat{z}$$

Now, note that given any point, if we take  $\hat{z}_{axis}$  to be the line from center of shell to the point, the point is always in the form  $(0,0,c)$

$$\therefore \vec{E} = \frac{q}{4\pi \epsilon_0 \cdot c^2} \hat{z}, \text{ where } c \text{ is distance from center to point, and } z\text{-axis is described as above}$$

b) Given a point charge, the field  $\vec{E} = \frac{q}{4\pi \epsilon_0 d^2} \hat{d}$ , where  $\hat{d}$  is the line from charge to the point in question

Is same as from a)