

QIP710: Assignment 5

Xiuzhe Luo (20812697)

Some properties of Shor's 9-qubit code

1.a

no, because a counterexample is that if there is a Z-error occurs at 2nd qubit, the syndrome is the same as occurring on 1st qubit, e.g if the input state is $\frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$, both single qubit Z-error on 1st qubit and 2nd qubit give the same syndrome 00100100.

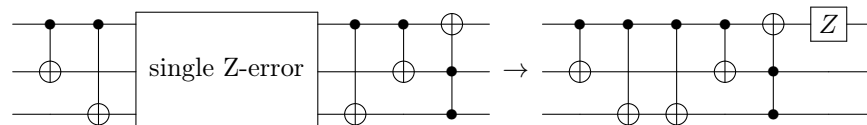
* I just check this by a script, the script is attached in Appendix.

But use the conclusion I prove in the next question, we can also show that the effect of Z-error in the same inner block on its syndrome is always to create a phase/or nothing on the syndrome instead of flip it, so we always get the same bitstring.

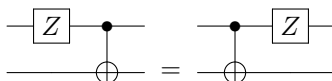
1.b

there can be one X-error occurs in the 3 blocks of the inner circuit, we pick 2 of the 3 which has 3 possibilities, then each of them can happen at 3 possible position for each block, thus $3 \times 3 \times 3 = 27$ weight 2 Pauli-X error in total.

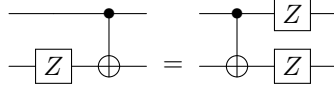
For Z error, it can only occur inside one block, this is because the inner circuit leave Z error intact and if there is a Z-error it always moves it to the first qubit (the protected qubit), when the ancilla qubits are $|00\rangle$ (note it is not equivalent, I omit other possible Z gate on 2, 3 wires since they don't change the protected qubit, so I'm not using $=$, but \rightarrow)



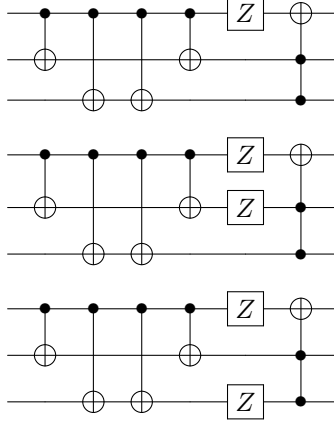
this is a statement from the slide, we can prove this easily by applying the swap rule between Pauli Z and CNOT, note we have



and



thus after we move the error gates from left to right, we always have one Pauli Z on the protected qubit.



when the ancilla qubits are $|00\rangle$, the Toffoli gate has no effect. thus if there are two Z error occurs in the inner block, we get two Z gate on the first qubit, which cancels. But if two Z error occurs at different inner block, then the outer block has to correct 2 Z-error on 1,4,7, which is not possible. Thus we can have $3 \times 3 = 9$ kinds of weight 2 Z errors.

Regarding to Y-errors, it can only happen when both Z and X error could happen, which means there cannot be double Y-error since X cannot be in the same block, while Z has to be in the same block.

Now the ZX-error can happen when single Z-error and single X-error can be corrected, which is anywhere on the qubits as long as they are not on the same qubit (since we are asking for weight 2 errors), thus we have $9 \times 8 = 72$ kinds of ZX-error.

ZY-error can only happen in the same block since single X-error can be corrected, but we cannot have two Z error in different block, thus we have $3 \times 6 = 18$ kinds of ZY-error.

XY-error can only happen in different block since single Z-error can be corrected, but we cannot have two X error in the same block, thus we have $6 \times 3 \times 3 = 54$ kinds of XY-error.

Thus we have $27 + 9 + 72 + 18 + 54 = 180$ kinds of weight 2 error.

1.c

since if Z error occurs at 2,3,5,6,8,9 which cancel itself on each protected qubit, and X error occurs at 1,4,7 which will be corrected by each inner block, the shor code corrects a weight 9 error.

Hadamard transform on uniform superposition of affine linear space

assume G is the $k \times n$ generator matrix for x , for $y \notin C^\perp$, when $Gy \neq 0$

$$\sum_{x \in C} (-1)^{x \cdot y} = \sum_{k \in \{0,1\}^k} (-1)^{aGy} = \sum_{k \in \{0,1\}^k} (-1)^{ab} = 0$$

since $y \notin C^\perp$, G can be viewed as the parity check matrix for C , thus $Gy \neq 0$. this also indicates that $|C||C^\perp| = 2^n$ So we have

$$\begin{aligned} H^{\otimes n} \left(\frac{1}{\sqrt{C}} \sum_{x \in C} |x + z\rangle \right) &= \frac{1}{\sqrt{C}} \sum_{x \in C} H^{\otimes n} |x + z\rangle \\ &= \frac{1}{\sqrt{C}} \frac{1}{2^{n/2}} \sum_{x \in C} \sum_{k \in \{0,1\}^n} (-1)^{x \cdot k} (-1)^{z \cdot k} |k\rangle \\ &= \frac{1}{\sqrt{|C^\perp|}} \sum_{k \in C^\perp} (-1)^{z \cdot k} |k\rangle + \frac{1}{\sqrt{C}} \frac{1}{2^{n/2}} \sum_{k \notin C^\perp} \sum_{x \in C} (-1)^{x \cdot k} |k\rangle \\ &= \frac{1}{\sqrt{|C^\perp|}} \sum_{k \in C^\perp} (-1)^{z \cdot k} |k\rangle \end{aligned}$$

Is the transpose a valid quantum operation?

3.a

this means $\text{tr}(\rho \rho^T) = 0$, $|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle)$ satisfy this condition, the output pure state is $\frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle)$

3.b, c

quantum channel always maps the density matrix which is a semi-definite matrix to another density matrix which is also a semi-definite matrix. However, if the transposition is a valid quantum channel, there exists a counterexample, let's apply this channel on the first qubit of a Bell state $|00\rangle + |11\rangle$:

$$\Lambda(|00\rangle\langle 00| + |11\rangle\langle 11| + |00\rangle\langle 11| + |11\rangle\langle 00|) = |00\rangle\langle 00| + |11\rangle\langle 11| + |10\rangle\langle 01| + |01\rangle\langle 10|$$

for convenience we omit the normalization factor, the eigenvalue of the output is not all non-negative (by calculating this numerically, there is one negative eigenvalue -1), thus it is not a density matrix, thus there is no such channel for all ρ .

Searching with good guessing algorithm

4.a

assume we sample k times, the probability of getting a satisfying assignment is

$$\begin{aligned} P(n) &= \sum_{k=1}^n p(1-p)^{k-1} = p \frac{1 - (1-p)^n}{p} = 1 - \left(\frac{3}{4}\right)^n \geq 1 - \epsilon \\ \epsilon &\geq \left(\frac{3}{4}\right)^n \\ \ln(\epsilon) &\geq n \ln\left(\frac{3}{4}\right) \\ n &\geq \frac{\ln(\epsilon)}{\ln(\frac{3}{4})} \end{aligned}$$

4.b

(note I'm using a different notation of x and y comparing to the question, the x in question is the k here, just for convenience)

denote summation the states satisfy f as $|x\rangle = 2 \sum_k A_k |k\rangle$, $\sum_k |A_k|^2 = 1/4$ and doesn't satisfy f : $|y\rangle = \frac{2}{\sqrt{3}} \sum_k B_k |k\rangle$, $\sum_k |B_k|^2 = 3/4$, we have

$$|\omega\rangle = U_g |0^n\rangle = \frac{1}{2} |x\rangle + \frac{\sqrt{3}}{2} |y\rangle$$

since $|x\rangle$ and $|y\rangle$ are on different basis, they are orthogonal, and we have

$$U_f |x\rangle |0\rangle = |x\rangle |1\rangle$$

$$U_f |x\rangle |1\rangle = |x\rangle |0\rangle$$

$$U_f |y\rangle |0\rangle = |y\rangle |0\rangle$$

$$U_f |y\rangle |1\rangle = |y\rangle |1\rangle$$

$$\begin{aligned} &U_f \left(\frac{1}{2} |x\rangle |-\rangle + \frac{\sqrt{3}}{2} |y\rangle |-\rangle \right) \\ &= -\frac{1}{2} |x\rangle |-\rangle + \frac{\sqrt{3}}{2} |y\rangle |-\rangle \end{aligned}$$

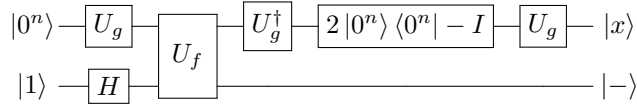
Thus U_f reflects the state on $|x\rangle$ axis. Now we create the other reflection operator $2|\omega\rangle\langle\omega| - I = U_g(2|0\rangle\langle 0| - I)U_g^\dagger$, and we have

$$\begin{aligned}
\langle \omega | x \rangle &= \frac{1}{2} \\
\langle \omega | y \rangle &= \frac{\sqrt{3}}{2} \\
(2 |\omega\rangle \langle \omega| - I) |x\rangle &= |\omega\rangle - |x\rangle = -\frac{1}{2} |x\rangle + \frac{\sqrt{3}}{2} |y\rangle \\
(2 |\omega\rangle \langle \omega| - I) |y\rangle &= \sqrt{3} |\omega\rangle - |y\rangle = \frac{\sqrt{3}}{2} |x\rangle + \frac{1}{2} |y\rangle
\end{aligned}$$

thus apply the reflection after oracle we have

$$\begin{aligned}
&(2 |\omega\rangle \langle \omega| - I) \left(-\frac{1}{2} |x\rangle + \frac{\sqrt{3}}{2} |y\rangle \right) |-\rangle \\
&= \left(\frac{1}{4} |x\rangle - \frac{\sqrt{3}}{4} |y\rangle + \frac{3}{4} |x\rangle + \frac{\sqrt{3}}{4} |y\rangle \right) |-\rangle \\
&= |x\rangle |-\rangle
\end{aligned}$$

which is the answer! The circuit looks like the following



it outputs the answer with only one query without any error.

Characterizing

5.a

assume we have some two qubit state below

$$|\psi\rangle = A_1 |00\rangle + A_2 |01\rangle + A_3 |10\rangle + A_4 |11\rangle$$

on Hadamard basis we have

$$\begin{aligned}
2H \otimes H |\psi\rangle &= (A_1 + A_2 + A_3 + A_4) |00\rangle + \\
&\quad (A_1 - A_2 + A_3 - A_4) |01\rangle + \\
&\quad (A_1 + A_2 - A_3 - A_4) |10\rangle + \\
&\quad (A_1 - A_2 - A_3 + A_4) |11\rangle
\end{aligned}$$

in order to have $a = b$ on both basis, we need $A_2 = A_3 = 0$ and $(A_1 - A_2 + A_3 - A_4) = (A_1 + A_2 - A_3 - A_4) = 0$ thus, $A_1 = A_4 = C$, so the state becomes $C(|00\rangle + |11\rangle)$

5.b

similarly we have

$$\begin{aligned} 2I \otimes H |\psi\rangle &= (A_1 + A_2) |00\rangle + (A_1 - A_2) |01\rangle + (A_3 + A_4) |10\rangle + (A_3 - A_4) |11\rangle \\ 2H \otimes I |\psi\rangle &= (A_1 + A_3) |00\rangle + (A_2 + A_4) |01\rangle + (A_1 - A_3) |10\rangle + (A_2 - A_4) |11\rangle \end{aligned}$$

thus $A_1 = A_2 = A_3 = C, A_4 = -C$, we get $\frac{1}{2}(|00\rangle + |01\rangle + |10\rangle - |11\rangle)$

You will need the latest develop version of Yao to run this, which can be installed with `add Yao#master` in Julia's `pkg` mode.

```
using Yao

shor(Es...) =shor(chain(Es...))

shor(E) =chain(9,
    cnot(1, 4),
    cnot(1, 7),

    put(1=>H), put(4=>H), put(7=>H),

    cnot(1,2), cnot(1,3),
    cnot(4,5), cnot(4,6),
    cnot(7,8), cnot(7,9),
    # error
    E,

    cnot(1,3), cnot(1,2),
    cnot((2, 3), 1),

    cnot(4,6), cnot(4,5),
    cnot((5, 6), 4),

    cnot(7,9), cnot(7,8),
    cnot((8, 9), 7),

    put(1=>H), put(4=>H), put(7=>H),
    cnot(1, 7), cnot(1, 4), cnot((4, 7), 1)
)

# aassignment answers
# 1.a
@vars αβ
r =α*ket"0" +β*ket"1" |>addbits!(8) |>shor(put(1=>Z)) |>partial_tr(1) |>expand
# output: α+ β|00100100>
r =α*ket"0" +β*ket"1" |>addbits!(8) |>shor(put(2=>Z)) |>partial_tr(1) |>expand
# α+ β|00100100>

# 1.c
r =α*ket"0" +β*ket"1" |>addbits!(8) |>
    shor(kron(1=>X, 2=>Z, 3=>Z, 4=>X, 5=>Z, 6=>Z, 7=>X, 8=>Z, 9=>Z)) |>
    partial_tr(2:9) |>expand
# output: α|0> + β|1>

# question 5
@vars A1 A2 A3 A4
psi =A1*ket"00" +A2 *ket"01" +A3 *ket"10" +A4 *ket"11"
# 5.a
copy(psi) |>kron(1=>H, 2=>H) |>expand
# 5.b
copy(psi) |>put(2, 1=>H) |>expand
copy(psi) |>put(2, 2=>H) |>expand
```