

# Multipartite Operator Entanglement & The AME Problem

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# Outline

- 1 Bipartite Operator Entanglement : A Review
- 2 Operator Entanglement Results
- 3 AME Results

# Motivation

- The study of the entangling capabilities of quantum operators was thrust into the limelight by the importance of entanglement.
- When dealing with operators & entanglement, the questions are two-fold.
  - ① **Entangling power** ( $e_p(U)$ ):- Average entanglement **created** by the operator when acting on product states. By convention, we take the average with respect to the Haar measure.
  - ② **Operator entanglement** :- Akin to the entanglement of states, operators live in the Hilbert-Schmidt space and thus, can be 'entangled' as well. Operator entanglement quantities are a measure of the '*entangled-ness*' of operators.
- Operator entanglement of the time-evolution operator characterizes information scrambling and delocalization. Thus, it can be used to characterize the many-body localized (MBL) phase and the random singlet phase (RSP)- **i.e. Operator entanglement is useful.**

# Bipartite Operator Entanglement : A Natural Extension of State Entanglement

- A 2-party state is entangled iff. it has more than 1 Schmidt coefficient.
- Binary test of entanglement.
- The same test for entanglement can be performed for bipartite operators through the Schmidt decomposition of an operator  $U_{AB} \in \mathcal{H}_{HS}^d \otimes \mathcal{H}_{HS}^d$ .

## Qualitative Test for Operator Entanglement

$$U_{AB} = \sum_{i=1}^{d^2} \sqrt{\lambda_i} m_i^A \otimes m_i^B, \text{ with } \sum_{i=1}^{d^2} \lambda_i = d^2.$$

$U_{AB}$  is unentangled iff.  $\exists$  **only one** Schmidt coefficient.

# Measures of Bipartite operator Entanglement

$\{\frac{\lambda_i}{d^2}\}$  forms a probability distribution and are local unitary invariant (LU-invariant).

## Quantitative Measures of Bipartite Operator Entanglement

$$E(U) = 1 - \frac{1}{d^4} \sum_{i=1}^{d^2} \lambda_i^2, E(U) \in [0, 1 - 1/d^2].$$

$$E(US) = 1 - \frac{1}{d^4} \sum_{i=1}^{d^2} \mu_i^2, E(US) \in [0, 1 - 1/d^2].$$

**Note:**—  $\{\lambda_i\}$  &  $\{\mu_i\}$  are local unitary invariant. The utility of these particular measures is in their connection to the *entangling power*.

$$e_p(U) = \frac{1}{E(S)} * (E(U) + E(US) - E(S)),$$

$$e_p \in [0, 1].$$

# The Choi-Jamialkowski Isomorphism (CJ Isomorphism)

## The CJ Isomorphism

An operator,  $X$  on  $\mathcal{H}_A^d$  can be mapped to a state  $|X\rangle \in \mathcal{H}_A^d \otimes \mathcal{H}_{A'}^d$ , as:—

$$\begin{aligned} |X\rangle_{AA'} &= \sum_{ij} X_{ij} |ij\rangle = \frac{1}{\sqrt{d}} \sum_{ij} \langle i|X|j\rangle |i\rangle |j\rangle = \frac{1}{\sqrt{d}} \sum_j X |jj\rangle \\ &= (X \otimes \mathbb{I}) |\Phi^+\rangle_{AA'}, \quad |\Phi^+\rangle := \frac{1}{\sqrt{d}} \sum_{i=0}^{d-1} |ii\rangle \quad \& \quad \langle i|X|j\rangle := \sqrt{d} \cdot X_{ij}. \end{aligned}$$

This isomorphism corresponds to the lexicographical vectorization of the operator,  $X$  (Row-wise vectorization).

# Illustration

## Lexicographic Vectorization

Consider  $\mathbb{I}_2$ . From Eq. 1,  $|X\rangle_{AA'} = |\Phi^+\rangle_{AA'}$ .

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \Rightarrow [1 \quad 0 \quad 0 \quad 1]$$

In general, for an operator  $X$  on  $\mathcal{H}_A^d$  :—

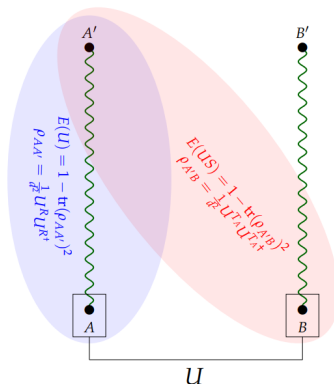
$$\begin{bmatrix} x_{11} & \cdots & x_{1d} \\ \vdots & \ddots & \vdots \\ x_{d1} & \cdots & x_{dd} \end{bmatrix} \Rightarrow [x_{11} \dots x_{1d} \quad x_{21} \dots x_{2d} \dots x_{d1} \dots x_{dd}]$$

# The Modified CJ Isomorphism for Bipartite Operators

## The CJ Isomorphism

Consider a bipartite unitary,  $U_{AB}$  on A & B of local dimension d. It will be isomorphic to a 4-party state in  $\mathcal{H}_A^d \otimes \mathcal{H}_{A'}^d \otimes \mathcal{H}_B^d \otimes \mathcal{H}_{B'}^d$ , defined as:-

$$|U\rangle_{AA'BB'} = (U_{AB} \otimes \mathbb{I}_{A'B'})|\Phi^+\rangle_{AA'}|\Phi^+\rangle_{BB'}.$$





# Restating Operator Entanglement Entropies Using State Bipartitions

$AB \mid A'B'$

$$\rho_{AB} = \frac{1}{d^2} UU^\dagger = \frac{1}{d^2} \mathbb{I}_{d^2} \rightarrow \text{Maximally Entangled Bipartition}$$

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$AA' \mid BB'$

$$\rho_{AA'} = \frac{1}{d^2} U^R U^{R\dagger},$$

$$S_2(\rho_{AA'}) = 1 - \frac{1}{d^4} \text{Tr} \left( (U^R U^{R\dagger})^2 \right) \equiv E(U),$$

where  $\langle i_1 j_1 | U^R | i_2 j_2 \rangle := \langle i_1 i_2 | U | j_1 j_2 \rangle$ .

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$A'B \mid AB'$

$$\begin{aligned} \rho_{A'B} &= \frac{1}{d^2} U^T U^{T\dagger} \\ S_2(\rho_{A'B}) &= 1 - \frac{\text{Tr} \left( (U^T U^{T\dagger})^2 \right)}{d^4} = E(US), \\ \text{where } \langle j_1 i_2 | U^T | i_1 j_2 \rangle &:= \langle i_1 i_2 | U | j_1 j_2 \rangle. \end{aligned}$$

# Duals, T-duals & 2-Unitaries : Definition & Maps

## Unitaries that Maximize the Linear Operator Entanglements

Dual Unitaries (Duals) : Unitary operators :  $E(U) = 1 - \frac{1}{d^2}$ .

T-Dual Unitaries (T-Duals) : Unitary operators :  $E(US) = 1 - \frac{1}{d^2}$ .

2-Unitaries : Unitary operators that maximize both  $E(U)$  &  $E(US)$ .

- ① Duals are period-2 fixed points of the realignment operation ( $U \mapsto U^{\mathcal{R}}$ ).
- ② T-duals are period-2 fixed points of the partial-trace operation ( $U \mapsto U^{\Gamma}$ ).
- ③ 2-unitaries are period-3 fixed points of the realignment-partial trace operation ( $U \mapsto (U^{\mathcal{R}})^{\Gamma}$ ).

The fact that duals, T-duals & 2-unitaries are fixed points of the corresponding operator reshaping is used to define non-linear maps that generate these unitary operators.

# Summary of Operator Entanglement Results

- The Multipartite Operator-State Isomorphism
- Connection Between the Reduced Density Operators & The Matrix Reshapings

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- Maximally Entangled Multipartite Unitary Operators
- The Bipartition Theorem for Operator Entanglement
- Multipartite Operator Entanglement Entropies
- Explorations of Operator Entanglement of Tripartite Unitary Operators
- Towards the Entangling Powers of Multipartite Unitary Operators



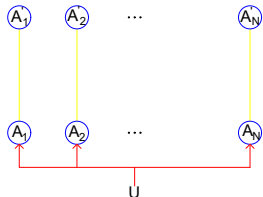
# The Multipartite Operator-State Isomorphism

- N-partite unitary,  $U_{A_1 A_2 \dots A_N}$ , defined over  $\mathcal{H}_{A_1}^d \otimes \mathcal{H}_{A_2}^d \otimes \dots \otimes \mathcal{H}_{A_N}^d$ .
- Isomorphic to a  $2N$ -party state in  $\mathcal{H}_{A_1}^d \otimes \mathcal{H}_{A'_1}^d \dots \otimes \mathcal{H}_{A_N}^d \otimes \mathcal{H}_{A'_N}^d$ .

$$U_{A_1 \dots A_N} \mapsto |U\rangle_{A_1 A'_1 \dots A_N A'_N}$$

$$|U\rangle_{A_1 A'_1 \dots A_N A'_N} := U_{A_1 A_2 \dots A_N} \otimes \mathbb{I}_{A'_1 A'_2 \dots A'_N} |\Phi^+\rangle_{A_1 A'_1} |\Phi^+\rangle_{A_2 A'_2} \dots |\Phi^+\rangle_{A_N A'_N} \quad (1)$$

$$\langle i_1 j_1 i_2 j_2 \dots i_N j_N | U \rangle_{A_1 A'_1 \dots A_N A'_N} = \frac{\langle i_1 i_2 \dots i_N | U_{A_1 \dots A_N} | j_1 j_2 \dots j_N \rangle}{d^{N/2}} \quad (2)$$



# Reduced Density Operators & Matrix Reshapings

- An N-partite unitary,  $U_{A_1 A_2 \dots A_N}$  isomorphic to  $|U\rangle_{A_1 A'_1 \dots A_N A'_N}$ .
- Consider the bipartition  $A_1 \dots A_p A'_{a_1} \dots A'_{a_q} | A_{p+1} \dots A_N A'_{a_{q+1}} \dots A'_{a_N}$ , where  $\{a_r\}_{r=1}^N$  is a permutation of  $\{r\}_{r=1}^N$ .

The Bipartition -  $A_1 \dots A_p A'_{a_1} \dots A'_{a_q} | A_{p+1} \dots A_N A'_{a_{q+1}} \dots A'_{a_N}$

$$\rho_{A_1 A'_1 \dots A_N A'_N} = |U\rangle\langle U|$$

$$\rho' := \text{Tr}_{A_{p+1} \dots A_N A'_{a_{q+1}} \dots A'_{a_N}} [\rho_{A_1 A'_1 \dots A_N A'_N}] = \frac{U^{(x)} \cdot U^{(x)\dagger}}{d^N}, \quad (3)$$

$$\langle i_1 \dots i_p j_{a_1} \dots j_{a_q} | U^{(x)} | i_{p+1} \dots i_N j_{a_{q+1}} \dots j_{a_N} \rangle = \langle i_1 \dots i_N | U | j_1 \dots j_N \rangle. \quad (4)$$

# k-Uniform Unitary Operators

## Definition

A unitary operator,  $U$  on  $(N,d)$  is termed a  $k$ -uniform unitary operator if its vectorization under the multipartite operator-state isomorphism,

$$|U\rangle_{A_1 A'_1 \dots A_N A'_N} = U_{A_1 A_2 \dots A_N} \otimes \mathbb{I}_{A'_1 A'_2 \dots A'_N} |\Phi^+\rangle_{A_1 A'_1} |\Phi^+\rangle_{A_2 A'_2} \cdots |\Phi^+\rangle_{A_N A'_N},$$

is a  $k$ -uniform state on  $2N$  parties.

- An  $\lfloor \frac{N}{2} \rfloor$ -uniform state of  $N$  parties is an AME state.
- An  $\lfloor \frac{N}{2} \rfloor$ -uniform unitary operator is a multi-unitary operator (Perfect tensor of  $(N,d)$ )
- To stress upon the generality of our definition of a  $k$ -uniform unitary, we conjecture that if  $\exists$  a multipartite operator-state isomorphism on  $U$  that produces a  $k'$ -uniform state,  $k' \leq k$ .

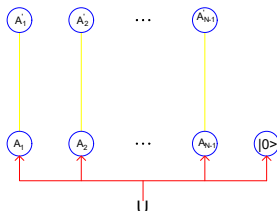
# Existence of k-Uniform States

## Construction of k-Uniform States for an Odd Number of Parties

$U$  is a  $k$ -uniform unitary operator on  $N$  parties, of local dimension  $d$ .

$$|U\rangle_{A_1 A'_1 \dots A_{N-1} A'_{N-1} A_N} = U_{A_1 \dots A_N} \otimes \mathbb{I}_{A'_1 \dots A'_{N-1}} |\Phi^+\rangle_{A_1 A'_1} \cdots |\Phi^+\rangle_{A_{N-1} A'_{N-1}} |0\rangle_{A_N} \quad (5)$$

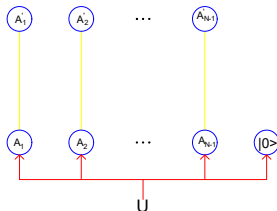
$\rightarrow (k-1)$ -uniform state on  $2N-1$  parties.



# Implications

## Implications

- If  $\exists$  a  $k$ -uniform state in  $(2N, d)$ ,  $\exists$  a  $(k-1)$ -uniform state in  $(2N-1, d)$ .
- If  $\nexists$  a  $(k-1)$ -uniform state in  $(2N-1, d)$ ,  $\nexists$  a  $k$ -uniform state in  $(2N, d)$ .
- If  $\text{AME}(2N, d) \exists$ ,  $\text{AME}(2N-1, d)$  exists.
- If  $\text{AME}(2N-1, d) \nexists$ ,  $\text{AME}(2N, d)$  does **not** exist.



# Maximally Entangled Multipartite Unitary Operators

Analogous to the definition of maximal entanglement in quantum states, we can define maximally entangled multipartite unitary operators.

## Maximal Operator Entanglement

$U$  is a unitary operator on  $N$  parties -  $A_1 \dots A_N$ .

$S \subset \{1, \dots, N\}$ .

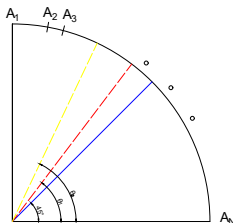
Then,  $U_{A_1 \dots A_N}$  is maximally entangled iff. the Schmidt coefficients of its decomposition into *any* partition -  $S|S^c$  - are **all equal** to  $d^{N-2|S|}$ .  
(wlog  $|S| \leq |S^c|$ )

- Similar to the bipartite case, the Schmidt values for any bipartite decomposition of  $U$  is related to the *singular values* of a reshaping of  $U$  of size  $d^{|S|} \times d^{2N-|S|}$ .
- This correspondence along with its derivation has been shown in my thesis.

# Bipartition Theorem for Operator Entanglement

## Bipartition Theorem

If  $U_{A_1 \dots A_N}$  is maximally entangled for a bipartition  $S|S^c$ , with  $|S| \leq |S^c|$ , then  $U$  is maximally entangled for the bipartition -  $S'|S'^c$ , for any  $S' \subset S$ .



# The Balanced Bipartition Theorem for Maximal Operator Entanglement

## Implication of the Bipartition Theorem

If  $U_{A_1 \dots A_N}$  is maximally entangled for all balanced bipartitions,  $S|S^c$  - i.e. where  $|S| = \lfloor N/2 \rfloor$ , then  $U$  is a maximally entangled unitary operator.

Thus, similar to AME states where only the entanglement in balanced bipartitions needs to be maximized to ensure that any reduction to  $\lfloor N/2 \rfloor$  parties or less should be maximally mixed, only the Schmidt decomposition of  $U$  wrt balanced bipartitions is important for maximal operator entanglement.



# Multipartite Operator Entanglement Entropies

For a bipartite unitary operator,  $U$ , the  $E(U)$  measure is a measure of the operator entanglement of  $U$ .

If  $E(U)$  is maximum, Schmidt coefficients are **equal** to 1.

As  $N \uparrow$ , the number of Schmidt decompositions possible  $\uparrow$ . Thus, we have to define as many operator entanglements of  $U$ . i.e.

## Operator Entanglement Entropy

$$U_{A_1 \dots A_N} = \sum_{i=1}^{d^{2*|S|}} \sqrt{\lambda_i} m_i^S \otimes m_i^{S^c},$$

$$E_S(U) := 1 - \frac{1}{d^N} * \sum_{i=1}^{d^{2*|S|}} \lambda_i^2$$

**Normalization Condition:-**  $\sum_{i=1}^{d^{2*|S|}} \lambda_i = d^N$ .

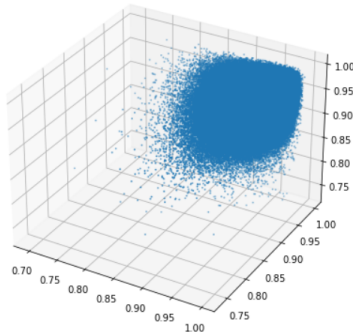
To obtain an exhaustive and complete set of operator entanglements, we must consider all **unique** Schmidt decompositions where the cardinality of  $S$  is  $\lfloor N/2 \rfloor$ .

# Future Directions

$N=3$  marks a transition from the entanglement features of bipartite unitary operators to that of multipartite unitary operators.

Within  $N=3$ ,  $d=2$  is the simplest non-trivial case. We study this in detail in the thesis.

For  $N=3$ , there are 3 possible bipartitions -  $A|BC$ ,  $B|AC$  &  $C|AB$  - and consequently, 3 operator entanglements. We study the distribution of these operator entanglements over CUE.



# Future Directions

- The multipartite linear operator entanglement of the unitary only characterizes part of its entanglement properties. We need to find the equivalences of  $E(US)$ ,  $e_p(U) \dots$  for a general operator on  $(N,d)$ .
- Determine the operator-state equivalences of the linear operator entanglement entropies —apart from  $E_x(U)$ .
- Co-opt the non-linear maps to generate maximally entangled multipartite operators.
- Explore unequal dimensions.



## Results Related to AME States

- Explored the efficiency of the non-linear maps.
- Found 2 classes of unitary operators in  $U(9)$  for which the  $M_{TR}$  map does not converge. (Convergence rate for CUE - 94%) [Shrigyan Brahmachari]
- Proved that **all** fixed points of the  $M_R$  &  $M_T$  maps are period-2. [Shrigyan Brahmachari]
- Developed a map to combine AME states of lower  $N$  to form  $k$ -uniform states of higher  $N$ .
- Proved a result on the separability of  $k$ -uniform states in terms of AME states.
- Used a modified version of the non-linear map to generate 3-unitary operators in  $(6,2)$ .
- Using the non-linear maps, constructed a general form of a 3-unitary in  $(6,2)$
- Attempts at constructing AME(8,4).
- Tests on the LU invariance of certain multi-unitary operators.

# References

- Creating ensembles of dual unitary and maximally entangling quantum evolutions (*Suhail et al.*)
- Entanglement measures of bipartite quantum gates and their thermalization under arbitrary interaction strength (*Bhargavi Jonnadula et al.*)
- Entanglement of Quantum Evolutions (*Paolo Zanardi*)
- Entanglement and quantum combinatorial designs (*Dardo Goyeneche*).
- Thirty Six Entangled Officers :A Quantum solution to a classically impossible problem (*Suhail et al.*)

# Acknowledgements

I would like to thank

- Prof. Arul for the opportunity, **immense** guidance and endless stream of open problems. Thank you for the snacks as well:)
- Suhail for his guidance, advice, support and the occasional *Usha* treat.
- Shrigyan my collaborator for his discussions and acting as a sounding-off board - and for keeping things running smoothly during the online semester.
- To the Quantum Information Research Group for providing us with a place to congregate and brainstorm.
- To Devcharan for providing me with a place to work after work hours.

# Thanks

