

Principles of Scientific Computing

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Preface

This book grew out of a one semester first course in Scientific Computing for graduate students at New York University. It represents our view of how advanced undergraduates or beginning graduate students should start learning the subject, assuming that they will eventually become professionals. It is a common foundation that we hope will serve people heading to one of the many areas that rely on computing. This generic class normally would be followed by more specialized work in a particular application area.

We started out to write a book that could be covered in an intensive one semester class. The present book is a little bigger than that, but it still benefits or suffers from many hard choices of material to leave out. Textbook authors serve students by selecting the few most important topics from very many important ones. Topics such as finite element analysis, constrained optimization, algorithms for finding eigenvalues, etc. are barely mentioned. In each case, we found ourselves unable to say enough about the topic to be helpful without crowding out the material here.

Scientific computing projects fail as often from poor software as from poor mathematics. Well-designed software is much more likely to *get the right answer* than naive “spaghetti code”. Each chapter of this book has a *Software* section that discusses some aspect of programming practice. Taken together, these form a short course on programming practice for scientific computing. Included are topics like modular design and testing, documentation, robustness, performance and cache management, and visualization and performance tools.

The exercises are an essential part of the experience of this book. Much important material is there. We have limited the number of exercises so that the instructor can reasonably assign all of them, which is what we do. In particular, each chapter has one or two major exercises that guide the student through turning the ideas of the chapter into software. These build on each other as students become progressively more sophisticated in numerical technique and software design. For example, the exercise for Chapter 6 draws on an LL^t factorization program written for Chapter 5 as well as software protocols from Chapter 3.

This book is part treatise and part training manual. We lay out the mathematical principles behind scientific computing, such as error analysis and condition number. We also attempt to train the student in how to think about computing problems and how to write good software. The experiences of scientific computing are as memorable as the theorems – a program running surprisingly faster than before, a beautiful visualization, a finicky failure prone computation suddenly becoming dependable. The programming exercises in this book aim to give the student this feeling for computing.

The book assumes a facility with the mathematics of quantitative modeling: multivariate calculus, linear algebra, basic differential equations, and elementary probability. There is some review and suggested references, but nothing that would substitute for classes in the background material. While sticking to the prerequisites, we use mathematics at a relatively high level. Students are expected to understand and manipulate asymptotic error expansions, to do perturbation theory in linear algebra, and to manipulate probability densities.

Most of our students have benefitted from this level of mathematics.

We assume that the student knows basic C++ and MATLAB. The C++ in this book is in a “C style”, and we avoid both discussion of object-oriented design and of advanced language features such as templates and C++ exceptions. We help students by providing partial codes (examples of what we consider good programming style) in early chapters. The training wheels come off by the end. We do not require a specific programming environment, but in some places we say how things would be done using Linux. Instructors may have to help students without access to Linux to do some exercises (install LAPACK in Chapter 4, use performance tools in Chapter 9). Some highly motivated students have been able learn programming as they go. The web site <http://www.math.nyu.edu/faculty/goodman/ScientificComputing/> has materials to help the beginner get started with C++ or Matlab.

Many of our views on scientific computing were formed during as graduate students. One of us (JG) had the good fortune to be associated with the remarkable group of faculty and graduate students at Serra House, the numerical analysis group of the Computer Science Department of Stanford University, in the early 1980’s. I mention in particular Marsha Berger, Petter Björstad, Bill Coughran, Gene Golub, Bill Gropp, Eric Grosse, Bob Higdon, Randy LeVeque, Steve Nash, Joe Oliger, Michael Overton, Robert Schreiber, Nick Trefethen, and Margaret Wright.

The other one (DB) was fortunate to learn about numerical technique from professors and other graduate students Berkeley in the early 2000s, including Jim Demmel, W. Kahan, Beresford Parlett, Yi Chen, Plamen Koev, Jason Riedy, and Rich Vuduc. I also learned a tremendous amount about making computations relevant from my engineering colleagues, particularly Sanjay Govindjee, Bob Taylor, and Panos Papadopoulos.

Colleagues at the Courant Institute who have influenced this book include Leslie Greengard, Gene Isaacson, Peter Lax, Charlie Peskin, Luis Reyna, Mike Shelley, and Olof Widlund. We also acknowledge the lovely book *Numerical Methods* by Germund Dahlquist and Åke Björk [2]. From an organizational standpoint, this book has more in common with *Numerical Methods and Software* by Kahaner, Moler, and Nash [13].

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Chapter 1

Introduction

Most problem solving in science and engineering uses scientific computing. A scientist might devise a system of differential equations to model a physical system, then use a computer to calculate their solutions. An engineer might develop a formula to predict cost as a function of several variables, then use a computer to find the combination of variables that minimizes that cost. A scientist or engineer needs to know science or engineering to make the models. He or she needs the principles of scientific computing to find out what the models predict.

Scientific computing is challenging partly because it draws on many parts of mathematics and computer science. Beyond this knowledge, it also takes discipline and practice. A problem-solving code is built and tested procedure by procedure. Algorithms and program design are chosen based on considerations of accuracy, stability, robustness, and performance. Modern software development tools include programming environments and debuggers, visualization, profiling, and performance tools, and high-quality libraries. The training, as opposed to just teaching, is in integrating all the knowledge and the tools and the habits to create high quality computing software “solutions.”

This book weaves together this knowledge and skill base through exposition and exercises. The bulk of each chapter concerns the mathematics and algorithms of scientific computing. In addition, each chapter has a *Software* section that discusses some aspect of programming practice or software engineering. The exercises allow the student to build small codes using these principles, not just program the algorithm du jour. Hopefully he or she will see that a little planning, patience, and attention to detail can lead to scientific software that is faster, more reliable, and more accurate.

One common theme is the need to understand what is happening “under the hood” in order to understand the accuracy and performance of our computations. We should understand how computer arithmetic works so we know which operations are likely to be accurate and which are not. To write fast code, we should know that adding is much faster if the numbers are in cache, that there is overhead in getting memory (using `new` in C++ or `malloc` in C), and that printing to the screen has even more overhead. It isn’t that we should not use dynamic memory or print statements, but using them in the wrong way can make a code much slower. State-of-the-art eigenvalue software will not produce accurate eigenvalues if the problem is ill-conditioned. If it uses dense matrix methods, the running time will scale as n^3 for an $n \times n$ matrix.

Doing the exercises also should give the student a feel for numbers. The exercises are calibrated so that the student will get a feel for run time by waiting for a run to finish (a moving target given hardware advances). Many exercises ask the student to comment on the sizes of numbers. We should have a feeling for whether 4.5×10^{-6} is a plausible roundoff error if the operands are of the order of magnitude of 50. Is it plausible to compute the inverse of an $n \times n$ matrix if $n = 500$ or $n = 5000$? How accurate is the answer likely to be? Is there enough memory? Will it take more than ten seconds? Is it likely that a Monte Carlo computation with $N = 1000$ samples gives .1% accuracy?

Many topics discussed here are treated superficially. Others are left out

altogether. Do not think the things left out are unimportant. For example, anyone solving ordinary differential equations must know the stability theory of Dalhquist and others, which can be found in any serious book on numerical solution of ordinary differential equations. There are many variants of the FFT that are faster than the simple one in Chapter 7, more sophisticated kinds of spline interpolation, etc. The same applies to things like software engineering and scientific visualization. Most high performance computing is done on parallel computers, which are not discussed here at all.

Chapter 2

Sources of Error

Scientific computing usually gives inexact answers. The code `x = sqrt(2)` produces something that is not the mathematical $\sqrt{2}$. Instead, `x` differs from $\sqrt{2}$ by an amount that we call the *error*. An *accurate* result has a small error. The goal of a scientific computation is rarely the exact answer, but a result that is as accurate as needed. Throughout this book, we use A to denote the exact answer to some problem and \hat{A} to denote the computed approximation to A . The error is $\hat{A} - A$.

There are four primary ways in which error is introduced into a computation:

- (i) *Roundoff error* from inexact computer arithmetic.
- (ii) *Truncation error* from approximate formulas.
- (iii) *Termination of iterations*.
- (iv) *Statistical error* in Monte Carlo.

This chapter discusses the first of these in detail and the others more briefly. There are whole chapters dedicated to them later on. What is important here is to understand the likely relative sizes of the various kinds of error. This will help in the design of computational algorithms. In particular, it will help us focus our efforts on reducing the largest sources of error.

We need to understand the various sources of error to debug scientific computing software. If a result is supposed to be A and instead is \hat{A} , we have to ask if the difference between A and \hat{A} is the result of a programming mistake. Some bugs are the usual kind – a mangled formula or mistake in logic. Others are peculiar to scientific computing. It may turn out that a certain way of calculating something is simply not accurate enough.

Error propagation also is important. A typical computation has several stages, with the results of one stage being the inputs to the next. Errors in the output of one stage most likely mean that the output of the next would be inexact even if the second stage computations were done exactly. It is unlikely that the second stage would produce the exact output from inexact inputs. On the contrary, it is possible to have *error amplification*. If the second stage output is very sensitive to its input, small errors in the input could result in large errors in the output; that is, the error will be amplified. A method with large error amplification is *unstable*.

The *condition number* of a problem measures the sensitivity of the answer to small changes in its input data. The condition number is determined by the problem, not the method used to solve it. The accuracy of a solution is limited by the condition number of the problem. A problem is called *ill-conditioned* if the condition number is so large that it is hard or impossible to solve it accurately enough.

A computational strategy is likely to be unstable if it has an ill-conditioned subproblem. For example, suppose we solve a system of linear differential equations using the eigenvector basis of the corresponding matrix. Finding eigenvectors of a matrix can be ill-conditioned, as we discuss in Chapter 4. This makes

the eigenvector approach to solving linear differential equations potentially unstable, even when the differential equations themselves are well-conditioned.

2.1 Relative error, absolute error, and cancellation

When we approximate A by \hat{A} , the *absolute error* is $e = \hat{A} - A$, and the *relative error* is $\epsilon = e/A$. That is,

$$\hat{A} = A + e \quad (\text{absolute error}) , \quad \hat{A} = A \cdot (1 + \epsilon) \quad (\text{relative error}). \quad (2.1)$$

For example, the absolute error in approximating $A = \sqrt{175}$ by $\hat{A} = 13$ is $e \approx .23$, and the relative error is $\epsilon \approx .017 < 2\%$.

If we say $e \approx .23$ and do not give A , we generally do not know whether the error is large or small. Whether an absolute error much less than one is “small” often depends entirely on how units are chosen for a problem. In contrast, relative error is dimensionless, and if we know \hat{A} is within 2% of A , we know the error is not too large. For this reason, relative error is often more useful than absolute error.

We often describe the accuracy of an approximation by saying how many decimal digits are correct. For example, Avogadro’s number with two digits of accuracy is $N_0 \approx 6.0 \times 10^{23}$. We write 6.0 instead of just 6 to indicate that Avogadro’s number is closer to 6×10^{23} than to 6.1×10^{23} or 5.9×10^{23} . With three digits the number is $N_0 \approx 6.02 \times 10^{23}$. The difference between $N_0 \approx 6 \times 10^{23}$ and $N_0 \approx 6.02 \times 10^{23}$ is 2×10^{21} , which may seem like a lot, but the relative error is about a third of one percent.

Relative error can grow through *cancellation*. For example, suppose $A = B - C$, with $B \approx \hat{B} = 2.38 \times 10^5$ and $C \approx \hat{C} = 2.33 \times 10^5$. Since the first two digits of B and C agree, then they *cancel* in the subtraction, leaving only one correct digit in A . Doing the subtraction exactly gives $\hat{A} = \hat{B} - \hat{C} = 5 \times 10^3$. The absolute error in A is just the sum of the absolute errors in B and C , and probably is less than 10^3 . But this gives \hat{A} a relative accuracy of less than 10%, even though the inputs \hat{B} and \hat{C} had relative accuracy a hundred times smaller. *Catastrophic cancellation* is losing many digits in one subtraction. More subtle is an accumulation of less dramatic cancellations over many steps, as illustrated in Exercise 3.

2.2 Computer arithmetic

For many tasks in computer science, all arithmetic can be done with integers. In scientific computing, though, we often deal with numbers that are not integers, or with numbers that are too large to fit into standard integer types. For this reason, we typically use *floating point* numbers, which are the computer version of numbers in scientific notation.

2.2.1 Bits and ints

The basic unit of computer storage is a *bit* (binary digit), which may be 0 or 1. Bits are organized into 32-bit or 64-bit *words*. There $2^{32} \approx$ four billion possible 32-bit words; a modern machine running at 2-3 GHz could enumerate them in a second or two. In contrast, there are $2^{64} \approx 1.8 \times 10^{19}$ possible 64-bit words; to enumerate them at the same rate would take more than a century.

C++ has several basic integer types: `short`, `int`, and `long int`. The language standard does not specify the sizes of these types, but most modern systems have a 16-bit `short`, and a 32-bit `int`. The size of a `long` is 32 bits on some systems and 64 bits on others. For portability, the C++ header file `cstdint` (or the C header `stdint.h`) defines types `int16_t`, `int32_t`, and `int64_t` that are exactly 8, 16, 32, and 64 bits.

An ordinary b -bit integer can take values in the range -2^{b-1} to $2^{b-1} - 1$; an unsigned b -bit integer (such as an `unsigned int`) takes values in the range 0 to $2^b - 1$. Thus a 32-bit integer can be between -2^{31} and $2^{31} - 1$, or between about -2 billion and +2 billion. Integer addition, subtraction, and multiplication are done exactly when the results are within the representable range, and integer division is rounded toward zero to obtain an integer result. For example, $(-7)/2$ produces -3.

When integer results are out of range (an *overflow*), the answer is not defined by the standard. On most platforms, the result will be wrap around. For example, if we set a 32-bit `int` to $2^{31} - 1$ and increment it, the result will usually be -2^{31} . Therefore, the loop

```
for (int i = 0; i < 2e9; ++i);
```

takes seconds, while the loop

```
for (int i = 0; i < 3e9; ++i);
```

never terminates, because the number `3e9` (three billion) is larger than any number that can be represented as an `int`.

2.2.2 Floating point basics

Floating point numbers are computer data-types that represent approximations to real numbers rather than integers. The *IEEE floating point standard* is a set of conventions for computer representation and processing of floating point numbers. Modern computers follow these standards for the most part. The standard has three main goals:

1. To make floating point arithmetic as accurate as possible.
2. To produce sensible outcomes in exceptional situations.
3. To standardize floating point operations across computers.

Floating point numbers are like numbers in ordinary scientific notation. A number in scientific notation has three parts: a *sign*, a *mantissa* in the interval $[1, 10)$, and an *exponent*. For example, if we ask MATLAB to display the number $-2752 = -2.572 \times 10^3$ in scientific notation (using `format short e`), we see

$$-2.7520\text{e}+03$$

For this number, the sign is negative, the mantissa is 2.7520, and the exponent is 3.

Similarly, a normal binary floating point number consists of a sign s , a mantissa $1 \leq m < 2$, and an exponent e . If \mathbf{x} is a floating point number with these three fields, then the value of \mathbf{x} is the real number

$$\text{val}(\mathbf{x}) = (-1)^s \times 2^e \times m. \quad (2.2)$$

For example, we write the number $-2752 = -2.752 \times 10^3$ as

$$\begin{aligned} -2752 &= (-1)^1 \times (2^{11} + 2^9 + 2^7 + 2^6) \\ &= (-1)^1 \times 2^{11} \times (1 + 2^{-2} + 2^{-4} + 2^{-5}) \\ &= (-1)^1 \times 2^{11} \times (1 + (.01)_2 + (.0001)_2 + (.00001)_2) \\ -2752 &= (-1)^1 \times 2^{11} \times (1.01011)_2. \end{aligned}$$

The bits in a floating point word are divided into three groups. One bit represents the sign: $s = 1$ for negative and $s = 0$ for positive, according to (2.2). There are $p - 1$ bits for the mantissa and the rest for the exponent. For example (see Figure 2.1), a 32-bit single precision floating point word has $p = 24$, so there are 23 mantissa bits, one sign bit, and 8 bits for the exponent.

Floating point formats allow a limited range of exponents ($e_{\min} \leq e \leq e_{\max}$). Note that in single precision, the number of possible exponents $\{-126, -125, \dots, 126, 127\}$, is 254, which is two less than the number of 8 bit combinations ($2^8 = 256$). The remaining two exponent bit strings (all zeros and all ones) have different interpretations described in Section 2.2.4. The other floating point formats, double precision and extended precision, also reserve the all zero and all one exponent bit patterns.

The mantissa takes the form

$$m = (1.b_1b_2b_3 \dots b_{p-1})_2,$$

where p is the total number of bits (binary digits)¹ used for the mantissa. In Figure 2.1, we list the exponent range for IEEE *single precision* (`float` in C/C++), IEEE *double precision* (`double` in C/C++), and the *extended precision* on the Intel processors (`long double` in C/C++).

Not every number can be exactly represented in binary floating point. For example, just as $1/3 = .33\overline{3}$ cannot be written exactly as a finite decimal fraction, $1/3 = (.0101\overline{01})_2$ also cannot be written exactly as a finite binary fraction.

¹ Because the first digit of a normal floating point number is always one, it is not stored explicitly.

Name	C/C++ type	Bits	p	$\epsilon_{\text{mach}} = 2^{-p}$	e_{min}	e_{max}
Single	<code>float</code>	32	24	$\approx 6 \times 10^{-8}$	-126	127
Double	<code>double</code>	64	53	$\approx 10^{-16}$	-1022	1023
Extended	<code>long double</code>	80	63	$\approx 5 \times 10^{-19}$	-16382	16383

Figure 2.1: Parameters for floating point formats.

If x is a real number, we write $\hat{x} = \text{round}(x)$ for the floating point number (of a given format) that is closest² to x . Finding \hat{x} is called *rounding*. The difference $\text{round}(x) - x = \hat{x} - x$ is *rounding error*. If x is in the range of normal floating point numbers ($2^{e_{\text{min}}} \leq x < 2^{e_{\text{max}}+1}$), then the closest floating point number to x has a relative error not more than $|\epsilon| \leq \epsilon_{\text{mach}}$, where the *machine epsilon* $\epsilon_{\text{mach}} = 2^{-p}$ is half the distance between 1 and the next floating point number.

The IEEE standard for arithmetic operations (addition, subtraction, multiplication, division, square root) is: *the exact answer, correctly rounded*. For example, the statement `z = x*y` gives `z` the value $\text{round}(\text{val}(x) \cdot \text{val}(y))$. That is: interpret the bit strings `x` and `y` using the floating point standard (2.2), perform the operation (multiplication in this case) exactly, then round the result to the nearest floating point number. For example, the result of computing `1/(float)3` in single precision is

$$(1.010101010101010101011)_2 \times 2^{-2}.$$

Some properties of floating point arithmetic follow from the above rule. For example, addition and multiplication are commutative: `x*y = y*x`. Division by powers of 2 is done exactly if the result is a normalized number. Division by 3 is rarely exact. Integers, not too large, are represented exactly. Integer arithmetic (excluding division and square roots) is done exactly. This is illustrated in Exercise 8.

Double precision floating point has smaller rounding errors because it has more mantissa bits. It has roughly 16 digit accuracy ($2^{-53} \sim 10^{-16}$, as opposed to roughly 7 digit accuracy for single precision. It also has a larger range of values. The largest double precision floating point number is $2^{1023} \sim 10^{307}$, as opposed to $2^{126} \sim 10^{38}$ for single. The hardware in many processor chips does arithmetic and stores intermediate results in extended precision, see below.

Rounding error occurs in most floating point operations. When using an unstable algorithm or solving a very sensitive problem, even calculations that would give the exact answer in *exact* arithmetic may give very wrong answers in *floating point* arithmetic. Being exactly right in exact arithmetic does not imply being approximately right in floating point arithmetic.

2.2.3 Modeling floating point error

Rounding error analysis models the generation and propagation of rounding

² If x is equally close to two floating point numbers, the answer is the number whose last bit is zero.

errors over the course of a calculation. For example, suppose x , y , and z are floating point numbers, and that we compute $\mathbf{fl}(x + y + z)$, where $\mathbf{fl}(\cdot)$ denotes the result of a floating point computation. Under IEEE arithmetic,

$$\mathbf{fl}(x + y) = \mathbf{round}(x + y) = (x + y)(1 + \epsilon_1),$$

where $|\epsilon_1| < \epsilon_{\text{mach}}$. A sum of more than two numbers must be performed pairwise, and usually from left to right. For example:

$$\begin{aligned} \mathbf{fl}(x + y + z) &= \mathbf{round}(\mathbf{round}(x + y) + z) \\ &= ((x + y)(1 + \epsilon_1) + z)(1 + \epsilon_2) \\ &= (x + y + z) + (x + y)\epsilon_1 + (x + y + z)\epsilon_2 + (x + y)\epsilon_1\epsilon_2 \end{aligned}$$

Here and below we use ϵ_1 , ϵ_2 , etc. to represent individual rounding errors.

It is often replace exact formulas by simpler approximations. For example, we neglect the product $\epsilon_1\epsilon_2$ because it is smaller than either ϵ_1 or ϵ_2 (by a factor of ϵ_{mach}). This leads to the useful approximation

$$\mathbf{fl}(x + y + z) \approx (x + y + z) + (x + y)\epsilon_1 + (x + y + z)\epsilon_2,$$

We also neglect higher terms in Taylor expansions. In this spirit, we have:

$$(1 + \epsilon_1)(1 + \epsilon_2) \approx 1 + \epsilon_1 + \epsilon_2 \quad (2.3)$$

$$\sqrt{1 + \epsilon} \approx 1 + \epsilon/2. \quad (2.4)$$

As an example, we look at computing the smaller root of $x^2 - 2x + \delta = 0$ using the quadratic formula

$$x = 1 - \sqrt{1 - \delta}. \quad (2.5)$$

The two terms on the right are approximately equal when δ is small. This can lead to catastrophic cancellation. We will assume that δ is so small that (2.4) applies to (2.5), and therefore $x \approx \delta/2$.

We start with the rounding errors from the $1 - \delta$ subtraction and square root. We simplify with (2.3) and (2.4):

$$\begin{aligned} \mathbf{fl}(\sqrt{1 - \delta}) &= \left(\sqrt{(1 - \delta)(1 + \epsilon_1)} \right) (1 + \epsilon_2) \\ &\approx (\sqrt{1 - \delta})(1 + \epsilon_1/2 + \epsilon_2) = (\sqrt{1 - \delta})(1 + \epsilon_d), \end{aligned}$$

where $|\epsilon_d| = |\epsilon_1/2 + \epsilon_2| \leq 1.5\epsilon_{\text{mach}}$. This means that relative error at this point is of the order of machine precision but may be as much as 50% larger.

Now, we account for the error in the second subtraction³, using $\sqrt{1 - \delta} \approx 1$ and $x \approx \delta/2$ to simplify the error terms:

$$\begin{aligned} \mathbf{fl}(1 - \sqrt{1 - \delta}) &\approx \left(1 - (\sqrt{1 - \delta})(1 + \epsilon_d) \right) (1 + \epsilon_3) \\ &= x \left(1 - \frac{\sqrt{1 - \delta}}{x} \epsilon_d + \epsilon_3 \right) \approx x \left(1 + \frac{2\epsilon_d}{\delta} + \epsilon_3 \right). \end{aligned}$$

³ For $\delta \leq 0.75$, this subtraction actually contributes *no* rounding error, since subtraction of floating point values within a factor of two of each other is *exact*. Nonetheless, we will continue to use our model of small relative errors in this step for the current example.

Therefore, for small δ we have

$$\hat{x} - x \approx x \frac{\epsilon_d}{x},$$

which says that the relative error from using the formula (2.5) is amplified from ϵ_{mach} by a factor on the order of $1/x$. The catastrophic cancellation in the final subtraction leads to a large relative error. In single precision with $x = 10^{-5}$, for example, we would have relative error on the order of $8\epsilon_{\text{mach}}/x \approx 0.2$. When we would only expect one or two correct digits in this computation.

In this case and many others, we can avoid catastrophic cancellation by rewriting the basic formula. In this case, we could replace (2.5) by the mathematically equivalent $x = \delta/(1 + \sqrt{1 - \delta})$, which is far more accurate in floating point.

2.2.4 Exceptions

The smallest normal floating point number in a given format is 2^{emin} . When a floating point operation yields a nonzero number less than 2^{emin} , we say there has been an *underflow*. The standard formats can represent some values less than the 2^{emin} as *denormalized* numbers. These numbers have the form

$$(-1)^s \times 2^{\text{emin}} \times (0.d_1d_2 \dots d_{p-1})_2.$$

Floating point operations that produce results less than about 2^{emin} in magnitude are rounded to the nearest denormalized number. This is called *gradual underflow*. When gradual underflow occurs, the relative error in the result may be greater than ϵ_{mach} , but it is much better than if the result were rounded to 0 or 2^{emin} .

With denormalized numbers, every floating point number except the largest in magnitude has the property that the distances to the two closest floating point numbers differ by no more than a factor of two. Without denormalized numbers, the smallest number to the right of 2^{emin} would be 2^{p-1} times closer than the largest number to the left of 2^{emin} ; in single precision, that is a difference of a factor of about eight billion! Gradual underflow also has the consequence that two floating point numbers are equal, $x = y$, if and only if subtracting one from the other gives exactly zero.

In addition to the normal floating point numbers and the denormalized numbers, the IEEE standard has encodings for $\pm\infty$ and Not a Number (*NaN*). When we print these values to the screen, we see “**inf**” and “**NaN**,” respectively.⁴ A floating point operation results in an **inf** if the exact result is larger than the largest normal floating point number (*overflow*), or in cases like $1/0$ or $\cot(0)$ where the exact result is infinite⁵. Invalid operations such as `sqrt(-1.)` and

⁴ The actual text varies from system to system. The Microsoft Visual Studio compilers print **Inf** rather than **NaN**, for example.

⁵ IEEE arithmetic distinguishes between positive and negative zero, so actually $1/0.0 = \text{inf}$ and $1/-0.0 = -\text{inf}$.

`log(-4.)` produce NaN. Any operation involving a NaN produces another NaN. Operations with `inf` are common sense: `inf + finite = inf`, `inf/inf = NaN`, `finite/inf = 0`, `inf + inf = inf`, `inf - inf = NaN`.

A floating point operation generates an *exception* if the exact result is not a normalized floating point number. The five types of exceptions are an inexact result (i.e. when there is a nonzero rounding error), an underflow, an overflow, an exact infinity, or an invalid operations. When one of these exceptions occurs, a flag is raised (i.e. a bit in memory is set to one). There is one flag for each type of exception. The C99 standard defines functions to raise, lower, and check the floating point exception flags.

2.3 Truncation error

Truncation error is the error in analytical approximations such as

$$f'(x) \approx \frac{f(x+h) - f(x)}{h} . \quad (2.6)$$

This is not an exact formula for f' , but it can be a useful approximation. We often think of truncation error as arising from truncating a Taylor series. In this case, the Taylor series formula,

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \cdots ,$$

is *truncated* by neglecting all the terms after the first two on the right. This leaves the approximation

$$f(x+h) \approx f(x) + hf'(x) ,$$

which can be rearranged to give (2.6). Truncation usually is the main source of error in numerical integration or solution of differential equations. The analysis of truncation error using Taylor series will occupy the next two chapters.

As an example, we take $f(x) = xe^x$, $x = 1$, and several h values. The truncation error is

$$e_{tr} = \frac{f(x+h) - f(x)}{h} - f'(x) .$$

In Chapter 3 we will see that (in exact arithmetic) e_{tr} roughly is proportional to h for small h . This is apparent in Figure 2.3. As h is reduced from 10^{-2} to 10^{-5} (a factor of 10^{-3}), the error decreases from 4.10×10^{-2} to 4.08×10^{-5} , approximately the same factor of 10^{-3} .

The numbers in Figure 2.3 were computed in double precision floating point arithmetic. The total error, e_{tot} , is a combination of truncation and roundoff error. Roundoff error is significant for the smallest h values: for $h = 10^{-8}$ the error is no longer proportional to h ; by $h = 10^{-10}$ the error has increased. Such small h values are rare in a practical calculation.

h	.3	.01	10^{-5}	10^{-8}	10^{-10}
$\widehat{f'}$	6.84	5.48	5.4366	5.436564	5.436562
e_{tot}	1.40	4.10×10^{-2}	4.08×10^{-5}	-5.76×10^{-8}	-1.35×10^{-6}

Figure 2.2: Estimates of $f'(x)$ using (2.6). The error is e_{tot} , which results from truncation and roundoff error. Roundoff error is apparent only in the last two columns.

n	1	3	6	10	20	67
x_n	1	1.46	1.80	1.751	1.74555	1.745528
e_n	-.745	-.277	5.5×10^{-2}	5.9×10^{-3}	2.3×10^{-5}	3.1×10^{-17}

Figure 2.3: Iterates of $x_{n+1} = \ln(y) - \ln(x_n)$ illustrating convergence to a limit that satisfies the equation $xe^x = y$. The error is $e_n = x_n - x$. Here, $y = 10$.

2.4 Iterative methods

A *direct method* computes the exact answer to a problem in finite time, assuming exact arithmetic. The only error in direct methods is roundoff error. However, many problems cannot be solved by direct methods. An example is finding A that satisfies an equation that has no explicit solution formula. An *iterative method* constructs a sequence of approximate solutions, A_n , for $n = 1, 2, \dots$. Hopefully, the approximations *converge* to the right answer: $A_n \rightarrow A$ as $n \rightarrow \infty$. In practice, we must stop the iteration for some large but finite n and accept A_n as the approximate answer.

For example, suppose we have a $y > 0$ and we want to find x such that $xe^x = y$. There is no formula for x , but we can write a program to carry out the iteration: $x_1 = 1$, $x_{n+1} = \ln(y) - \ln(x_n)$. The numbers x_n are *iterates*. The limit $x = \lim_{n \rightarrow \infty} x_n$ (if it exists), is a *fixed point* of the iteration, i.e. $x = \ln(y) - \ln(x)$, which implies that $xe^x = y$. Figure 2.3 demonstrates the convergence of the iterates in this case with $y = 10$. The *initial guess* is $x_1 = 1$. After 20 iterations, we have $x_{20} \approx 1.75$. The relative error is $e_{20} \approx 2.3 \times 10^{-5}$, which might be small enough, depending on the application.

After 67 iterations, the relative error is $(x_{67} - x)/x \approx 3.1 \times 10^{-17}/1.75 \approx 1.8 \times 10^{-17}$, which means that x_{67} is the correctly rounded result in double precision with $\epsilon_{\text{mach}} \approx 1.1 \times 10^{-16}$. This shows that supposedly approximate iterative methods can be as accurate as direct methods in floating point arithmetic. For example, if the equation were $e^x = y$, the direct method would be $x = \log(y)$. It is unlikely that the relative error resulting from the direct formula would be significantly smaller than the error we could get from an iterative method.

Chapter 6 explains many features of Figure 2.3. The *convergence rate* is the rate of decrease of e_n . We will be able to calculate it and to find methods that converge faster. In some cases an iterative method may fail to converge even though the solution is a fixed point of the iteration.

n	10	100	10^4	10^6	10^6
\hat{A}	.603	.518	.511	.5004	.4991
error	.103	1.8×10^{-2}	1.1×10^{-2}	4.4×10^{-4}	-8.7×10^{-4}

Figure 2.4: Statistical errors in a demonstration Monte Carlo computation.

2.5 Statistical error in Monte Carlo

Monte Carlo means using random numbers as a computational tool. For example, suppose⁶ $A = E[X]$, where X is a random variable with some known distribution. *Sampling* X means using the computer random number generator to create independent random variables X_1, X_2, \dots , each with the distribution of X . The simple Monte Carlo method would be to generate n such samples and calculate the *sample mean*:

$$A \approx \hat{A} = \frac{1}{n} \sum_{k=1}^n X_k.$$

The difference between \hat{A} and A is *statistical error*. A theorem in probability, the *law of large numbers*, implies that $\hat{A} \rightarrow A$ as $n \rightarrow \infty$. Monte Carlo statistical errors typically are larger than roundoff or truncation errors. This makes Monte Carlo a method of last resort, to be used only when other methods are not practical.

Figure 2.4 illustrates the behavior of this Monte Carlo method for the random variable $X = \frac{3}{2}U^2$ with U uniformly distributed in the interval $[0, 1]$. The exact answer is $A = E[X] = \frac{3}{2}E[U^2] = .5$. The value $n = 10^6$ is repeated to illustrate the fact that statistical error is random (see Chapter 9 for a clarification of this). The errors even with a million samples are much larger than those in the right columns of Figures 2.3 and 2.3.

2.6 Error propagation and amplification

Errors can grow as they *propagate* through a computation. For example, consider the divided difference (2.6):

```
f1 = . . . ; // approx of f(x)
f2 = . . . ; // approx of f(x+h)
fPrimeHat = ( f2 - f1 ) / h ; // approx of derivative
```

There are three contributions to the final error in f' :

$$\hat{f}' = f'(x)(1 + \epsilon_{\text{pr}})(1 + \epsilon_{\text{tr}})(1 + \epsilon_r) \approx f'(x)(1 + \epsilon_{\text{pr}} + \epsilon_{\text{tr}} + \epsilon_r). \quad (2.7)$$

⁶ $E[X]$ is the *expected value* of X . Chapter 9 has some review of probability.

It is unlikely that $f_1 = \widehat{f(x)} \approx f(x)$ is exact. Many factors may contribute to the errors $e_1 = f_1 - f(x)$ and $e_2 = f_2 - f(x+h)$, including inaccurate x values and roundoff in the code to evaluate f . The *propagated error* comes from using inexact values of $f(x+h)$ and $f(x)$:

$$\frac{f_2 - f_1}{h} = \frac{f(x+h) - f(x)}{h} \left(1 + \frac{e_2 - e_1}{f_2 - f_1} \right) = \frac{f(x+h) - f(x)}{h} (1 + \epsilon_{\text{pr}}). \quad (2.8)$$

The truncation error in the difference quotient approximation is

$$\frac{f(x+h) - f(x)}{h} = f'(x)(1 + \epsilon_{\text{tr}}). \quad (2.9)$$

Finally, there is roundoff error in evaluating $(f_2 - f_1) / h$:

$$\widehat{f'} = \frac{f_2 - f_1}{h} (1 + \epsilon_r). \quad (2.10)$$

If we multiply the errors from (2.8)–(2.10) and simplify, we get (2.7).

Most of the error in this calculation comes from truncation error and propagated error. The subtraction and the division in the evaluation of the divided difference each have relative error of at most ϵ_{mach} ; thus, the roundoff error ϵ_r is at most about $2\epsilon_{\text{mach}}$, which is relatively small⁷. We noted in Section 2.3 that truncation error is roughly proportional to h . The propagated error ϵ_{pr} in the outputs is roughly proportional to the absolute input errors e_1 and e_2 *amplified* by a factor of h^{-1} :

$$\epsilon_{\text{pr}} = \frac{e_2 - e_1}{f_2 - f_1} \approx \frac{e_2 - e_1}{f'(x)h}.$$

Even if $\epsilon_1 = e_1/f(x)$ and $\epsilon_2 = e_2/f(x+h)$ are small, ϵ_{pr} may be quite large. This increase in relative error by a large factor in one step is another example of catastrophic cancellation, which we described in Section 2.1. If the numbers $f(x)$ and $f(x+h)$ are nearly equal, the difference can have much less relative accuracy than the numbers themselves. More subtle is *gradual error growth* over many steps. Exercise 2.15 has an example in which the error roughly doubles at each stage. Starting from double precision roundoff level, the error after 30 steps is negligible, but the error after 60 steps is larger than the answer.

An algorithm is *unstable* if its error mainly comes from amplification. This *numerical instability* can be hard to discover by standard debugging techniques that look for the first place something goes wrong, particularly if there is gradual error growth.

In scientific computing, we use *stability theory*, or the study of propagation of small changes by a process, to search for error growth in computations. In a typical stability analysis, we focus on propagated error only, ignoring the original sources of error. For example, Exercise 8 involves the backward recurrence $f_{k-1} = f_{k+1} - f_k$. In our stability analysis, we assume that the subtraction is performed exactly and that the error in f_{k-1} is entirely due to errors in f_k

⁷ If h is a power of two and f_2 and f_1 are within a factor of two of each other, then $\epsilon_r = 0$.

and f_{k+1} . That is, if $\hat{f}_k = f_k + e_k$ is the computer approximation, then the e_k satisfy the *error propagation equation* $e_{k-1} = e_{k+1} - e_k$. We then would use the theory of recurrence relations to see whether the e_k can grow relative to the f_k as k decreases. If this error growth is possible, it will usually happen.

2.7 Condition number and ill-conditioned problems

The *condition number* of a problem measures the sensitivity of the answer to small changes in the data. If κ is the condition number, then we expect relative error at least $\kappa\epsilon_{\text{mach}}$, regardless of the algorithm. A problem with large condition number is *ill-conditioned*. For example, if $\kappa > 10^7$, then there probably is no algorithm that gives anything like the right answer in single precision arithmetic. Condition numbers as large as 10^7 or 10^{16} can and do occur in practice.

The definition of κ is simplest when the answer is a single number that depends on a single scalar variable, x : $A = A(x)$. A change in x causes a change in A : $\Delta A = A(x + \Delta x) - A(x)$. The condition number measures the relative change in A caused by a small relative change of x :

$$\left| \frac{\Delta A}{A} \right| \approx \kappa \left| \frac{\Delta x}{x} \right|. \quad (2.11)$$

Any algorithm that computes $A(x)$ must round x to the nearest floating point number, \hat{x} . This creates a relative error (assuming x is within the range of normalized floating point numbers) of $|\Delta x/x| = |(\hat{x} - x)/x| \sim \epsilon_{\text{mach}}$. If the rest of the computation were done exactly, the computed answer would be $\hat{A}(x) = A(\hat{x})$ and the relative error would be (using (2.11))

$$\left| \frac{\hat{A}(x) - A(x)}{A(x)} \right| = \left| \frac{A(\hat{x}) - A(x)}{A(x)} \right| \approx \kappa \left| \frac{\Delta x}{x} \right| \sim \kappa \epsilon_{\text{mach}}. \quad (2.12)$$

If A is a differentiable function of x with derivative $A'(x)$, then, for small Δx , $\Delta A \approx A'(x)\Delta x$. With a little algebra, this gives (2.11) with

$$\kappa = \left| A'(x) \cdot \frac{x}{A(x)} \right|. \quad (2.13)$$

An algorithm is *unstable* if relative errors in the output are much larger than relative errors in the input. This analysis argues that any computation for an ill-conditioned problem must be unstable. Even if $A(x)$ is evaluated exactly, relative errors in the input of size ϵ are amplified by a factor of κ . The formulas (2.12) and (2.13) represent an lower bound for the accuracy of any algorithm. An ill-conditioned problem is not going to be solved accurately, unless it can be reformulated to improve the conditioning.

A *backward stable* algorithm is as stable as the condition number allows. This sometimes is stated as $\hat{A}(x) = A(\tilde{x})$ for some \tilde{x} that is within on the order

of ϵ_{mach} of x , $|\tilde{x} - x|/|x| \lesssim C\epsilon_{\text{mach}}$ where C is a modest constant. That is, the computer produces a result, \tilde{A} , that is the exact correct answer to a problem that differs from the original one only by roundoff error in the input. Many algorithms in computational linear algebra in Chapter 5 are backward stable in roughly this sense. This means that they are unstable only when the underlying problem is ill-conditioned

The condition number (2.13) is dimensionless because it measures relative sensitivity. The extra factor $x/A(x)$ removes the units of x and A . Absolute sensitivity is just $A'(x)$. Note that both sides of our starting point (2.11) are dimensionless with dimensionless κ .

As an example consider the problem of evaluating $A(x) = R \sin(x)$. The condition number formula (2.13) gives

$$\kappa(x) = \left| \cos(x) \cdot \frac{x}{\sin(x)} \right|.$$

Note that the problem remains well-conditioned (κ is not large) as $x \rightarrow 0$, even though $A(x)$ is small when x is small. For extremely small x , the calculation could suffer from underflow. But the condition number blows up as $x \rightarrow \pi$, because small relative changes in x lead to much larger relative changes in A . This illustrates quirk of the condition number definition: typical values of A have the order of magnitude R and we can evaluate A with error much smaller than this, but certain individual values of A may not be computed to high relative precision. In most applications that would not be a problem.

There is no perfect definition of condition number for problems with more than one input or output. Suppose at first that the single output $A(x)$ depends on n inputs $x = (x_1, \dots, x_n)$. Of course A may have different sensitivities to different components of x . For example, $\Delta x_1/x_1 = 1\%$ may change A much more than $\Delta x_2/x_2 = 1\%$. If we view (2.11) as saying that $|\Delta A/A| \approx \kappa \epsilon$ for $|\Delta x/x| = \epsilon$, a worst case multicomponent generalization could be

$$\kappa = \frac{1}{\epsilon} \max \left| \frac{\Delta A}{A} \right|, \quad \text{where} \quad \left| \frac{\Delta x_k}{x_k} \right| \leq \epsilon \text{ for all } k.$$

We seek the worst case⁸ Δx . For small ϵ , we write

$$\Delta A \approx \sum_{k=1}^n \frac{\partial A}{\partial x_k} \Delta x_k,$$

then maximize subject to the constraint $|\Delta x_k| \leq \epsilon |x_k|$ for all k . The maximum occurs at $\Delta x_k = \pm \epsilon x_k$, which (with some algebra) leads to one possible generalization of (2.13):

$$\kappa = \sum_{k=1}^n \left| \frac{\partial A}{\partial x_k} \cdot \frac{x_k}{A} \right|. \quad (2.14)$$

⁸ As with rounding, typical errors tend to be on the order of the worst case error.

This formula is useful if the inputs are known to similar relative accuracy, which could happen even when the x_k have different orders of magnitude or different units. Condition numbers for multivariate problems are discussed using matrix norms in Section 4.3. The analogue of (2.13) is (4.28).

2.8 Software

Each chapter of this book has a *Software* section. Taken together they form a mini-course in software for scientific computing. The material ranges from simple tips to longer discussions of bigger issues. The programming exercises illustrate the chapter's software principles as well as the mathematical material from earlier sections.

Scientific computing projects fail because of bad software as often as they fail because of bad algorithms. The principles of scientific software are less precise than the mathematics of scientific computing, but are just as important. Like other programmers, scientific programmers should follow general software engineering principles of modular design, documentation, and testing. Unlike other programmers, scientific programmers must also deal with the approximate nature of their computations by testing, analyzing, and documenting the error in their methods, and by composing modules in a way that does not needlessly amplify errors. Projects are handsomely rewarded for extra efforts and care taken to do the software “right.”

2.8.1 General software principles

This is not the place for a long discussion of general software principles, but it may be worthwhile to review a few that will be in force even for the small programs called for in the exercises below. In our classes, students are required to submit the program source as well the output, and points are deducted for failing to follow basic software protocols listed here.

Most software projects are collaborative. Any person working on the project will write code that will be read and modified by others. He or she will read code that will be read and modified by others. In single person projects, that “other” person may be you, a few months later, so you will end up helping yourself by coding in a way friendly to others.

- Make code easy for others to read. Choose helpful variable names, not too short or too long. Indent the bodies of loops and conditionals. Align similar code to make it easier to understand visually. Insert blank lines to separate “paragraphs” of code. Use descriptive comments liberally; these are part of the documentation.
- Code documentation includes comments and well chosen variable names. Larger codes need separate documents describing how to use them and how they work.

```

void do_something(double tStart, double tFinal, int n)
{
    double dt = ( tFinal - tStart ) / n; // Size of each step
    for (double t = tStart; t < tFinal; t += dt) {
        // Body of loop
    }
}

```

Figure 2.5: A code fragment illustrating a pitfall of using a floating point variable to regulate a for loop.

- Design large computing tasks into modules that perform sub-tasks.
- A code design includes a plan for building it, which includes a plan for testing. Plan to create separate routines that test the main components of your project.
- Know and use the tools in your programming environment. This includes code specific editors, window based debuggers, formal or informal version control, and `make` or other building and scripting tools.

2.8.2 Coding for floating point

Programmers who forget the inexactness of floating point may write code with subtle bugs. One common mistake is illustrated in Figure 2.5. For `tFinal > tStart`, this code would give the desired n iterations in exact arithmetic. Because of rounding error, though, the variable `t` might be above or below `tFinal`, and so we do not know whether this code while execute the loop body n times or $n + 1$ times. The routine in Figure 2.5 has other issues, too. For example, if `tFinal <= tStart`, then the loop body will never execute. The loop will also never execute if either `tStart` or `tFinal` is a `NaN`, since any comparison involving a `NaN` is false. If `n = 0`, then `dt` will be `Inf`. Figure 2.6 uses exact integer arithmetic to guarantee n executions of the for loop.

Floating point results may depend on how the code is compiled and executed. For example, `w = x + y + z` is executed as if it were `A = x + y; w = A + z`. The new *anonymous variable*, `A`, may be computed in extended precision or in the precision of `x`, `y`, and `z`, depending on factors often beyond control or the programmer. The results will be as accurate as they should be, but they will be exactly identical. For example, on a Pentium running Linux⁹, the statement

```

double eps = 1e-16;
cout << 1 + eps - 1 << endl;

```

⁹ This is a 64-bit Pentium 4 running Linux 2.6.18 with GCC version 4.1.2 and Intel C version 10.0.

```

void do_something(double tStart, double tFinal, int n)
{
    double dt = ( tFinal - tStart ) / n;
    double t = tStart;
    for (int i = 0; i < n; ++i) {
        // Body of loop
        t += dt;
    }
}

```

Figure 2.6: A code fragment using an integer variable to regulate the loop of Figure 2.5. This version also is more robust in that it runs correctly if $t_{\text{Final}} \leq t_{\text{Start}}$ or if $n = 0$.

prints $1\text{e-}16$ when compiled with the Intel C compiler, which uses 80-bit intermediate precision to compute $1 + \text{eps}$, and 0 when compiled with the GNU C compiler, which uses double precision for $1 + \text{eps}$.

Some compilers take what seem like small liberties with the IEEE floating point standard. This is particularly true with high levels of optimization. Therefore, the output may change when the code is recompiled with a higher level of optimization. This is possible even if the code is absolutely correct.

2.8.3 Plotting

Visualization tools range from simple plots and graphs to more sophisticated surface and contour plots, to interactive graphics and movies. They make it possible to explore and understand data more reliably and faster than simply looking at lists of numbers. We discuss only simple graphs here, with future software sections discussing other tools. Here are some things to keep in mind.

Learn your system and your options. Find out what visualization tools are available or easy to get on your system. Choose a package designed for scientific visualization, such as MATLAB (or Gnuplot, R, Python, etc.), rather than one designed for commercial presentations such as Excel. Learn the options such as line style (dashes, thickness, color, symbols), labeling, etc. Also, learn how to produce figures that are easy to read on the page as well as on the screen. Figure 2.7 shows a MATLAB script that sets the graphics parameters so that printed figures are easier to read.

Use scripting and other forms of automation. You will become frustrated typing several commands each time you adjust one detail of the plot. Instead, assemble the sequence of plot commands into a script.

Choose informative scales and ranges. Figure 2.8 shows the convergence of the fixed-point iteration from Section 2.4 by plotting the residual $r_i = |x_i +$

```
% Set MATLAB graphics parameters so that plots are more readable
% in print. These settings are taken from the front matter in
% L.N. Trefethen's book 'Spectral Methods in MATLAB' (SIAM, 2000),
% which includes many beautifully produced MATLAB figures.

set(0, 'defaultaxesfontsize', 12, ...
      'defaultaxeslinewidth', .7, ...
      'defaultlinelength', .8, ...
      'defaultpatchlinewidth', .7);
```

Figure 2.7: A MATLAB script to set graphics parameters so that printed figures are easier to read.

$\log(x_i) - y$ against the iteration number. On a linear scale (first plot), all the residual values after the tenth step are indistinguishable *to plotting accuracy* from zero. The semi-logarithmic plot shows how big each of the iterates is. It also makes clear that the $\log(r_i)$ is nearly proportional to i ; this is the hallmark of *linear convergence*, which we will discuss more in Chapter 6.

Combine curves you want to compare into a single figure. Stacks of graphs are as frustrating as arrays of numbers. You may have to scale different curves differently to bring out the relationship they have to each other. If the curves are supposed to have a certain slope, include a line with that slope. If a certain x or y value is important, draw a horizontal or vertical line to mark it in the figure. Use a variety of line styles to distinguish the curves. Exercise 9 illustrates some of these points.

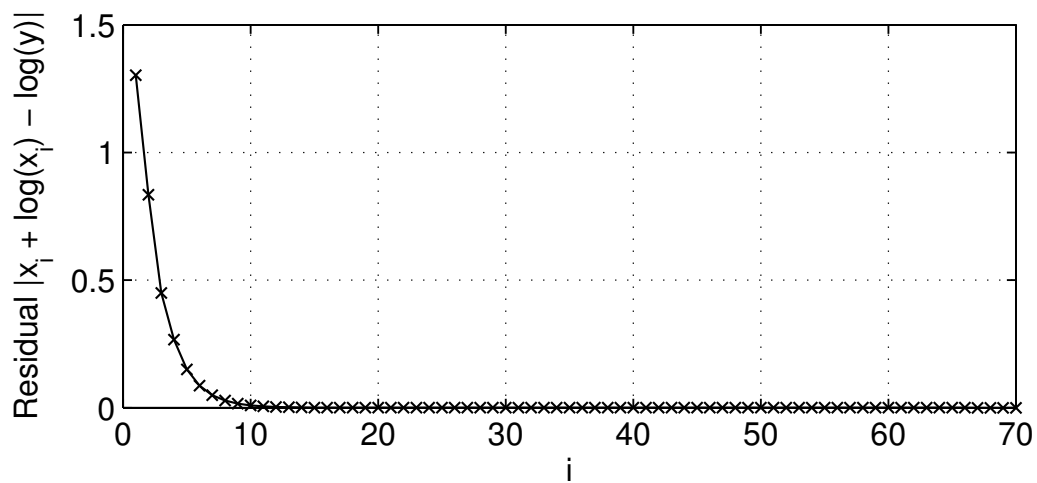
Make plots self-documenting. Figure 2.8 illustrates mechanisms in MATLAB for doing this. The horizontal and vertical axes are labeled with values and text. Parameters from the run, in this case x_1 and y , are embedded in the title.

2.9 Further reading

We were inspired to start our book with a discussion of sources of error by the book *Numerical Methods and Software* by David Kahaner, Cleve Moler, and Steve Nash [13]. Another interesting version is in *Scientific Computing* by Michael Heath [9]. A much more thorough description of error in scientific computing is Higham's *Accuracy and Stability of Numerical Algorithms* [10]. Acton's *Real Computing Made Real* [1] is a jeremiad against common errors in computation, replete with examples of the ways in which computations go awry – and ways those computations may be fixed.

A classic paper on floating point arithmetic, readily available online, is Goldberg's "What Every Computer Scientist Should Know About Floating-Point Arithmetic" [6]. Michael Overton has written a nice short book *IEEE Floating*

Convergence of $x_{i+1} = \log(y) - \log(x_i)$, with $x_1 = 1$, $y = 10$



Convergence of $x_{i+1} = \log(y) - \log(x_i)$, with $x_1 = 1$, $y = 10$

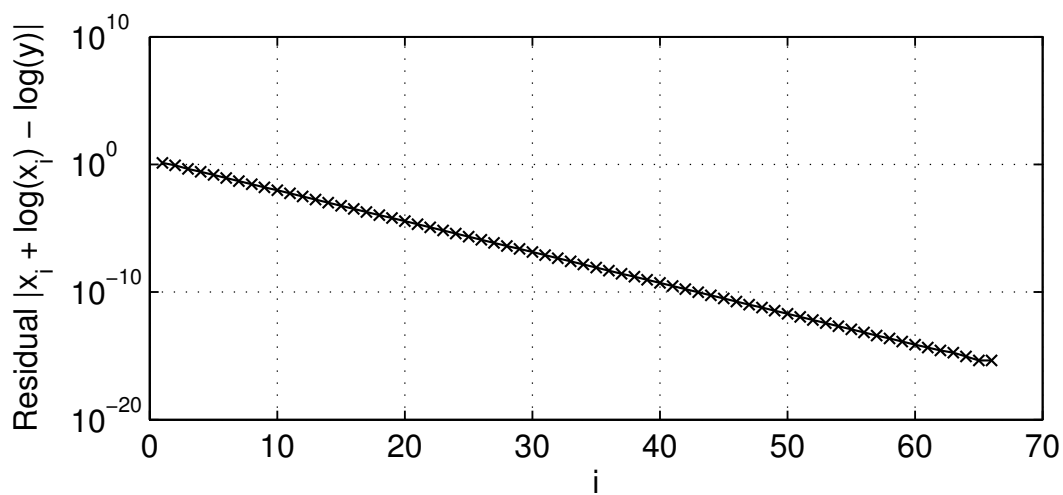


Figure 2.8: Plots of the convergence of $x_{i+1} = \log(y) - \log(x_i)$ to a fixed point on linear and logarithmic scales.

```

% sec2_8b(x1, y, n, fname)
%
% Solve the equation  $x \exp(x) = y$  by the fixed-point iteration
%    $x(i+1) = \log(y) - \log(x(i))$ ;
% and plot the convergence of the residual  $|x + \log(x) - \log(y)|$  to zero.
%
% Inputs:
%   x1    - Initial guess (default: 1)
%   y     - Right side (default: 10)
%   n     - Maximum number of iterations (default: 70)
%   fname - Name of an eps file for printing the convergence plot.

function sec2_8b(x, y, n, fname)

    % Set default parameter values
    if nargin < 1, x = 1; end
    if nargin < 2, y = 10; end
    if nargin < 3, n = 70; end

    % Compute the iterates
    for i = 1:n-1
        x(i+1) = log(y) - log(x(i));
    end

    % Plot the residual error vs iteration number on a log scale
    f = x + log(x) - log(y);
    semilogy(abs(f), 'x-');

    % Label the x and y axes
    xlabel('i');
    ylabel('Residual  $|x_i + \log(x_i) - \log(y)|$ ');

    % Add a title. The sprintf command lets us format the string.
    title(sprintf(['Convergence of  $x_{i+1} = \log(y) - \log(x_i)$ , ...',
                  'with  $x_1 = %g$ ,  $y = %g$ '], x(1), y));
    grid; % A grid makes the plot easier to read

    % If a filename is provided, print to that file (sized for book)
    if nargin == 4
        set(gcf, 'PaperPosition', [0 0 6 3]);
        print('-deps', fname);
    end

```

Figure 2.9: MATLAB code to plot convergence of the iteration from Section 2.4.

Point Arithmetic [18] that goes into more detail.

There are many good books on software design, among which we recommend *The Pragmatic Programmer* by Hunt and Thomas [11] and *The Practice of Programming* by Kernighan and Pike [14]. There are far fewer books that specifically address *numerical* software design, but one nice exception is *Writing Scientific Software* by Oliveira and Stewart [16].

2.10 Exercises

1. It is common to think of $\pi^2 = 9.87$ as approximately ten. What are the absolute and relative errors in this approximation?
2. If `x` and `y` have type `double`, and `(fabs(x - y) >= 10)` evaluates to `TRUE`, does that mean that `y` is not a good approximation to `x` in the sense of relative error?
3. Show that $f_{jk} = \sin(x_0 + (j - k)\pi/3)$ satisfies the recurrence relation

$$f_{j,k+1} = f_{j,k} - f_{j+1,k} . \quad (2.15)$$

We view this as a formula that computes the f values on level $k + 1$ from the f values on level k . Let \hat{f}_{jk} for $k \geq 0$ be the floating point numbers that come from implementing $f_{j0} = \sin(x_0 + j\pi/3)$ and (2.15) (for $k > 0$) in double precision floating point. If $|\hat{f}_{jk} - f_{jk}| \leq \epsilon$ for all j , show that $|\hat{f}_{j,k+1} - f_{j,k+1}| \leq 2\epsilon$ for all j . Thus, if the level k values are very accurate, then the level $k + 1$ values still are pretty good.

Write a program (C/C++ or MATLAB) that computes $e_k = \hat{f}_{0k} - f_{0k}$ for $1 \leq k \leq 60$ and $x_0 = 1$. Note that f_{0n} , a single number on level n , depends on $f_{0,n-1}$ and $f_{1,n-1}$, two numbers on level $n - 1$, and so on down to n numbers on level 0. Print the e_k and see whether they grow monotonically. Plot the e_k on a linear scale and see that the numbers seem to go bad suddenly at around $k = 50$. Plot the e_k on a log scale. For comparison, include a straight line that would represent the error if it were exactly to double each time.

4. What are the possible values of `k` after the `for` loop is finished?

```
float x = 100.0*rand() + 2;
int n = 20, k = 0;
float dy = x/n;
for (float y = 0; y < x; y += dy)
    k++; /* Body does not change x, y, or dy */
```

5. We wish to evaluate the function $f(x)$ for x values around 10^{-3} . We expect f to be about 10^5 and f' to be about 10^{10} . Is the problem too ill-conditioned for single precision? For double precision?

6. Use the fact that the floating point sum $x + y$ has relative error $\epsilon < \epsilon_{\text{mach}}$ to show that the absolute error in the sum S computed below is no worse than $(n - 1)\epsilon_{\text{mach}} \sum_{k=0}^{n-1} |x_k|$:

```
double compute_sum(double x[], int n)
{
    double S = 0;
    for (int i = 0; i < n; ++i)
        S += x[i];
    return S;
}
```

7. Starting with the declarations

```
float x, y, z, w;
const float oneThird = 1/ (float) 3;
const float oneHalf  = 1/ (float) 2;
// const means these never are reassigned
```

we do lots of arithmetic on the variables x , y , z , w . In each case below, determine whether the two arithmetic expressions result in the same floating point number (down to the last bit) as long as no NaN or inf values or denormalized numbers are produced.

(a)

```
( x * y ) + ( z - w )
( z - w ) + ( y * x )
```

(b)

```
( x + y ) + z
x + ( y + z )
```

(c)

```
x * oneHalf + y * oneHalf
( x + y ) * oneHalf
```

(d)

```
x * oneThird + y * oneThird
( x + y ) * oneThird
```

8. The *Fibonacci numbers*, f_k , are defined by $f_0 = 1$, $f_1 = 1$, and

$$f_{k+1} = f_k + f_{k-1} \quad (2.16)$$

for any integer $k > 1$. A small perturbation of them, the *piB numbers* (“p” instead of “f” to indicate a perturbation), p_k , are defined by $p_0 = 1$, $p_1 = 1$, and

$$p_{k+1} = c \cdot p_k + p_{k-1}$$

for any integer $k > 1$, where $c = 1 + \sqrt{3}/100$.

- (a) Plot the f_n and p_n in one together on a log scale plot. On the plot, mark $1/\epsilon_{\text{mach}}$ for single and double precision arithmetic. This can be useful in answering the questions below.
- (b) Rewrite (2.16) to express f_{k-1} in terms of f_k and f_{k+1} . Use the computed f_n and f_{n-1} to recompute f_k for $k = n-2, n-3, \dots, 0$. Make a plot of the difference between the original $f_0 = 1$ and the recomputed \hat{f}_0 as a function of n . What n values result in no accuracy for the recomputed f_0 ? How do the results in single and double precision differ?
- (c) Repeat b. for the pib numbers. Comment on the striking difference in the way precision is lost in these two cases. Which is more typical? *Extra credit:* predict the order of magnitude of the error in recomputing p_0 using what you may know about recurrence relations and what you should know about computer arithmetic.

9. The binomial coefficients, $a_{n,k}$, are defined by

$$a_{n,k} = \binom{n}{k} = \frac{n!}{k!(n-k)!}$$

To compute the $a_{n,k}$, for a given n , start with $a_{n,0} = 1$ and then use the recurrence relation $a_{n,k+1} = \frac{n-k}{k+1} a_{n,k}$.

- (a) For a range of n values, compute the $a_{n,k}$ this way, noting the largest $a_{n,k}$ and the accuracy with which $a_{n,n} = 1$ is computed. Do this in single and double precision. Why is roundoff not a problem here as it was in problem 8? Find n values for which $\hat{a}_{n,n} \approx 1$ in double precision but not in single precision. How is this possible, given that roundoff is not a problem?
- (b) Use the algorithm of part (a) to compute

$$E(k) = \frac{1}{2^n} \sum_{k=0}^n k a_{n,k} = \frac{n}{2} . \quad (2.17)$$

Write a program without any safeguards against overflow or zero divide (this time only!)¹⁰. Show (both in single and double precision) that the computed answer has high accuracy as long as the intermediate results are within the range of floating point numbers. As with (a), explain how the computer gets an accurate, small, answer when the intermediate numbers have such a wide range of values. Why is cancellation not a problem? Note the advantage of a wider range of values: we can compute $E(k)$ for much larger n in double precision. Print $E(k)$ as computed by (2.17) and $M_n = \max_k a_{n,k}$. For large n , one should be `inf` and the other `NaN`. Why?

¹⁰One of the purposes of the IEEE floating point standard was to allow a program with overflow or zero divide to run and print results.

- (c) For fairly large n , plot $a_{n,k}/M_n$ as a function of k for a range of k chosen to illuminate the interesting “bell shaped” behavior of the $a_{n,k}$ near $k = n/2$. Combine the curves for $n = 10$, $n = 20$, and $n = 50$ in a single plot. Choose the three k ranges so that the curves are close to each other. Choose different line styles for the three curves.

Chapter 3

Local Analysis

Among the most common computational tasks are differentiation, interpolation, and integration. The simplest methods used for these operations are *finite difference* approximations for derivatives, *polynomial* interpolation, and *panel method* integration. Finite difference formulas, integration rules, and interpolation form the core of most scientific computing projects that involve solving differential or integral equations.

Finite difference formulas range from simple *low order* approximations (3.14a) – (3.14c) to more complicated *high order* methods such as (3.14e). High order methods can be far more accurate than low order ones. This can make the difference between getting useful answers and not in serious applications. We will learn to design highly accurate methods rather than relying on simple but often inefficient low order ones.

Many methods for these problems involve a *step size*, h . For each h there is an approximation¹ $\hat{A}(h) \approx A$. We say \hat{A} is *consistent* if² $\hat{A}(h) \rightarrow A$ as $h \rightarrow 0$. For example, we might estimate $A = f'(x)$ using the finite difference formula (3.14a): $\hat{A}(h) = (f(x+h) - f(x))/h$. This is consistent, as $\lim_{h \rightarrow 0} \hat{A}(h)$ is the definition of $f'(x)$. The accuracy of the approximation depends on f , but the *order of accuracy* depends only on how many derivatives f has.³ The approximation is *first order accurate* if the error is nearly proportional to h for small enough h . It is *second order* if the error goes like h^2 . When h is small, $h^2 \ll h$, so approximations with a higher order of accuracy can be much more accurate.

The design of difference formulas and integration rules is based on *local analysis* using approximations to a function f about a *base point* x . These approximations consist of the first few terms of the *Taylor series* expansion of f about x . The *first order* approximation is

$$f(x+h) \approx f(x) + hf'(x). \quad (3.1)$$

The *second order* approximation is more complicated and more accurate:

$$f(x+h) \approx f(x) + f'(x)h + \frac{1}{2}f''(x)h^2. \quad (3.2)$$

Figure 3.1 illustrates the first and second order approximations. *Truncation error* is the difference between $f(x+h)$ and one of these approximations. For example, the truncation error for the first order approximation is

$$f(x) + f'(x)h - f(x+h).$$

To see how Taylor series are used, substitute the approximation (3.1) into the finite difference formula $f'(x) \approx (f(x+h) - f(x))/h$ (given in (3.14a)). We find

$$\hat{A}(h) \approx f'(x) = A. \quad (3.3)$$

¹In the notation of Chapter 2, \hat{A} is an estimate of the desired answer, A .

²Here, and in much of this chapter, we implicitly ignore rounding errors. Truncation errors are larger than rounding errors in the majority of applications.

³The nominal order of accuracy may only be achieved if f is smooth enough, a point that is important in many applications.

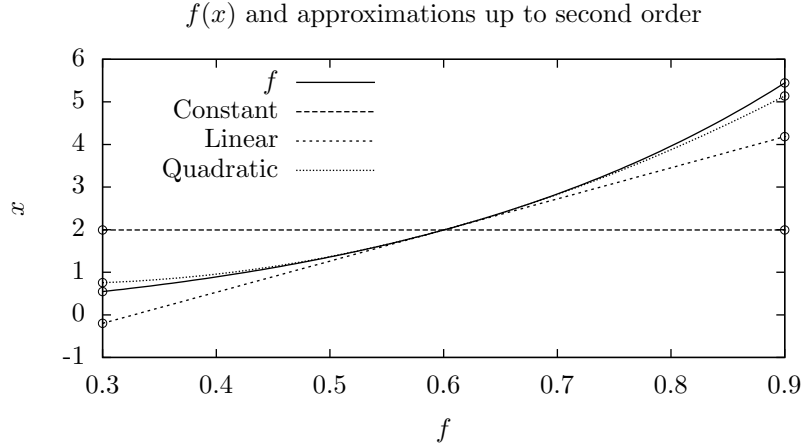


Figure 3.1: Plot of $f(x) = xe^{2x}$ together with Taylor series approximations of order zero, one, and two. The base point is $x = .6$ and h ranges from $-.3$ to $.3$. Notice the convergence for fixed h as the order increases. Also, the higher order curves make closer contact with $f(x)$ as $h \rightarrow 0$.

The more accurate Taylor approximation (3.2) allows us to estimate the error in (3.14a). Substituting (3.2) into (3.14a) gives

$$\hat{A}(h) \approx A + A_1 h, \quad A_1 = \frac{1}{2} f''(x). \quad (3.4)$$

This *asymptotic error expansion* is an estimate of $\hat{A} - A$ for a given f and h . It shows that the error is roughly proportional to h , for small h . This understanding of truncation error leads to more sophisticated computational strategies. *Richardson extrapolation* combines $\hat{A}(h)$ and $\hat{A}(2h)$ to create higher order estimates with much less error. *Adaptive methods* are automatic ways to find h such that $|\hat{A}(h) - A| \leq e$, where e is a specified level of accuracy. Error expansions like (3.4) are the basis of many adaptive methods.

This chapter focuses on truncation error and mostly ignores roundoff. In most practical computations that have truncation error, including numerical solution of differential equations or integral equations, the truncation error is much larger than roundoff. In the example shown in Figure 2.3, we saw that for values of h ranging from 10^{-2} to 10^{-5} , the computed error is roughly $4.1h$, as (3.4) suggests it should be. Roundoff in this case begins to dominate only around $h = 10^{-8}$. More sophisticated high-order approximations reach roundoff sooner, which can be an issue in testing, but rarely is an issue in production runs.

3.1 Taylor series and asymptotic expansions

The Taylor series expansion of a function f about the point x is

$$f(x+h) = \sum_{n=0}^{\infty} \frac{1}{n!} f^{(n)}(x) h^n . \quad (3.5)$$

The notation $f^{(n)}(x)$ refers to the n^{th} derivative of f evaluated at x . The *partial sum* of order p is a degree p polynomial in h :

$$F_p(x, h) = \sum_{n=0}^p \frac{1}{n!} f^{(n)}(x) h^n . \quad (3.6)$$

The partial sum F_p is the Taylor approximation to $f(x+h)$ of order p . It is a polynomial of order p in the variable h . Increasing p makes the approximation more complicated and more accurate. The order $p=0$ partial sum is simply $F_0(x, h) = f(x)$. The first and second order approximations are (3.1) and (3.2) respectively.

The Taylor series sum *converges* if the partial sums converge to f :

$$\lim_{p \rightarrow \infty} F_p(x, h) = f(x+h) .$$

If there is a positive h_0 so that the series converges whenever $|h| < h_0$, then f is *analytic* at x . A function probably is analytic at most points if there is a formula for it, or it is the solution of a differential equation. Figure 3.1 plots a function $f(x) = xe^{2x}$ together with the Taylor approximations of order zero, one, and two. The symbols at the ends of the curves illustrate the convergence of this Taylor series when $h = \pm 3$. When $h = .3$, the series converges monotonically: $F_0 < F_1 < F_2 < \dots \rightarrow f(x+h)$. When $h = -.3$, there are approximants on both sides of the answer: $F_0 > F_2 > f(x+h) > F_1$.

For our purposes, we will often use the partial sums F_p without letting p go to infinity. To analyze how well the approximations F_p approximate f near the base point x , it is useful to introduce *big-O notation*. If $B_1(h) > 0$ for $h \neq 0$, we write $B_2(h) = O(B_1(h))$ (as $h \rightarrow 0$ ⁴ to mean there is a C such that $|B_1(h)| \leq CB_2(h)$ in an open neighborhood of $h = 0$. A common misuse of the notation, which we will follow, is to write $B_2(h) = O(B_1(h))$ when B_1 can be negative, rather than $B_2(h) = O(|B_1(h)|)$. We say B_2 is an order smaller than B_1 (or an order of approximation smaller) if $B_2(h) = O(hB_1(h))$. If B_2 is an order smaller than B_1 , then we have $|B_2(h)| < |B_1(h)|$ for small enough h ($h < 1/C$), and reducing h further makes B_2 much smaller than B_1 .

An *asymptotic expansion* is a sequence of approximations with increasing order of accuracy, like the partial sums $F_p(x, h)$. We can make an analogy

⁴ Big-O notation is also used to describe the asymptotic behavior of functions as they go to infinity; we say $f(t) = O(g(t))$ for $g \geq 0$ as $t \rightarrow \infty$ if $|f(t)|/g(t)$ is bounded for all $t > T$. It will usually be clear from context whether we are looking at limits going to zero or going to infinity.

between asymptotic expansions of functions and decimal expansions of numbers if we think of orders of magnitude rather than orders of accuracy. The expansion

$$\pi = 3.141592 \cdots = 3 + 1 \cdot .1 + 4 \cdot (.1)^2 + 1 \cdot (.1)^3 + 5 \cdot (.1)^4 + \cdots \quad (3.7)$$

is a sequence of approximations

$$\begin{aligned} \hat{A}_0 &\approx 3 \\ \hat{A}_1 &\approx 3 + 1 \cdot .1 = 3.1 \\ \hat{A}_2 &\approx 3 + 1 \cdot .1 + 4 \cdot (.1)^2 = 3.14 \\ &\text{etc.} \end{aligned}$$

The approximation \hat{A}_p is an order of magnitude more accurate than \hat{A}_{p-1} . The error $\hat{A}_p - \pi$ is of the order of magnitude $(.1)^{p+1}$, which also is the order of magnitude of the first neglected term. The error in $\hat{A}_3 = 3.141$ is $\hat{A}_3 - \pi \approx 6 \cdot 10^{-4}$. This error is approximately the same as the next term, $5 \cdot (.1)^4$. Adding the next term gives an approximation whose error is an order of magnitude smaller:

$$\hat{A}_3 + 5 \cdot (.1)^4 - \pi = \hat{A}_4 - \pi \approx -9 \cdot 10^{-5}.$$

We change notation when we view the Taylor series as an asymptotic expansion, writing

$$f(x+h) \sim f(x) + f'(x) \cdot h + \frac{1}{2} f''(x) h^2 + \cdots \quad (3.8)$$

This means that the right side is an *asymptotic series* that may or may not converge. It represents $f(x+h)$ in the sense that the partial sums $F_p(x, h)$ are a family of approximations of increasing order of accuracy:

$$|F_p(x, h) - f(x+h)| = O(h^{p+1}). \quad (3.9)$$

The asymptotic expansion (3.8) is much like the decimal expansion (3.7). The term in (3.8) of order $O(h^2)$ is $\frac{1}{2} f''(x) \cdot h^2$. The term in (3.7) of order of magnitude 10^{-2} is $4 \cdot (.1)^2$. The error in a p term approximation is roughly the first neglected term, since all other neglected terms are at least one order smaller.

Figure 3.1 illustrates the asymptotic nature of the Taylor approximations. The lowest order approximation is $F_0(x, h) = f(x)$. The graph of F_0 touches the graph of f when $h = 0$ but otherwise has little in common. The graph of $F_1(x, h) = f(x) + f'(x)h$ not only touches the graph of f when $h = 0$, but the curves are tangent. The graph of $F_2(x, h)$ not only is tangent, but has the same curvature and is a better fit for small and not so small h .

3.1.1 Technical points

In this subsection, we present two technical points for mathematically minded readers. First, we prove the basic fact that underlies most of the analysis in this

chapter, that the Taylor series (3.5) is an asymptotic expansion. Second, we give an example of an asymptotic expansion that converges to the wrong answer, and another example of an asymptotic expansion that does not converge at all.

The asymptotic expansion property of Taylor series comes from the Taylor series *remainder theorem*.⁵ If the derivatives of f up to order $p + 1$ exist and are continuous in the interval $[x, x + h]$, then there is a $\xi \in [x, x + h]$ so that

$$f(x + h) - F_p(x, h) = \frac{1}{(p + 1)!} f^{(p+1)}(\xi) h^{p+1}. \quad (3.10)$$

If we take

$$C = \frac{1}{(p + 1)!} \max_{y \in [x, x+h]} |f^{(p+1)}(y)|,$$

then we find that

$$|F_p(x, h) - f(x + h)| \leq C \cdot h^{p+1}.$$

This is the proof of (3.9), which states that the Taylor series is an asymptotic expansion.

The approximation $F_p(x, h)$ includes terms in the sum (3.8) through order p . The first neglected term is $\frac{1}{(p+1)!} f^{(p+1)}(x)$, the term of order $p + 1$. This also is the difference $F_{p+1} - F_p$. It differs from the the Taylor series remainder (3.10) only in ξ being replaced by x . Since $\xi \in [x, x + h]$, this is a small change if h is small. Therefore, the error in the F_p is nearly equal to the first neglected term.

An asymptotic expansion can converge to the wrong answer or not converge at all. We give an example of each. These are based on the fact that exponentials *beat* polynomials in the sense that, for any n ,

$$t^n e^{-t} \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

If we take $t = a/|x|$ (because x may be positive or negative), this implies that

$$\frac{1}{x^n} e^{-a/|x|} \rightarrow 0 \quad \text{as } x \rightarrow 0. \quad (3.11)$$

Consider the function $f(x) = e^{-1/|x|}$. This function is continuous at $x = 0$ if we define $f(0) = 0$. The derivative at zero is (using (3.11))

$$f'(0) = \lim_{h \rightarrow 0} \frac{f(h) - f(0)}{h} = \lim_{h \rightarrow 0} \frac{e^{-1/|h|}}{h} = 0.$$

When $x \neq 0$, we calculate $f'(x) = \pm \frac{1}{|x|^2} e^{-1/|x|}$. The first derivative is continuous at $x = 0$ because (3.11) implies that $f'(x) \rightarrow 0 = f'(0)$ as $x \rightarrow 0$. Continuing in this way, one can see that each of the higher derivatives vanishes at $x = 0$ and is continuous. Therefore $F_p(0, h) = 0$ for any p , as $f^{(n)}(0) = 0$ for all n . Thus clearly $F_p(0, h) \rightarrow 0$ as $p \rightarrow \infty$. But $f(h) = e^{1/h} > 0$ if $h \neq 0$, so the Taylor series, while asymptotic, converges to the wrong answer.

⁵ See any good calculus book for a derivation and proof.

What goes wrong here is that the derivatives $f^{(p)}$ are zero at $x = 0$, but they are large for x close zero. The remainder theorem (3.10) implies that $F_p(x, h) \rightarrow f(x + h)$ as $p \rightarrow \infty$ if

$$M_p = \frac{h^p}{p!} \max_{x \leq \xi \leq x+h} |f^{(p)}(\xi)| \rightarrow 0 \quad \text{as } p \rightarrow \infty.$$

Taking $x = 0$ and any $h > 0$, function $f(x) = e^{-1/|x|}$ has $M_p \rightarrow \infty$ as $p \rightarrow \infty$.

Here is an example of an asymptotic Taylor series that does not converge at all. Consider

$$f(h) = \int_0^{1/2} e^{-x/h} \frac{1}{1-x} dx. \quad (3.12)$$

The integrand goes to zero *exponentially* as $h \rightarrow 0$ for any fixed x . This suggests⁶ that most of the integral comes from values of x near zero and that we can approximate the integral by approximating the integrand near $x = 0$. Therefore, we write $1/(1-x) = 1 + x + x^2 + \dots$, which converges for all x in the range of integration. Integrating separately gives

$$f(h) = \int_0^{1/2} e^{-x/h} dx + \int_0^{1/2} e^{-x/h} x dx + \int_0^{1/2} e^{-x/h} x^2 dx + \dots$$

We get a simple formula for the integral of the general term $e^{-x/h} x^n$ if we change the upper limit from $1/2$ to ∞ . For any fixed n , changing the upper limit of integration makes an exponentially small change in the integral, see problem (6). Therefore the n^{th} term is (for any $p > 0$)

$$\begin{aligned} \int_0^{1/2} e^{-x/h} x^n dx &= \int_0^{\infty} e^{-x/h} x^n dx + O(h^p) \\ &= n! h^{n+1} + O(h^p). \end{aligned}$$

Assembling these gives

$$f(h) \sim h + h^2 + 2h^3 + \dots + (n-1)! \cdot h^n + \dots \quad (3.13)$$

This is an asymptotic expansion because the partial sums are asymptotic approximations:

$$|h + h^2 + 2h^3 + \dots + (p-1)! \cdot h^p - f(h)| = O(h^{p+1}).$$

But the infinite sum does not converge; for any $h > 0$ we have $n! \cdot h^{n+1} \rightarrow \infty$ as $n \rightarrow \infty$.

In these examples, the higher order approximations have smaller ranges of validity. For (3.12), the three term approximation $f(h) \approx h + h^2 + 2h^3$ is reasonably accurate when $h = .3$ but the six term approximation is less accurate, and the ten term “approximation” is 4.06 for an answer less than .5. The ten term approximation is very accurate when $h = .01$ but the fifty term “approximation” is astronomical.

⁶A more precise version of this intuitive argument is in exercise 6.

3.2 Numerical Differentiation

One basic numerical task is estimating the derivative of a function from function values. Suppose we have a smooth function, $f(x)$, of a single variable, x . The problem is to combine several values of f to estimate f' . These *finite difference*⁷ approximations are useful in themselves, and because they underlie methods for solving differential equations. Several common finite difference approximations are

$$f'(x) \approx \frac{f(x+h) - f(x)}{h} \quad (3.14a)$$

$$f'(x) \approx \frac{f(x) - f(x-h)}{h} \quad (3.14b)$$

$$f'(x) \approx \frac{f(x+h) - f(x-h)}{2h} \quad (3.14c)$$

$$f'(x) \approx \frac{-f(x+2h) + 4f(x+h) - 3f(x)}{2h} \quad (3.14d)$$

$$f'(x) \approx \frac{-f(x+2h) + 8f(x+h) - 8f(x-h) + f(x+2h)}{12h} \quad (3.14e)$$

Formulas (3.14a)-(3.14c) have simple geometric interpretations as the slopes of lines connecting nearby points on the graph of $f(x)$. A carefully drawn figure shows that (3.14c) is more accurate than (3.14a). We give an analytical explanation of this below. Formulas (3.14d) and (3.14e) are more technical. The formulas (3.14a), (3.14b), and (3.14d) are *one sided* because they use values only on one side of x . The formulas (3.14c) and (3.14e) are *centered* because they use points symmetrical about x and with opposite weights.

The Taylor series expansion (3.8) allows us to calculate the accuracy of each of these approximations. Let us start with the simplest, the first-order formula (3.14a). Substituting (3.8) into the right side of (3.14a) gives

$$\frac{f(x+h) - f(x)}{h} \sim f'(x) + h \frac{f''(x)}{2} + h^2 \frac{f'''(x)}{6} + \dots \quad (3.15)$$

This may be written:

$$\frac{f(x+h) - f(x)}{h} = f'(x) + E_a(h) \ ,$$

where

$$E_a(h) \sim \frac{1}{2}f''(x) \cdot h + \frac{1}{6}f'''(x) \cdot h^2 + \dots \quad (3.16)$$

In particular, this shows that $E_a(h) = O(h)$, which means that the one sided two point finite difference approximation is first order accurate. Moreover,

$$E_a(h) = \frac{1}{2}f''(x) \cdot h + O(h^2) \ , \quad (3.17)$$

⁷ Isaac Newton thought of the differential dx as something infinitely small and yet not zero. He called such quantities *infinitesimals*. The infinitesimal difference was $df = f(x+dx) - f(x)$. The *finite* difference occurs when the change in x is finitely small but not infinitely small.

which is to say that, to *leading order*, the error is proportional to h and given by $\frac{1}{2}f''(x)h$.

Taylor series analysis applied to the two point centered difference approximation (3.14c) leads to

$$f'(x) = \frac{f(x+h) - f(x-h)}{2h} + E_c(h)$$

where

$$\begin{aligned} E_c(h) &\sim \frac{1}{6}f'''(x) \cdot h^2 + \frac{1}{24}f^{(5)}(x) \cdot h^4 + \dots \\ &= \frac{1}{6}f'''(x) \cdot h^2 + O(h^4) \end{aligned} \quad (3.18)$$

This centered approximation is second order accurate, $E_c(h) = O(h^2)$. This is one order more accurate than the one sided approximations (3.14a) and (3.14b). Any centered approximation such as (3.14c) or (3.14e) must be at least second order accurate because of the symmetry relation $\hat{A}(-h) = \hat{A}(h)$. Since $A = f'(x)$ is independent of h , this implies that $E(h) = \hat{A}(h) - A$ is symmetric. If $E(h) = c \cdot h + O(h^2)$, then

$$E(-h) = -c \cdot h + O(h^2) = E(h) + O(h^2) \approx -E(h) \quad \text{for small } h,$$

which contradicts $E(-h) = E(h)$. The same kind of reasoning shows that the $O(h^3)$ term in (3.18) must be zero.

A Taylor series analysis shows that the three point one sided formula (3.14d) is second order accurate, while the four point centered approximation (3.14e) is fourth order. Sections 3.3.1 and 3.5 give two ways to find the coefficients 4, -3, and 8 achieve these higher orders of accuracy.

Figure 3.2 illustrates many of these features. The first is that the higher order formulas (3.14c), (3.14d), and (3.14e) are more accurate when h is small. For $h = .5$, the first order two point one sided difference formula is more accurate than the second order accurate three point formula, but their proper asymptotic ordering is established by $h = .01$. For $h \leq 10^{-5}$ with the fourth order centered difference formula and $h = 10^{-7}$ with the second order formula, double precision roundoff error makes the results significantly different from what they would be in exact arithmetic. The rows labeled \hat{E} give the leading order Taylor series estimate of the error. For the first order formula, (3.17) shows that this is $\hat{E}(h) = \frac{1}{2}f''(x)h$. For the second order centered formula, (3.18) gives leading order error $\hat{E}_c(h) = \frac{1}{6}f'''(x)h^2$. For the three point one sided formula, the coefficient of $f'''(x)h^2$ is $\frac{1}{3}$, twice the coefficient for the second order centered formula. For the fourth order formula, the coefficient of $f^{(5)}(x)h^4$ is $\frac{1}{30}$. The table shows that \hat{E} is a good predictor of E , if h is at all small, until roundoff gets in the way. The smallest error⁸ in the table comes from the fourth order

⁸The error would have been -3×10^{-19} rather than -6×10^{-12} , seven orders of magnitude smaller, in exact arithmetic. The best answer comes despite some catastrophic cancellation, but not completely catastrophic.

h	(3.14a)	(3.14c)	(3.14d)	(3.14e)
.5	3.793849	0.339528	7.172794	0.543374
E	2.38e+00	-1.08e+00	5.75e+00	-8.75e-01
\widehat{E}	5.99e+00	-1.48e+00	-2.95e+00	-1.85e+00
.01	2.533839	1.359949	1.670135	1.415443
E	1.12e+00	-5.84e-02	2.52e-01	-2.87e-03
\widehat{E}	1.20e+00	-5.91e-02	-1.18e-01	-2.95e-03
5×10^{-3}	1.999796	1.403583	1.465752	1.418128
E	5.81e-01	-1.47e-02	4.74e-02	-1.83e-04
\widehat{E}	5.99e-01	-1.48e-02	-2.95e-02	-1.85e-04
10^{-3}	1.537561	1.417720	1.419642	1.418311
E	1.19e-01	-5.91e-04	1.33e-03	-2.95e-07
\widehat{E}	1.20e-01	-5.91e-04	-1.18e-03	-2.95e-07
10^{-5}	1.418431	1.418311	1.418311	1.418311
E	1.20e-04	-5.95e-10	1.16e-09	-6.05e-12
\widehat{E}	1.20e-04	-5.91e-10	-1.18e-09	-2.95e-19
10^{-7}	1.418312	1.418311	1.418311	1.418311
E	1.20e-06	2.76e-10	3.61e-09	8.31e-10
\widehat{E}	1.20e-06	-5.91e-14	-1.18e-13	-2.95e-27

Figure 3.2: Estimates of $f'(x)$ with $f(x) = \sin(5x)$ and $x = 1$ using formulas (3.14a), (3.14c), (3.14d), and (3.14e). Each group of three rows corresponds to one h value. The top row gives the finite difference estimate of $f'(x)$, the middle row gives the error $E(h)$, and the third row is $\widehat{E}(h)$, the leading Taylor series term in the error formula. All calculations were done in double precision floating point arithmetic.

formula and $h = 10^{-5}$. It is impossible to have an error this small with a first or second order formula no matter what the step size. Note that the error in the (3.14e) column increased when h was reduced from 10^{-5} to 10^{-7} because of roundoff.

A difference approximation may not achieve its expected order of accuracy if the requisite derivatives are infinite or do not exist. As an example of this, let $f(x)$ be the function

$$f(x) = \begin{cases} 0 & \text{if } x \leq 0 \\ x^2 & \text{if } x \geq 0 \end{cases}.$$

If we want $f'(0)$, the formulas (1c) and (1e) are only first order accurate despite their higher accuracy for smoother functions. This f has a mild singularity, a discontinuity in its second derivative. Such a singularity is hard to spot on a graph, but may have a drastic effect on the numerical analysis of the function.

We can use finite differences to approximate higher derivatives such as

$$\frac{f(x+h) - 2f(x) + f(x-h)}{h^2} = f''(x) + \frac{h^2}{12}f^{(4)} + O(h^4) ,$$

and to estimate partial derivatives of functions depending on several variables, such as

$$\frac{f(x+h, y) - f(x-h, y)}{2h} \sim \frac{\partial}{\partial x}f(x, y) + \frac{h^2}{3} \frac{\partial^3 f}{\partial x^3}(x, y) + \cdots .$$

3.2.1 Mixed partial derivatives

Several new features arise only when evaluating mixed partial derivatives or sums of partial derivatives in different variables. For example, suppose we want to evaluate⁹ $f_{xy} = \partial_x \partial_y f(x, y)$. Rather than using the same h for both¹⁰ x and y , we use step size Δx for x and Δy for y . The first order one sided approximation for f_y is

$$f_y \approx \frac{f(x, y + \Delta y) - f(x, y)}{\Delta y} .$$

We might hope this, and

$$f_y(x + \Delta x, y) \approx \frac{f(x + \Delta x, y + \Delta y) - f(x + \Delta x, y)}{\Delta y} ,$$

are accurate enough so that

$$\begin{aligned} \partial_x(\partial_y f) &\approx \frac{f_y(x + \Delta x, y) - f_y(x, y)}{\Delta x} \\ &\approx \frac{\frac{f(x + \Delta x, y + \Delta y) - f(x + \Delta x, y)}{\Delta y} - \frac{f(x, y + \Delta y) - f(x, y)}{\Delta y}}{\Delta x} \\ f_{xy} &\approx \frac{f(x + \Delta x, y + \Delta y) - f(x + \Delta x, y) - f(x, y + \Delta y) + f(x, y)}{\Delta x \Delta y} \end{aligned} \quad (3.19)$$

is consistent¹¹.

To understand the error in (3.19), we need the Taylor series for functions of more than one variable. The rigorous remainder theorem is more complicated,

⁹ We abbreviate formulas by denoting partial derivatives by subscripts, $\partial_x f = f_x$, etc., and by leaving out the arguments if they are (x, y) , so $f(x + \Delta x, y) - f(x, y) = f(x + \Delta x, y) - f \approx \Delta x f_x(x, y) = \Delta x f_x$.

¹⁰ The expression $f(x+h, y+h)$ does not even make sense if x and y have different physical units.

¹¹ The same calculation shows that the right side of (3.19) is an approximation of $\partial_y(\partial_x f)$. This is one proof that $\partial_y \partial_x f = \partial_x \partial_y f$.

but it suffices here to use all of the “first” neglected terms. The expansion is

$$\begin{aligned}
f(x + \Delta x, y + \Delta y) &\sim f + \Delta x f_x + \Delta y f_y \\
&\quad + \frac{1}{2} \Delta x^2 f_{xx} + \Delta x \Delta y f_{xy} + \frac{1}{2} \Delta y^2 f_{yy} \\
&\quad + \frac{1}{6} \Delta x^3 f_{xxx} + \frac{1}{2} \Delta x^2 \Delta y f_{xxy} + \frac{1}{2} \Delta x \Delta y^2 f_{xyy} + \frac{1}{6} \Delta y^3 f_{yyy} \\
&\quad + \dots \\
&\quad + \frac{1}{p!} \sum_{k=0}^p \binom{p}{k} \Delta x^{p-k} \Delta y^k \partial_x^{p-k} \partial_y^k f + \dots .
\end{aligned}$$

If we keep just the terms on the top row on the right, the second order terms on the second row are the first neglected terms, and (using the inequality $\Delta x \Delta y \leq \Delta x^2 + \Delta y^2$):

$$f(x + \Delta x, y + \Delta y) = f + \Delta x f_x + \Delta y f_y + O(\Delta x^2 + \Delta y^2) .$$

Similarly,

$$\begin{aligned}
f(x + \Delta x, y + \Delta y) &= f + \Delta x f_x + \Delta y f_y + \frac{1}{2} \Delta x^2 f_{xx} + \Delta x \Delta y f_{xy} + \frac{1}{2} \Delta y^2 f_{yy} \\
&\quad + O(\Delta x^3 + \Delta y^3) .
\end{aligned}$$

Of course, the one variable Taylor series is

$$f(x + \Delta x, y) = f + \Delta x f_x + \frac{1}{2} \Delta x^2 f_{xx} + O(\Delta x^3) , \text{ etc.}$$

Using all these, and some algebra, gives

$$\begin{aligned}
&\frac{f(x + \Delta x, y + \Delta y) - f(x + \Delta x, y) - f(x, y + \Delta y) + f}{\Delta x \Delta y} \\
&= f_{xy} + O\left(\frac{\Delta x^3 + \Delta y^3}{\Delta x \Delta y}\right) . \tag{3.20}
\end{aligned}$$

This shows that the approximation (3.19) is first order, at least if Δx is roughly proportional to Δy . The second order Taylor expansion above gives a quantitative estimate of the error:

$$\begin{aligned}
&\frac{f(x + \Delta x, y + \Delta y) - f(x + \Delta x, y) - f(x, y + \Delta y) + f}{\Delta x \Delta y} - f_{xy} \\
&\approx \frac{1}{2} (\Delta x f_{xxy} + \Delta y f_{xyy}) . \tag{3.21}
\end{aligned}$$

This formula suggests (and it is true) that in exact arithmetic we could let $\Delta x \rightarrow 0$, with Δy fixed but small, and still have a reasonable approximation to f_{xy} . The less detailed version (3.20) suggests that might not be so.

A partial differential equation may involve a *differential operator* that is a sum of partial derivatives. One way to approximate a differential operator is to approximate each of the terms separately. For example, the *Laplace operator* (or *Laplacian*), which is $\Delta = \partial_x^2 + \partial_y^2$ in two dimensions, may be approximated by

$$\begin{aligned}\Delta f(x, y) &= \partial_x^2 f + \partial_y^2 f \\ &\approx \frac{f(x + \Delta x, y) - 2f + f(x - \Delta x, y)}{\Delta x^2} \\ &\quad + \frac{f(x, y + \Delta y) - 2f + f(x, y - \Delta y)}{\Delta y^2} .\end{aligned}$$

If $\Delta x = \Delta y = h$ (x and y have the same units in the Laplace operator), then this becomes

$$\Delta f \approx \frac{1}{h^2} (f(x + h, y) + f(x - h, y) + f(x, y + h) + f(x, y - h) - 4f) . \quad (3.22)$$

This is the *standard* five point approximation (seven points in three dimensions). The leading error term is

$$\frac{h^2}{12} (\partial_x^4 f + \partial_y^4 f) . \quad (3.23)$$

The simplest *heat equation* (or *diffusion equation*) is $\partial_t f = \frac{1}{2} \partial_x^2 f$. The *space* variable, x , and the *time* variable, t have different units. We approximate the differential operator using a first order forward difference approximation in time and a second order centered approximation in space. This gives

$$\partial_t f - \frac{1}{2} \partial_x^2 f \approx \frac{f(x, t + \Delta t) - f}{\Delta t} - \frac{f(x + \Delta x, t) - 2f + f(x - \Delta x, t)}{2\Delta x^2} . \quad (3.24)$$

The leading order error is the sum of the leading errors from time differencing ($\frac{1}{2} \Delta t \partial_t^2 f$) and space differencing ($\frac{\Delta x^2}{24} \partial_x^4 f$), which is

$$\frac{1}{2} \Delta t \partial_t^2 f - \frac{\Delta x^2}{24} \partial_x^4 f . \quad (3.25)$$

For many reasons, people often take Δt proportional to Δx^2 . In the simplest case of $\Delta t = \Delta x^2$, the leading error becomes

$$\Delta x^2 \left(\frac{1}{2} \partial_t^2 f - \frac{1}{24} \partial_x^4 f \right) .$$

This shows that the overall approximation (3.24) is second order accurate if we take the time step to be the square of the space step.

3.3 Error Expansions and Richardson Extrapolation

The error expansions (3.16) and (3.18) above are instances of a common situation that we now describe more systematically and abstractly. We are trying

to compute A and there is an approximation with

$$\hat{A}(h) \rightarrow A \text{ as } h \rightarrow 0 .$$

The error is $E(h) = \hat{A}(h) - A$. A general *asymptotic error expansion* in powers of h has the form

$$\hat{A}(h) \sim A + h^{p_1} A_1 + h^{p_2} A_2 + \cdots , \quad (3.26)$$

or, equivalently,

$$E(h) \sim h^{p_1} A_1 + h^{p_2} A_2 + \cdots .$$

As with Taylor series, the expression (3.26) does not imply that the series on the right converges to $\hat{A}(h)$. Instead, the asymptotic relation (3.26) means that, as $h \rightarrow 0$,

$$\left. \begin{aligned} \hat{A}(h) - (A + h^{p_1} A_1) &= O(h^{p_2}) & (a) \\ \hat{A}(h) - (A + h^{p_1} A_1 + h^{p_2} A_2) &= O(h^{p_3}) & (b) \\ \text{and so on.} \end{aligned} \right\} \quad (3.27)$$

It goes without saying that $0 < p_1 < p_2 < \cdots$. The statement (3.27a) says not only that $A + A_1 h^{p_1}$ is a good approximation to $\hat{A}(h)$, but that the error has the same order as the first neglected term, $A_2 h^{p_2}$. The statement (3.27b) says that including the $O(h^{p_2})$ term improves the approximation to $O(h^{p_3})$, and so on.

Many asymptotic error expansions arise from Taylor series manipulations. For example, the two point one sided difference formula error expansion (3.15) gives $p_1 = 1$, $A_1 = \frac{1}{2} f''(x)$, $p_2 = 2$, $A_2 = \frac{1}{6} f'''(x)$, etc. The error expansion (3.18) for the two point centered difference formula implies that $p_1 = 2$, $p_2 = 4$, $A_1 = \frac{1}{6} f'''(x)$, and $A_2 = \frac{1}{24} f^{(5)}(x)$. The three point one sided formula has $p_1 = 2$ because it is second order accurate, but $p_2 = 3$ instead of $p_2 = 4$. The fourth order formula has $p_1 = 4$ and $p_2 = 6$.

It is possible that an approximation is p^{th} order accurate in the big O sense, $|E(h)| \leq Ch^p$, without having an asymptotic error expansion of the form (3.26). Figure 3.4 has an example showing that this can happen when the function $f(x)$ is not sufficiently smooth. Most of the extrapolation and debugging tricks described here do not apply in those cases.

We often work with asymptotic error expansions for which we know the powers p_k but not the coefficients, A_k . For example, in finite difference approximations, the A_k depend on the function f but the p_k do not. Two techniques that use this information are *Richardson extrapolation* and *convergence analysis*. Richardson extrapolation combines $\hat{A}(h)$ approximations for several values of h to produce a new approximation that has greater order of accuracy than $\hat{A}(h)$. Convergence analysis is a debugging method that tests the order of accuracy of numbers produced by a computer code.

3.3.1 Richardson extrapolation

Richardson extrapolation increases the order of accuracy of an approximation provided that the approximation has an asymptotic error expansion of the form (3.26) with known p_k . In its simplest form, we compute $\hat{A}(h)$ and $\hat{A}(2h)$ and then form a linear combination that eliminates the leading error term. Note that

$$\begin{aligned}\hat{A}(2h) &= A + (2h)^{p_1} A_1 + (2h)^{p_2} A_2 + \cdots \\ &= A + 2^{p_1} h^{p_1} A_1 + 2^{p_2} h^{p_2} A_2 + \cdots ,\end{aligned}$$

so

$$\frac{2^{p_1} \hat{A}(h) - \hat{A}(2h)}{2^{p_1} - 1} = A + \frac{2^{p_1} - 2^{p_2}}{2^{p_1} - 1} h^{p_2} A_2 + \frac{2^{p_3} - 2^{p_2}}{2^{p_1} - 1} h^{p_3} A_3 + \cdots .$$

In other words, the *extrapolated* approximation

$$\hat{A}^{(1)}(h) = \frac{2^{p_1} \hat{A}(h) - \hat{A}(2h)}{2^{p_1} - 1} \quad (3.28)$$

has order of accuracy $p_2 > p_1$. It also has an asymptotic error expansion,

$$\hat{A}^{(1)}(h) = A + h^{p_2} A_2^{(1)} + h^{p_3} A_3^{(1)} + \cdots ,$$

where $A_2^{(1)} = \frac{2^{p_1} - 2^{p_2}}{2^{p_1} - 1} A_2$, and so on.

Richardson extrapolation can be repeated to remove more asymptotic error terms. For example,

$$\hat{A}^{(2)}(h) = \frac{2^{p_2} \hat{A}^{(1)}(h) - \hat{A}^{(1)}(2h)}{2^{p_2} - 1}$$

has order p_3 . Since $\hat{A}^{(1)}(h)$ depends on $\hat{A}(h)$ and $\hat{A}(2h)$, $\hat{A}^{(2)}(h)$ depends on $\hat{A}(h)$, $\hat{A}(2h)$, and $\hat{A}(4h)$. It is not necessary to use powers of 2, but this is natural in many applications. Richardson extrapolation will not work if the underlying approximation, $\hat{A}(h)$, has accuracy of order h^p in the $O(h^p)$ sense without at least one term of an asymptotic expansion.

Richardson extrapolation allows us to derive higher order difference approximations from low order ones. Start, for example, with the first order one sided approximation to $f'(x)$ given by (3.14a). Taking $p_1 = 1$ in (3.28) leads to the second order approximation

$$\begin{aligned}f'(x) &\approx 2 \frac{f(x+h) - f(x)}{h} - \frac{f(x+2h) - f(x)}{2h} \\ &= \frac{-f(x+2h) + 4f(x+h) - 3f(x)}{2h} ,\end{aligned}$$

which is the second order three point one sided difference approximation (3.14d). Starting with the second order centered approximation (3.14c) (with $p_1 = 2$ and

$p_2 = 4$) leads to the fourth order approximation (3.14e). The second order one sided formula has $p_1 = 2$ and $p_2 = 3$. Applying Richardson extrapolation to it gives a one sided formula that uses $f(x + 4h)$, $f(x + 2h)$, $f(x + h)$, and $f(x)$ to give a third order approximation. A better third order one sided approximation would use $f(x + 3h)$ instead of $f(x + 4h)$. Section 3.5 explains how to do this.

Richardson extrapolation may also be applied to the output of a complex code. Run it with step size h and $2h$ and apply (3.28) to the output. This is sometimes applied to differential equations as an alternative to making up high order schemes from scratch, which can be time consuming and intricate.

3.3.2 Convergence analysis

We can test a code, and the algorithm it is based on, using ideas related to Richardson extrapolation. A naive test would be to do runs with decreasing h values to check whether $\hat{A}(h) \rightarrow A$ as $h \rightarrow 0$. A *convergence analysis* based on asymptotic error expansions can be better. For one thing, we might not know A . Even if we run a test case where A is known, it is common that a code with mistakes limps to convergence, but not as accurately or reliably as the correct code would. If we bother to write a code that is more than first order accurate, we should test that we are getting the order of accuracy we worked for.

There are two cases, the case where A is known and the case where A is not known. While we would not write a code solely to solve problems for which we know the answers, such problems are useful test cases. Ideally, a code should come with a suite of tests, and some of these tests should be nontrivial. For example, the fourth-order approximation (3.14e) gives the exact answer for any polynomial of degree less than five, so a test suite of only low-order polynomials would be inappropriate for this rule.

If A is known, we can run the code with step size h and $2h$ and, from the resulting approximations, $\hat{A}(h)$ and $\hat{A}(2h)$, compute

$$\begin{aligned} E(h) &\approx A_1 h^{p_1} + A_2 h^{p_2} + \dots, \\ E(2h) &\approx 2^{p_1} A_1 h^{p_1} + 2^{p_2} A_2 h^{p_2} + \dots. \end{aligned}$$

For small h the first term is a good enough approximation so that the ratio should be approximately the characteristic value

$$R(h) = \frac{E(2h)}{E(h)} \approx 2^{p_1}. \quad (3.29)$$

Figure 3.3 illustrates this phenomenon. As $h \rightarrow 0$, the ratios converge to the expected result $2^{p_1} = 2^3 = 8$. Figure 3.4 shows what may happen when we apply this convergence analysis to an approximation that is second order accurate in the big O sense without having an asymptotic error expansion. The error gets very small but the error ratio does not have simple behavior as in Figure 3.3. The difference between the convergence in these two cases appears clearly in the log-log error plots shown in Figure 3.5.

h	Error: $E(h)$	Ratio: $E(h)/E(h/2)$
.1	4.8756e-04	3.7339e+00
.05	1.3058e-04	6.4103e+00
.025	2.0370e-05	7.3018e+00
.0125	2.7898e-06	7.6717e+00
6.2500e-03	3.6364e-07	7.8407e+00
3.1250e-03	4.6379e-08	7.9215e+00
1.5625e-03	5.8547e-09	7.9611e+00
7.8125e-04	7.3542e-10	—————

Figure 3.3: Convergence study for a third order accurate approximation. As $h \rightarrow 0$, the ratio converges to $2^3 = 8$. The h values in the left column decrease by a factor of two from row to row.

h	Error: $E(h)$	Ratio: $E(h)/E(h/2)$
.1	1.9041e-02	2.4014e+00
.05	7.9289e-03	1.4958e+01
.025	5.3008e-04	-1.5112e+00
.0125	-3.5075e-04	3.0145e+00
6.2500e-03	-1.1635e-04	1.9880e+01
3.1250e-03	-5.8529e-06	-8.9173e-01
1.5625e-03	6.5635e-06	2.8250e+00
7.8125e-04	2.3233e-06	—————

Figure 3.4: Convergence study for an approximation that is second order accurate in the sense that $|E(h)| = O(h^2)$ but that has no asymptotic error expansion. The h values are the same as in Figure 3.3. The errors decrease in an irregular fashion.

Convergence analysis can be applied even when A is not known. In this case we need three approximations, $\hat{A}(4h)$, $\hat{A}(2h)$, and $\hat{A}(h)$. Again assuming the existence of an asymptotic error expansion (3.26), we get, for small h ,

$$R'(h) = \frac{\hat{A}(4h) - \hat{A}(2h)}{\hat{A}(2h) - \hat{A}(h)} \approx 2^{p_1} . \quad (3.30)$$

3.4 Integration

Numerical integration means finding approximations for quantities such as

$$I = \int_a^b f(x) dx .$$

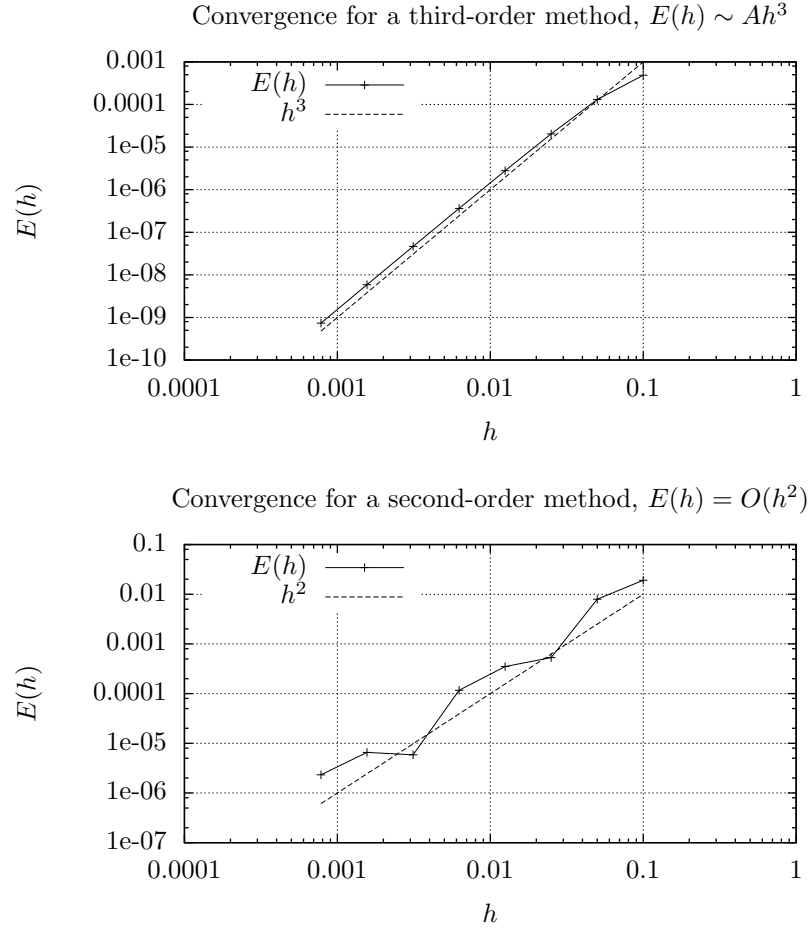


Figure 3.5: Log-log plots of the convergence study data in Figures 3.3 (top) and 3.4 (bottom). The existence of an asymptotic error expansion in the first example shows up graphically as convergence to a straight line with slope 3.

Rectangle	$\hat{I}_k = h_k f(x_k)$	1 st order
Trapezoid	$\hat{I}_k = \frac{h_k}{2} (f(x_k) + f(x_{k+1}))$	2 nd order
Midpoint	$\hat{I}_k = h_k f(x_{k+1/2})$	2 nd order
Simpson	$\hat{I}_k \approx \frac{h_k}{6} (f(x_k) + 4f(x_{k+1/2}) + f(x_{k+1}))$	4 th order
2 point GQ	$\hat{I}_k = \frac{h_k}{2} (f(x_{k+1/2} - h_k \xi) + f(x_{k+1/2} + h_k \xi))$	4 th order
3 point GQ	$\hat{I}_k = \frac{h_k}{18} (5f(x_{k+1/2} - h_k \eta) + 8f(x_{k+1/2}) + 5f(x_{k+1/2} + h_k \eta))$	6 th order

Figure 3.6: Common panel integration rules. The last two are Gauss quadrature (Gauss – Legendre to be precise) formulas. The definitions are $\xi = \frac{1}{2\sqrt{3}}$ and $\eta = \frac{1}{2}\sqrt{\frac{3}{5}}$.

We discuss only *panel methods* here, though there are other elegant methods. In a panel method, the integration interval, $[a, b]$, is divided into n subintervals, or *panels*, $P_k = [x_k, x_{k+1}]$, where $a = x_0 < x_1 < \dots < x_n = b$. If the panel P_k is small, we can get an accurate approximation to

$$I_k = \int_{P_k} f(x) dx = \int_{x_k}^{x_{k+1}} f(x) dx \quad (3.31)$$

using a few evaluations of f inside P_k . Adding these approximations gives an approximation to I :

$$\hat{I} = \sum_{k=0}^{n-1} \hat{I}_k. \quad (3.32)$$

Some common panel integral approximations are given in Figure 3.6, where we write $x_{k+1/2} = (x_{k+1} + x_k)/2$ for the midpoint of the panel and $h_k = x_{k+1} - x_k$ is the width. Note that x_k is the left endpoint of P_k and the right endpoint of P_{k-1} . In the trapezoid rule and Simpson's rule, we need not evaluate $f(x_k)$ twice.

For our error analysis, we assume that all the panels are the same size

$$h = \Delta x = |P_k| = x_{k+1} - x_k \text{ for all } k.$$

Given this restriction, not every value of h is allowed because $b - a = nh$ and n is an integer. When we take $h \rightarrow 0$, we will assume that h only takes allowed values $h = (b - a)/n$. The *local truncation error* is the integration error over one panel. The *global error* is the sum of the local truncation errors in all the panels. The global error usually is one power of h larger than the local truncation error. If the error per panel is $O(h^q)$, then the total error will be of the order of nh^q , where n is the number of panels. Since $n = (b - a)/h$, this suggests that the global error will be of order $h^q \cdot (b - a)/h = O(h^{q-1})$.

For the local truncation error analysis, let $P = [x_*, x_* + h]$ be a generic panel. The panel integration rule approximates the panel integral

$$I_P = \int_P f(x)dx = \int_{x_*}^{x_*+h} f(x)dx$$

with the approximation, \hat{I}_P . For example, the rectangle rule (top row of Figure 3.6) has panel integration rule

$$\int_{x_*}^{x_*+h} f(x)dx \approx \hat{I}_P(h) = hf(x_*) .$$

To estimate the difference between I_P and $\hat{I}_P(h)$, we expand f in a Taylor series about x_* :

$$f(x) \sim f(x_*) + f'(x_*)(x - x_*) + \frac{1}{2}f''(x_*)(x - x_*)^2 + \cdots .$$

Integrating this term by term leads to

$$\begin{aligned} I_P &\sim \int_P f(x_*)dx + \int_P f'(x_*)(x - x_*)dx + \cdots \\ &= f(x_*)h + \frac{1}{2}f'(x_*)h^2 + \frac{1}{6}f''(x_*)h^3 + \cdots . \end{aligned}$$

The error in integration over this panel then is

$$E(P, h) = \hat{I}_P(h) - I_P \sim -\frac{1}{2}f'(x_*)h^2 - \frac{1}{6}f''(x_*)h^3 - \cdots . \quad (3.33)$$

This shows that the local truncation error for the rectangle rule is $O(h^2)$ and identifies the leading error coefficient.

$$\begin{aligned} E &= \hat{I} - I \\ &= \sum_{k=0}^{n-1} \hat{I}_k - I_k \\ E &\sim -\sum_{k=0}^{n-1} \frac{1}{2}f'(x_k)h^2 - \sum_{k=0}^{n-1} \frac{1}{6}f''(x_k)h^3 - \cdots . \end{aligned} \quad (3.34)$$

We sum over k and use simple inequalities to get the order of magnitude of the global error:

$$\begin{aligned} |E| &< \approx \frac{1}{2} \sum_{k=0}^{n-1} |f'(x_k)| h^2 \\ &\leq n \cdot \frac{1}{2} \max_{a \leq x \leq b} |f'(x)| \cdot h^2 \\ &= \frac{b-a}{h} O(h^2) \\ &= O(h) . \end{aligned}$$

This shows that the rectangle rule is first order accurate overall.

Looking at the global error in more detail leads to an asymptotic error expansion. Applying the rectangle rule error bound to another function, $g(x)$, we have

$$\sum_{k=0}^{n-1} g(x_k)h = \int_a^b g(x)dx + O(h) .$$

Taking $g(x) = f'(x)$ gives

$$\sum_{k=0}^{n-1} f'(x_k)h = \int_a^b f'(x)dx + O(h) = f(b) - f(a) + O(h) .$$

From (3.34) we have

$$\begin{aligned} E &\approx - \left(\sum_{k=0}^{n-1} f'(x_k)h \right) \frac{h}{2} \\ &\approx - \left(\int_a^b f'(x)dx \right) \frac{h}{2} \\ E &\approx -\frac{1}{2} (f(b) - f(a)) h . \end{aligned} \tag{3.35}$$

This gives the first term in the asymptotic error expansion. It shows that the leading error not only is bounded by h , but roughly is proportional to h . It also demonstrates the curious fact that if f is differentiable then the leading error term is determined by the values of f at the endpoints and is independent of the values of f between. This is not true if f has a discontinuity in the interval $[a, b]$.

To get the next term in the error expansion, apply (3.34) to the error itself, i.e.

$$\begin{aligned} \sum_{k=0}^{n-1} f'(x_k)h &= \int_a^b f'(x)dx - \frac{h}{2} (f'(b) - f'(a)) + O(h^2) \\ &= f(b) - f(a) - \frac{h}{2} (f'(b) - f'(a)) + O(h^2) . \end{aligned}$$

In the same way, we find that

$$\sum_{k=0}^{n-1} f''(x_k) \frac{h^3}{6} = (f'(b) - f'(a)) \frac{h^2}{6} + O(h^3) .$$

Combining all these gives the first two terms in the error expansion:

$$E(h) \sim -\frac{1}{2} (f(b) - f(a)) h + \frac{1}{12} (f'(b) - f'(a)) h^2 + \dots . \tag{3.36}$$

It is clear that this procedure can be used to continue the expansion as far as we want, but you would have to be very determined to compute, for example,

n	Computed Integral	Error	Error/h	$(E - A_1 h)/h^2$	$(E - A_1 h - A_2 h^2)/h^3$
10	3.2271	-0.2546	-1.6973	0.2900	-0.7250
20	3.3528	-0.1289	-1.7191	0.2901	-0.3626
40	3.4168	-0.0649	-1.7300	0.2901	-0.1813
80	3.4492	-0.0325	-1.7354	0.2901	-0.0907
160	3.4654	-0.0163	-1.7381	0.2901	-0.0453

Figure 3.7: Experiment illustrating the asymptotic error expansion for rectangle rule integration.

n	Computed Integral	Error	Error/h	$(E - A_1 h)/h^2$
10	7.4398e-02	-3.1277e-02	-3.1277e-01	-4.2173e-01
20	9.1097e-02	-1.4578e-02	-2.9156e-01	-4.1926e-01
40	9.8844e-02	-6.8314e-03	-2.7326e-01	-1.0635e-01
80	1.0241e-01	-3.2605e-03	-2.6084e-01	7.8070e-01
160	1.0393e-01	-1.7446e-03	-2.7914e-01	-1.3670e+00
320	1.0482e-01	-8.5085e-04	-2.7227e-01	-5.3609e-01
640	1.0526e-01	-4.1805e-04	-2.6755e-01	1.9508e+00
1280	1.0546e-01	-2.1442e-04	-2.7446e-01	-4.9470e+00
2560	1.0557e-01	-1.0631e-04	-2.7214e-01	-3.9497e+00
5120	1.0562e-01	-5.2795e-05	-2.7031e-01	1.4700e+00

Figure 3.8: Experiment illustrating the breakdown of the asymptotic expansion for a function with a continuous first derivative but discontinuous second derivative.

the coefficient of h^4 . An elegant and more systematic discussion of this error expansion is carried out in the book of Dahlquist and Bjork. The resulting error expansion is called the *Euler Maclaurin* formula. The coefficients $1/2$, $1/12$, and so on, are related to the *Bernoulli numbers*.

The error expansion (3.36) will not be valid if the integrand, f , has singularities inside the domain of integration. Suppose, for example, that $a = 0$, $b = 1$, $u = 1/\sqrt{2}$, and $f(x) = 0$ for $x \leq u$ and $f(x) = \sqrt{x - u}$ for $x \geq u$. In this case the error expansion for the rectangle rule approximation to $\int_0^1 f(x)dx$ has one valid term only. This is illustrated in Figure 3.8. The “Error/h” column shows that the first coefficient, A_1 , exists. Moreover, A_1 is given by the formula (3.36). The numbers in the last column do not tend to a limit. This shows that the coefficient A_2 does not exist. The error expansion does not exist beyond the first term.

The analysis of the higher order integration methods listed in Figure 3.6 is easier if we use a symmetric basic panel. From now on, the panel of length h will have x_* in the center, rather at the left end, that is

$$P = [x_* - h/2, x_* + h/2] \quad .$$

If we now expand $f(x)$ in a Taylor series about x_* and integrate term by term,

we get

$$\int_P f(x)dx = \int_{x=x_*-\frac{h}{2}}^{x_*+\frac{h}{2}} f(x)dx \sim f(x_*)h + \frac{f''(x_*)}{24}h^3 + \frac{f^{(4)}(x_*)}{1920}h^5 + \dots$$

For the midpoint rule, this leads to a global error expansion in even powers of h , $E \approx A_1 h^2 + A_2 h^4 + \dots$, with $A_1 = (f'(b) - f'(a))/24$. Each of the remaining panel methods is symmetric about the center of the panel. This implies that each of them has local truncation error containing only odd powers of h and global error containing only even powers of h .

The leading power of h in the error expansion is the order of accuracy. It can be determined by a simple observation: the order of the local truncation error is one more than the degree of the lowest monomial that is not integrated exactly by the panel method. For example, the rectangle rule integrates $f(x) = x^0 \equiv 1$ exactly but gets $f(x) = x^1 \equiv x$ wrong. The order of the lowest monomial not integrated exactly is 1 so the local truncation error is $O(h^2)$ and the global error is $O(h)$. The midpoint rule integrates x^0 and x^1 correctly but gets x^2 wrong. The order of the lowest monomial not integrated exactly is 2 so the local truncation error is $O(h^3)$ and the global error is $O(h^2)$. If the generic panel has x_* in the center, then

$$\int_P (x - x_*)^n dx$$

is always done exactly if n is odd. This is because both the exact integral and its panel method approximation are zero by symmetry.

To understand why this rule works, think of the Taylor expansion of $f(x)$ about the midpoint, x_* . This approximates f by a sum of monomials. Applying the panel integral approximation to f is the same as applying the approximation to each monomial and summing the results. Moreover, the integral of a monomial $(x - x_*)^n$ over P is proportional to h^{n+1} , as is the panel method approximation to it, regardless of whether the panel method is exact or not. The first monomial that is not integrated exactly contributes something proportional to h^{n+1} to the error.

Using this rule it is easy to determine the accuracy of the approximations in Figure 3.6. The trapezoid rule integrates constants and linear functions exactly, but it gets quadratics wrong. This makes the local truncation error third order and the global error second order. The Simpson's rule coefficients $1/6$ and $2/3$ are designed exactly to integrate constants and quadratics exactly, which they do. Simpson's rule integrates cubics exactly (by symmetry) but gets quartics wrong. This gives Simpson's rule fourth order global accuracy. The two point Gauss quadrature also does constants and quadratics correctly but quartics wrong (check this!). The three point Gauss quadrature rule does constants, quadratics, and quartics correctly but gets $(x - x_*)^6$ wrong. That makes it sixth order accurate.

3.5 The method of undetermined coefficients

The method of undetermined coefficients is a general way to find an approximation formula of a desired type. Suppose we want to estimate some A in terms of given data $g_1(h)$, $g_2(h)$, \dots . The method is to assume a linear estimation formula of the form

$$\widehat{A}(h) = a_1(h)g_1(h) + a_2(h)g_2(h) + \dots, \quad (3.37)$$

then determine the unknown coefficients $a_k(h)$ by matching Taylor series up to the highest possible order. The coefficients often take the form of a constant times some power of h : $a_k(h) = a_k h^{p_k}$. The algebra is simpler if we guess or figure out the powers first. The estimator is *consistent* if $\widehat{A}(h) - A \rightarrow 0$ as $h \rightarrow \infty$. Generally (but not always), being consistent is the same as being at least first order accurate. At the end of our calculations, we may discover that there is no consistent estimator of the desired type.

We illustrate the method in a simple example: estimate $f'(x)$ from $g_1 = f(x)$ and $g_2 = f(x+h)$. As above, we will leave out the argument x whenever possible and write f for $f(x)$, f' for $f'(x)$, etc. The estimator is (dropping the x argument)

$$f' \approx \widehat{A} = a_1(h)f + a_2(h)f(x+h).$$

Now expand in Taylor series:

$$f(x+h) = f + f'h + \frac{1}{2}f''h^2 + \dots$$

The estimator is

$$\widehat{A} = a_1(h)f + a_2(h)f + a_2(h)f'h + a_2(h)f''h^2 + \dots \quad (3.38)$$

Looking at the right hand side, we see various coefficients, f , f' , and so on. Since the relation is supposed to work whatever the values of f , f' , etc. may be, we choose a_1 and a_2 so that the coefficient of f is zero. From (3.38), this leads to

$$0 = a_1(h) + a_2(h).$$

To make the estimator consistent, we try

$$1 = a_2(h)h.$$

These two conditions lead to

$$a_2 = \frac{1}{h}, \quad a_1 = \frac{-1}{h}, \quad (3.39)$$

so the estimate is

$$\begin{aligned} f'(x) \approx \widehat{A} &= \frac{-1}{h}f(x) + \frac{1}{h}f(x+h) \\ &= \frac{f(x+h) - f(x)}{h}. \end{aligned}$$

This is the first order one sided difference approximation we saw earlier. Plugging the values (3.39) into (3.38) shows that the estimator satisfies $\hat{A} = f' + O(h)$, which is the first order accuracy we found before.

A more complicated problem is to estimate $f'(x)$ from $f(x)$, $f(x-h)$, $f(x+h)$, $f(x+2h)$. This is not centered nor is it completely one sided, but it is biased to one side. It has proven useful in high accuracy wave simulations. This time we guess that all the coefficients have a power $1/h$, as all the estimates of f' so far have this property. Thus assume the form:

$$f' \approx \hat{A} = \frac{1}{h} (a_{-1}f(x-h) + a_0f + a_1f(x+h) + a_2f(x+2h)) .$$

The Taylor series expansions are

$$\begin{aligned} f(x-h) &= f - f'h + \frac{f''}{2}h^2 - \frac{f'''}{6}h^3 + \frac{f^{(4)}}{24}h^4 + \dots \\ f(x+h) &= f + f'h + \frac{f''}{2}h^2 + \frac{f'''}{6}h^3 + \frac{f^{(4)}}{24}h^4 + \dots \\ f(x+2h) &= f + 2f'h + 2f''h^2 + \frac{4f'''}{3}h^3 + \frac{2f^{(4)}}{3}h^4 + \dots \end{aligned}$$

Equating powers of h turns out to be the same as equating the coefficients of f , f' , etc. from both sides:

$$\begin{aligned} f, \quad O(h^{-1}) : \quad & 0 = a_{-1} + a_0 + a_1 + a_2 \\ f', \quad O(h^0) : \quad & 1 = -a_{-1} + a_1 + 2a_2 \\ f'', \quad O(h^1) : \quad & 0 = \frac{1}{2}a_{-1} + \frac{1}{2}a_1 + 2a_2 \\ f''', \quad O(h^2) : \quad & 0 = \frac{-1}{6}a_{-1} + \frac{1}{6}a_1 + \frac{4}{3}a_2 \end{aligned} \tag{3.40}$$

We could compute the $O(h^3)$ equation but already we have four equations for the four unknown coefficients. If we would use the $O(h^3)$ equation in place of the $O(h^2)$ equation, we loose an order of accuracy in the resulting approximation.

These are a system of four linear equations in the four unknowns a_{-1} through a_2 , which we solve in an ad hoc way. Notice that the combination $b = -a_{-1} + a_1$ appears in the second and fourth equations. If we substitute b , these equations are

$$\begin{aligned} 1 &= b + 2a_2, \\ 0 &= \frac{1}{6}b + \frac{4}{3}a_2. \end{aligned}$$

which implies that $b = -8a_2$ and then that $a_2 = -\frac{1}{6}$ and $b = \frac{4}{3}$. Then, since $-4a_2 = \frac{2}{3}$, the third equation gives $a_{-1} + a_1 = \frac{2}{3}$. Since $b = \frac{4}{3}$ is known, we get two equations for a_{-1} and a_1 :

$$\begin{aligned} a_1 - a_{-1} &= \frac{4}{3}, \\ a_1 + a_{-1} &= \frac{2}{3}. \end{aligned}$$

The solution is $a_1 = 1$ and $a_{-1} = -\frac{1}{3}$. With these, the first equation leads to $a_0 = \frac{-1}{2}$. Finally, our approximation is

$$f'(x) = \frac{1}{h} \left(\frac{-1}{3}f(x-h) - \frac{1}{2}f(x) + f(x+h) - \frac{1}{6}f(x+2h) \right) + O(h^3) .$$

Note that the first step in this derivation was to approximate f by its Taylor approximation of order 3, which would be exact if f were a polynomial of order 3. The derivation has the effect of making \hat{A} exact on polynomials of degree 3 or less. The four equations (3.40) arise from asking \hat{A} to be exact on constants, linear functions, quadratics, and cubics. We illustrate this approach with the problem of estimating $f''(x)$ as accurately as possible from $f(x)$, $f(x+h)$, $f'(x)$ and $f'(x+h)$. The estimator we seek has the form

$$f'' \approx \hat{A} = af + bf(x+h) + cf' + df'(x+h) .$$

We can determine the four unknown coefficients a , b , c , and d by requiring the approximation to be exact on constants, linears, quadratics, and cubics. It does not matter what x value we use, so let us take $x = 0$. This gives, respectively, the four equations:

$$\begin{aligned} 0 &= a + b && (\text{constants, } f = 1) , \\ 0 &= bh + c + d && (\text{linears, } f = x) , \\ 1 &= b\frac{h^2}{2} + dh && (\text{quadratics, } f = x^2/2) , \\ 0 &= b\frac{h^3}{6} + d\frac{h^2}{2} && (\text{cubics, } f = x^3/6) . \end{aligned}$$

Solving these gives

$$a = \frac{-6}{h^2} , \quad b = \frac{6}{h^2} , \quad c = \frac{-4}{h} , \quad d = \frac{-2}{h} .$$

and the approximation

$$f''(x) \approx \frac{6}{h^2}(-f(x) + f(x+h)) - \frac{2}{h}(2f'(x) + f'(x+h)) .$$

A Taylor series calculation shows that this is second order accurate.

3.6 Adaptive parameter estimation

In most real computations, the computational strategy is not fixed in advance, but is adjusted *adaptively* as the computation proceeds. If we are using one of the approximations $\hat{A}(h)$, we might not know an appropriate h when we write the program, and the user might not have the time or expertise to choose h for each application. For example, exercise 12 involves hundreds or thousands of numerical integrations. It is out of the question for the user to experiment manually to find a good h for each one. We need instead a systematic procedure for finding an appropriate step size.

Suppose we are computing something about the function f , a derivative or an integral. We want a program that takes f , and a desired level of accuracy¹²,

¹²This is absolute error. We also could seek a bound on relative error: $|\hat{A} - A| / |A| \leq \epsilon$.

e , and returns \hat{A} with $|\hat{A} - A| \leq e$ with a high degree of confidence. We have $\hat{A}(h)$ that we can evaluate for any h , and we want an automatic way to choose h so that $|\hat{A}(h) - A| \leq e$. A natural suggestion would be to keep reducing h until the answer stops changing. We seek a quantitative version of this.

Asymptotic error expansions of Section 3.3 give one approach. For example, if $\hat{A}(h)$ is a second order accurate approximation to an unknown A and h is small enough we can estimate the error using the leading term:

$$E(h) = \hat{A}(h) - A \approx A_1 h^2 .$$

We can estimate $A_1 h^2$ from $\hat{A}(h)$ and $\hat{A}(2h)$ using the ideas that give (3.28). The result is the Richardson extrapolation error estimate

$$E(h) \approx A_1 h^2 \approx \frac{1}{3} (\hat{A}(2h) - \hat{A}(h)) . \quad (3.41)$$

The adaptive strategy would be to keep reducing h by a factor of two until the estimated error (3.41) is within the tolerance¹³:

```
double adaptive1(double h,      // Step size
                 double eps)   // Desired error bound
{
    double Ah  = A(h);
    double Ah2 = A(h/2);
    while (fabs(Ah2 - Ah) > 3*eps) {
        h  = h/2;
        Ah = Ah2;
        Ah2 = A(h/2);
    }
    return Ah2;
}
```

A natural strategy might be to stop when $|\hat{A}(2h) - \hat{A}(h)| \leq e$. Our quantitative asymptotic error analysis shows that this strategy is off by a factor of 3. We achieve accuracy roughly e when we stop at $|\hat{A}(2h) - \hat{A}(h)| \leq 3e$. This is because $\hat{A}(h)$ is more accurate than $\hat{A}(2h)$.

We can base reasonably reliable software on refinements of the basic strategy of `adaptive1`. Some drawbacks of `adaptive1` are that

1. It needs an *initial guess*, a starting value of h .
2. It may be an infinite loop.
3. It might terminate early if the initial h is outside the *asymptotic range* where error expansions are accurate.

¹³ In the increment part we need not evaluate $\hat{A}(2h)$ because this is what we called $\hat{A}(h)$ before we replaced h with $h/2$.

4. If $\hat{A}(h)$ does not have an asymptotic error expansion, the program will not detect this.
5. It does not return the best possible estimate of A .

A plausible initial guess, h_0 , will depend on the scales (length or time, etc.) of the problem. For example 10^{-10} meters is natural for a problem in atomic physics but not in airplane design. The programmer or the user should supply h_0 based on understanding of the problem. The programmer can take $h_0 = 1$ if he or she thinks the user will use natural units for the problem (Ångströms for atomic physics, meters for airplanes). It might happen that you need $h = h_0/1000$ to satisfy `adaptive1`, but you should give up if $h = h_0\epsilon_{\text{mach}}$. For integration we need an initial $n = (b - a)/h$. It might be reasonable to take $n_0 = 10$, so that $h_0 = (b - a)/10$.

Point 2 says that we need some criterion for giving up. As discussed more in Section 3.7, we should anticipate the ways our software can fail and report failure. When to give up should depend on the problem. For numerical differentiation, we can stop when roundoff or propagated error from evaluating f (see Chapter 2, Section?) creates an error as big as the answer. For integration limiting the number of refinements to 20, would limit the number of panels to $n_0 \cdot 2^{20} \approx n_0 \cdot 10^6$. The revised program might look like

```
const int HMIN_REACHED = -1;

int adaptive2(double h,          // Step size
              double eps,       // Desired error bound
              double& result)    // Final A estimate (output)
{
    double Ah    = A(h);
    double Ah2   = A(h/2);
    double hmin = 10*DBL_EPSILON*h; // DBL_EPSILON is in cfloat (float.h)
    while (fabs(Ah2 - Ah) > 3*eps) {
        h = h/2;
        if (h <= hmin) {
            result = (4*Ah-Ah2)/3; // Extrapolated best result
            return HMIN_REACHED;  // Return error code;
        }
        Ah = Ah2;
        Ah2 = A(h/2);
    }
    result = (4*Ah2-Ah)/3; // Extrapolated best result
    return 0;
}
```

We cannot have perfect protection from point 3, though *premature termination* is unlikely if h_0 is sensible and e (the desired accuracy) is small enough. A more cautious programmer might do more convergence analysis, for example

asking that the $\hat{A}(4h)$ and $\hat{A}(2h)$ error estimate be roughly 2^p times larger than the $\hat{A}(2h)$ and $\hat{A}(h)$ estimate. There might be irregularities in $f(x)$, possibly jumps in some derivative, that prevent the asymptotic error analysis but do not prevent convergence. It would be worthwhile returning an error flag in this case, as some commercial numerical software packages do.

Part of the risk in point 4 comes from the possibility that $\hat{A}(h)$ converges more slowly than the hoped for order of accuracy suggests. For example if $\hat{A}(h) \approx A + A_1 h^{1/2}$, then the error is three times that suggested by `adaptive1`. The extra convergence analysis suggested above might catch this.

Point 5 is part of a paradox afflicting many error estimation strategies. We estimate the size of $E(h) = \hat{A}(h) - A$ by estimating the value of $E(h)$. This leaves us a choice. We could ignore the error estimate (3.41) and report $\hat{A}(h)$ as the approximate answer, or we could subtract out the estimated error and report the more accurate $\hat{A}(h) - \hat{E}(h)$. This is the same as applying one level of Richardson extrapolation to \hat{A} . The corrected approximation probably is more accurate, but we have no estimate of its error. The only reason to be dissatisfied with this is that we cannot report an answer with error less than e until the error is far less than e .

3.7 Software

Scientific programmers have many opportunities to err. Even in short computations, a sign error in a formula or an error in program logic may lead to an utterly wrong answer. More intricate computations provide even more opportunities for bugs. Programs that are designed without care become more and more difficult to read as they grow in size, until

Things fall apart; the centre cannot hold;
Mere anarchy is loosed upon the world.¹⁴

Scientific programmers can do many things to make codes easier to debug and more reliable. We write *modular* programs composed of short procedures so that we can understand and test our calculations one piece at a time. We design these procedures to be *flexible* so that we can re-use work from one problem in solving another problem, or a problem with different parameters. We program *defensively*, checking for invalid inputs and nonsensical results and reporting errors rather than failing silently. In order to uncover hidden problems, we also test our programs thoroughly, using techniques like *convergence analysis* that we described in Section 3.3.2.

3.7.1 Flexibility and modularity

Suppose you want to compute $I = \int_0^2 f(x) dx$ using a panel method. You could write the following code using the rectangle rule:

¹⁴ From “The Second Coming,” with apologies to William Butler Yeats.

```
double I = 0;
for (int k = 0; k < 100; ++k)
    I += .02*f(.02*k);
```

In this code, the function name, the domain of integration, and the number of panels are all *hard-wired*. If you later decided you needed 300 panels for sufficient accuracy, it would be easy to introduce a bug by changing the counter in line 2 without changing anything in line 3. It would also be easy to introduce a bug by replacing .02 with the integer expression 2/300 (instead of 2.0/300).

A more flexible version would be:

```
double rectangle_rule(double (*f)(double), // The integrand
                     double a, double b,   // Left and right ends
                     unsigned n)           // Number of panels
{
    double sum = 0;
    double h = (b-a)/n;
    for (unsigned k = 0; k < n; ++k) {
        double xk = ( a*(n-k) + b*k )/n;
        sum += f(xk);
    }
    return sum*h;
}
```

In this version of the code, the function, domain of integration, and number of panels are all *parameters*. We could use $n = 100$ or $n = 300$ and still get a correct result. This extra flexibility costs only a few extra lines of code that take just a few seconds to type.

This subroutine also *documents* what we were trying to compute by giving names to what were previously just numeric constants. We can tell immediately that the call `rectangle_rule(f, 0, 2, 100)` should integrate over $[0, 2]$, while we could only tell that in the original version by reading the code carefully and multiplying $.02 \times 100$.

What if we are interested in integrating a function that has parameters in addition to the variable of integration? We could pass the extra parameters via a global variable, but this introduces many opportunities for error. If we need to deal with such integrands, we might make the code even more flexible by adding another argument:

```
// An integrand_t is a type of function that takes a double and a double*
// and returns a double. The second argument is used to pass parameters.
typedef double (*integrand_t)(double, double*);

double rectangle_rule(integrand_t f,          // The integrand
                     double* fparams,        // Integrand parameters
                     double a, double b,     // Left and right ends
                     unsigned n)             // Number of panels
```

```

{
    double sum = 0;
    double h = (b-a)/n;
    for (unsigned k = 0; k < n; ++k) {
        double xk = ( a*(n-k) + b*k )/n;
        sum += f(xk, fparams);
    }
    return sum*h;
}

```

We could make this routine even more flexible by making `fparams` a `void*`, or by writing in a style that used more C++ language features¹⁵.

The `rectangle_rule` function illustrates how designing with subroutines improves code flexibility and documentation. When we call a procedure, we care about the *interface*: what data do we have to pass in, and what results do we get out? But we can use the procedure without knowing all the details of the *implementation*. Because the details of the implementation are hidden, when we call a procedure we can keep our focus on the big picture of our computation.

To design robust software, we need to design a testing plan, and another advantage of well-designed procedures is that they make testing easier. For example, we might write test a generic adaptive Richardson extrapolation routine as in Section 3.6, and separately write and test an integration routine like `rectangle_rule`. If the procedures work correctly separately, they have a good chance of working correctly together.

3.7.2 Error checking and failure reports

Most scientific routines can fail, either because the user provides nonsensical input or because it is too difficult to solve the problem to the requested accuracy. Any module that can fail must have a way to report these failures. The simplest way to handle an error is to report an error message and exit; but while this is fine for some research software, it is generally not appropriate for commercial software. A more sophisticated programmer might write errors to a log file, return a flag that indicates when an error has occurred, or throw an exception.¹⁶

For example, let's consider our `rectangle_rule` procedure. We have implicitly assumed in this case that the input arguments are reasonable: the function pointer cannot be `NULL`, there is at least one panel, and the number of panels is not larger than some given value of `MAX_PANELS`. These are *preconditions* for the

¹⁵ This routine is written in C-style C++, which we will use throughout the book. A “native” C++ solution would probably involve additional language features, such as function templates or virtual methods.

¹⁶ We will avoid discussing C++ exception handling with `try/catch`, except to say that it exists. While the `try/catch` model of exception handling is conceptually simple, and while it's relatively simple in languages like Java and MATLAB that have garbage collection, using exceptions well in C++ involves some points that are beyond the scope of this course. For a thorough discussion of C++ exception handling and its implications, we recommend the discussion in *Effective C++* by Scott Meyers.

routine. We can catch errors in these preconditions by using the C++ `assert` macro defined in `assert.h` or `cassert`:

```
const int MAX_PANELS = 10000000;
typedef double (*integrand_t)(double, double*);

double rectangle_rule(integrand_t f,          // The integrand
                     double* fparams,        // Integrand parameters
                     double a, double b,      // Left and right ends
                     unsigned n)              // Number of panels
{
    assert(f != NULL);
    assert(a <= b);
    assert(n != 0 && n < MAX_PANELS);
    ...
}
```

When an assertion fails, the system prints a diagnostic message that tells the programmer where the problem is (file and line number), then exits. Such a hard failure is probably appropriate if the user passes a `NULL` pointer into our `rectangle_rule` routine, but what if the user just uses too large a value of `n`? In that case, there is a reasonable default behavior that would probably give an adequate answer. For a kinder failure mode, we can return our best estimate and an indication that there might be a problem by using an information flag:

```
const int MAX_PANELS = 10000000;
typedef double (*integrand_t)(double, double*);

/* Integrates f on [a,b] via an n-panel rectangle rule.
 * We assume [a,b] and [b,a] are the same interval (no signed measure).
 * If an error occurs, the output argument info takes a negative value:
 *   info == -1 -- n = 0; computed with n = 1
 *   info == -2 -- n was too large; computed with n = MAX_PANELS
 * If no error occurs, info is set to zero.
 */
double rectangle_rule(integrand_t f,          // The integrand
                     double* fparams,        // Integrand parameters
                     double a, double b,      // Left and right ends
                     unsigned n,              // Number of panels
                     int& info)               // Status code
{
    assert(f != NULL);
    info = 0;                                // 0 means "success"
    if (n == 0) {
        n = 1;
        info = -1;                           // -1 means "n too small"
    } else if (n > MAX_PANELS) {
```

```

        n = MAX_PANELS;
        info = -2;                // -2 means "n too big"
    }
    if (b < a)
        swap(a,b);                // Allow bounds in either order
    ...
}

```

Note that we describe the possible values of the `info` output variable in a comment before the function definition. Like the return value and the arguments, the types of failures a routine can experience are an important part of the interface, and they should be documented accordingly¹⁷.

In addition to checking the validity of the input variables, we also want to be careful about our assumptions involving intermediate calculations. For example, consider the loop

```

while (error > targetError) {
    ... refine the solution ...;
}

```

If `targetError` is too small, this could be an infinite loop. Instead, we should at least put in an iteration counter:

```

const int max_iter = 1000; // Stop runaway loop
int iter = 0;
while (error > targetError) {
    if (++iter > max_iter) {
        ... report error and quit ...;
    }
    ... make the solution more accurate ...;
}

```

3.7.3 Unit testing

A *unit test* is a test to make sure that a particular routine does what it is supposed to do. In general, a unit test suite should include problems that exercise every line of code, including failures from which the routine is supposedly able to recover. As a programmer finds errors during development, it also makes sense to add relevant test cases to the unit tests so that those bugs do not recur. In many cases, it makes sense to design the unit tests *before* actually writing a routine, and to use the unit tests to make sure that various revisions to a routine are correct.

¹⁷ There is an unchecked error in this function: the arguments to the integrand could be invalid. In this C-style code, we might just allow `f` to indicate that it has been called with invalid arguments by returning NaN. If we were using C++ exception handling, we could allow `f` to throw an exception when it was given invalid arguments, which would allow more error information to propagate back to the top-level program without any need to redesign the `rectangle_rule` interface.

In addition, high-quality numerical codes are designed with test problems that probe the accuracy and stability of the computation. These test cases should include some trivial problems that can be solved exactly (except possibly for rounding error) as well as more difficult problems.

3.8 References and further reading

For a review of one variable calculus, I recommend the Schaum outline. The chapter on Taylor series explains the remainder estimate clearly.

There several good old books on classical numerical analysis. Two favorites are *Numerical Methods* by Germund Dahlquist and Åke Björk [2], and *Analysis of Numerical Methods* by Gene Isaacson and Herb Keller [12]. Particularly interesting subjects are the symbolic calculus of finite difference operators and Gaussian quadrature.

There are several applications of convergent Taylor series in scientific computing. One example is the *fast multipole method* of Leslie Greengard and Vladimir Rokhlin.

3.9 Exercises

1. Verify that (3.23) represents the leading error in the approximation (3.22). *Hint*, this does not require multidimensional Taylor series. Why?
2. Use multidimensional Taylor series to show that the *rotated* five point operator

$$\frac{1}{2h^2} (f(x+h, y+h) + f(x+h, y-h) + f(x-h, y+h) + f(x-h, y-h) - 4f)$$

is a consistent approximation to Δf . Use a symmetry argument to show that the approximation is at least second order accurate. Show that the leading error term is

$$\frac{h^2}{12} (\partial_x^4 f + 6\partial_x^2 \partial_y^2 f + \partial_y^4 f) .$$

3. For $a \in [0.5, 2]$, consider the following recurrence defining $x_k(a) = x_k$:

$$x_{k+1} = \frac{a + x_k x_{k-1}}{x_k + x_{k-1}}$$

with $x_0 = x_1 = 1 + (a - 1)/2$. This recurrence converges to a function $f(a)$ which is smooth for $a \in [0.5, 2]$.

- (a) Define δ_k to be the largest value of $|x_{k+1}(a) - x_k(a)|$ over twenty evenly spaced values starting at $a = 0.5$ and ending with $a = 2.0$. Plot the largest value of δ_k versus k on a semi-logarithmic plot from

$k = 1$ up to $k = 10$ and comment on the apparent behavior. About how accurate do you think $x_4(a)$ is as an approximation to $f(x)$? About how accurate do you think $x_6(a)$ is?

- (b) Write a program to compute $\widehat{x'_k}(a; h) = (x_k(a + h) - x_k(a - h))/2h$. For $k = 4$ and $a = 0.5625$, show a log-log plot of $\widehat{x'_k}(a; h) - \widehat{x'_k}(a; h/2)$ for $h = 2^{-1}, 2^{-2}, \dots, 2^{-20}$. For h not too small, the plot should be roughly linear on a log-log scale; what is the slope of the line?
- (c) For $a = 0.5625$, the exact value of $f'(a)$ is $2/3$; on a single log-log plot, show the difference between $\widehat{x'_k}(0.5625; h)$ and $2/3$ for $h = 2^{-1}, 2^{-2}, \dots, 2^{-20}$ for $k = 4, 6$. Comment on the accuracy for $k = 4$ in light of parts (a) and (b).
4. Find a formula that estimates $f''(x)$ using the four values $f(x)$, $f(x + h)$, $f(x + 2h)$, and $f(x + 3h)$ with the highest possible order of accuracy. What is this order of accuracy? For what order polynomials does the formula give the exact answer?
5. Suppose we have panels P_k as in (3.31) and panel averages $F_k = \int_{P_k} f(x)dx / (x_{k+1} - x_k)$.
- (a) What is the order of accuracy of F_k as an estimate of $f((x_k + x_{k+1})/2) = f(x_{k+1/2})$?
- (b) Assuming the panels all have size h , find a higher order accurate estimate of $f(x_{k+1/2})$ using F_k , F_{k-1} , and F_{k+1} .
6. In this exercise, we computationally explore the accuracy of the asymptotic series for the function (3.12).
- (a) Define

$$\begin{aligned} L_n(h) &= \sum_{k=0}^{n-1} \int_{k/n}^{(k+1)/n} \frac{e^{-x/h}}{1-a} dx \\ R_n(h) &= \sum_{k=0}^{n-1} \int_{k/n}^{(k+1)/n} \frac{e^{-x/h}}{1-b} dx \\ f(h) &= \int_0^{0.5} \frac{e^{-x/h}}{1-x} dx \end{aligned}$$

and show that for any $n > 1$, $h > 0$,

$$L_n(h) \leq f(h) \leq R_n(h)$$

- (b) Compute $R_{200}(0.3)$ and $L_{200}(0.3)$, and give a bound on the relative error in approximating $f(0.3)$ by $R_{200}(0.3)$. Let $\hat{f}_k(h)$ be the asymptotic series for $f(h)$ expanded through order k , and use the result above to make a semilogarithmic plot of the (approximate) relative error for $h = 0.3$ and $k = 1$ through 30. For what value of k does $\hat{f}_k(0.3)$ best approximate $f(0.3)$, and what is the relative error?

7. An application requires accurate values of $f(x) = e^x - 1$ for x very close to zero.¹⁸
- (a) Show that the problem of evaluating $f(x)$ is well conditioned for small x .
 - (b) How many digits of accuracy would you expect from the code `f = exp(x) - 1`; for $x \sim 10^{-5}$ and for $x \sim 10^{-10}$ in single and in double precision?
 - (c) Let $f(x) = \sum_{n=1}^{\infty} f_n x^n$ be the Taylor series about $x = 0$. Calculate the first three terms, the terms involving x , x^2 , and x^3 . Let $p(x)$ be the degree three Taylor approximation of $f(x)$ about $x = 0$.
 - (d) Assuming that x_0 is so small that the error is nearly equal to the largest neglected term, estimate $\max |f(x) - p(x)|$ when $|x| \leq x_0$.
 - (e) We will evaluate $f(x)$ using

```
if ( abs(x) > x0 ) f = exp(x) - 1;
else               f = p3(x); // given by part c.
```

What x_0 should we choose to maximize the accuracy of $f(x)$ for $|x| < 1$ assuming double precision arithmetic and that the exponential function is evaluated to full double precision accuracy (exact answer correctly rounded)?

- 8. Suppose that $f(x)$ is a function that is evaluated to full machine precision but that there is ϵ_{mach} rounding error in evaluating $\hat{A} = (f(x+h) - f(x))/h$. What value of h minimizes the total absolute error including both rounding error¹⁹ and truncation error? This will be $h_*(\epsilon_{\text{mach}}) \sim \epsilon_{\text{mach}}^q$. Let $e_*(\epsilon_{\text{mach}})$ be the error in the resulting best estimate of $f'(x)$. Show that $e_* \sim \epsilon_{\text{mach}}^r$ and find r .
- 9. Repeat Exercise 8 with the two point centered difference approximation to $f'(x)$. Show that the best error possible with centered differencing is much better than the best possible with the first order approximation. This is one of the advantages of higher order finite difference approximations.
- 10. Verify that the two point Gauss quadrature formula of Figure 3.6 is exact for monomials of degree less than six. This involves checking the functions $f(x) = 1$, $f(x) = x^2$, and $f(x) = x^4$ because the odd order monomials are exact by symmetry. Check that the three point Gauss quadrature formula is exact for monomials of degree less than 8.
- 11. Find the replacement to adaptive halting criterion (3.41) for a method of order p .

¹⁸ The C standard library function `expm1` evaluates $e^x - 1$ to high relative accuracy even for x close to zero.

¹⁹ Note that both x and $f(x)$ will generally be rounded. Recall from the discussion in Section 2.6 that the subtraction and division in the divided difference formula usually commit negligible, or even zero, additional rounding error.

12. In this exercise, we will write a simple adaptive integration routine in order to illustrate some of the steps that go into creating robust numerical software. We will then use the routine to explore the convergence of an asymptotic series. You will be given much of the software; your job is to fill in the necessary pieces, as indicated below (and by `TODO` comments in the code you are given).

The method we will use is not very sophisticated, and in practice, you would usually use a library routine from `QUADPACK`, or possibly a quadrature routine in `MATLAB`, to compute the integrals in this example. However, you will use the same software design ideas we explore here when you attack more complicated problems on your own.

- (a) Write a procedure based on Simpson's rule to estimate the definite integral

$$\int_a^b f(x) dx$$

using panel integration method with uniformly-sized panels. For your implementation of Simpson's rule, the values $f(x_k)$ should be computed only once. Your procedure should run cleanly against a set of unit tests that ensure that you check the arguments correctly, that you exactly integrate low-order polynomials. Your integration procedures should have the same signature as the last version of `rectangle_rule` from the software section:

```
typedef (*integrand_t)(double x, double* fargs);

double simpson_rule(integrand_t f, double* fargs
                    double a, double b, unsigned n, int& info);
double gauss3_rule(integrand_t f, double* fargs
                   double a, double b, unsigned n, int& info);
```

If there is an error, the `info` flag should indicate which argument was a problem (-1 for the first argument, -1 for the second, etc). On success, `info` should return the number of function evaluations performed.

- (b) For each of the quadrature rules you have programmed, use (3.30) to compute the apparent order of accuracy of

$$I_1(h) \approx \int_0^1 e^x dx$$

and

$$I_2(h) \approx \int_0^1 \sqrt{x} dx.$$

The order of accuracy for I_1 will be computed as part of your standard test suite. You will need to add the computation of the order of

accuracy of I_2 yourself. Do your routines obtain the nominal order of accuracy for I_1 ? For I_2 ?

- (c) Write code that does one step of Richardson extrapolation on Simpson's rule in order. Test your code using the test harness from the first part. This procedure should be sixth-order accurate; why is it sixth order rather than fifth? Repeat the procedure for the three-point Gauss rule: write a code to do one step of extrapolation and test to see that it is eighth order.
- (d) We want to know how the function²⁰

$$f(t) = \int_0^1 \cos(tx^2) dx \quad (3.42)$$

behaves for large t .

Write a procedure based on the three-point Gauss quadrature rule with Richardson estimation to find n that gives $f(t)$ to within a specified tolerance (at least 10^{-9}). The procedure should work by repeatedly doubling n until the estimated error, based on comparing approximations, is less than this tolerance. This routine should be robust enough to quit and report failure if it is unable to achieve the requested accuracy.

- (e) The approximation,

$$f(t) \sim \sqrt{\frac{\pi}{8t}} + \frac{1}{2t} \sin(t) - \frac{1}{4t^2} \cos(t) - \frac{3}{8t^3} \sin(t) + \cdots, \quad (3.43)$$

holds for large t ²¹. Use the procedure developed in the previous part to estimate the error in using one, two, and three terms on the right side of (3.43) for t in the range $1 \leq t \leq 1000$. Show the errors versus t on a log-log plot. In all cases we want to evaluate f so accurately that the error in our f value is much less than the error of the approximation (3.43). Note that even for a fixed level of accuracy, more points are needed for large t . Plot the integrand to see why.

²⁰ The function $f(t)$ is a rescaling of the *Fresnel integral* $C(t)$.

²¹ This asymptotic expansion is described, for example, in Exercise 3.1.4 of Olver's *Asymptotics and Special Functions*.

Chapter 4

Linear Algebra I, Theory and Conditioning

4.1 Introduction

Linear algebra and calculus are the basic tools of quantitative science. The operations of linear algebra include solving systems of equations, finding subspaces, solving least squares problems, factoring matrices, and computing eigenvalues and eigenvectors. There is good, publicly available software to perform these computations, and in most cases this software is faster and more accurate than code you write yourself. Chapter 5 outlines some of the basic algorithms of computational linear algebra. This chapter discusses more basic material.

Conditioning is a major concern in many linear algebra computations. Easily available linear algebra software is *backward stable*, which essentially¹ means that the results are as accurate as the conditioning of the problem allows. Even a backward stable method produces large errors if the condition number is of the order of $1/\epsilon_{\text{mach}}$. For example, if the condition number is 10^{18} , even double precision calculations are likely to yield a completely wrong answer. Unfortunately, such condition numbers occur in problems that are not terribly large or rare.

If a computational method for a well-conditioned problem is unstable (much less accurate than its conditioning allows), it is likely because one of the subproblems is ill-conditioned. For example, the problem of computing the matrix exponential, e^A , may be well-conditioned while the problem of computing the eigenvectors of A is ill-conditioned. A stable algorithm for computing e^A (see Exercise 12) in that case must avoid using the eigenvectors of A .

The condition number of a problem (see Section 2.7) measures how small perturbations in the data affect the answer. This is called *perturbation theory*. Suppose A is a matrix² and $f(A)$ is the solution of a linear algebra problem involving A , such as x that satisfies $Ax = b$, or λ and v that satisfy $Av = \lambda v$. Perturbation theory seeks to estimate $\Delta f = f(A + \Delta A) - f(A)$ when ΔA is small. Usually, this amounts to calculating the derivative of f with respect to A .

We simplify the results of perturbation calculations using bounds that involve vector or matrix *norms*. For example, suppose we want to say that all the entries in ΔA or Δv are small. For a vector, v , or a matrix, A , the norm, $\|v\|$ or $\|A\|$, is a number that characterizes the size of v or A . Using norms, we can say that the relative size of a perturbation in A is $\|\Delta A\|/\|A\|$.

The condition number of a problem involving A depends on the problem as well as on A . The condition number of $f(A) = A^{-1}b$ (i.e. solving the system of linear equations $Ax = b$) is very different from the problem of finding the eigenvalues of A . There are matrices that have well conditioned eigenvalues but poorly conditioned eigenvectors. What is commonly called “the” condition number of A is the worst case condition number of solving $Ax = b$, taking the worst possible b .

¹The precise definition of backward stability is in Chapter 5.

² This notation replaces our earlier $A(x)$. In linear algebra, A always is a matrix and x never is a matrix.

4.2 Review of linear algebra

This section reviews some linear algebra that we will use later. It is not a substitute for a course on linear algebra. We assume that most of the topics are familiar to the reader. People come to scientific computing with vastly differing perspectives on linear algebra, and will know some of the concepts we describe by different names and notations. This section should give everyone a common language.

Much of the power of linear algebra comes from this interaction between the abstract and the concrete. Our review connects the abstract language of vector spaces and linear transformations to the concrete language of matrix algebra. There may be more than one concrete representation corresponding to any abstract linear algebra problem. We will find that different representations often lead to numerical methods with very different properties.

4.2.1 Vector spaces

A *vector space* is a set of elements that may be added and multiplied by *scalars*³ (either real or complex numbers, depending on the application). Vector addition is commutative ($u + v = v + u$) and associative ($((u + v) + w = u + (v + w))$). Multiplication by scalars is distributive over vector addition ($a(u + v) = au + av$ and $(a + b)u = au + bu$ for scalars a and b and vectors u and v). There is a unique zero vector, 0 , with $0 + u = u$ for any vector u .

The standard vector spaces are \mathbb{R}^n (or \mathbb{C}^n), consisting of column vectors

$$u = \begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_n \end{pmatrix},$$

where the *components*, u_k , are arbitrary real (or complex) numbers. Vector addition and scalar multiplication are done componentwise.

A subset V' of a vector space V is a *subspace* of V if sums and scalar multiples of elements in V' remain in V' . That is, V' is *closed* under vector addition and scalar multiplication. This means V' is also a vector space under the vector addition and scalar multiplication operations of V . For example, suppose $V = \mathbb{R}^n$ and V' consists of all vectors whose components sum to zero ($\sum_{k=1}^n u_k = 0$). If we add two such vectors or multiply by a scalar, the result also has the zero sum property. On the other hand, the set of vectors whose components sum to one ($\sum_{k=1}^n u_k = 1$) is not closed under vector addition or scalar multiplication.

³ Physicists use the word “scalar” in a different way. For them, a scalar is a number that is the same in any coordinate system. The components of a vector in a particular basis are not scalars in this sense.

The *span* of a set of vectors $\text{span}(f_1, f_2, \dots, f_n) \subset V$ is the subspace of V consisting of *linear combination* of the vectors f_j :

$$u = u_1 f_1 + \dots + u_n f_n, \quad (4.1)$$

where u_k are scalar coefficients. We say f_1, \dots, f_n are *linearly independent* if $u = 0$ implies that $u_k = 0$ for all k in (4.1). Recall that the f_j are linearly independent if and only if the representation (4.1) uniquely determines the *expansion coefficients*, u_k . A theorem of linear algebra states that if the f_j are not linearly independent, then it is possible to find a subset of them with the same span. If $V = \text{span}(f_1, \dots, f_n)$ and the f_j are linearly independent, then the f_j are a *basis* for V .

The standard vector spaces \mathbb{R}^n and \mathbb{C}^n have standard bases $\{e_1, \dots, e_n\}$, where e_k is the vector with all zero components but for a single 1 in position k . This is a basis because

$$u = \begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_n \end{pmatrix} = u_1 \begin{pmatrix} 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix} + u_2 \begin{pmatrix} 0 \\ 1 \\ \vdots \\ 0 \end{pmatrix} + \dots + u_n \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 1 \end{pmatrix} = \sum_{k=1}^n u_k e_k.$$

In view of this, there is little distinction between coordinates, components, and expansion coefficients, all of which are denoted u_k . If V has a basis with n elements, we say the *dimension* of V is n . It is possible to make this definition because of the theorem that states that every basis of V has the same number of elements. A vector space that does not have a finite basis is called *infinite-dimensional*⁴.

An *inner product space* is a vector space that has an *inner product* $\langle \cdot, \cdot \rangle$, which is a scalar function of two vector arguments with the following properties:

1. $\langle u, av_1 + bv_2 \rangle = a\langle u, v_1 \rangle + b\langle u, v_2 \rangle$;
2. $\langle u, v \rangle = \overline{\langle v, u \rangle}$, where \bar{z} refers to the complex conjugate of z ;
3. $\langle u, u \rangle \geq 0$;
4. $\langle u, u \rangle = 0$ if and only if $u = 0$.

When u and v are vectors with $\langle u, v \rangle = 0$, we say u and v are *orthogonal*. If u and v are n component column vectors ($u \in \mathbb{C}^n$, $v \in \mathbb{C}^n$), their *standard inner product* (sometimes called the *dot product*) is

$$\langle u, v \rangle = \sum_{k=1}^n \bar{u}_k v_k. \quad (4.2)$$

The complex conjugates are not needed when the entries of u and v are real.

⁴ An infinite dimensional vector space might have an infinite basis.

Spaces of polynomials are interesting examples of vector spaces. A polynomial in the variable x is a linear combination of powers of x , such as $2 + 3x^4$, or 1, or $\frac{1}{3}(x-1)^2(x^3-3x)^6$. We could multiply out the last example to write it as a linear combination of powers of x . The *degree* of a polynomial is the highest power that it contains. The product $\frac{1}{3}(x-1)^2(x^3-3x)^6$ has degree 20. The vector space P_d is the set of all polynomials of degree at most d . This space has a basis consisting of $d+1$ elements:

$$f_0 = 1, \quad f_1 = x, \quad \dots, \quad f_d = x^d. \quad (4.3)$$

The *power basis* (4.3) is one basis for P_3 (with $d = 3$, so P_3 has dimension 4). Another basis consists of the first four *Hermite* polynomials

$$H_0 = 1, \quad H_1 = x, \quad H_2 = x^2 - 1, \quad H_3 = x^3 - 3x.$$

The Hermite polynomials are orthogonal with respect to a certain inner product⁵:

$$\langle p, q \rangle = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} p(x)q(x)e^{-x^2/2} dx. \quad (4.4)$$

Hermite polynomials are useful in probability because if X is a standard normal random variable, then they are uncorrelated:

$$E[H_j(X)H_k(X)] = \langle H_j, H_k \rangle = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} H_j(x)H_k(x)e^{-x^2/2} dx = 0 \quad \text{if } j \neq k.$$

Still another basis of P_3 consists of Lagrange interpolating polynomials for the points 1, 2, 3, and 4:

$$\begin{aligned} l_1 &= \frac{(x-2)(x-3)(x-4)}{(1-2)(1-3)(1-4)}, & l_2 &= \frac{(x-1)(x-3)(x-4)}{(2-1)(2-3)(2-4)}, \\ l_3 &= \frac{(x-1)(x-2)(x-4)}{(3-1)(3-2)(3-4)}, & l_4 &= \frac{(x-1)(x-2)(x-3)}{(4-1)(4-2)(4-3)}. \end{aligned}$$

These are useful for interpolation because, for example, $l_1(1) = 1$ while $l_2(1) = l_3(1) = l_4(1) = 0$. If we want $u(x)$ to be a polynomial of degree 3 taking specified values $u(1) = u_1$, $u(2) = u_2$, $u(3) = u_3$, and $u(4) = u_4$, the answer is

$$u(x) = u_1 l_1(x) + u_2 l_2(x) + u_3 l_3(x) + u_4 l_4(x).$$

The Lagrange interpolating polynomials are linearly independent because if $0 = u(x) = u_1 l_1(x) + u_2 l_2(x) + u_3 l_3(x) + u_4 l_4(x)$ for all x then in particular $u(x) = 0$ at $x = 1, 2, 3$, and 4 , so $u_1 = u_2 = u_3 = u_4 = 0$. Let $V' \subset P_3$ be the set of polynomials $p \in P_3$ that satisfy $p(2) = 0$ and $p(3) = 0$. This is a subspace of P_3 . A basis for it consists of l_1 and l_4 .

If $V' \subset V$ is a subspace of dimension m of a vector space of dimension n , then it is possible to find a basis of V consisting of vectors f_k so that the first m

⁵The reader can verify that the formula (4.4) defines an inner product on the vector space P_3 .

of the f_k form a basis of V' . For example, if $V = P_3$ and V' is the polynomials that vanish at $x = 2$ and $x = 3$, we can take

$$f_1 = l_1, \quad f_2 = l_4, \quad f_3 = l_2, \quad f_4 = l_3.$$

In general, the dimension of a subspace V' is the dimension of V minus the number of linearly independent conditions that define V' . If there are any nontrivial constraints that define V' , then V' is a *proper* subspace of V ; that is, there is some $u \in V$ that is not in V' , $m < n$. One common task in computational linear algebra is finding a *well-conditioned* basis for a subspace.

4.2.2 Matrices and linear transformations

Suppose V and W are vector spaces. A function L from V to W is *linear* if $L(v_1 + v_2) = L(v_1) + L(v_2)$ for any vectors $v_1, v_2 \in V$, and $L(av) = aL(v)$ for any scalar a and vector $v \in V$. Linear functions are also called *linear transformations*. By convention, we write Lv instead of $L(v)$, even though L represents a function from V to W . This makes algebra with linear transformations look just like matrix algebra, deliberately blurring the distinction between linear transformations and matrices. The simplest example is $V = \mathbb{R}^n$, $W = \mathbb{R}^m$, and $Lu = A \cdot u$ for some $m \times n$ matrix A . The notation $A \cdot u$ refers to the product of the matrix A and the vector u . Most of the time we just write Au .

Any linear transformation between finite dimensional vector spaces may be represented by a matrix. Suppose f_1, \dots, f_n is a basis for V , and g_1, \dots, g_m is a basis for W . For each k , the linear transformation of f_k is an element of W and may be written as a linear combination of the g_j :

$$Lf_k = \sum_{j=1}^m a_{jk} g_j.$$

Because the transformation is linear, we can calculate what happens to a vector $u \in V$ in terms of its expansion $u = \sum_k u_k f_k$. Let $w \in W$ be the *image* of u , $w = Lu$, written as $w = \sum_j w_j g_j$. We find

$$w_j = \sum_{k=1}^n a_{jk} u_k,$$

which is ordinary matrix-vector multiplication.

The matrix that represents L depends on the basis. For example, suppose $V = P_3$, $W = P_2$, and L represents differentiation:

$$L(p_0 + p_1x + p_2x^2 + p_3x^3) = \frac{d}{dx}(p_0 + p_1x + p_2x^2 + p_3x^3) = p_1 + 2p_2x + 3p_3x^2.$$

If we take the basis $1, x, x^2, x^3$ for V , and $1, x, x^2$ for W , then the matrix is

$$\begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 3 \end{pmatrix}.$$

The matrix would be different if we used the Hermite polynomial basis for V (see Exercise 1).

Conversely, an $m \times n$ matrix, A , represents a linear transformation from \mathbb{R}^n to \mathbb{R}^m (or from \mathbb{C}^n to \mathbb{C}^m). We denote this transformation also by A . If $v \in \mathbb{R}^n$ is an n -component column vector, then the matrix-vector product $w = Av$ is a column vector with m components. As before, the notation deliberately is ambiguous. The matrix A is the matrix that represents the linear transformation A using standard bases of \mathbb{R}^n and \mathbb{R}^m .

A matrix also may represent a change of basis within the same space V . If f_1, \dots, f_n , and g_1, \dots, g_n are different bases of V , and u is a vector with expansions $u = \sum_k v_k f_k$ and $u = \sum_j w_j g_j$, then we may write

$$\begin{pmatrix} v_1 \\ \cdot \\ v_n \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & \cdot & a_{1n} \\ a_{21} & a_{22} & \cdot & a_{2n} \\ \cdot & \cdot & \cdot & \cdot \\ a_{n1} & a_{n2} & \cdot & a_{nn} \end{pmatrix} \begin{pmatrix} w_1 \\ \cdot \\ w_n \end{pmatrix}.$$

As before, the matrix elements a_{jk} are the expansion coefficients of g_j with respect to the f_k basis⁶. For example, suppose $u \in P_3$ is given in terms of Hermite polynomials or simple powers: $u = \sum_{j=0}^3 v_j H_j(x) = \sum_{k=0}^3 w_k x^k$, then

$$\begin{pmatrix} v_0 \\ v_1 \\ v_2 \\ v_3 \end{pmatrix} = \begin{pmatrix} 1 & 0 & -1 & 0 \\ 0 & 1 & 0 & -3 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} w_0 \\ w_1 \\ w_2 \\ w_3 \end{pmatrix}.$$

We may reverse the change of basis by using the inverse matrix:

$$\begin{pmatrix} w_1 \\ \cdot \\ w_n \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & \cdot & a_{1n} \\ a_{21} & a_{22} & \cdot & a_{2n} \\ \cdot & \cdot & \cdot & \cdot \\ a_{n1} & a_{n2} & \cdot & a_{nn} \end{pmatrix}^{-1} \begin{pmatrix} v_1 \\ \cdot \\ v_n \end{pmatrix}.$$

Two bases must have the same number of elements because only square matrices can be invertible.

Composition of linear transformations corresponds to matrix multiplication. If L is a linear transformation from V to W , and M is a linear transformation from W to a third vector space, Z , then ML is the *composite* transformation that takes V to Z . The composite of L and M is defined if the *target* (or *range*) of L , is the same as the *source* (or *domain*) of M , W in this case. If A is an $m \times n$ matrix and B is $p \times q$, then the target of A is $W = \mathbb{R}^m$ and the source of B is \mathbb{R}^q . Therefore, the composite AB is defined if $n = p$. This is the condition for the matrix product $A \cdot B$ (usually written without the dot) to be defined. The result is a transformation from $V = \mathbb{R}^q$ to $Z = \mathbb{R}^m$, i.e., the $m \times q$ matrix AB .

⁶We write a_{jk} for the (j, k) entry of A .

For vector spaces V and W , the set of linear transformations from V to W forms a vector space. We can add two linear transformations and multiply a linear transformation by a scalar. This applies in particular to $m \times n$ matrices, which represent linear transformations from \mathbb{R}^n to \mathbb{R}^m . The entries of $A + B$ are $a_{jk} + b_{jk}$. An $n \times 1$ matrix has a single column and may be thought of as a column vector. The product Au is the same whether we consider u to be a column vector or an $n \times 1$ matrix. A $1 \times n$ matrix has a single row and may be thought of as a *row vector*. If r is such a row vector, the product rA makes sense, but Ar does not. It is useful to distinguish between row and column vectors although both are determined by n components. The product ru is a 1×1 matrix $((1 \times n) \cdot (n \times 1))$ gives 1×1 , i.e., a scalar.

If the source and targets are the same, $V = W$, or $n = m = p = q$, then both composites LM and ML both are defined, but they probably are not equal. Similarly, if A and B are $n \times n$ matrices, then AB and BA are defined, but $AB \neq BA$ in general. That is, composition and matrix multiplication are *noncommutative*. If A , B , and C are matrices so that the products AB and BC both are defined, then the products $A \cdot (BC)$ and $(AB) \cdot C$ also are defined. The *associative* property of matrix multiplication is the fact that these are equal: $A(BC) = (AB)C$. In practice, there may be good reasons for computing BC then multiplying by A , rather than finding AB and multiplying by C .

If A is an $m \times n$ matrix with real entries a_{jk} , the *transpose* of A , written A^T , is the $n \times m$ matrix whose (j, k) entry is $(a^T)_{jk} = a_{kj}$. If A has complex entries, then A^* , the *adjoint* of A , is the $n \times m$ matrix with entries $(a^*)_{jk} = \bar{a}_{kj}$ (\bar{a} is the complex conjugate of a). If A is real, then $A^T = A^*$. The transpose or adjoint of a column vector is a row vector with the same number of entries, and vice versa. A square matrix ($m = n$) is *symmetric* if $A = A^T$, and *self-adjoint* (or *Hermitian*) if $A = A^*$. In the case of real matrices, symmetry and self-adjointness are the same.

For a linear operator L on an inner product space, the adjoint L^* is an operator that satisfies $\langle L^*u, v \rangle = \langle u, Lv \rangle$ for all u and v . The reader should check that the matrix adjoint described above satisfies the definition of an operator when the matrix A is interpreted as an operator from \mathbb{R}^m to \mathbb{R}^n with the standard inner product.

4.2.3 Vector norms

The *norm* of a vector u , written $\|u\|$, is a single number that describes the magnitude of u . There are different ways to measure overall size and therefore different vector norms. We say $\|u\|$ (technically, a real number associated to the vector u), is a norm if it satisfies the axioms:

1. *Positive definiteness*: $\|u\| \geq 0$ with $\|u\| = 0$ only when $u = 0$;
2. *Homogeneity*: $\|au\| = |a|\|u\|$ for any scalar a ;
3. *Triangle inequality*: $\|u + v\| \leq \|u\| + \|v\|$ for any vectors u and v .

There are several simple norms for \mathbb{R}^n or \mathbb{C}^n that are useful in scientific computing. One is the l^1 norm

$$\|u\|_1 = \|u\|_{l^1} = \sum_{k=1}^n |u_k| .$$

Another is the l^∞ norm, also called the *sup* norm or the *max* norm⁷:

$$\|u\|_\infty = \|u\|_{l^\infty} = \max_{k=1,\dots,n} |u_k| .$$

Another is the l^2 norm, also called the *Euclidean* norm

$$\|u\|_2 = \|u\|_{l^2} = \left(\sum_{k=1}^n |u_k|^2 \right)^{1/2} = \langle u, u \rangle^{1/2} .$$

The l^2 norm is natural, for example, for vectors representing positions or velocities in three dimensional space. If the components of $u \in \mathbb{R}^n$ represent probabilities, the l^1 norm might be more appropriate. In some cases we may have a norm defined indirectly or with a definition that is hard to turn into a number. For example in the vector space P_3 of polynomials of degree 3, we can define a norm

$$\|p\| = \max_{a \leq x \leq b} |p(x)| . \quad (4.5)$$

There is no simple formula for $\|p\|$ in terms of the coefficients of p .

An appropriate choice of norms is not always obvious. For example, what norm should we use for the two-dimensional subspace of P_3 consisting of polynomials that vanish at $x = 2$ and $x = 3$? In other cases, we might be concerned with vectors whose components have very different magnitudes, perhaps because they are associated with measurements in different units. This might happen, for example, if the components of u represent different factors (or variables) in a linear regression. The first factor, u_1 , might be age of a person, the second, u_2 , income, the third the number of children. In units of years and dollars, we might get

$$u = \begin{pmatrix} 45 \\ 50000 \\ 2 \end{pmatrix} . \quad (4.6)$$

However, most people would consider a five dollar difference in annual salary to be small, while a five-child difference in family size is significant. In situations like these we can define for example, a *dimensionless* version of the l^1 norm:

$$\|u\| = \sum_{k=1}^n \frac{1}{s_k} \cdot |u_k| ,$$

⁷The name l^∞ comes from a generalization of the l^2 norm below. The l^p norm is $\|u\|_p = (\sum \|u_k\|^2)^{1/p}$. One can show that $\|u\|_p \rightarrow \|u\|_\infty$ as $p \rightarrow \infty$.

where the *scale factor* s_k^{-1} is a typical value of a quantity with the units of u_k in the problem at hand. In the example above, we might use $s_1 = 40$ (years), $s_2 = 60000$ (dollars per year), and $s_3 = 2.3$ (children). This is equivalent to using the l^1 norm for the problem expressed in a different basis, $\{s_k^{-1}e_k\}$. In many computations, it makes sense to change to an appropriately scaled basis before turning to the computer.

4.2.4 Norms of matrices and linear transformations

Suppose L is a linear transformation from V to W . If we have norms for the spaces V and W , we can define a corresponding norm of L , written $\|L\|$, as the largest amount by which it stretches a vector:

$$\|L\| = \max_{u \neq 0} \frac{\|Lu\|}{\|u\|} . \quad (4.7)$$

The norm definition (4.7) implies that for all u ,

$$\|Lu\| \leq \|L\| \cdot \|u\| . \quad (4.8)$$

Moreover, $\|L\|$ is the *sharp constant* in the inequality (4.8) in the sense that if $\|Lu\| \leq C \cdot \|u\|$ for all u , then $C \geq \|L\|$. Thus, (4.7) is equivalent to saying that $\|L\|$ is the sharp constant in (4.8).

The different vector norms give rise to different matrix norms. The matrix norms corresponding to certain standard vector norms are written with corresponding subscripts, such as

$$\|L\|_{l^2} = \max_{u \neq 0} \frac{\|Lu\|_{l^2}}{\|u\|_{l^2}} . \quad (4.9)$$

For $V = W = \mathbb{R}^n$, it turns out that (for the linear transformation represented in the standard basis by A)

$$\|A\|_{l^1} = \max_k \sum_j |a_{jk}| ,$$

and

$$\|A\|_{l^\infty} = \max_j \sum_k |a_{jk}| .$$

Thus, the l^1 matrix norm is the maximum *column sum* while the max norm is the maximum *row sum*. Other norms, such as the l^2 matrix norm, are hard to compute explicitly in terms of the entries of A .

Any norm defined by (4.7) in terms of vector norms has several properties derived from corresponding properties of vector norms. One is homogeneity: $\|xL\| = |x| \|L\|$. Another is that $\|L\| \geq 0$ for all L , with $\|L\| = 0$ only for $L = 0$. The triangle inequality for vector norms implies that if L and M are two linear transformations from V to W , then $\|L + M\| \leq \|L\| + \|M\|$. Finally, we have

$\|LM\| \leq \|L\| \|M\|$. This is because the composite transformation stretches no more than the product of the individual maximum stretches:

$$\|M(Lu)\| \leq \|M\| \|Lu\| \leq \|M\| \|L\| \|u\| .$$

Of course, all these properties hold for matrices of the appropriate sizes.

All of these norms have uses in the theoretical parts of scientific computing, the l^1 and l^∞ norms because they are easy to compute and the l^2 norm because it is invariant under orthogonal transformations such as the discrete Fourier transform. The norms are not terribly different from each other. For example, $\|A\|_{l^1} \leq n \|A\|_{l^\infty}$ and $\|A\|_{l^\infty} \leq n \|A\|_{l^1}$. For $n \leq 1000$, this factor of n may not be so important if we are asking, for example, about catastrophic ill-conditioning.

4.2.5 Eigenvalues and eigenvectors

Let A be an $n \times n$ matrix, or a linear transformation from V to V . If

$$Ar = \lambda r .$$

and $r \neq 0$, we say that λ is an *eigenvalue*, and that r is the corresponding *eigenvector*⁸ of A .

Eigenvalues and eigenvectors are useful in understanding dynamics related to A . For example, the differential equation $\frac{du}{dt} = \dot{u} = Au$ has solutions $u(t) = e^{\lambda t}r$. If the differential equation describes something oscillating, A will have at least one complex eigenvalue. In general, eigenvalues and eigenvectors may be complex even though A is real. This is one reason people work with complex vectors in \mathbb{C}^n , even for applications that seem to call for \mathbb{R}^n .

The special case of the *symmetric eigenvalue problem* (A symmetric for real A or self-adjoint for complex A), is vastly different from the general, or *unsymmetric* problem. One of the differences is conditioning. The set of eigenvalues of a self-adjoint matrix is always a well-conditioned function of the matrix – a rare example of uniform good fortune. By contrast, the eigenvalues of an unsymmetric matrix may be so ill-conditioned, even for n as small as 20, that they are not computable in double precision arithmetic. Eigenvalues of unsymmetric matrices are too useful to ignore, but we can get into trouble if we forget their potential ill-conditioning. Eigenvectors, even for self-adjoint matrices, may be ill-conditioned.

We return to the unsymmetric eigenvalue problem. An $n \times n$ matrix may have as many as n eigenvalues, denoted λ_k , $k = 1, \dots, n$. If all the eigenvalues are distinct, the corresponding eigenvectors, denoted r_k , with $Ar_k = \lambda_k r_k$ must be linearly independent, and therefore form a basis for \mathbb{C}^n . These n linearly independent vectors can be assembled to form the columns of an $n \times n$ eigenvector

⁸ Note that “the” eigenvector is at best unique up to scaling: if r is an eigenvector, then so is ar for any scalar a . Those who fail to understand this fact often complain needlessly when a computed eigenvector is scaled differently from the one they had in mind.

matrix that we call the *right eigenvector* matrix.

$$R = \begin{pmatrix} \vdots & & & \vdots \\ r_1 & \cdot & \cdot & r_n \\ \vdots & & & \vdots \end{pmatrix} . \quad (4.10)$$

We also consider the diagonal eigenvalue matrix with the eigenvalues on the diagonal and zeros in all other entries:

$$\Lambda = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix} .$$

The eigenvalue – eigenvector relations may be expressed in terms of these matrices as

$$AR = R\Lambda . \quad (4.11)$$

To see this, check that multiplying R by A is the same as multiplying each of the columns of R by A . Since these are eigenvectors, we get

$$\begin{aligned} A \begin{pmatrix} \vdots & & & \vdots \\ r_1 & \cdot & \cdot & r_n \\ \vdots & & & \vdots \end{pmatrix} &= \begin{pmatrix} \vdots & & & \vdots \\ \lambda_1 r_1 & \cdot & \cdot & \lambda_n r_n \\ \vdots & & & \vdots \end{pmatrix} \\ &= \begin{pmatrix} \vdots & & & \vdots \\ r_1 & \cdot & \cdot & r_n \\ \vdots & & & \vdots \end{pmatrix} \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix} \\ &= R\Lambda . \end{aligned}$$

Since the columns of R are linearly independent, R is invertible, we can multiply (4.11) from the right and from the left by R^{-1} to get

$$R^{-1}ARR^{-1} = R^{-1}R\Lambda R^{-1} ,$$

then cancel the $R^{-1}R$ and RR^{-1} , and define⁹ $L = R^{-1}$ to get

$$LA = \Lambda L .$$

This shows that the k^{th} row of L is an eigenvector of A if we put the A on the right:

$$l_k A = \lambda_k l_k .$$

Of course, the λ_k are the same: there is no difference between “right” and “left” eigenvalues.

⁹ Here L refers to a matrix, not a general linear transformation.

The matrix equation we used to define L , $LR = I$, gives useful relations between left and right eigenvectors. The (j, k) entry of LR is the product of row j of L with column k of R . When $j = k$ this product should be a diagonal entry of I , namely one. When $j \neq k$, the product should be zero. That is

$$\left. \begin{aligned} l_k r_k &= 1 \\ l_j r_k &= 0 \quad \text{if } j \neq k. \end{aligned} \right\} \quad (4.12)$$

These are called *biorthogonality* relations. For example, r_1 need not be orthogonal to r_2 , but it is orthogonal to l_2 . The set of vectors r_k is not orthogonal, but the two sets l_j and r_k are biorthogonal. The left eigenvectors are sometimes called *adjoint eigenvectors* because their transposes form right eigenvectors for the adjoint of A :

$$A^* l_j^* = \lambda_j l_j^*.$$

Still supposing n distinct eigenvalues, we may take the right eigenvectors to be a basis for \mathbb{R}^n (or \mathbb{C}^n if the entries are not real). As discussed in Section 4.2.2, we may express the action of A in this basis. Since $Ar_j = \lambda_j r_j$, the matrix representing the linear transformation A in this basis will be the diagonal matrix Λ . In the framework of Section 4.2.2, this is because if we expand a vector $v \in \mathbb{R}^n$ in the r_k basis, $v = v_1 r_1 + \cdots + v_n r_n$, then $Av = \lambda_1 v_1 r_1 + \cdots + \lambda_n v_n r_n$. For this reason finding a complete set of eigenvectors and eigenvalues is called *diagonalizing* A . A matrix with n linearly independent right eigenvectors is *diagonalizable*.

If A does not have n distinct eigenvalues then there may be no basis in which the action of A is diagonal. For example, consider the matrix

$$J = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}.$$

Clearly, $J \neq 0$ but $J^2 = 0$. A diagonal or diagonalizable matrix cannot have this property: if $\Lambda^2 = 0$ then $\Lambda = 0$, and if the relations $A \neq 0$, $A^2 = 0$ in one basis, they hold in any other basis. In general a *Jordan block* with eigenvalue λ is a $k \times k$ matrix with λ on the diagonal, 1 on the *superdiagonal* and zeros elsewhere:

$$\begin{pmatrix} \lambda & 1 & 0 & \cdots & 0 \\ 0 & \lambda & 1 & 0 & 0 \\ \vdots & 0 & \ddots & \ddots & \vdots \\ \vdots & \vdots & \ddots & \lambda & 1 \\ 0 & 0 & \cdots & 0 & \lambda \end{pmatrix}.$$

If a matrix has fewer than n distinct eigenvalues, it might or might not be diagonalizable. If A is not diagonalizable, a theorem of linear algebra states that there is a basis in which A has *Jordan form*. A matrix has Jordan form if it is *block diagonal* with Jordan blocks of various sizes and eigenvalues on the diagonal. It can be hard to compute the Jordan form of a matrix numerically, as we will see.

If A has a Jordan block, a basis of eigenvectors will not exist; and even if A is diagonalizable, transforming to an eigenvector basis may be very sensitive. For this reason, most software for the unsymmetric eigenvalue problem actually computes a *Schur form*:

$$AQ = QT,$$

where T is an upper triangular matrix with the eigenvalues on the diagonal,

$$T = \begin{pmatrix} \lambda_1 & t_{12} & t_{13} & \dots & t_{1n} \\ 0 & \lambda_2 & t_{23} & \dots & t_{2n} \\ 0 & 0 & \lambda_3 & \dots & t_{3n} \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \dots & 0 & \lambda_n \end{pmatrix},$$

and the columns q_k of Q are *orthonormal*, i.e.

$$q_j^* q_k = \begin{cases} 1, & j = k \\ 0, & j \neq k. \end{cases}$$

Equivalently, we can say that Q is an *orthogonal matrix*, which means that $Q^*Q = I$. In many applications, the Schur form is an acceptable substitute for an eigenvalue – eigenvector decomposition. When eigenvectors are needed, they are computed from the Schur form.

The eigenvalue – eigenvector problem for self-adjoint matrices is different and in many ways simpler than the general nonsymmetric eigenvalue problem. The eigenvalues are real. The left eigenvectors are adjoints of the right eigenvectors: $l_k = r_k^*$. There are no Jordan blocks. Every self-adjoint matrix is diagonalizable even if the number of distinct eigenvalues is less than n . A complete set of eigenvectors of a symmetric matrix forms an orthonormal basis; that is, R is orthogonal. The matrix form of the eigenvalue relation (4.11) may be written $R^*AR = \Lambda$, or $A = R\Lambda R^*$, or $R^*A = \Lambda R^*$. The latter shows (yet again) that the rows of R^* , which are r_k^* , are left eigenvectors of A .

4.2.6 Differentiation and perturbation theory

The main technique in the perturbation theory of Section 4.3 is implicit differentiation. We use the formalism of infinitesimal *virtual perturbations*, which is related to tangent vectors in differential geometry. It may seem roundabout at first, but it makes actual calculations quick.

Suppose $f(x)$ represents m functions, $f_1(x), \dots, f_m(x)$ of n variables, x_1, \dots, x_n . The Jacobian matrix¹⁰, $f'(x)$, is the $m \times n$ matrix of partial derivatives $f'_{jk}(x) = \partial_{x_k} f_j(x)$. If f is differentiable (and f' is Lipschitz continuous), then the first derivative approximation is (writing x_0 for x to clarify some discussion below)

$$f(x_0 + \Delta x) - f(x_0) = \Delta f = f'(x_0)\Delta x + O(\|\Delta x\|^2). \quad (4.13)$$

¹⁰ See any good book on multivariate calculus.

Here Δf and Δx are column vectors.

Suppose s is a scalar “parameter” and $x(s)$ is a differentiable curve in \mathbb{R}^n with $x(0) = x_0$. The function $f(x)$ then defines a curve in \mathbb{R}^m with $f(x(0)) = f(x_0)$. We define two vectors, called *virtual perturbations*,

$$\dot{x} = \left. \frac{d}{ds} \right|_{s=0} x(s) \quad , \quad \dot{f} = \left. \frac{d}{dx} \right|_{s=0} f(x(s)) \quad .$$

The multivariate calculus chain rule implies the virtual perturbation formula

$$\dot{f} = f'(x_0)\dot{x} \quad . \quad (4.14)$$

This formula is nearly the same as (4.13). The virtual perturbation strategy is to calculate the linear relationship (4.14) between virtual perturbations and use it to find the approximate relation (4.13) between actual perturbations. For this, it is important that any $\dot{x} \in \mathbb{R}^n$ can be the virtual perturbation corresponding to some curve: just take the straight “curve” $x(s) = x_0 + s\dot{x}$.

The Leibniz rule (product rule) for matrix multiplication makes virtual perturbations handy in linear algebra. Suppose $A(s)$ and $B(s)$ are differentiable curves of matrices and compatible for matrix multiplication. Then the virtual perturbation of AB is given by the product rule

$$\left. \frac{d}{ds} \right|_{s=0} AB = \dot{A}B + A\dot{B} \quad . \quad (4.15)$$

To see this, note that the (jk) entry of $A(s)B(s)$ is $\sum_l a_{jl}(s)b_{lk}(s)$. Differentiating this using the ordinary product rule then setting $s = 0$ yields

$$\sum_l \dot{a}_{jl}b_{lk} + \sum_l a_{jl}\dot{b}_{lk} \quad .$$

These terms correspond to the terms on the right side of (4.15). We can differentiate products of more than two matrices in this way.

As an example, consider perturbations of the inverse of a matrix, $B = A^{-1}$. The variable x in (4.13) now is the matrix A , and $f(A) = A^{-1}$. Apply implicit differentiation to the formula $AB = I$, use the fact that I is constant, and we get $\dot{A}B + A\dot{B} = 0$. Then solve for \dot{B} and use $A^{-1} = B$, and get

$$\dot{B} = -A^{-1}\dot{A}A^{-1} \quad .$$

The corresponding actual perturbation formula is

$$\Delta(A^{-1}) = -A^{-1}(\Delta A)A^{-1} + O(\|\Delta A\|^2) \quad . \quad (4.16)$$

This is a generalization of the fact that the derivative of $1/x$ is $-1/x^2$, so $\Delta(1/x) \approx (-1/x^2)\Delta x$. When x is replaced by A and ΔA does not commute with A , we have to worry about the order of the factors. The correct order is (4.16), and not, for example, $\dot{A}A^{-1}A^{-1} = \dot{A}A^{-2}$.

For future reference we comment on the case $m = 1$, which is the case of one function of n variables. The $1 \times n$ Jacobian matrix may be thought of as a row vector. We often write this as ∇f , and calculate it from the fact that $\dot{f} = \nabla f(x) \cdot \dot{x}$ for all \dot{x} . In particular, x is a stationary point of f if $\nabla f(x) = 0$, which is the same as $\dot{f} = 0$ for all \dot{x} . For example, suppose $f(x) = x^*Ax$ for some $n \times n$ matrix A . This is a product of the $1 \times n$ matrix x^* with A with the $n \times 1$ matrix x . The Leibniz rule (4.15) gives, if A is constant,

$$\dot{f} = \dot{x}^*Ax + x^*A\dot{x}.$$

Since the 1×1 real matrix \dot{x}^*Ax is equal to its transpose, this is

$$\dot{f} = x^*(A + A^*)\dot{x}.$$

This implies that (both sides are row vectors)

$$\nabla (x^*Ax) = x^*(A + A^*). \quad (4.17)$$

If $A^* = A$, we recognize this as a generalization of the $n = 1$ formula $\frac{d}{dx}(ax^2) = 2ax$.

4.2.7 Variational principles for the symmetric eigenvalue problem

A *variational principle* is a way to find something by solving a maximization or minimization problem. The *Rayleigh quotient* for an $n \times n$ matrix is

$$Q(x) = \frac{x^*Ax}{x^*x} = \frac{\langle x, Ax \rangle}{\langle x, x \rangle}. \quad (4.18)$$

If x is complex, x^* is the adjoint. In either case, the denominator is $x^*x = \sum_{k=1}^n |x_k|^2 = \|x\|_{l^2}^2$. The Rayleigh quotient is defined for $x \neq 0$. A vector r is a *stationary point* if $\nabla Q(r) = 0$. If r is a stationary point, the corresponding value $\lambda = Q(r)$ is a *stationary value*.

Theorem 1 *Suppose A is self-adjoint. A vector $x \neq 0$ is an eigenvector if and only if it is a stationary point of the Rayleigh quotient; and if x is an eigenvector, the Rayleigh quotient is the corresponding eigenvalue.*

Proof: The underlying principle is the calculation (4.17). If $A^* = A$ (this is where symmetry matters) then $\nabla (x^*Ax) = 2x^*A$. Since (4.18) is a quotient, we differentiate using the quotient rule. We already know $\nabla(x^*Ax) = 2x^*A$. Also, $\nabla x^*x = 2x^*$ (the rule with $A = I$). The result is

$$\nabla Q = 2 \left(\frac{1}{x^*x} \right) x^*A - 2 \left(\frac{x^*Ax}{(x^*x)^2} \right) x^* = \frac{2}{x^*x} (x^*A - x^*Q(x)).$$

If x is a stationary point ($\nabla Q = 0$), then $x^*A = \left(\frac{x^*Ax}{x^*x} \right) x^*$, or, taking the adjoint,

$$Ax = \left(\frac{x^*Ax}{x^*x} \right) x.$$

This shows that x is an eigenvector with

$$\lambda = \frac{x^* Ax}{x^* x} = Q(x)$$

as the corresponding eigenvalue. Conversely if $Ar = \lambda r$, then $Q(r) = \lambda$ and the calculations above show that $\nabla Q(r) = 0$. This proves the theorem. \square

A simple observation shows that there is at least one stationary point of Q for Theorem 1 to apply to. If α is a real number, then $Q(\alpha x) = Q(x)$. We may choose α so that¹¹ $\|\alpha x\| = 1$. This shows that

$$\max_{x \neq 0} Q(x) = \max_{\|x\|=1} Q(x) = \max_{\|x\|=1} x^* Ax .$$

A theorem of analysis states that if $Q(x)$ is a continuous function on a compact set, then there is an r so that $Q(r) = \max_x Q(x)$ (the max is attained). The set of x with $\|x\| = 1$ (the *unit sphere*) is compact and Q is continuous. Clearly if $Q(r) = \max_x Q(x)$, then $\nabla Q(r) = 0$, so r is a stationary point of Q and an eigenvector of A .

Now suppose we have found m orthogonal eigenvectors r_1, \dots, r_m . If x is orthogonal to r_j , i.e. $r_j^* x = 0$, then so is Ax :

$$r_j^*(Ax) = (r_j^* A)x = \lambda_j r_j^* x = 0.$$

Therefore, the subspace $V_m \subset \mathbb{C}^n$ of all x that are orthogonal to r_1, \dots, r_m is an *invariant subspace*: if $x \in V_m$, then $Ax \in V_m$. Thus A defines a linear transformation from V_m to V_m , which we call A_m . Chapter 5 gives a proof that A_m is symmetric in a suitable basis. Therefore, Theorem 1 implies that A_m has at least one eigenvector, r_{m+1} , with eigenvalue λ_{m+1} . Since $r_{m+1} \in V_m$, the action of A and A_m on r_{m+1} is the same, which means that $Ar_{m+1} = \lambda_{m+1} r_{m+1}$. After finding r_{m+1} , we can repeat the procedure to find r_{m+2} , and continue until we eventually find a full set of n orthogonal eigenvectors.

4.2.8 Least squares

Suppose A is an $m \times n$ matrix representing a linear transformation from \mathbb{R}^n to \mathbb{R}^m , and we have a vector $b \in \mathbb{R}^m$. If $n < m$ the linear equation system $Ax = b$ is *overdetermined* in the sense that there are more equations than variables to satisfy them. If there is no x with $Ax = b$, we may seek an x that minimizes the *residual*

$$r = Ax - b . \tag{4.19}$$

This *linear least squares* problem

$$\min_x \|Ax - b\|_{l^2} , \tag{4.20}$$

¹¹In this section and the next, $\|x\| = \|x\|_{l^2}$.

is the same as finding x to minimize the *sum of squares*

$$\|r\|_{l^2}^2 = SS = \sum_{j=1}^n r_j^2 .$$

Linear least squares problems arise through *linear regression* in statistics. A *linear regression model* models the response, b , as a linear combination of m *explanatory* vectors, a_k , each weighted by a *regression coefficient*, x_k . The residual, $R = (\sum_{k=1}^m a_k x_k) - b$, is modeled as a Gaussian random variable¹² with mean zero and variance σ^2 . We do n *experiments*. The explanatory variables and response for experiment j are a_{jk} , for $k = 1 \dots, m$, and b_j , and the residual (for given regression coefficients) is $r_j = \sum_{k=1}^m a_{jk} x_k - b_j$. The *log likelihood* function is $f(x) = -\sigma^2 \sum_{j=1}^n r_j^2$. Finding regression coefficients to maximize the log likelihood function leads to (4.20).

The *normal equations* give one approach to least squares problems. We calculate:

$$\begin{aligned} \|r\|_{l^2}^2 &= r^* r \\ &= (Ax - b)^* (Ax - b) \\ &= x^* A^* Ax - 2x^* A^* b + b^* b . \end{aligned}$$

Setting the gradient to zero as in the proof of Theorem 1 leads to the normal equations

$$A^* Ax = A^* b , \quad (4.21)$$

which can be solved by

$$x = (A^* A)^{-1} A^* b . \quad (4.22)$$

The matrix $M = A^* A$ is the *moment matrix* or the *Gram matrix*. It is symmetric, and *positive definite* if A has rank m , so the *Choleski decomposition* of M (see Chapter 5) is a good way to solve (4.21). The matrix $(A^* A)^{-1} A^*$ is the *pseudoinverse* of A . If A were square and invertible, it would be A^{-1} (check this). The normal equation approach is the fastest way to solve dense linear least squares problems, but it is not suitable for the subtle ill-conditioned problems that arise often in practice.

The singular value decomposition in Section 4.2.9 and the *QR* decomposition from Section 5.5 give better ways to solve ill-conditioned linear least squares problems.

4.2.9 Singular values and principal components

Let A be an $m \times n$ matrix that represents a linear transformation from \mathbb{R}^n to \mathbb{R}^m . The *right singular vectors*, $v_k \in \mathbb{R}^n$ form an orthonormal basis for \mathbb{R}^n . The *left*

¹² See any good book on statistics for definitions of Gaussian random variable and the log likelihood function. What is important here is that a systematic statistical procedure, the *maximum likelihood* method, tells us to minimize the sum of squares of residuals.

singular vectors, $u_k \in \mathbb{R}^m$, form an orthonormal basis for \mathbb{R}^m . Corresponding to each v_k and u_k pair is a non-negative *singular value*, σ_k with

$$Av_k = \sigma_k u_k . \quad (4.23)$$

By convention these are ordered so that $\sigma_1 \geq \sigma_2 \geq \cdots \geq 0$. If $n < m$ we interpret (4.23) as saying that $\sigma_k = 0$ for $k > n$. If $n > m$ we say $\sigma_k = 0$ for $k > m$.

The non-zero singular values and corresponding singular vectors may be found one by one using variational and orthogonality principles similar to those in Section 4.2.7. We suppose A is not the zero matrix (not all entries equal to zero). The first step is the variational principle:

$$\sigma_1 = \max_{x \neq 0} \frac{\|Ax\|}{\|x\|} . \quad (4.24)$$

As in Section 4.2.7, the maximum is achieved, and $\sigma_1 > 0$. Let $v_1 \in \mathbb{R}^n$ be a maximizer, normalized to have $\|v_1\| = 1$. Because $\|Av_1\| = \sigma_1$, we may write $Av_1 = \sigma_1 u_1$ with $\|u_1\| = 1$. This is the relation (4.23) with $k = 1$.

The optimality condition calculated as in the proof of Theorem 1 implies that

$$u_1^* A = \sigma_1 v_1^* . \quad (4.25)$$

Indeed, since $\sigma_1 > 0$, (4.24) is equivalent to¹³

$$\begin{aligned} \sigma_1^2 &= \max_{x \neq 0} \frac{\|Ax\|^2}{\|x\|^2} \\ &= \max_{x \neq 0} \frac{(Ax)^*(Ax)}{x^*x} \\ \sigma_1^2 &= \max_{x \neq 0} \frac{x^*(A^*A)x}{x^*x} . \end{aligned} \quad (4.26)$$

Theorem 1 implies that the solution to the maximization problem (4.26), which is v_1 , satisfies $\sigma^2 v_1 = A^* A v_1$. Since $Av_1 = \sigma u_1$, this implies $\sigma_1 v_1 = A^* u_1$, which is the same as (4.25).

The analogue of the eigenvalue orthogonality principle is that if $x^* v_1 = 0$, then $(Ax)^* u_1 = 0$. This is true because

$$(Ax)^* u_1 = x^* (A^* u_1) = x^* \sigma_1 v_1 = 0 .$$

Therefore, if we define $V_1 \subset \mathbb{R}^n$ by $x \in V_1$ if $x^* v_1 = 0$, and $U_1 \subset \mathbb{R}^m$ by $y \in U_1$ if $y^* u_1 = 0$, then A also defines a linear transformation (called A_1) from V_1 to U_1 . If A_1 is not identically zero, we can define

$$\sigma_2 = \max_{\substack{x \in V_1 \\ x \neq 0}} \frac{\|Ax\|^2}{\|x\|^2} = \max_{\substack{x^* v_1 = 0 \\ x \neq 0}} \frac{\|Ax\|^2}{\|x\|^2} ,$$

¹³These calculations make constant use of the associativity of matrix multiplication, even when one of the matrices is a row or column vector.

and get $Av_2 = \sigma_2 u_2$ with $v_2^* v_1 = 0$ and $u_2^* u_1 = 0$. This is the second step in constructing orthonormal bases satisfying (4.23). Continuing in this way, we can continue finding orthonormal vectors v_k and u_k that satisfy (4.23) until we reach $A_k = 0$ or $k = m$ or $k = n$. After that point, we may complete the v and u bases arbitrarily as in Chapter 5 with remaining singular values being zero.

The *singular value decomposition* (SVD) is a matrix expression of the relations (4.23). Let U be the $m \times m$ matrix whose columns are the left singular vectors u_k (as in (4.10)). The orthonormality relations¹⁴ $u_j^* u_k = \delta_{jk}$ are equivalent to U being an orthogonal matrix: $U^* U = I$. Similarly, we can form the orthogonal $n \times n$ matrix, V , whose columns are the right singular vectors v_k . Finally, the $m \times n$ matrix, Σ , has the singular values on its diagonal (with somewhat unfortunate notation), $\sigma_{jj} = \sigma_j$, and zeros off the diagonal, $\sigma_{jk} = 0$ if $j \neq k$. With these definitions, the relations (4.23) are equivalent to $AV = U\Sigma$, which more often is written

$$A = U\Sigma V^* . \quad (4.27)$$

This is the singular value decomposition. Any matrix may be factored, or decomposed, into the product of the form (4.27) where U is an $m \times m$ orthogonal matrix, Σ is an $m \times n$ diagonal matrix with nonnegative entries, and V is an $n \times n$ orthogonal matrix.

A calculation shows that $A^* A = V\Sigma^* \Sigma V^* = V\Lambda V^*$. This implies that the eigenvalues of $A^* A$ are given by $\lambda_j = \sigma_j^2$ and the right singular vectors of A are the eigenvectors of $A^* A$. Section 4.3 explains that the condition number associated with solving linear systems may be taken to be $\kappa_{l^2}(A) = \sigma_1(A)/\sigma_n(A)$. This implies that $\kappa_{l^2}(A^* A) = \kappa_{l^2}(A)^2$. This means that the condition number of solving the normal equations (4.21) is the square of the condition number of the original least squares problem (4.20). If the condition number of a least squares problem is $\kappa_{l^2}(A) = 10^5$, even the best solution algorithm can amplify errors by a factor of 10^5 . Solving using the normal equations can amplify rounding errors by a factor of 10^{10} .

Many people call singular vectors u_k and v_k *principal components*. They refer to the singular value decomposition as *principal component analysis*, or *PCA*. One application is *clustering*, in which you have n objects, each with m measurements, and you want to separate them into two clusters, say “girls” and “boys”. You assemble the data into a matrix, A , and compute, say, the largest two singular values and corresponding left singular vectors, $u_1 \in \mathbb{R}^m$ and $u_2 \in \mathbb{R}^m$. The data for object k is $a_k \in \mathbb{R}^m$, and you compute $z_k \in \mathbb{R}^2$ by $z_{k1} = u_1^* a_k$ and $z_{k2} = u_2^* a_k$, the components of a_k in the principal component directions. You then plot the n points z_k in the plane and look for clusters, or maybe just a line that separates one group of points from another. Surprising as may seem, this simple procedure does identify clusters in practical problems.

¹⁴ Here δ_{jk} is the *Kronecker delta*, which is equal to one when $j = k$ and zero otherwise.

4.3 Condition number

Ill-conditioning can be a serious problem in numerical solution of problems in linear algebra. We take into account possible ill-conditioning when we choose computational strategies. For example, the matrix exponential $\exp(A)$ (see Exercise 12) can be computed using the eigenvectors and eigenvalues of A . We will see in Section 4.3.3 that the eigenvalue problem may be ill conditioned even when the problem of computing $\exp(A)$ is fine. In such cases we need a way to compute $\exp(A)$ that does not use the eigenvectors and eigenvalues of A .

As we said in Section 2.7 (in slightly different notation), the condition number is the ratio of the change in the answer to the change in the problem data, with (i) both changes measured in relative terms, and (ii) the change in the problem data being small. Norms provide a way to extend this definition to functions and data with more than one component. Let $f(x)$ represent m functions of n variables, with Δx being a change in x and $\Delta f = f(x + \Delta x) - f(x)$ the corresponding change in f . The size of Δx , relative to x is $\|\Delta x\| / \|x\|$, and similarly for Δf . In the multivariate case, the size of Δf depends not only on the size of Δx , but also on the direction. The norm-based condition number seeks the worst case Δx , which leads to

$$\kappa(x) = \lim_{\epsilon \rightarrow 0} \max_{\|\Delta x\| = \epsilon} \frac{\|f(x + \Delta x) - f(x)\| / \|f(x)\|}{\|\Delta x\| / \|x\|} . \quad (4.28)$$

Except for the maximization over directions Δx with $\|\Delta x\| = \epsilon$, this is the same as the earlier definition 2.11.

Still following Section 2.7, we express (4.28) in terms of derivatives of f . We let $f'(x)$ represent the $m \times n$ Jacobian matrix of first partial derivatives of f , as in Section 4.2.6, so that, $\Delta f = f'(x)\Delta x + O(\|\Delta x\|^2)$. Since $O(\|\Delta x\|^2) / \|\Delta x\| = O(\|\Delta x\|)$, the ratio in (4.28) may be written

$$\frac{\|\Delta f\|}{\|\Delta x\|} \cdot \frac{\|x\|}{\|f\|} = \frac{\|f'(x)\Delta x\|}{\|\Delta x\|} \cdot \frac{\|x\|}{\|f\|} + O(\|\Delta x\|) .$$

The second term on the right disappears as $\Delta x \rightarrow 0$. Maximizing the first term on the right over Δx yields the norm of the matrix $f'(x)$. Altogether, we have

$$\kappa(x) = \|f'(x)\| \cdot \frac{\|x\|}{\|f(x)\|} . \quad (4.29)$$

This differs from the earlier condition number definition (2.13) in that it uses norms and maximizes over Δx with $\|\Delta x\| = \epsilon$.

In specific calculations we often use a slightly more complicated way of stating the definition (4.28). Suppose that P and Q are two positive quantities and there is a C so that $P \leq C \cdot Q$. We say that C is the *sharp* constant if there is no C' with $C' < C$ so that $P \leq C' \cdot Q$. For example, we have the inequality $\sin(2\epsilon) \leq 3\epsilon$ for all $\epsilon > 0$. But $C = 3$ is not the sharp constant because the inequality also is true with $C' = 2$, which is sharp.

This sharp constant idea is not exactly what we want because it is required to hold for all ϵ (large or small), and because the inequality you might want for small ϵ is not exactly true. For example, there is no inequality $\tan(\epsilon) \leq C\epsilon$ that applies for all $\epsilon > 0$. As $\epsilon \rightarrow 0$, the constant seems to be $C = 1$, but this is not strictly true, since $\tan(\epsilon) > \epsilon$ for $0 < \epsilon < \frac{\pi}{2}$. Therefore we write

$$P(\epsilon) \lesssim CQ(\epsilon) \quad \text{as } \epsilon \rightarrow 0, \quad (4.30)$$

if $P(\epsilon) > 0$, $Q(\epsilon) > 0$, and

$$\lim_{\epsilon \rightarrow 0} \frac{P(\epsilon)}{Q(\epsilon)} \leq C.$$

This is the sharp constant, the smallest C so that

$$P(\epsilon) \leq C \cdot Q(\epsilon) + O(\epsilon) \quad \text{as } \epsilon \rightarrow 0.$$

The definition (4.28) is precisely that $\kappa(x)$ is the sharp constant in the inequality

$$\frac{\|\Delta f\|}{\|f\|} \lesssim \kappa(x) \frac{\|\Delta x\|}{\|x\|} \quad \text{as } \|x\| \rightarrow 0. \quad (4.31)$$

4.3.1 Linear systems, direct estimates

We start with the condition number of calculating $b = Au$ in terms of u with A fixed. This fits into the general framework above, with u playing the role of x , and Au of $f(x)$. Of course, A is the Jacobian of the function $u \rightarrow Au$, so (4.29) gives

$$\kappa(A, u) = \|A\| \cdot \frac{\|u\|}{\|Au\|}. \quad (4.32)$$

This shows that computing Au is ill-conditioned if $\|A\|$ is much larger than the ratio $\|Au\| / \|u\|$. The norm of A is the maximum A can stretch any vector ($\max \|Av\| / \|v\|$). Computing Au is ill-conditioned if it stretches some vector u much more than it stretches u .

The condition number of solving a linear system $Au = b$ (finding u as a function of b) is the same as the condition number of computing $u = A^{-1}b$. The formula (4.32) gives this as

$$\kappa(A^{-1}, b) = \|A^{-1}\| \cdot \frac{\|b\|}{\|A^{-1}b\|} = \|A^{-1}\| \cdot \frac{\|Au\|}{\|u\|}.$$

This is large if there is some vector that is stretched much less than u . Of course, “stretching” factors can be less than one. For future reference, note that $\kappa(A^{-1}, b)$ is not the same as (4.32).

The traditional definition of the condition number of the Au computation takes the worst case relative error not only over perturbations Δu but also over vectors u . Taking the maximum over Δu led to (4.32), so we need only maximize it over u :

$$\kappa(A) = \|A\| \cdot \max_{u \neq 0} \frac{\|u\|}{\|Au\|}. \quad (4.33)$$

Since $A(u + \Delta u) - Au = A\Delta u$, and u and Δu are independent variables, this is the same as

$$\kappa(A) = \max_{u \neq 0} \frac{\|u\|}{\|Au\|} \cdot \max_{\Delta u \neq 0} \frac{\|A\Delta u\|}{\|\Delta u\|} . \quad (4.34)$$

To evaluate the maximum, we suppose A^{-1} exists.¹⁵ Substituting $Au = v$, $u = A^{-1}v$, gives

$$\max_{u \neq 0} \frac{\|u\|}{\|Au\|} = \max_{v \neq 0} \frac{\|A^{-1}v\|}{\|v\|} = \|A^{-1}\| .$$

Thus, (4.33) leads to

$$\kappa(A) = \|A\| \|A^{-1}\| \quad (4.35)$$

as the worst case condition number of the forward problem.

The computation $b = Au$ with

$$A = \begin{pmatrix} 1000 & 0 \\ 0 & 10 \end{pmatrix}$$

illustrates this discussion. The error amplification relates $\|\Delta b\| / \|b\|$ to $\|\Delta u\| / \|u\|$. The worst case would make $\|b\|$ small relative to $\|u\|$ and $\|\Delta b\|$ large relative to $\|\Delta u\|$: amplify u the least and Δu the most. This is achieved by taking $u = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ so that $Au = \begin{pmatrix} 0 \\ 10 \end{pmatrix}$ with amplification factor 10, and $\Delta u = \begin{pmatrix} \epsilon \\ 0 \end{pmatrix}$ with $A\Delta u = \begin{pmatrix} 1000\epsilon \\ 0 \end{pmatrix}$ and amplification factor 1000. This makes $\|\Delta b\| / \|b\|$ 100 times larger than $\|\Delta u\| / \|u\|$. For the condition number of calculating $u = A^{-1}b$, the worst case is $b = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ and $\Delta b = \begin{pmatrix} \epsilon \\ 0 \end{pmatrix}$, which amplifies the error by the same factor of $\kappa(A) = 100$.

The informal condition number (4.33) has advantages and disadvantages over the more formal one (4.32). At the time we design a computational strategy, it may be easier to estimate the informal condition number than the formal one, as we may know more about A than u . If we have no idea what u will come up, we have a reasonable chance of getting something like the worst one. Also, by coincidence, $\kappa(A)$ defined by (4.33) determines the convergence rate of some iterative methods for solving linear systems involving A . On the other hand, in solving differential equations $\kappa(A)$ often is much larger than $\kappa(A, u)$. In such cases, the error in solving $Au = b$ is much less than the condition number $\kappa(A)$ would suggest. For example, in Exercise 11, $\kappa(A)$ is on the order of n^2 , where n is the number of unknowns. The truncation error for the second order discretization is on the order of $1/n^2$. A naive estimate using (4.33) might suggest that solving the system amplifies the $O(n^{-2})$ truncation error by a factor of n^2 to be on the same order as the solution itself. This does not happen because the u we seek is smooth, and not like the worst case.

¹⁵See Exercise 8 for a the l^2 condition number of the $u \rightarrow Au$ problem with singular or non-square A .

The informal condition number (4.35) also has the strange feature than $\kappa(A) = \kappa(A^{-1})$, since $(A^{-1})^{-1} = A$. This gives the mistaken problem that solving a *forward* problem (computing Au from u) has the same stability properties as solving the *inverse* problem, (computing u from $b = Au$). For example, solving the heat equation forward in time is far different from the inverse problem of solving it backward in time.¹⁶ Again, the more precise definition (4.32) does not have this drawback.

4.3.2 Linear systems, perturbation theory

If $Au = b$, we can study the dependence of u on A through perturbation theory. The starting point is the perturbation formula (4.16). Taking norms gives

$$\|\Delta u\| \lesssim \|A^{-1}\| \|\Delta A\| \|u\|, \quad (\text{for small } \Delta A), \quad (4.36)$$

so

$$\frac{\|\Delta u\|}{\|u\|} \lesssim \|A^{-1}\| \|A\| \cdot \frac{\|\Delta A\|}{\|A\|} \quad (4.37)$$

This shows that the condition number satisfies $\kappa \leq \|A^{-1}\| \|A\|$. The condition number is equal to $\|A^{-1}\| \|A\|$ if the inequality (4.36) is (approximately for small ΔA) sharp, which it is because we can take $\Delta A = \epsilon I$ and u to be a maximum stretch vector for A^{-1} . Note that the condition number formula (4.35) applies to the problem of solving the linear system $Au = b$ both when we consider perturbations in b and in A , though the derivations here are different.

4.3.3 Eigenvalues and eigenvectors

The eigenvalue relationship is $Ar_j = \lambda_j r_j$. Perturbation theory allows to estimate the changes in λ_j and r_j that result from a small ΔA . These perturbation results are used throughout science and engineering. The symmetric or self-adjoint case is often called *Rayleigh-Schödinger* perturbation theory.¹⁷ Using the virtual perturbation method of Section 4.2.6, differentiating the eigenvalue relation using the product rule yields

$$\dot{A}r_j + A\dot{r}_j = \dot{\lambda}_j r_j + \lambda_j \dot{r}_j. \quad (4.38)$$

Multiply this from the left by r_j^* and use the fact that r_j^* is a left eigenvector gives

$$r_j^* \dot{A}r_j = \dot{\lambda}_j r_j^* r_j.$$

If r_j is normalized so that $r_j^* r_j = \|r_j\|_{l^2}^2 = 1$, then the right side is just $\dot{\lambda}_j$. Trading virtual perturbations for actual small perturbations turns this into the famous formula

$$\Delta \lambda_j = r_j^* \Delta A r_j + O(\|\Delta A\|^2). \quad (4.39)$$

¹⁶See, for example, the book by Fritz John on partial differential equations.

¹⁷Lord Rayleigh used it to study vibrational frequencies of plates and shells. Later Erwin Schrödinger used it to compute energies (which are eigenvalues) in quantum mechanics.

We get a condition number estimate by recasting this in terms of relative errors on both sides. The important observation is that $\|r_j\|_{l^2} = 1$, so $\|\Delta A \cdot r_j\|_{l^2} \leq \|\Delta A\|_{l^2}$ and finally

$$|r_j^* \Delta A r_j| \lesssim \|\Delta A\|_{l^2} .$$

This inequality is sharp because we can take $\Delta A = \epsilon r_j r_j^*$, which is an $n \times n$ matrix with (see Exercise 7) $\|\epsilon r_j r_j^*\|_{l^2} = |\epsilon|$. Putting this into (4.39) gives the also sharp inequality,

$$\left| \frac{\Delta \lambda_j}{\lambda_j} \right| \leq \frac{\|A\|_{l^2}}{|\lambda_j|} \frac{\|\Delta A\|_{l^2}}{\|A\|_{l^2}} .$$

We can put this directly into the abstract condition number definition (4.28) to get the conditioning of λ_j :

$$\kappa_j(A) = \frac{\|A\|_{l^2}}{|\lambda_j|} = \frac{|\lambda|_{max}}{|\lambda_j|} \quad (4.40)$$

Here, $\kappa_j(A)$ denotes the condition number of the problem of computing λ_j , which is a function of the matrix A , and $\|A\|_{l^2} = |\lambda|_{max}$ refers to the eigenvalue of largest absolute value.

The condition number formula (4.40) predicts the sizes of errors we get in practice. Presumably λ_j depends in some way on all the entries of A and the perturbations due to roundoff will be on the order of the entries themselves, multiplied by the machine precision, ϵ_{mach} , which are on the order of $\|A\|$. Only if λ_j is very close to zero, by comparison with $|\lambda_{max}|$, will it be hard to compute with high relative accuracy. All of the other eigenvalue and eigenvector problems have much worse condition number difficulties.

The eigenvalue problem for non-symmetric matrices can be much more sensitive. To derive the analogue of (4.39) for non-symmetric matrices we start with (4.38) and multiply from the left with the corresponding left eigenvector, l_j , and using the normalization condition $l_j r_j = 1$. After simplifying, the result is

$$\dot{\lambda}_j = l_j \dot{A} r_j , \quad \Delta \lambda_j = l_j \Delta A r_j + O(\|\Delta A\|^2) . \quad (4.41)$$

In the non-symmetric case, the eigenvectors need not be orthogonal and the eigenvector matrix R need not be well conditioned. For this reason, it is possible that l_k , which is a row of R^{-1} is very large. Working from (4.41) as we did for the symmetric case leads to

$$\left| \frac{\Delta \lambda_j}{\lambda_j} \right| \leq \|l_j^*\| \|r_j\| \frac{\|A\|}{|\lambda_j|} \frac{\|\Delta A\|}{\|A\|} .$$

Therefore, the condition number of the non-symmetric eigenvalue problem is (again because the inequalities are sharp)

$$\kappa_j(A) = \|l_j\| \|r_j\| \frac{\|A\|}{|\lambda_j|} . \quad (4.42)$$

A useful upper bound is $\|l_j\|\|r_j\| \leq \kappa_{LS}(R)$, where $\kappa_{LS}(R) = \|R^{-1}\|\|R\|$ is the linear systems condition number of the right eigenvector matrix. Since A is not symmetric, we cannot replace $\|A\|$ by $|\lambda|_{max}$ as we did for (4.40). In the symmetric case, the only reason for ill-conditioning is that we are looking for a (relatively) tiny number. For non-symmetric matrices, it is also possible that the eigenvector matrix is ill-conditioned. It is possible to show that if a family of matrices approaches a matrix with a Jordan block, the condition number of R approaches infinity. For a symmetric or self-adjoint matrix, R is orthogonal or unitary, so that $\|R\|_{l_2} = \|R^*\|_{l_2} = 1$ and $\kappa_{LS}(R) = 1$.

The eigenvector perturbation theory uses the same ideas, with the extra trick of expanding the derivatives of the eigenvectors in terms of the eigenvectors themselves. We expand the virtual perturbation \dot{r}_j in terms of the eigenvectors r_k . Call the expansion coefficients m_{jk} , and we have

$$\dot{r}_j = \sum_{l=1}^n m_{jl} r_l .$$

For the symmetric eigenvalue problem, if all the eigenvalues are distinct, the formula follows from multiplying (4.38) from the left by r_k^* :

$$m_{jk} = \frac{r_k^* \dot{A} r_j}{\lambda_j - \lambda_k} ,$$

so that

$$\Delta r_j = \sum_{k \neq j} \frac{r_k^* \Delta A r_j}{\lambda_j - \lambda_k} + O(\|\Delta A\|^2) .$$

(The term $j = k$ is omitted because $m_{jj} = 0$: differentiating $r_j^* r_j = 1$ gives $r_j^* \dot{r}_j = 0$.) This shows that the eigenvectors have condition number “issues” even when the eigenvalues are well-conditioned, if the eigenvalues are close to each other. Since the eigenvectors are not uniquely determined when eigenvalues are equal, this seems plausible. The unsymmetric matrix perturbation formula is

$$m_{kj} = \frac{l_j \dot{A} r_k}{\lambda_j - \lambda_k} .$$

Again, we have the potential problem of an ill-conditioned eigenvector basis, combined with the possibility of close eigenvalues. The conclusion is that the eigenvector conditioning can be problematic, even though the eigenvalues are fine, for closely spaced eigenvalues.

4.4 Software

4.4.1 Software for numerical linear algebra

There have been major government-funded efforts to produce high quality software for numerical linear algebra. This culminated in the public domain software package *LAPACK*. LAPACK is a combination and extension of earlier

packages *EISPACK*, for solving eigenvalue problems, and *LINPACK*, for solving systems of equations. LAPACK is used in many mathematical packages, including MATLAB.

Our advice is to use LAPACK for dense linear algebra computations whenever possible¹⁸, either directly or through an environment like MATLAB. These routines are extensively tested, and are much less likely to have subtle bugs than codes you developed yourself or copied from a textbook. The LAPACK software also includes condition estimators that are often more sophisticated than the basic algorithms themselves, and includes methods such as *equilibration* to improve the conditioning of problems.

LAPACK is built upon a library of *Basic Linear Algebra Subroutines* (BLAS). A BLAS library includes routines that perform such tasks as computing dot products and multiplying matrices. Different systems often have their own specially-tuned BLAS libraries, and these tuned libraries have much better performance than the reference implementation. Computations using LAPACK with an optimized BLAS can be faster than computations with the reference BLAS by an order of magnitude or more.

On many platforms, LAPACK and a tuned BLAS are packaged as a pre-compiled library. On machines running OS X, for example, LAPACK and a tuned BLAS are provided as part of the vecLib framework. On Windows and Linux machines, optimized LAPACK and BLAS implementations are part of Intel's Math Kernel Libraries (MKL), which are sold commercially but with a free version under Windows, and as part of the AMD Core Math Library (ACML), which is freely available from AMD. The [LAPACK Frequently Asked Questions](http://www.netlib.org/lapack) list from <http://www.netlib.org/lapack> includes as its first item a list of places where you can get LAPACK and the BLAS pre-packaged as part of a vendor library.

LAPACK is written in Fortran 90¹⁹, not C++. It is not too difficult to call Fortran routines from C++, but the details vary from platform to platform. One alternative to mixed-language programming is to use CLAPACK, which is an automatic translation of LAPACK into the C language. CLAPACK provides a Fortran-style interface to LAPACK, but because it is written entirely in C, it is often easier to link. As of this writing, there is a standardized C interface to the BLAS (the CBLAS), but there is no standard C interface to LAPACK. One of the advantages of using LAPACK via a vendor-provided library is that many of these libraries provide a native C-style interface to the LAPACK routines.

4.4.2 Linear algebra in Matlab

The MATLAB system uses LAPACK and BLAS for most of its dense linear algebra operation. This includes the `*` command to multiply matrices, the

¹⁸ LAPACK has routines for *dense* linear algebra. *Sparse* matrices, such as those with only a few nonzero entries, should generally be solved using other packages. We will discuss sparse matrices briefly in Chapter 5

¹⁹ Most of LAPACK was developed in Fortran 77. However, the most recent version (LAPACK 3.2) uses some Fortran 90 features to provide extended precision routines.

```

N = 500;
A = rand(N);
B = rand(N);

tic; C = A*B; t1 = toc;

tic;
C2 = zeros(N);
for i = 1:N
    for j = 1:N
        for k = 1:N
            C2(i,j) = C2(i,j) + A(i,k)*B(k,j);
        end
    end
end
t2 = toc;

fprintf('For N = %d: built-in = %f sec; ours = %f sec\n', ...
        N, t1, t2);

```

Figure 4.1: Example MATLAB program that multiplies two matrices with a built-in primitive based on the BLAS and with a hand-written loop.

backslash command to solve linear systems and least squares problems (\backslash), the `eig` command to compute eigenvalues and eigenvectors, and the `svd` command to compute singular values and vectors.

Where MATLAB provides a built-in routine based on LAPACK and the BLAS, it will usually be *much* more efficient than a corresponding routine that you would write yourself. For example, consider the code shown in Figure 4.1 to multiply two matrices using the built-in `*` operator and a hand-written loop. On a desktop Pentium 4 system running version 7.6 of MATLAB, we can see the difference between the two calls clearly:

```
For N = 500: built-in = 0.061936 sec; ours = 5.681781 sec
```

Using the MATLAB built-in functions that call the BLAS is almost a hundred times faster than writing our own matrix multiply routine! In general, the most efficient MATLAB programs make heavy use of linear algebra primitives that call BLAS and LAPACK routines.

Where LAPACK provides very different routines for different matrix structures (e.g. different algorithms for symmetric and unsymmetric eigenvalue problems), MATLAB tests the matrix structure and automatically chooses an appropriate routine. It's possible to see this indirectly by looking at the timing of

```

A = rand(600);      % Make a medium-sized random matrix
B = (A+A')/2;      % Take the symmetric part of A
Bhat = B+eps*A;    % Make B just slightly nonsymmetric

tic; eig(B); t1 = toc; % Time eig on a symmetric problem
tic; eig(Bhat); t2 = toc; % Time eig on a nonsymmetric problem

fprintf('Symmetric: %f sec; Nonsymmetric: %f sec\n', t1, t2);

```

Figure 4.2: Example MATLAB program that times the computation of eigenvalues for a symmetric matrix and a slightly nonsymmetric matrix.

the following MATLAB script in Figure 4.2. On a desktop Pentium 4 system running version 7.6 of MATLAB, we can see the difference between the two calls clearly:

```
Symmetric: 0.157029 sec; Nonsymmetric: 1.926543 sec
```

4.4.3 Mixing C++ and Fortran

Like an enormous amount of scientific software, LAPACK is written in Fortran. Because of this, it is useful to call Fortran code from C. Unfortunately, many systems do not include a Fortran compiler; and even when there is a Fortran compiler, it may be difficult to figure out which libraries are needed to link together the Fortran and C codes. The `f2c` translator, available from <http://www.netlib.org/f2c>, translates Fortran codes into C. In this section, we will describe how to use `f2c`-translated codes in a C++ program. As an example, we will translate and use the reference version of the BLAS dot product routine `ddot`, which is available at <http://www.netlib.org/blas/ddot.f>. In our description, we assume a command line shell of the sort provided by OS X, Linux, or the Cygwin environment under Windows. We also assume that `f2c` has already been installed somewhere on the system.

The first step is to actually translate the code from Fortran to C:

```
f2c -C++ ddot.f
```

This command generates a function `ddot.c` with the translated routine. The `-C++` option says that the translation should be done to C++-compatible C. In Fortran, `ddot` has the signature

```

      DOUBLE PRECISION FUNCTION DDOT(N,DX,INCX,DY,INCY)
*      .. Scalar Arguments ..
      INTEGER INCX,INCY,N
*      ..

```

```

#include <iostream>
#include "f2c.h"

extern "C"
double ddot_(const integer& N, const double* dx,
             const integer& incX, const double* dy,
             const integer& incY);

int main()
{
    using namespace std;
    double xy[] = {1, 2};
    double result = ddot_(2, xy, 1, xy, 1);
    cout << "Result = " << result << endl; // Result = 5
    return 0;
}

```

Figure 4.3: Example C++ program that calls an f2c-translated routine.

```

*      .. Array Arguments ..
      DOUBLE PRECISION DX(*),DY(*)

```

The corresponding C function `ddot_` has the signature

```

double ddot_(integer* n, double* dx, integer* incx,
             double* dy, integer* incy);

```

We note three things about the signature for this translated routine:

- The C name of the routine is all lower case, with an underscore appended at the end: `ddot_`.
- Fortran uses *call-by-reference* semantics; that is, all arguments are passed by reference (through a pointer in C) rather than by value.
- An `integer` is a typedef defined in the `f2c.h` header file. Depending on the system, a Fortran `integer` may correspond to a C `int` or `long`.

Once we have translated the function to C, we need to be able to use it in another program. We will use as an example the code in Figure 4.3, which computes the dot product of a two-element vector with itself. We note a few things about this program:

- The `extern "C"` directive tells the C++ compiler to use C-style linkage²⁰,

²⁰ C++ linkage involves *name mangling*, which means that the compiler adds type information into its internal representation for a function name. In C-style linkage, we keep only the function name.

which is what `f2c` requires. This is important when using C++ with any other language.

- Rather than using C-style pointers to pass the vector length `N` and the parameters `incX` and `incY`, we use C++-style references²¹. We tell the compiler that the `ddot` routine uses these values purely as input parameters by declaring that these are references to *constant* values. If we do this, the C++ compiler will allow us to pass in literal constants like 2 and 1 when we call `ddot_`, rather than requiring that we write

```
// If we had ddot_(integer* N, ...)
integer two = 2;
integer one = 1;
ddot_(&two, xy, &one, xy, &one);
```

Finally, we compile the program:

```
c++ -o example.exe -I f2cdir example.cc ddot.c -L f2cdir -lf2c
```

Here `-I f2cdir` tells the compiler to search the directory `f2cdir` for the header file `f2c.h`; and the flag `-L f2cdir` tells the compiler where to find the support library `libf2c`.

The steps we used to incorporate the Fortran `ddot` routine in our C++ code using `f2c` are similar to the steps we would use to incorporate a Fortran library into any C++ code:

1. Follow the directions to compile the library (or to translate it using `f2c`, in our case). If you can find a pre-built version of the library, you may be able to skip this step.
2. Write a C++ program that includes a prototype for the library routine. Remember that the C name will probably be slightly different from the Fortran name (a trailing underscore is common), and that all arguments in Fortran are passed by reference.
3. Compile your C++ program and link it against the Fortran library and possibly some Fortran support libraries.

4.5 Resources and further reading

For a practical introduction to linear algebra written by a numerical analyst, try the books by Gilbert Strang [23, 24]. More theoretical treatments may be found in the book by Peter Lax [15] or the one by Paul Halmos [8]. The linear algebra book by Peter Lax also has a beautiful discussion of eigenvalue perturbation

²¹ The C++ compiler would complain at us if it were keeping type information on the arguments to `ddot_`, but those are discarded because the compiler is using C linkage. References and pointers are implemented nearly identically, so we can get away with this in practice.

theory and some of its applications. More applications may be found in the book *Theory of Sound* by Lord Rayleigh (reprinted by Dover Press) and in any book on quantum mechanics [20].

There are several good books that focus exclusively on topics in numerical linear algebra. The textbook by Lloyd N. Trefethen [25] provides a good survey of numerical linear algebra, while the textbook by James Demmel [4] goes into more detail. Both books provide excellent discussion of condition numbers. Nicholas Higham's book *Accuracy and Stability of Numerical Algorithms* [10] includes a much more detailed description of conditioning and related issues. The standard reference book for topics in numerical linear algebra is by Gene Golub and Charles Van Loan [7].

Beresford Parlett has a book on the theory and computational methods for the symmetric eigenvalue problem [19]; though it does not include some of the most recent methods, it remains an excellent and highly readable book. G.W. Stewart has devoted the second volume of his *Matrix Algorithms* book entirely to eigenvalue problems, including both the symmetric and the nonsymmetric case [22].

The software repository *Netlib*, <http://netlib.org>, is a source for LAPACK and for many other numerical codes.

4.6 Exercises

1. Let L be the differentiation operator that takes P_3 to P_2 described in Section 4.2.2. Let $f_k = H_k(x)$ for $k = 0, 1, 2, 3$ be the Hermite polynomial basis of P_3 and $g_k = H_k(x)$ for $k = 0, 1, 2$ be the Hermite basis of P_2 . What is the matrix, A , that represents this L in these bases?
2. Suppose L is a linear transformation from V to V and that f_1, \dots, f_n , and g_1, \dots, g_n are two bases of V . Any $u \in V$ may be written in a unique way as $u = \sum_{k=1}^n v_k f_k$, or as $u = \sum_{k=1}^n w_k g_k$. There is an $n \times n$ matrix, R that relates the f_k expansion coefficients v_k to the g_k coefficients w_k by $v_j = \sum_{k=1}^n r_{jk} w_k$. If v and w are the column vectors with components v_k and w_k respectively, then $v = Rw$. Let A represent L in the f_k basis and B represent L in the g_k basis.
 - (a) Show that $B = R^{-1}AR$.
 - (b) For $V = P_3$, and $f_k = x^k$, and $g_k = H_k$, find R .
 - (c) Let L be the linear transformation $Lp = q$ with $q(x) = \partial_x(xp(x))$. Find the matrix, A , that represents L in the monomial basis f_k .
 - (d) Find the matrix, B , that represents L in the Hermite polynomial basis H_k .
 - (e) Multiply the matrices to check explicitly that $B = R^{-1}AR$ in this case.

3. If A is an $n \times m$ matrix and B is an $m \times l$ matrix, then AB is an $n \times l$ matrix. Show that $(AB)^* = B^*A^*$. Note that the incorrect suggestion A^*B^* in general is not compatible for matrix multiplication.
4. Let $V = \mathbb{R}^n$ and M be an $n \times n$ real matrix. This exercise shows that $\|u\| = (u^*Mu)^{1/2}$ is a vector norm whenever M is *positive definite* (defined below).
 - (a) Show that $u^*Mu = u^*M^*u = u^* \left(\frac{1}{2}(M + M^*) \right) u$ for all $u \in V$. This means that as long as we consider functions of the form $f(u) = u^*Mu$, we may assume M is symmetric. For the rest of this question, assume M is symmetric. Hint: u^*Mu is a 1×1 matrix and therefore equal to its transpose.
 - (b) Show that the function $\|u\| = (u^*Mu)^{1/2}$ is homogeneous: $\|au\| = |a| \|u\|$.
 - (c) We say M is positive definite if $u^*Mu > 0$ whenever $u \neq 0$. Show that if M is positive definite, then $\|u\| \geq 0$ for all u and $\|u\| = 0$ only for $u = 0$.
 - (d) Show that if M is symmetric and positive definite (SPD), then $|u^*Mv| \leq \|u\| \|v\|$. This is the *Cauchy-Schwarz inequality*. Hint (a famous old trick): $\phi(t) = (u + tv)^*M(u + tv)$ is a quadratic function of t that is non-negative for all t if M is positive definite. The Cauchy-Schwarz inequality follows from requiring that the minimum value of ϕ is not negative, assuming $M^* = M$.
 - (e) Use the Cauchy-Schwarz inequality to verify the triangle inequality in its squared form $\|u + v\|^2 \leq \|u\|^2 + 2\|u\| \|v\| + \|v\|^2$.
 - (f) Show that if $M = I$ then $\|u\|$ is the l^2 norm of u .
5. Verify that $\|p\|$ defined by (4.5) on $V = P_3$ is a norm as long as $a < b$.
6. Suppose A is the $n \times n$ matrix that represents a linear transformation from \mathbb{R}^n to \mathbb{R}^n in the standard basis e_k . Let B be the matrix of the same linear transformation in the scaled basis (??).
 - (a) Find a formula for the entries b_{jk} in terms of the a_{jk} and \bar{u}_k .
 - (b) Find a matrix formula for B in terms of A and the *diagonal scaling* matrix $W = \text{diag}(\bar{u}_k)$ (defined by $w_{kk} = \bar{u}_k$, $w_{jk} = 0$ if $j \neq k$) and W^{-1} .
7. Show that if $u \in \mathbb{R}^m$ and $v \in \mathbb{R}^n$ and $A = uv^*$, then $\|A\|_{l^2} = \|u\|_{l^2} \cdot \|v\|_{l^2}$. Hint: Note that $Aw = bu$ where b is a scalar, so $\|Aw\|_{l^2} = |b| \cdot \|u\|_{l^2}$. Also, be aware of the *Cauchy-Schwarz* inequality: $|v^*w| \leq \|v\|_{l^2} \|w\|_{l^2}$.
8. Suppose that A is an $n \times n$ invertible matrix. Show that

$$\|A^{-1}\| = \max_{u \neq 0} \frac{\|u\|}{\|Au\|} = \left(\min_{u \neq 0} \frac{\|Au\|}{\|u\|} \right)^{-1}.$$

9. The *symmetric part* of the real $n \times n$ matrix is $M = \frac{1}{2}(A + A^*)$. Show that $\nabla \left(\frac{1}{2}x^*Ax \right) = Mx$.
10. The *informal* condition number of the problem of computing the action of A is

$$\kappa(A) = \max_{x \neq 0, \Delta x \neq 0} \frac{\frac{\|A(x+\Delta x) - Ax\|}{\|Ax\|}}{\frac{\|x+\Delta x\|}{\|x\|}}.$$

Alternatively, it is the sharp constant in the estimate

$$\frac{\|A(x + \Delta x) - Ax\|}{\|Ax\|} \leq C \cdot \frac{\|\Delta x\|}{\|x\|},$$

which bounds the worst case relative change in the answer in terms of the relative change in the data. Show that for the l^2 norm,

$$\kappa = \sigma_{\max}/\sigma_{\min},$$

the ratio of the largest to smallest singular value of A . Show that this formula holds even when A is not square.

11. We wish to solve the *boundary value problem* for the differential equation

$$\frac{1}{2}\partial_x^2 u = f(x) \quad \text{for } 0 < x < 1, \quad (4.43)$$

with *boundary conditions*

$$u(0) = u(1) = 0. \quad (4.44)$$

We *discretize* the interval $[0, 1]$ using a uniform *grid* of points $x_j = j\Delta x$ with $n\Delta x = 1$. The $n - 1$ unknowns, U_j , are approximations to $u(x_j)$, for $j = 1, \dots, n - 1$. If we use a second order approximation to $\frac{1}{2}\partial_x^2 u$, we get discrete equations

$$\frac{1}{2} \frac{1}{\Delta x^2} (U_{j+1} - 2U_j + U_{j-1}) = f(x_j) = F_j. \quad (4.45)$$

Together with boundary conditions $U_0 = U_n = 0$, this is a system of $n - 1$ linear equations for the vector $U = (U_1, \dots, U_{n-1})^*$ that we write as $AU = F$.

- (a) Check that there are $n - 1$ distinct eigenvectors of A having the form $r_{kj} = \sin(k\pi x_j)$. Here r_{kj} is the j component of eigenvector r_k . Note that $r_{k,j+1} = \sin(k\pi x_{j+1}) = \sin(k\pi(x_j + \Delta x))$, which can be evaluated in terms of r_{kj} using trigonometric identities.
- (b) Use the eigenvalue information from part (a) to show that $\|A^{-1}\| \rightarrow 2/\pi^2$ as $n \rightarrow \infty$ and $\kappa(A) = O(n^2)$ (in the informal sense) as $n \rightarrow \infty$. All norms are l^2 .

- (c) Suppose $\tilde{U}_j = u(x_j)$ where $u(x)$ is the exact but unknown solution of (4.43), (4.44). Show that if $u(x)$ is smooth then the *residual*²², $R = A\tilde{U} - F$, satisfies $\|R\| = O(\Delta x^2) = O(1/n^2)$. For this to be true we have to adjust the definition of $\|U\|$ to be consistent with the L^2 integral $\|u\|_{L^2}^2 = \int_{x=0}^1 u^2(x)dx$. The discrete approximation is $\|U\|_{l^2}^2 = \Delta x \sum_{k=1}^n U_j^2$.
- (d) Show that $A(\tilde{U} - U) = R$. Use part (b) to show that $\|U - \tilde{U}\| = O(\Delta x^2)$ (with the Δx modified $\|\cdot\|$).
- (e) (harder) Create a fourth order five point central difference approximation to $\partial_x^2 u$. You can do this using Richardson extrapolation from the second order three point formula. Use this to find an A so that solving $AU = F$ leads to a fourth order accurate U . The hard part is what to do at $j = 1$ and $j = n - 1$. At $j = 1$ the five point approximation to $\partial_x^2 u$ involves U_0 and U_{-1} . It is fine to take $U_0 = u(0) = 0$. Is it OK to take $U_{-1} = -U_1$?
- (f) Write a program in Matlab to solve $AU = F$ for the second order method. The matrix A is symmetric and *tridiagonal* (has nonzeros only on three *diagonals*, the *main diagonal*, and the immediate sub and super diagonals). Use the Matlab matrix operation appropriate for symmetric positive definite tridiagonal matrices. Do a convergence study to show that the results are second order accurate.
- (g) (extra credit) Program the fourth order method and check that the results are fourth order when $f(x) = \sin(\pi x)$ but not when $f(x) = \max(0, .15 - (x - .5)^2)$. Why are the results different?
12. This exercise explores conditioning of the non-symmetric eigenvalue problem. It shows that although the problem of computing the fundamental solution is well-conditioned, computing it using eigenvalues and eigenvectors can be an unstable algorithm because the problem of computing eigenvalues and eigenvectors is ill-conditioned. For parameters $0 < \lambda < \mu$, there is a *Markov chain transition rate* matrix, A , whose entries are $a_{jk} = 0$ if $|j - k| > 1$. If $1 \leq j \leq n - 2$, $a_{j,j-1} = \mu$, $a_{jj} = -(\lambda + \mu)$, and $a_{j,j+1} = \lambda$ (taking j and k to run from 0 to $n - 1$). The other cases are $a_{00} = -\lambda$, $a_{01} = \lambda$, $a_{n-1,n-1} = -\mu$, and $a_{n-1,n-2} = \mu$. This matrix describes a continuous time Markov process with a random walker whose position at time t is the integer $X(t)$. Transitions $X \rightarrow X + 1$ happen with rate λ and transitions $X \rightarrow X - 1$ have rate μ . The transitions $0 \rightarrow -1$ and $n - 1 \rightarrow n$ are not allowed. This is the *M/M/1 queue* used in operations research to model queues ($X(t)$ is the number of customers in the queue at time t , λ is the rate of arrival of new customers, μ is the service rate. A customer arrival is an $X \rightarrow X + 1$ transition.). For each t , we can consider the row vector $p(t) = (p_1(t), \dots, p_n(t))$ where $p_j(t) = \text{Prob}(X(t) = j)$.

²²Residual refers to the extent to which equations are not satisfied. Here, the equation is $AU = F$, which \tilde{U} does not satisfy, so $R = A\tilde{U} - F$ is the residual.

These probabilities satisfy the differential equation $\dot{p} = \frac{d}{dt}p = pA$. The solution can be written in terms of the *fundamental solution*, $S(t)$, which in an $n \times n$ matrix that satisfies $\dot{S} = SA$, $S(0) = I$.

- (a) Show that if $\dot{S} = SA$, $S(0) = I$, then $p(t) = p(0)S(t)$.
- (b) The *matrix exponential* may be defined through the Taylor series $\exp(B) = \sum_{k=0}^{\infty} \frac{1}{k!} B^k$. Use matrix norms and the fact that $\|B^k\| \leq \|B\|^k$ to show that the infinite sum of matrices converges.
- (c) Show that the fundamental solution is given by $S(t) = \exp(tA)$. To do this, it is enough to show that $\exp(tA)$ satisfies the differential equation $\frac{d}{dt} \exp(tA) = \exp(tA)A$ using the infinite series, and show $\exp(0A) = I$.
- (d) Suppose $A = R\Lambda L$ is the eigenvalue and eigenvector decomposition of A , show that $\exp(tA) = R \exp(t\Lambda) L$, and that $\exp(t\Lambda)$ is the obvious diagonal matrix.
- (e) Use the Matlab function `[R,Lam] = eig(A)`; to calculate the eigenvalues and right eigenvector matrix of A . Let r_k be the k^{th} column of R . For $k = 1, \dots, n$, print r_k , Ar_k , $\lambda_k r_k$, and $\|\lambda_k r_k - Ar_k\|$ (you choose the norm). Mathematically, one of the eigenvectors is a multiple of the vector $\mathbf{1}$ defined in part h. The corresponding eigenvalue is $\lambda = 0$. The computed eigenvalue is not exactly zero. Take $n = 4$ for this, but do not hard wire $n = 4$ into the Matlab code. Set $\lambda = 1$ and $\mu = 4$.
- (f) Let $L = R^{-1}$, which can be computed in Matlab using `L=R\(-1)`; . Let l_k be the k^{th} row of L , check that the l_k are left eigenvectors of A as in part e. Corresponding to $\lambda = 0$ is a left eigenvector that is a multiple of p_{∞} from part h. Check this.
- (g) Write a program in Matlab to calculate $S(t)$ using the eigenvalues and eigenvectors of A as above. Compare the results to those obtained using the Matlab built in function `S = expm(t*A)`; . Use the values $\lambda = 1$, $\mu = 4$, $t = 1$, and n ranging from $n = 4$ to $n = 80$. Compare the two computed $\hat{S}(t)$ (one using eigenvalues, the other just using `expm`) using the l^1 matrix norm. Use the Matlab routine `cond(R)` to compute the condition number of the eigenvector matrix, R . Print three columns of numbers, n , error, condition number. Comment on the quantitative relation between the error and the condition number.
- (h) Here we figure out which of the answers is correct. To do this, use the known fact that $\lim_{t \rightarrow \infty} S(t) = S_{\infty}$ has the simple form $S_{\infty} = \mathbf{1}p_{\infty}$, where $\mathbf{1}$ is the column vector with all ones, and p_{∞} is the row vector with $p_{\infty,j} = ((1 - \rho)/(1 - \rho^n))r^j$, with $\rho = \lambda/\mu$. Take $t = 3 * n$ (which is close enough to $t = \infty$ for this purpose) and the same values of n and see which version of $S(t)$ is correct. What can you say about the stability of computing the matrix exponential using the ill conditioned eigenvalue/eigenvector problem?

13. This exercise explores eigenvalue and eigenvector perturbation theory for the matrix A defined in exercise 12. Let B be the $n \times n$ matrix with $b_{jk} = 0$ for all (j, k) except $b_{00} = -1$ and $b_{1, n-1} = 1$ (as in exercise 12, indices run from $j = 0$ to $j = n - 1$). Define $A(s) = A + sB$, so that $A(0) = A$ and $\frac{dA(s)}{ds} = B$ when $s = 0$.
- (a) For $n = 20$, print the eigenvalues of $A(s)$ for $s = 0$ and $s = .1$. What does this say about the condition number of the eigenvalue eigenvector problem? All the eigenvalues of a real tridiagonal matrix are real²³ but that $A(s = .1)$ is not tridiagonal and its eigenvalues are not real.
 - (b) Use first order eigenvalue perturbation theory to calculate $\dot{\lambda}_k = \frac{d}{ds}\lambda_k$ when $s = 0$. What size s do you need for $\Delta\lambda_k$ to be accurately approximated by $s\dot{\lambda}_k$? Try $n = 5$ and $n = 20$. Note that first order perturbation theory always predicts that eigenvalues stay real, so $s = .1$ is much too large for $n = 20$.

²³It is easy to see that if A is tridiagonal then there is a diagonal matrix, W , so that WAW^{-1} is symmetric. Therefore, A has the same eigenvalues as the symmetric matrix WAW^{-1} .

Chapter 5

Linear Algebra II, Algorithms

5.1 Introduction

This chapter discusses some of the algorithms of computational linear algebra. For routine applications it probably is better to use these algorithms as software packages rather than to recreate them. Still, it is important to know what they do and how they work. Some applications call for variations on the basic algorithms here. For example, Chapter 6 refers to a modification of the Cholesky decomposition.

Many algorithms of numerical linear algebra may be formulated as ways to calculate matrix factorizations. This point of view gives conceptual insight and often suggests alternative algorithms. This is particularly useful in finding variants of the algorithms that run faster on modern processor hardware. Moreover, computing matrix factors explicitly rather than implicitly allows the factors, and the work in computing them, to be re-used. For example, the work to compute the LU factorization of an $n \times n$ matrix, A , is $O(n^3)$, but the work to solve $Au = b$ is only $O(n^2)$ once the factorization is known. This makes it much faster to re-solve if b changes but A does not.

This chapter does not cover the many factorization algorithms in great detail. This material is available, for example, in the book of Golub and Van Loan [7] and many other places. Our aim is to make the reader aware of what the computer does (roughly), and how long it should take. First we explain how the classical Gaussian elimination algorithm may be viewed as a matrix factorization, the LU factorization. The algorithm presented is not the practical one because it does not include *pivoting*. Next, we discuss the Cholesky (LL^*) decomposition, which is a natural version of LU for symmetric positive definite matrices. Understanding the details of the Cholesky decomposition will be useful later when we study optimization methods and still later when we discuss sampling multivariate normal random variables with correlations. Finally, we show how to compute matrix factorizations, such as the QR decomposition, that involve orthogonal matrices.

5.2 Counting operations

A simple way to estimate running time is to count the number of floating point operations, or *flops* that a particular algorithm performs. We seek to estimate $W(n)$, the number of flops needed to solve a problem of size n , for large n . Typically,

$$W(n) \approx Cn^p, \quad (5.1)$$

for large n . Most important is the power, p . If $p = 3$, then $W(2n) \approx 8W(n)$. This is the case for most factorization algorithms in this chapter.

One can give work estimates in the “big O” notation from Section 3.1. We say that $W(n) = O(n^p)$ if there is a C so that $W(n) \leq Cn^p$ for all $n > 0$. This is a less precise and less useful statement than (5.1). It is common to say $W(n) = O(n^p)$ when one really means (5.1). For example, we say that matrix

multiplication takes $O(n^3)$ operations and back substitution takes $O(n^2)$, when we really mean that they satisfy (5.1) with $p = 3$ and $p = 2$ respectively¹.

Work estimates like (5.1) often are derived by using integral approximations to sums. As an example of this idea, consider the sum

$$S_1(n) = \sum_{k=1}^n k = \frac{n(n+1)}{2} = \frac{1}{2}n^2 + \frac{1}{2}n \approx \frac{1}{2}n^2. \quad (5.2)$$

We would like to pass directly to the useful approximation at the end without the complicated exact formula in the middle. As an example, the simple fact (draw a picture)

$$k^2 = \int_{k-1}^k x^2 dx + O(k).$$

implies that (using (5.2) to add up $O(k)$)

$$S_2(n) = \sum_{k=1}^n k^2 = \int_0^n x^2 dx + \sum_{k=1}^n O(k) = \frac{1}{3}n^3 + O(n^2).$$

The integral approximation to (5.2) is easy to picture in a graph. The sum represents the area under a staircase in a large square. The integral represents the area under the diagonal. The error is the smaller area between the staircase and the diagonal.

Consider the problem of computing $C = AB$, where A and B are general $n \times n$ matrices. The standard method uses the formulas

$$c_{jk} = \sum_{l=1}^n a_{jl}b_{lk}. \quad (5.3)$$

The right side requires n multiplies and about n adds for each j and k . That makes $2n \times n^2 = 2n^3$ multiplies and adds in total. The formulas (5.3) have (almost) exactly $2n^3$ flops in this sense. More generally, suppose A is $n \times m$ and B is $m \times p$. Then AB is $n \times p$, and it takes (approximately) m flops to calculate each entry. Thus, the entire AB calculation takes about $2nmp$ flops.

Now consider a matrix triple product ABC , where the matrices are not square but are compatible for multiplication. Say A is $n \times m$, B is $m \times p$, and C is $p \times q$. If we do the calculation as $(AB)C$, then we first compute AB and then multiply by C . The work to do it this way is $2(nmp + npq)$ (because AB is $n \times p$). On the other hand doing it as $A(BC)$ has total work $2(mpq + nmq)$. These numbers are not the same. Depending on the shapes of the matrices, one could be much smaller than the other. The associativity of matrix multiplication allows us to choose the order of the operations to minimize the total work.

¹ The actual relation can be made precise by writing $W = \Theta(n^2)$, for example, which means that $W = O(n^2)$ and $n^2 = O(W)$.

5.3 Gauss elimination and LU decomposition

5.3.1 A 3×3 example

Gauss elimination is a simple systematic way to solve systems of linear equations. For example, suppose we have the system of equations

$$\begin{array}{rrcr} 4x_1 & + & 4x_2 & + & 2x_3 & = & 2 \\ 4x_1 & + & 5x_2 & + & 3x_3 & = & 3 \\ 2x_1 & + & 3x_2 & + & 3x_3 & = & 5. \end{array}$$

Now we *eliminate* x_1 from the second two equations, first subtracting the first equation from the second equation, then subtracting half the first equation from the third equation.

$$\begin{array}{rrcr} 4x_1 & + & 4x_2 & + & 2x_3 & = & 2 \\ & & x_2 & + & x_3 & = & 1 \\ & & x_2 & + & 2x_3 & = & 4. \end{array}$$

Now we eliminate x_2 from the last equation by subtracting the second equation:

$$\begin{array}{rrcr} 4x_1 & + & 4x_2 & + & 2x_3 & = & 2 \\ & & x_2 & + & x_3 & = & 1 \\ & & & & x_3 & = & 3. \end{array}$$

Finally, we solve the last equation for $x_3 = 3$, then *back-substitute* into the second equation to find $x_2 = -2$, and then back-substitute into the first equation to find $x_1 = 1$.

We gain insight into the elimination procedure by writing it in matrix terms. We begin with the matrix equation $Ax = b$:

$$\begin{pmatrix} 4 & 4 & 2 \\ 4 & 5 & 3 \\ 2 & 3 & 3 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} 2 \\ 3 \\ 5 \end{pmatrix}.$$

The operation of eliminating x_1 can be written as an invertible linear transformation. Let M_1 be the transformation that subtracts the first component from the second component and half the first component from the third component:

$$M_1 = \begin{pmatrix} 1 & 0 & 0 \\ -1 & 1 & 0 \\ -0.5 & 0 & 1 \end{pmatrix}. \quad (5.4)$$

The equation $Ax = b$ is equivalent to $M_1Ax = M_1b$, which is

$$\begin{pmatrix} 4 & 4 & 2 \\ 0 & 1 & 1 \\ 0 & 1 & 2 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} 2 \\ 1 \\ 4 \end{pmatrix}.$$

We apply a second linear transformation to eliminate x_2 :

$$M_2 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & -1 & 1 \end{pmatrix}.$$

The equation $M_2 M_1 A x = M_2 M_1 b$ is

$$\begin{pmatrix} 4 & 4 & 2 \\ 0 & 1 & 1 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} 2 \\ 1 \\ 3 \end{pmatrix}.$$

The matrix $U = M_2 M_1 A$ is *upper triangular*, and so we can apply back substitution.

We can rewrite the equation $U = M_2 M_1 A$ as $A = M_1^{-1} M_2^{-1} U = LU$. The matrices M_1 and M_2 are *unit lower triangular*: that is, all of the diagonal elements are one, and all the elements above the diagonal are zero. All unit lower triangular matrices have inverses that are unit lower triangular, and products of unit lower triangular matrices are also unit lower triangular, so $L = M_1^{-1} M_2^{-1}$ is unit lower triangular². The inverse of M_2 corresponds to undoing the last elimination step, which subtracts the second component from the third, by adding the third component back to the second. The inverse of M_1 can be constructed similarly. The reader should verify that in matrix form we have

$$L = M_1^{-1} M_2^{-1} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 1 & 1 & 0 \\ 0.5 & 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 1 & 1 & 0 \\ 0.5 & 1 & 1 \end{pmatrix}.$$

Note that the subdiagonal elements of L form a record of the elimination procedure: when we eliminated the i th variable, we subtracted l_{ij} times the i th equation from the j th equation. This turns out to be a general pattern.

One advantage of the LU factorization interpretation of Gauss elimination is that it provides a framework for organizing the computation. We seek lower and upper triangular matrices L and U so that $LU = A$:

$$\begin{pmatrix} 1 & 0 & 0 \\ l_{21} & 1 & 0 \\ l_{31} & l_{32} & 1 \end{pmatrix} \begin{pmatrix} u_{11} & u_{12} & u_{13} \\ 0 & u_{22} & u_{23} \\ 0 & 0 & u_{33} \end{pmatrix} = \begin{pmatrix} 4 & 4 & 2 \\ 4 & 5 & 3 \\ 2 & 3 & 3 \end{pmatrix}. \quad (5.5)$$

If we multiply out the first row of the product, we find that

$$1 \cdot u_{1j} = a_{1j}$$

for each j ; that is, the first row of U is the same as the first row of A . If we multiply out the first column, we have

$$l_{i1} u_{11} = a_{i1}$$

² The unit lower triangular matrices actually form an *algebraic group*, a subgroup of the group of invertible $n \times n$ matrices.

```

function [L,U] = mylu(A)

    n = length(A);    % Get the dimension of A
    L = eye(n);        % Initialize the diagonal elements of L to 1

    for i = 1:n-1      % Loop over variables to eliminate
        for j = i+1:n  % Loop over equations to eliminate from

            % Subtract L(j,i) times the ith row from the jth row
            L(j,i) = A(j,i) / A(i,i);
            for k = i:n
                A(j,k) = A(j,k) - L(j,i)*A(i,k);
            end

        end
    end

    U = A;              % A is now transformed to upper triangular

```

Figure 5.1: Example MATLAB program to compute $A = LU$.

or $l_{i1} = a_{i1}/u_{11} = a_{i1}/a_{11}$; that is, the first column of L is the first column of A divided by a scaling factor.

Finding the first column of L and U corresponds to finding the multiples of the first row we need to subtract from the other rows to do the elimination. Actually doing the elimination involves replacing a_{ij} with $a_{ij} - l_{i1}u_{1j} = a_{ij} - a_{i1}a_{11}^{-1}a_{1j}$ for each $i > 1$. In terms of the LU factorization, this gives us a reduced factorization problem with a smaller matrix:

$$\begin{pmatrix} 1 & 0 \\ l_{32} & 1 \end{pmatrix} \begin{pmatrix} u_{22} & u_{23} \\ 0 & u_{33} \end{pmatrix} = \begin{pmatrix} 5 & 3 \\ 3 & 3 \end{pmatrix} - \begin{pmatrix} l_{21} \\ l_{31} \end{pmatrix} \begin{pmatrix} u_{21} & u_{31} \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 1 & 2 \end{pmatrix}.$$

We can continue onward in this way to compute the rest of the elements of L and U . These calculations show that the LU factorization, if it exists, is unique (remembering to put ones on the diagonal of L).

5.3.2 Algorithms and their cost

The basic Gauss elimination algorithm transforms the matrix A into its upper triangular factor U . As we saw in the previous section, the scale factors that appear when we eliminate the i th variable from the j th equation can be interpreted as subdiagonal entries of a unit lower triangular matrix L such that $A = LU$. We show a straightforward MATLAB implementation of this idea in Figure 5.1. These algorithms lack *pivoting*, which is needed for stability. How-

```
function x = mylusolve(L,U,b)

    n = length(b);
    y = zeros(n,1);
    x = zeros(n,1);

    % Forward substitution: L*y = b
    for i = 1:n

        % Subtract off contributions from other variables
        rhs = b(i);
        for j = 1:i-1
            rhs = rhs - L(i,j)*y(j);
        end

        % Solve the equation L(i,i)*y(i) = 1*y(i) = rhs
        y(i) = rhs;

    end

    % Back substitution: U*x = y
    for i = n:-1:1

        % Subtract off contributions from other variables
        rhs = y(i);
        for j = i+1:n
            rhs = rhs - U(i,j)*x(j);
        end

        % Solve the equation U(i,i)*x(i) = rhs
        x(i) = rhs / U(i,i);

    end

end
```

Figure 5.2: Example MATLAB program that performs forward and backward substitution with LU factors (without pivoting)

ever, pivoting does not significantly affect the estimates of operation counts for the algorithm, which we will estimate now.

The total number of arithmetic operations for the LU factorization is given by

$$W(n) = \sum_{i=1}^{n-1} \sum_{j=i+1}^n \left(1 + \sum_{k=i}^n 2 \right).$$

We get this work expression by converting loops in Figure 5.1 into sums and counting the number of operations in each loop. There is one division inside the loop over j , and there are two operations (a multiply and an add) inside the loop over k . There are far more multiplies and adds than divisions, so let us count only these operations:

$$W(n) \approx \sum_{i=1}^{n-1} \sum_{j=i+1}^n \sum_{k=i}^n 2.$$

If we approximate the sums by integrals, we have

$$W(n) \approx \int_0^n \int_x^n \int_x^n 2 \, dz \, dy \, dx = \frac{2}{3} n^3.$$

If we are a little more careful, we find that $W(n) = \frac{2}{3} n^3 + O(n^2)$.

We can compute the LU factors of A without knowing the right hand side b , and writing the LU factorization alone does not solve the system $Ax = b$. Once we know $A = LU$, though, solving $Ax = b$ becomes a simple two-stage process. First, *forward substitution* gives y such that $Ly = b$; then *back substitution* gives x such that $Ux = y$. Solving by forward and backward substitution is equivalent to writing $x = U^{-1}y = U^{-1}L^{-1}b$. Figure 5.2 gives a MATLAB implementation of the forward and backward substitution loops.

For a general $n \times n$ system, forward and backward substitution each take about $\frac{1}{2}n^2$ multiplies and the same number of adds. Therefore, it takes a total of about $2n^2$ multiply and add operations to solve a linear system once we have the LU factorization. This costs much less time than the LU factorization, which takes about $\frac{2}{3}n^3$ multiplies and adds.

The factorization algorithms just described may fail or be numerically unstable even when A is well conditioned. To get a stable algorithm, we need to introduce *pivoting*. In the present context this means adaptively reordering the equations or the unknowns so that the elements of L do not grow. Details are in the references.

5.4 Cholesky factorization

Many applications call for solving linear systems of equations with a symmetric and positive definite A . An $n \times n$ matrix is *positive definite* if $x^*Ax > 0$ whenever $x \neq 0$. Symmetric positive definite (SPD) matrices arise in many applications.

If B is an $m \times n$ matrix with $m \geq n$ and $\text{rank}(B) = n$, then the product $A = B^*B$ is SPD. This is what happens when we solve a linear least squares problem using the normal equations, see Section 4.2.8). If $f(x)$ is a scalar function of $x \in R^n$, the *Hessian* matrix of second partials has entries $h_{jk}(x) = \partial^2 f(x) / \partial x_j \partial x_k$. This is symmetric because $\partial^2 f / \partial x_j \partial x_k = \partial^2 f / \partial x_k \partial x_j$. The minimum of f probably is taken at an x_* with $H(x_*)$ positive definite, see Chapter 6. Solving elliptic and parabolic partial differential equations often leads to large sparse SPD linear systems. The variance/covariance matrix of a multivariate random variable is symmetric, and positive definite except in degenerate cases.

We will see that A is SPD if and only if $A = LL^*$ for a lower triangular matrix L . This is the *Cholesky factorization*, or *Cholesky decomposition* of A . As with the LU factorization, we can find the entries of L from the equations for the entries of $LL^* = A$ one at a time, in a certain order. We write it out:

$$\begin{pmatrix} l_{11} & 0 & 0 & \cdots & 0 \\ l_{21} & l_{22} & 0 & \cdots & \vdots \\ l_{31} & l_{32} & l_{33} & \ddots & \\ \vdots & \vdots & & \ddots & 0 \\ l_{n1} & l_{n2} & \cdots & & l_{nn} \end{pmatrix} \cdot \begin{pmatrix} l_{11} & l_{21} & l_{31} & \cdots & l_{n1} \\ 0 & l_{22} & l_{32} & \cdots & l_{n2} \\ 0 & 0 & l_{33} & \ddots & \vdots \\ \vdots & \vdots & & \ddots & \\ 0 & 0 & \cdots & & l_{nn} \end{pmatrix} \\ = \begin{pmatrix} a_{11} & a_{21} & a_{31} & \cdots & a_{n1} \\ a_{21} & a_{22} & a_{32} & \cdots & a_{n2} \\ a_{31} & a_{32} & a_{33} & \ddots & \vdots \\ \vdots & \vdots & & \ddots & \\ a_{n1} & a_{n2} & \cdots & & a_{nn} \end{pmatrix} .$$

Notice that we have written, for example, a_{32} for the $(2, 3)$ entry because A is symmetric. We start with the top left corner. Doing the matrix multiplication gives

$$l_{11}^2 = a_{11} \implies l_{11} = \sqrt{a_{11}} .$$

The square root is real because $a_{11} > 0$ because A is positive definite and³ $a_{11} = e_1^* A e_1$. Next we match the $(2, 1)$ entry of A . The matrix multiplication gives:

$$l_{21} l_{11} = a_{21} \implies l_{21} = \frac{1}{l_{11}} a_{21} .$$

The denominator is not zero because $l_{11} > 0$ because $a_{11} > 0$. We could continue in this way, to get the whole first column of L . Alternatively, we could match $(2, 2)$ entries to get l_{22} :

$$l_{21}^2 + l_{22}^2 = a_{22} \implies l_{22} = \sqrt{a_{22} - l_{21}^2} .$$

³Here e_1 is the vector with one as its first component and all the rest zero. Similarly $a_{kk} = e_k^* A e_k$.

It is possible to show (see below) that if the square root on the right is not real, then A was not positive definite. Given l_{22} , we can now compute the rest of the second column of L . For example, matching $(3, 2)$ entries gives:

$$l_{31} \cdot l_{21} + l_{32} \cdot l_{22} = a_{32} \implies l_{32} = \frac{1}{l_{22}} (a_{32} - l_{31} \cdot l_{21}) .$$

Continuing in this way, we can find all the entries of L . It is clear that if L exists and if we always use the positive square root, then all the entries of L are uniquely determined.

A slightly different discussion of the Cholesky decomposition process makes it clear that the Cholesky factorization exists whenever A is positive definite. The algorithm above assumed the existence of a factorization and showed that the entries of L are uniquely determined by $LL^* = A$. Once we know the factorization exists, we know the equations are solvable, in particular, that we never try to take the square root of a negative number. This discussion represents L as a product of simple lower triangular matrices, a point of view that will be useful in constructing the QR decomposition (Section 5.5).

Suppose we want to apply Gauss elimination to A and find an elementary matrix of the type (5.4) to set $a_{i1} = a_{1i}$ to zero for $i = 2, 3, \dots, n$. The matrix would be

$$M_1 = \begin{pmatrix} 1 & 0 & 0 & \dots & 0 \\ -\frac{a_{21}}{a_{11}} & 1 & 0 & \dots & 0 \\ -\frac{a_{31}}{a_{11}} & 0 & 1 & \dots & 0 \\ \vdots & & & \ddots & \vdots \\ -\frac{a_{n1}}{a_{11}} & 0 & 0 & \dots & 1 \end{pmatrix}$$

Because we only care about the nonzero pattern, let us agree to write a star as a placeholder for possibly nonzero matrix entries that we might not care about in detail. Multiplying out $M_1 A$ gives:

$$M_1 A = \begin{pmatrix} a_{11} & a_{12} & a_{13} & \dots \\ 0 & * & * & \dots \\ 0 & * & * & \\ \vdots & & & \end{pmatrix} .$$

Only the entries in row two have changed, with the new values indicated by primes. Note that $M_1 A$ has lost the symmetry of A . We can restore this symmetry by multiplying from the right by M_1^* . This has the effect of subtracting $\frac{a_{1i}}{a_{11}}$ times the first column of $M_1 A$ from column i for $i = 2, \dots, n$. Since the top row of A has not changed, this has the effect of setting the $(1, 2)$ through $(1, n)$ entries to zero:

$$M_1 A M_1^* = \begin{pmatrix} a_{11} & 0 & 0 & \dots \\ 0 & * & * & \dots \\ 0 & * & * & \\ \vdots & & & \end{pmatrix} .$$

Continuing in this way, elementary matrices E_{31} , etc. will set to zero all the elements in the first row and top column except a_{11} . Finally, let D_1 be the diagonal matrix which is equal to the identity except that $d_{11} = 1/\sqrt{a_{11}}$. All in all, this gives ($D_1^* = D_1$):

$$D_1 M_1 A M_1^* D_1^* = \begin{pmatrix} 1 & 0 & 0 & \cdots \\ 0 & a'_{22} & a'_{23} & \cdots \\ 0 & a'_{23} & a'_{33} & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}. \quad (5.6)$$

We define L_1 to be the lower triangular matrix

$$L_1 = D_1 M_1,$$

so the right side of (5.6) is $A_1 = L_1 A L_1^*$ (check this). Note that the trailing submatrix elements a'_{ik} satisfy

$$a_{ik'} = a_{ik} - (L_1)_{i1}(L_1)_{k1}.$$

The matrix L_1 is nonsingular since D_1 and M_1 are nonsingular. To see that A_1 is positive definite, simply define $y = L^*x$, and note that $y \neq 0$ if $x \neq 0$ (L_1 being nonsingular), so $x^* A_1 x = x^* L A L^* x = y^* A y > 0$ since A is positive definite. In particular, this implies that $a'_{22} > 0$ and we may find an L_2 that sets a'_{22} to one and all the a_{2k} to zero.

Eventually, this gives $L_{n-1} \cdots L_1 A L_1^* \cdots L_{n-1}^* = I$. Solving for A by reversing the order of the operations leads to the desired factorization:

$$A = L_1^{-1} \cdots L_{n-1}^{-1} L_{n-1}^{-*} \cdots L_1^{-*},$$

where we use the common convention of writing B^{-*} for the inverse of the transpose of B , which is the same as the transpose of the inverse. Clearly, L is given by $L = L_1^{-1} \cdots L_{n-1}^{-1}$.

As in the case of Gaussian elimination, the product matrix L turns out to take a rather simple form:

$$l_{ij} = \begin{cases} -(L_j)_i, & i > j \\ (L_i)_i, & i = j. \end{cases}$$

Alternately, the Cholesky factorization can be seen as the end result of a sequence of transformations to a lower triangular form, just as we thought about the U matrix in Gaussian elimination as the end result of a sequence of transformations to upper triangular form.

Once we have the Cholesky decomposition of A , we can solve systems of equations $Ax = b$ using forward and back substitution, as we did for the LU factorization.

5.5 Least squares and the QR factorization

Many problems in linear algebra call for linear transformations that do not change the l^2 norm:

$$\|Qx\|_{l^2} = \|x\|_{l^2} \quad \text{for all } x \in R^n. \quad (5.7)$$

A real matrix satisfying (5.7) is orthogonal, because⁴

$$\|Qx\|_{l^2}^2 = (Qx)^* Qx = x^* Q^* Qx = x^* x = \|x\|_{l^2}^2.$$

(Recall that $Q^*Q = I$ is the definition of orthogonality for square matrices.)

Just as we were able to transform a matrix A into upper triangular form by a sequence of lower triangular transformations (Gauss transformations), we are also able to transform A into upper triangular form by a sequence of orthogonal transformations. The *Householder reflector* for a vector v is the orthogonal matrix

$$H(v) = I - 2 \frac{vv^*}{v^*v}.$$

The Householder reflector corresponds to reflection through the plane normal to v . We can reflect any given vector w into a vector $\|w\|e_1$ by reflecting through a plane normal to the line connecting w and $\|w\|e_1$

$$H(w - \|w\|e_1)w = \|w\|e_1.$$

In particular, this means that if A is an $m \times n$ matrix, we can use a Householder transformation to make all the subdiagonal entries of the first column into zeros:

$$H_1 A = H(Ae_1 - \|Ae_1\|e_1)A = \begin{pmatrix} \|Ae_1\| & a'_{12} & a'_{13} & \dots & \\ 0 & a'_{22} & a'_{23} & \dots & \\ 0 & a'_{32} & a'_{33} & \dots & \vdots \\ & & & \ddots & \vdots \end{pmatrix}$$

Continuing in the way, we can apply a sequence of Householder transformations in order to reduce A to an upper triangular matrix, which we usually call R :

$$H_n \dots H_2 H_1 A = R.$$

Using the fact that $H_j = H_j^{-1}$, we can rewrite this as

$$A = H_1 H_2 \dots H_n R = QR,$$

where Q is a product of Householder transformations.

The QR factorization is a useful tool for solving least squares problems. If A is a rectangular matrix with $m > n$ and $A = QR$, then the invariance of the l^2 norm under orthogonal transformations implies

$$\|Ax - b\|_{l^2} = \|Q^*Ax - Q^*b\|_{l^2} = \|Rx - b'\|_{l^2}.$$

⁴This shows an orthogonal matrix satisfies (5.7). Exercise 8 shows that a matrix satisfying (5.7) must be orthogonal.

Therefore, minimizing $\|Ax - b\|$ is equivalent to minimizing $\|Rx - b'\|$, where R is upper triangular⁵:

$$\begin{pmatrix} r_{11} & r_{12} & \cdots & r_{1n} \\ 0 & r_{22} & \cdots & \vdots \\ \vdots & \ddots & \ddots & \\ 0 & \cdots & 0 & r_{nn} \\ \hline 0 & \cdots & 0 & \\ \vdots & & & \vdots \\ 0 & \cdots & 0 & \end{pmatrix} \cdot \begin{pmatrix} x_1 \\ x_2 \\ \vdots \\ x_n \end{pmatrix} - \begin{pmatrix} b'_1 \\ b'_2 \\ \vdots \\ b'_n \\ \hline b'_{n+1} \\ \vdots \\ b'_m \end{pmatrix} = \begin{pmatrix} r'_1 \\ r'_2 \\ \vdots \\ r'_n \\ \hline r'_{n+1} \\ \vdots \\ r'_m \end{pmatrix}$$

Assuming none of the diagonals r_{jj} is zero, the best we can do is to choose x so that the first n components of r' are set to zero. Clearly x has no effect on the last $m - n$ components of r'

5.6 Software

5.6.1 Representing matrices

Computer memories are typically organized into a *linear address space*: that is, there is a number (an address) associated with each location in memory. A one-dimensional array maps naturally to this memory organization: the address of the i th element in the array is distance i (measured in appropriate units) from the *base address* for the array. However, we usually refer to matrix elements not by a single array offset, but by row and column indices. There are different ways that a pair of indices can be converted into an array offset. Languages like FORTRAN and MATLAB use *column major order*, so that a 2×2 matrix would be stored in memory a column at a time: `a11`, `a21`, `a12`, `a22`. In contrast, C/C++ use *row major order*, so that if we wrote

```
double A[2][2] = { {1, 2},
                   {3, 4} };
```

the order of the elements in memory would be 1, 2, 3, 4.

C++ has only very primitive built-in support for multi-dimensional arrays. For example, the first dimension of a built-in two-dimensional array in C++ must be known at compile time. However, C++ allows programmers to write *classes* that describe new data types, like matrices, that the basic language does not fully support. The uBLAS library and the Matrix Template Library are examples of widely-used libraries that provide matrix classes.

An older, C-style way of writing matrix algorithms is to manually manage the mappings between row and column indices and array offsets. This can

⁵The upper triangular part of the LU decomposition is called U while the upper triangular part of the QR decomposition is called R . The use of r_{jk} for the entries of R and r_j for the entries of the residual is not unique to this book.

```

/*
 * Solve the n-by-n lower triangular system Lx = b.
 * L is assumed to be in column-major form.
 */
void forward_substitute(const double* Ldata, double* x,
                       const double* b, int n)
{
#define L(i,j) Ldata[(j)*n+(i)]
    for (int i = 0; i < n; ++i) {
        x[i] = b[i];
        for (int j = 0; j < i; ++j)
            x[i] -= L(j,i)*x[j];
        x[i] /= L(i,i);
    }
#undef L
}

```

Figure 5.3: C++ code to solve a lower triangular system using forward substitution. Note the macro definition (and corresponding undefinition) used to simplify array accesses.

be made a little more convenient using C macros. For example, consider the forward substitution code in Figure 5.3. The line

```
#define L(i,j) Ldata[(j)*n+(i)]
```

defines `L` to be a macro with two arguments. Any time the compiler encounters `L(i,j)` while the macro is defined, it will substitute `Ldata[(j)*n+(i)]`. For example, if we wanted to look at an entry of the first superdiagonal, we might write `L(i,i+1)`, which the compiler would translate to `Ldata[(i+1)*n+(i)]`. Note that the parentheses in the macro definition are important – this is *not* the same as looking at `Ldata[i+1*n+i]`! Also note that we have followed the C convention of zero-based indexing, so the first matrix entry is at `L(0,0) = Ldata[0]`.

In addition to the column-major and row-major layouts, there are many other ways to represent matrices in memory. Often, it is possible to represent an $n \times n$ matrix in a way that uses far less than the n^2 numbers of a standard row-major or column-major layout. A matrix that can be represented efficiently with much fewer than n^2 memory entries is called *sparse*; other matrices are called *dense*. Much of the discussion here and in Chapter 5 applies mainly to dense matrices. A modern (2009) desktop computer with 2 GB (2^{31} bytes) of memory can store at most 2^{28} double precision numbers, so a square matrix with $n > 2^{14} \approx 16000$ variables would not even fit in memory. Sparse matrix

methods can handle larger problems and often give faster methods even for problems that can be handled using dense matrix methods. For example, finite element computations often lead to sparse matrices with hundreds of thousands of variables that can be solved in minutes.

One way a matrix can be sparse is for most of its entries to be zero. For example, discretizations of the Laplace equation in three dimensions have as few as seven non-zero entries per row, so that $7/n$ is the fraction of entries of A that are not zero. Sparse matrices in this sense also arise in circuit problems, where a non-zero entry in A corresponds to a direct connection between two elements in the circuit. Such matrices may be stored in *sparse matrix format*, in which we keep lists noting which entries are not zero and the values of the non-zero elements. Computations with such sparse matrices try to avoid *fill in*. For example, they would avoid explicit computation of A^{-1} because most of its entries are not zero. Sparse matrix software has heuristics that often do very well in avoiding fill in. The interested reader should consult the references.

In some cases it is possible to compute the *matrix vector product* $y = Ax$ for a given x efficiently without calculating the entries of A explicitly. One example is the *discrete Fourier transform* (DFT) described in Chapter 1. This is a *full* matrix with n^2 non-zero entries, but the FFT (*fast Fourier transform*) algorithm computes $y = Ax$ in $O(n \log(n))$ operations. Another example is the *fast multipole method* that computes forces from mutual electrostatic interaction of n charged particles with b bits of accuracy in $O(nb)$ work. Many finite element packages never *assemble* the stiffness matrix, A .

Computational methods can be *direct* or *iterative*. A direct method have only rounding errors. They would get the exact answer in exact arithmetic using a predetermined number of arithmetic operations. For example, Gauss elimination computes the *LU* factorization of A using $O(n^3)$ operations. Iterative methods produce a sequence of approximate solutions that converge to the exact answer as the number of iterations goes to infinity. They often are faster than direct methods for very large problems, particularly when A is sparse.

5.6.2 Performance and caches

In scientific computing, *performance* refers to the time it takes to run the program that does the computation. Faster computers give more performance, but so do better programs. To write high performance software, we should know what happens inside the computer, something about the compiler and about the hardware. This and several later Software sections explore performance-related issues.

When we program, we have a mental model of the operations that the computer is doing. We can use this model to estimate how long a computation will take. For example, we know that Gaussian elimination on an $n \times n$ matrix takes about $\frac{2}{3}n^3$ flops, so on a machine that can execute at R flop/s, the elimination procedure will take at least $\frac{2n^3}{3R}$ seconds. Desktop machines that run at a rate of 2-3 gigahertz (billions of cycles per second) can often execute two floating

point operations per cycle, for a peak performance of 4-5 gigaflops. On a desktop machine capable of 4 gigaflop/s (billions of flop/s), we would expect to take at least about a sixth of a second to factor a matrix with $n = 1000$. In practice, though, only very carefully written packages will run anywhere near this fast. A naive implementation of LU factorization is likely to run $10\times$ to $100\times$ slower than we would expect from a simple flop count.

The failure of our work estimates is not due to a flaw in our formula, but in our model of the computer. A computer can only do arithmetic operations as fast as it can fetch data from memory, and often the time required to fetch data is much greater than the time required to perform arithmetic. In order to try to keep the processor supplied with data, modern machines have a *memory hierarchy*. The memory hierarchy consists of different size memories, each with a different *latency*, or time taken between when memory is requested and when the first byte can be used by the processor⁶. At the top of the hierarchy are a few *registers*, which can be accessed immediately. The register file holds on the order of a hundred bytes. The *level 1 cache* holds tens of kilobytes – 32 and 64 kilobyte cache sizes are common at the time of this writing – within one or two clock cycles. The *level 2 cache* holds a couple megabytes of information, and usually has a latency of around ten clock cycles. *Main memories* are usually one or two gigabytes, and might take around fifty cycles to access. When the processor needs a piece of data, it will first try the level 1 cache, then resort to the level 2 cache and then to the main memory only if there is a *cache miss*. The programmer has little or no control over the details of how a processor manages its caches. However, cache policies are designed so that codes with good *locality* suffer relatively few cache misses, and therefore have good performance. There are two versions of cache locality: temporal and spatial.

Temporal locality occurs when we use the same variables repeatedly in a short period of time. If a computation involves a *working set* that is not so big that it cannot fit in the level 1 cache, then the processor will only have to fetch the data from memory once, after which it is kept in cache. If the working set is too large, we say that the program *thrashes the cache*. Because dense linear algebra routines often operate on matrices that do not fit into cache, they are often subject to cache thrashing. High performance linear algebra libraries organize their computations by partitioning the matrix into blocks that fit into cache, and then doing as much work as possible on each block before moving onto the next.

In addition to temporal locality, most programs also have some *spatial locality*, which means that when the processor reads one memory location, it is likely to read nearby memory locations in the immediate future. In order to get the best possible spatial locality, high performance scientific codes often access their data with *unit stride*, which means that matrix or vector entries are accessed one after the other in the order in which they appear in memory.

⁶ Memories are roughly characterized by their *latency*, or the time taken to respond to a request, and the *bandwidth*, or the rate at which the memory supplies data after a request has been initiated.

5.6.3 Programming for performance

Because of memory hierarchies and other features of modern computer architecture, simple models fail to accurately predict performance. The picture is even more subtle in most high-level programming languages because of the natural, though mistaken, inclination to assume that program statements take the about the same amount of time just because they can be expressed with the same number of keystrokes. For example, printing two numbers to a file is several hundred times slower than adding two numbers. A poorly written program may even spend more time allocating memory than it spends computing results⁷.

Performance tuning of scientific codes often involves hard work and difficult trade-offs. Other than using a good compiler and turning on the optimizer⁸, there are few easy fixes for speeding up the performance of code that is too slow. The source code for a highly tuned program is often much more subtle than the textbook implementation, and that subtlety can make the program harder to write, test, and debug. Fast implementations may also have different stability properties from their standard counterparts. The first goal of scientific computing is to get an answer which is accurate enough for the purpose at hand, so we do not consider it progress to have a tuned code which produces the wrong answer in a tenth the time that a correct code would take. We can try to prevent such mis-optimizations from affecting us by designing a thorough set of unit tests *before* we start to tune.

Because of the potential tradeoffs between speed, clarity, and numerical stability, performance tuning should be approached carefully. When a computation is slow, there is often one or a few critical *bottlenecks* that take most of the time, and we can best use our time by addressing these first. It is not always obvious where bottlenecks occur, but we can usually locate them by *timing* our codes. A *profiler* is an automated tool to determine where a program is spending its time. Profilers are useful both for finding bottlenecks both in compiled languages like C and in high-level languages like MATLAB⁹.

Some bottlenecks can be removed by calling fast libraries. Rewriting a code to make more use of libraries also often clarifies the code – a case in which improvements to speed and clarity go hand in hand. This is particularly the case in MATLAB, where many problems can be solved most concisely with a few calls to fast built-in factorization routines that are much faster than hand-written loops. In other cases, the right way to remedy a bottleneck is to change algorithms or data structures. A classic example is a program that solves a sparse linear system – a tridiagonal matrix, for example – using a general-purpose dense routine. One good reason for learning about matrix factorization algorithms is that you will then know which of several possible factorization-based solution methods is fastest or most stable, even if you have not written

⁷ We have seen this happen in actual codes that we were asked to look at.

⁸ Turning on compiler optimizations is one of the simplest things you can do to improve the performance of your code.

⁹ The MATLAB profiler is a beautifully informative tool – type `help profile` at the MATLAB prompt to learn more about it

the factorization codes yourself. One often hears the statement that increasing computer power makes it unnecessary to find faster algorithms. We believe the opposite: the greater the computer power, the larger the problems one can attempt, and the greater the difference a good algorithm can make.

Some novice scientific programmers (and even some experienced programmers) write codes with bottlenecks that have nothing to do with the main part of their computation. For example, consider the following fragment of MATLAB code:

```
n = 1000;
A = [];
for i = 1:n
    A(i,i) = 1;
end
```

On one of our desktop machines, the time to execute this loop went from about six seconds to under two milliseconds just by changing the second statement to `A = zeros(n)`. The problem with the original code is that at each step MATLAB is forced to enlarge the matrix, which it does by allocating a larger block of memory, copying the old matrix contents to the new location, and only then writing the new element. Therefore, the original code takes $O(i^2)$ time to run step i , and the overall cost of the loop scales like $O(n^3)$. Fortunately, this sort of blunder is relatively easy to fix. It is always worth timing a code before trying to tune it, just to make sure that the bottlenecks are where you think they are, and that your code is not wasting time because of a programming blunder.

5.7 References and resources

The algorithms of numerical linear algebra for dense matrices are described in great detail in the book by Charles Van Loan and Gene Golub [7] and in the book by James Demmel [4]. The book *Direct Methods for Sparse Linear Systems* by Tim Davis describes computational methods for matrices stored in sparse matrix format [3]. Still larger problems are solved by *iterative methods*. Generally speaking, iterative methods are not very effective unless the user can concoct a good *preconditioner*, which is an approximation to the inverse. Effective preconditioners usually depend in physical understanding of the problem and are problem specific. The book *Templates for the Solution of Linear Systems* provides a concise description of many different iterative methods and preconditioners [1].

While the topic is generally covered in courses in computer architecture, there are relatively few textbooks on performance optimization that seem suitable for scientific programmers without a broad CS background. The book *Performance Optimization of Numerically Intensive Codes* by Stefan Goedecker and Adolfo Hoesie [1] is one noteworthy exception. We also recommend *High Performance Computing* by Kevin Down and Charles Severance [21]. Truly

high performance computing is done on computers with more than one processor, which is called *parallel computing*. There are many specialized algorithms and programming techniques for parallel computing.

The LAPACK software package is designed to make the most of memory hierarchies. The performance of LAPACK depends on fast Basic Linear Algebra Subroutines (BLAS). There is a package called *ATLAS* that produces automatically tuned BLAS libraries for different architectures, as well as choosing good parameters (block sizes, etc.) for LAPACK depending on the cache performance of your particular processor. The LAPACK manual is published by SIAM, the *Society for Industrial and Applied Mathematics*.

5.8 Exercises

1. The solution to $Au = b$ may be written $u = A^{-1}b$. This can be a good way to analyze algorithms involving linear systems (see Sections 4.3.1 and 6.3). But we try to avoid forming A^{-1} explicitly in computations because it is more than twice as expensive as solving the linear equations. A good way to form $B = A^{-1}$ is to solve the matrix equation $AB = I$. Gauss elimination applied to A gives $A = LU$, where the entries of L are the pivots used in elimination.
 - (a) Show that about $\frac{2}{3}n^3$ work reduces $AB = I$ to $UB = L^{-1}$, where the entries of U and L^{-1} are known.
 - (b) Show that computing the entries of B from $UB = L^{-1}$ takes about $\frac{1}{2}n^3$ work. Hint: It takes one flop per element for each of the n elements of the bottom row of B , then two flops per element of the $n-1$ row of B , and so on to the top. The total is $n \times (1 + 2 + \cdots + n)$.
 - (c) Use this to verify the claim that computing A^{-1} is more than twice as expensive as solving $Au = b$.
2. Show that a symmetric $n \times n$ real matrix is positive definite if and only if all its eigenvalues are positive. Hint: If R is a right eigenvector matrix, then, for a symmetric matrix we may normalize so that $R^{-1} = R^*$ and $A = R\Lambda R^*$, where Λ is a diagonal matrix containing the eigenvalues (See Sections 4.2.5 and 4.2.7). Then $x^*Ax = x^*R\Lambda R^*x = (x^*R)\Lambda(R^*x) = y^*\Lambda y$, where $y = R^*x$. Show that $y^*\Lambda y > 0$ for all $y \neq 0$ if and only if all the diagonal entries of Λ are positive. Show that if A is positive definite, then there is a $C > 0$ so that $x^*Ax > C\|x\|_{l^2}^2$ for all x . (Hint: $\|x\|_{l^2} = \|Rx\|_{l^2}$, $C = \lambda_{\min}$.)
3. Write a program to compute the LL^* decomposition of an SPD matrix A . Your procedure should have as arguments the dimension, n , and the matrix A . The output should be the Cholesky factor, L . Your procedure must detect and report a matrix that is not positive definite and should not perform the operation `sqrtc` if $c < 0$. Write another procedure that has n and L as arguments and returns the product LL^* . Hand in: (i) printouts

of the two procedures and the driving program, (ii) a printout of results showing that the testing routine reports failure when $LL^* \neq A$, (iii) a printout showing that the Cholesky factoring procedure reports failure when A is not positive definite, (iv) a printout showing that the Cholesky factoring procedure works correctly when applied to a SPD matrix, proven by checking that $LL^* = A$.

4. A square matrix A has *bandwidth* $2k + 1$ if $a_{jp} = 0$ whenever $|j - p| > k$. A *subdiagonal* or *superdiagonal* is a set of matrix elements on one side of the main diagonal (below for sub, above for super) with $j - k$, the distance to the diagonal, fixed. The bandwidth is the number of nonzero bands. A bandwidth 3 matrix is *tridiagonal*, bandwidth 5 makes *pentadiagonal*, etc.
 - (a) Show that a SPD matrix with bandwidth $2k + 1$ has a Cholesky factor with nonzeros only on the diagonal and up to k bands below.
 - (b) Show that the Cholesky decomposition algorithm computes this L in work proportional to $k^2 n$ (if we skip operations on entries of A outside its nonzero bands).
 - (c) Write a procedure for Cholesky factorization of tridiagonal SPD matrices, and apply it to the matrix of Exercise 11, compare the running time with this dense matrix factorizer and the one from Exercise 5.4. Of course, check that the answer is the same, up to roundoff.
5. Suppose v_1, \dots, v_m is an orthonormal basis for a vector space $V \subseteq R^n$. Let L be a linear transformation from V to V . Let A be the matrix that represents L in this basis. Show that the entries of A are given by

$$a_{jk} = v_j^* L v_k. \quad (5.8)$$

Hint: Show that if $y \in V$, the representation of y in this basis is $y = \sum_j y_j v_j$, where $y_j = v_j^* y$. In physics and theoretical chemistry, inner products of the form (5.8) are called *matrix elements*. For example, the eigenvalue perturbation formula (4.39) (in physicist terminology) simply says that the perturbation in an eigenvalue is (nearly) equal to the appropriate matrix element of the perturbation in the matrix.

6. Suppose A is an $n \times n$ symmetric matrix and $V \subset R^n$ is an *invariant subspace* for A (i.e. $Ax \in V$ if $x \in V$). Show that A defines a linear transformation from V to V . Show that there is a basis for V in which this linear transformation (called *A restricted to V*) is represented by a symmetric matrix. Hint: construct an orthonormal basis for V .
7. If Q is an $n \times n$ matrix, and $(Qx)^* Qy = x^* y$ for all x and y , show that Q is an orthogonal matrix. Hint: If $(Qx)^* Qy = x^* (Q^* Q)y = x^* y$, we can explore the entries of $Q^* Q$ by choosing particular vectors x and y .
8. If $\|Qx\|_{l_2} = \|x\|_{l_2}$ for all x , show that $(Qx)^* Qy = x^* y$ for all x and y . Hint (*polarization*): If $\|Q(x + sy)\|_{l_2}^2 = \|x + sy\|_{l_2}^2$ for all s , then $(Qx)^* Qy = x^* y$.

Chapter 6

Nonlinear Equations and Optimization

6.1 Introduction

This chapter discusses two related computational problems. One is *root finding*, or solving systems of nonlinear equations. This means that we seek values of n variables, $(x_1, \dots, x_n) = x \in R^n$, to satisfy n nonlinear equations $f(x) = (f_1(x), \dots, f_n(x)) = 0$. We assume that $f(x)$ is a smooth function of x . The other problem is *smooth optimization*, or finding the minimum (or maximum¹) value of a smooth *objective function*, $V(x)$. These problems are closely related. Optimization algorithms use the gradient of the objective function, solving the system of equations $g(x) = \nabla V(x) = 0$. Conversely, root finding can be seen as minimizing $\|f(x)\|^2$.

The theory here is for *black box* methods. These are algorithms that do not depend on details of the definitions of the functions $f(x)$ or $V(x)$. The code doing the root finding will learn about f only by “user-supplied” procedures that supply values of f or V and their derivatives. The person writing the root finding or optimization code need not “open the box” to see how the user procedure works. This makes it possible for specialists to create general purpose optimization and root finding software that is efficient and robust, without knowing all the problems it may be applied to.

There is a strong incentive to use derivative information as well as function values. For root finding, we use the $n \times n$ Jacobian matrix, $f'(x)$, with entries $f'(x)_{jk} = \partial_{x_k} f_j(x)$. For optimization, we use the gradient and the $n \times n$ *Hessian* matrix of second partials $H(x)_{jk} = \partial_{x_j} \partial_{x_k} V(x)$. It may seem like too much extra work to go from the n components of f to the n^2 entries of f' , but algorithms that use f' often are much faster and more reliable than those that do not.

There are drawbacks to using general-purpose software that treats each specific problem as a black box. Large-scale computing problems usually have specific features that have a big impact on how they should be solved. Reformulating a problem to fit into a generic $f(x) = 0$ or $\min_x V(x)$ form may increase the condition number. Problem-specific solution strategies may be more effective than the generic Newton’s method. In particular, the Jacobian or the Hessian may be sparse in a way that general purpose software cannot take advantage of. Some more specialized algorithms are in Exercise 6c (Marquart-Levenberg for nonlinear least squares), and Section ?? (Gauss-Seidel iteration for large systems).

The algorithms discussed here are *iterative* (see Section 2.4). They produce a sequence of approximations, or *iterates*, that should converge to the desired solution, x_* . In the simplest case, each *iteration* starts with a *current iterate*, \bar{x} , and produces a *successor iterate*, $x' = \Phi(\bar{x})$. The algorithm starts from an *initial guess*², x_0 , then produces a sequence of iterates $x_{k+1} = \Phi(x_k)$. The algorithm succeeds if the iterates converge to the solution: $x_k \rightarrow x_*$ as $k \rightarrow \infty$.

¹Optimization refers either to minimization or maximization. But finding the maximum of $V(x)$ is the same as finding the minimum of $-V(x)$.

²Here, the subscript denotes the iteration number, not the component. In n dimensions, iterate x_k has components $x_k = (x_{k1}, \dots, x_{kn})$.

An iterative method fails if the iterates fail to converge or converge to the wrong answer. For an algorithm that succeeds, the *convergence rate* is the rate at which $\|x_k - x_*\| \rightarrow 0$ as $k \rightarrow \infty$.

An iterative method is *locally convergent* if it succeeds whenever the initial guess is close enough to the solution. That is, if there is an $R > 0$ so that if $\|x_0 - x_*\| \leq R$ then $x_k \rightarrow x_*$ as $k \rightarrow \infty$. The algorithms described here, mostly variants of *Newton's method*, all are locally convergent if the problem is non-degenerate (terminology below). An iterative method is *globally convergent* if it finds the answer from any initial guess. Between these extremes are algorithms that are more or less *robust*. The algorithms described here consist of a relatively simple locally convergent method, usually Newton's method, enhanced with *safeguards* that guarantee that some progress is made toward the solution from any \bar{x} . We will see that safeguards based on mathematical analysis and reasoning are more effective than heuristics.

All iterative methods need some kind of *convergence criterion* (more properly, *halting criterion*). One natural possibility is to stop when the relative change in x is small enough: $\|x_{k+1} - x_k\| / \|x_k\| \leq \epsilon$. It also makes sense to check that the *residuals*, the components of $f(x)$ or $\nabla V(x)$, are small. Even without roundoff error, an iterative method would be very unlikely to get the exact answer. However, as we saw in Section 2.4, good algorithms and well conditioned problems still allow essentially optimal accuracy: $\|x_k - x_*\| / \|x_*\| \sim \epsilon_{\text{mach}}$.

The final section of this chapter is on methods that do not use higher derivatives. The discussion applies to linear or nonlinear problems. For optimization, it turns out that the rate of convergence of these methods is determined by the condition number of H for solving linear systems involving H , see Section 4.3.1 and the condition number formula (4.35). More precisely, the number of iterations needed to reduce the error by a factor of 2 is proportional to $\kappa(H) = \lambda_{\max}(H) / \lambda_{\min}(H)$. This, more than linear algebra roundoff, explains our fixation on condition number. The condition number $\kappa(H) = 10^4$ could arise in a routine partial differential equation problem. This bothers us not so much because it makes us lose 4 out of 16 double precision digits of accuracy, but because it takes tens of thousands of iterations to solve the problem with a naive method.

6.2 Solving a single nonlinear equation

The simplest problem is that of solving a single equation in a single variable: $f(x) = 0$. Single variable problems are easier than multi-variable problems. There are simple criteria that guarantee a solution exists. Some algorithms for one dimensional problems, Newton's method in particular, have analogues for higher dimensional problems. Others, such as bisection, are strictly one dimensional. Algorithms for one dimensional problems are components for algorithms for higher dimensional problems.

6.2.1 Bisection

Bisection search is a simple, robust way to find a zero of a function on one variable. It does not require $f(x)$ to be differentiable, but merely continuous. It is based on a simple topological fact called the *intermediate value theorem*: if $f(x)$ is a continuous real-valued function of x on the interval $a \leq x \leq b$ and $f(a) < 0 < f(b)$, then there is at least one $x_* \in (a, b)$ with $f(x_*) = 0$. A similar theorem applies in the case $b < a$ or $f(a) > 0 > f(b)$.

The bisection search algorithm consists of repeatedly bisecting an interval in which a root is known to lie. Suppose we have an interval³ $[\bar{a}, \bar{b}]$ with $f(\bar{a}) < 0$ and $f(\bar{b}) > 0$. The intermediate value theorem tells us that there is a root of f in $[\bar{a}, \bar{b}]$. The uncertainty in the location of this root is the length of the interval $|\bar{b} - \bar{a}|$. To cut that uncertainty in half, we bisect the interval. The midpoint is $\bar{c} = (\bar{a} + \bar{b})/2$. We determine the sign of $f(\bar{c})$, probably by evaluating it. If $f(\bar{c}) > 0$ then we know there is a root of f in the sub-interval $[\bar{a}, \bar{c}]$. In this case, we take the new interval to be $[a', b']$, with $a' = \bar{a}$ and $b' = \bar{c}$. In the other case, $f(\bar{c}) < 0$, we take $a' = \bar{c}$ and $b' = \bar{b}$. In either case, f changes sign over the half size interval $[a', b']$.

To start the bisection algorithm, we need an initial interval $[a_0, b_0]$ over which f changes sign. Running the bisection procedure then produces intervals $[a_k, b_k]$ whose size decreases at an exponential rate:

$$|b_k - a_k| = 2^{-k} |b_0 - a_0| .$$

To get a feeling for the convergence rate, use the approximate formula $2^{10} = 10^3$. This tells us that we get three decimal digits of accuracy for each ten iterations. This may seem good, but Newton's method is much faster, when it works. Moreover, Newton's method generalizes to more than one dimension while there is no useful multidimensional analogue of bisection search. Exponential convergence often is called *linear convergence* because of the linear relationship $|b_{k+1} - a_{k+1}| = \frac{1}{2} |b_k - a_k|$. Newton's method is faster than this.

Although the bisection algorithm is robust, it can fail if the computed approximation to $f(x)$ has the wrong sign. The user of bisection should take into account the accuracy of the function approximation as well as the interval length when evaluating the accuracy of a computed root.

6.2.2 Newton's method for a nonlinear equation

As in the previous section, we want to find a value, x_* , that solves a single nonlinear equation $f(x_*) = 0$. We have procedures that return $f(x)$ and $f'(x)$ for any given x . At each iteration, we have a current iterate, \bar{x} and we want to find an x' that is closer to x_* . Suppose that \bar{x} is close to x_* . The values $f(\bar{x})$ and $f'(\bar{x})$ determine the tangent line to the graph of $f(x)$ at the point \bar{x} . The new iterate, x' , is the point where this tangent line crosses the x axis. If $f(x)$ is

³The interval notation $[a, b]$ used here is not intended to imply that $a < b$. For example, the interval $[5, 2]$ consists of all numbers between 5 and 2, endpoints included.

close to zero, then x' should be close to \bar{x} and the tangent line approximation should be close to f at x' , which suggests that $f(x')$ should be small.

More analytically, the tangent line approximation (See Section 3.1) is

$$f(x) \approx F^{(1)}(x) = f(\bar{x}) + f'(\bar{x}) \cdot (x - \bar{x}) . \quad (6.1)$$

Finding where the line crosses the x axis is the same as setting $F^{(1)}(x) = 0$ and solving for x' :

$$x' = \bar{x} - f'(\bar{x})^{-1} f(\bar{x}) . \quad (6.2)$$

This is the basic Newton method.

The local convergence rate of Newton's method is governed by the error in the approximation (6.1). The analysis assumes that the root x_* is *non-degenerate*, which means that $f'(x_*) \neq 0$. The convergence for degenerate roots is different, see Exercise 1. For a non-degenerate root, we will have $f'(\bar{x}) \neq 0$ for \bar{x} close enough to x_* . Assuming this, (6.2) implies that $|x' - \bar{x}| = O(|f(\bar{x})|)$. This, together with the Taylor series error bound

$$f(x') - F^{(1)}(x') = O(|x' - \bar{x}|^2) ,$$

and the *Newton equation* $F^{(1)}(x') = 0$, implies that

$$|f(x')| = O(|f(\bar{x})|^2) .$$

This means that there is a $C > 0$ so that

$$|f(x')| \leq C \cdot |f(\bar{x})|^2 . \quad (6.3)$$

This manifestation of *local quadratic convergence* says that the residual at the next iteration is roughly proportional⁴ to the square of the residual at the current iterate.

Quadratic convergence is very fast. In a typical problem, once $x_k - x_*$ is moderately small, the residual will be at roundoff levels in a few more iterations. For example, suppose that⁵ $C = 1$ in (6.3) and that $|x_k - x_*| = .1$. Then $|x_{k+1} - x_*| \leq .01$, $|x_{k+2} - x_*| \leq 10^{-4}$, and $|x_{k+4} - x_*| \leq 10^{-16}$. The number of correct digits doubles at each iteration. By contrast, a *linearly convergent* iteration with $|x_{k+1} - x_*| \leq .1 \cdot |x_k - x_*|$ gains one digit of accuracy per iteration and takes 15 iterations rather than 4 to go from .1 to 10^{-16} . Bisection search needs about 50 iterations to reduce the error by a factor of 10^{15} ($10^{15} = (10^3)^5 \approx (2^{10})^5 = 2^{50}$).

Unfortunately, the quadratic convergence of Newton's method is local. There is no guarantee that $x_k \rightarrow x_*$ as $k \rightarrow \infty$ if the initial guess is not close to x_* . A program for finding x_* must take this possibility into account. See Section 3.7.2 for some ideas on how to do this.

⁴Strictly speaking, (6.3) is just a bound, not an estimate. However, Exercise 2 shows that $f(x')$ really is approximately proportional to $f(\bar{x})^2$.

⁵This does not make sense on dimensional grounds. It would be more accurate and more cumbersome to describe this stuff in terms of relative error.

6.3 Newton's method in more than one dimension

Newton's method applies also to solving systems of nonlinear equations. The linear approximation (6.1) applies in dimensions $n > 1$ if $f'(\bar{x})$ is the Jacobian matrix evaluated at \bar{x} , and f and $(x - \bar{x})$ are column vectors. We write the *Newton step* as $x' - \bar{x} = z$, so $x' = \bar{x} + z$. Newton's method determines z by replacing the nonlinear equations, $f(\bar{x} + z) = 0$, with the linear approximation,

$$0 = f(\bar{x}) + f'(\bar{x})z . \quad (6.4)$$

To carry out one step of Newton's method, we must evaluate the function $f(\bar{x})$, the Jacobian, $f'(\bar{x})$, then solve the linear system of equations (6.4). We may write this as

$$z = -(f'(\bar{x}))^{-1} f(\bar{x}) , \quad (6.5)$$

which is a natural generalization of the one dimensional formula (6.2). In computational practice (see Exercise 1) it usually is more expensive to form $(f'(\bar{x}))^{-1}$ than to solve (6.4).

Newton's method for systems of equations also has quadratic (very fast) local convergence to a non-degenerate solution x_* . As in the one dimensional case, this is because of the error bound in the linear approximation (6.1). For $n > 1$, we write the Taylor approximation error bound in terms of norms:

$$\|f(\bar{x} + z) - F^{(1)}\| = \|f(\bar{x} + z) - \{ f(\bar{x}) + f'(\bar{x})z \}\| = O(\|z\|^2) .$$

We see from (6.5) that⁶

$$\|z\| \leq C \|f(\bar{x})\| .$$

Together, these inequalities imply that if $\bar{x} - x_*$ is small enough then

$$\|f(x')\| = \|f(\bar{x} + z)\| \leq C \|f(\bar{x})\|^2 ,$$

which is quadratic convergence, exactly as in the one dimensional case.

In practice, Newton's method can be frustrating for its lack of robustness. The user may need some ingenuity to find an x_0 close enough to x_* to get convergence. In fact, it often is hard to know whether a system of nonlinear equations has a solution at all. There is nothing as useful as the intermediate value theorem from the one dimensional case, and there is no multi-dimensional analogue of the robust but slow bisection method in one dimension.

While Newton's method can suffer from extreme ill conditioning, it has a certain robustness against ill conditioning that comes from its *affine invariance*. Affine invariance states Newton's method is invariant under *affine transformations*. An affine transformation is a mapping $x \rightarrow Ax + b$ (it would be linear

⁶The definition of a non-degenerate solution is that $f'(x_*)$ is nonsingular. If \bar{x} is close enough to x_* , then $f'(\bar{x})$ will be close enough to $f'(x_*)$ that it also will be nonsingular (See (4.16)). Therefore $\|z\| \leq \|(f'(\bar{x}))^{-1}\| \|f(\bar{x})\| \leq C \|f(\bar{x})\|$.

without the b). An affine transformation⁷ of $f(x)$ is $g(y) = Af(By)$, where A and B are invertible $n \times n$ matrices. The affine invariance is that if we start from corresponding initial guesses: $x_0 = By_0$, and create iterates y_k by applying Newton's method to $g(y)$ and x_k by applying Newton's method to $f(x)$, then the iterates also correspond: $x_k = By_k$. This means that Newton's method works exactly as well on the original equations $f(x) = 0$ as on the transformed equations $g(y) = 0$. For example, we can imagine changing from x to variables y in order to give each of the unknowns the same units. If $g(y) = 0$ is the best possible rescaling of the original equations $f(x) = 0$, then applying Newton's method to $f(x) = 0$ gives equivalent iterates.

This argument can be restated informally as saying that Newton's method makes sense on dimensional grounds and therefore is natural. The variables x_1, \dots, x_n may have different units, as may the functions f_1, \dots, f_n . The n^2 entries $f'(x)$ all may have different units, as may the entries of $(f')^{-1}$. The matrix vector product that determines the components of the Newton step (see (6.4)), $z = -(f')^{-1} f(\bar{x})$, involves adding a number of contributions (entries in the matrix $(f')^{-1}$ multiplying components of f) that might seem likely to have a variety of units. Nevertheless, each of the n terms in the sum implicit in the matrix-vector product (6.5) defining a component z_j has the same units as the corresponding component, x_j . See Section 6.6 for a more detailed discussion of this point.

6.3.1 Quasi-Newton methods

Local quadratic convergence is the incentive for evaluating the Jacobian matrix. Evaluating the Jacobian matrix may not be so expensive if much of the work in evaluating f can be re-used in calculating f' . There are other situations where the Jacobian is nearly impossible to evaluate analytically. One possibility would be to estimate f' using finite differences. Column k of $f'(\bar{x})$ is $\partial_{x_k} f(\bar{x})$. The cheapest and least accurate approximation to this is the first-order one-sided difference formula⁸: $(f(\bar{x} + \Delta x_k e_k) - f(\bar{x})) / \Delta x_k$. Evaluating all of f' in this way would take n extra evaluations of f per iteration, which may be so expensive that it outweighs the fast local convergence.

Quasi-Newton methods replace the true $f'(\bar{x})$ in the Newton equations (6.4) by estimates of $f'(\bar{x})$ built up from function values over a sequence of iterations. If we call this approximate Jacobian A_k , the quasi-Newton equations are

$$0 = f(x_k) + A_k z_k . \quad (6.6)$$

The simplest such method is the *secant method* for one-dimensional root finding. Using the current x_k and $f(x_k)$, and the previous x_{k-1} and $f(x_{k-1})$, we use the slope of the line connecting the current $(x_k, f(x_k))$ to the previous

⁷Actually, this is a linear transformation. It is traditional to call it affine though the constant terms are missing.

⁸ e_k is the unit vector in the x_k direction and Δx_k is a step size in that direction. Different components of x may have different units and therefore require different step sizes.

$(x_{k-1}, f(x_{k-1}))$ to estimate the slope of the tangent line at x_k . The result is

$$A_k = \frac{f(x_k) - f(x_{k-1})}{x_k - x_{k-1}}, \quad x_{k+1} = x_k - f(x_k)/A_k. \quad (6.7)$$

The local convergence rate of the secant method (6.7) is better than linear ($|x_{k+1} - x_*| \leq C|x_k - x_*|$) and worse than quadratic.

Many multidimensional quasi-Newton methods work by updating A_k at each iteration so that $A_{k+1}z_k = f(x_{k+1}) - f(x_k)$. In higher dimensions, this does not determine A_{k+1} completely, because it involves n equations, which is not enough to define the n^2 elements of A_{k+1} . The references give several suggestions for update formulas. The good ones have the property that if you apply them to linear equations, you find the exact $A = f'$ in n steps. It is not clear that such a property makes quasi-Newton methods better than ordinary Newton's method with finite difference approximations to the elements of the Jacobian.

6.4 One variable optimization

Suppose $n = 1$ and we wish to find the minimum of the function of a single variable, $V(x)$. Please bear with the following long list of definitions. We say that x_* is a *local minimum* of V if there is an $R > 0$ so that $V(x_*) \leq V(x)$ whenever $|x - x_*| \leq R$. We say that x_* is a *strict local minimum* if $V(x) > V(x_*)$ whenever $x \neq x_*$ and $|x - x_*| \leq R$. We say that x_* is a *global minimum* if $V(x_*) \leq V(x)$ for all x for which $V(x)$ is defined, and a *strict global minimum* if $V(x_*) < V(x)$ for all $x \neq x_*$ for which V is defined. Finally, x_* is a *nondegenerate local minimum* if $V''(x_*) > 0$. The Taylor series remainder theorem implies that if $V'(x_*) = 0$ and $V''(x_*) > 0$, then x_* is at least a strict local minimum. The function $V(x)$ is *convex* if ⁹ $\alpha V(x) + \beta V(y) > V(\alpha x + \beta y)$ whenever $\alpha \geq 0$, $\beta \geq 0$, and $\alpha + \beta = 1$. The function is *strictly convex* if $V''(x) > 0$ for all x . A strictly convex function is convex, but the function $V(x) = x^4$ is not strictly convex, because $V''(0) = 0$. This function has a strict but degenerate global minimum at $x_* = 0$.

For the curious, there is an analogue of bisection search in one variable optimization called *golden section search*. It applies to any continuous function that is *unimodal*, meaning that V has a single global minimum and no local minima. The *golden mean*¹⁰ is $r = (1 + \sqrt{5})/2 \approx 1.62$. At each stage of bisection search we have an interval $[\bar{a}, \bar{b}]$ in which there must be at least one root. At each stage of golden section search we have an interval $[\bar{a}, \bar{c}]$ and a third point $\bar{b} \in [\bar{a}, \bar{c}]$ with

$$|\bar{a} - \bar{c}| = r |\bar{a} - \bar{b}|. \quad (6.8)$$

⁹The reader should check that this is the same as the geometric condition that the line segment connecting the points $(x, V(x))$ and $(y, V(y))$ lies above the graph of V .

¹⁰This number comes up in many ways. From Fibonacci numbers it is $r = \lim_{k \rightarrow \infty} f_{k+1}/f_k$. If $(\alpha + \beta)/\alpha = \alpha/\beta$ and $\alpha > \beta$, then $\alpha/\beta = r$. This has the geometric interpretation that if we remove an $\alpha \times \alpha$ square from one end of an $\alpha \times (\alpha + \beta)$ rectangle, then the remaining smaller $\beta \times \alpha$ rectangle has the same aspect ratio as the original $\alpha \times (\alpha + \beta)$ rectangle. Either of these leads to the equation $r^2 = r + 1$.

As with our discussion of bisection search, the notation $[\bar{a}, \bar{c}]$ does not imply that $\bar{a} < \bar{c}$. In bisection, we assume that $f(\bar{a}) \cdot f(\bar{b}) < 0$. Here, we assume that $f(\bar{b}) < f(\bar{a})$ and $f(\bar{b}) < f(\bar{c})$, so that there must be a local minimum within $[\bar{a}, \bar{c}]$. Now (this is the clever part), consider a fourth point in the larger subinterval, $d = (1 - \frac{1}{r})\bar{a} + \frac{1}{r}\bar{b}$. Evaluate $f(\bar{d})$. If $f(\bar{d}) < f(\bar{b})$, take $a' = \bar{a}$, $b' = \bar{d}$, and $c' = \bar{b}$. Otherwise, take $a' = \bar{c}$, $b' = \bar{b}$, and $c' = \bar{d}$, reversing the sense of the interval. In either case, $|a' - c'| = r|a' - b'|$, and $f(b') < f(a')$ and $f(b') < f(c')$, so the iteration can continue. Each stage reduces the uncertainty in the minimizer by a factor of $\frac{1}{r}$, since $|a' - c'| = \frac{1}{r}|\bar{a} - \bar{c}|$.

6.5 Newton's method for local optimization

Most of the properties listed in Section 6.4 are the same for multi-variable optimization. We denote the gradient as $g(x) = \nabla V(x)$, and the Hessian matrix of second partials as $H(x)$. An x_* with $g(x_*) = 0$ and $H(x_*)$ positive definite (see Section 5.4 and Exercise 2) is called non-degenerate, a natural generalization of the condition $V'' > 0$ for one variable problems. Such a point is at least a local minimum because the Taylor series with error bound is

$$V(x_* + z) - V(x_*) = \frac{1}{2}z^*H(x_*)z + O(\|z\|^3).$$

Exercise 2 shows that the first term on the right is positive and larger than the second if $H(x_*)$ is positive definite and $\|z\|$ is small enough. If $H(x_*)$ is negative definite (obvious definition), the same argument shows that x_* is at least a local maximum. If $H(x_*)$ has some positive and some negative eigenvalues (and $g(x_*) = 0$) then x_* is neither a local minimum nor a local maximum, but is called a *saddle point*. In any case, a local minimum must satisfy $g(x_*) = 0$ if V is differentiable.

We can use Newton's method from Section 6.3 to seek a local minimum by solving the equations $g(x) = 0$, but we must pay attention to the difference between row and column vectors. We have been considering x , the Newton step, z , etc. to be column vectors while $\nabla V(x) = g(x)$ is a row vector. For this reason, we consider applying Newton's method to the column vector of equations $g^*(x) = 0$. The Jacobian matrix of the column vector function $g^*(x)$ is the Hessian H (check this). Therefore, the locally convergent Newton method is

$$x' = \bar{x} + z,$$

where the step z is given by the Newton equations

$$H(\bar{x})z = -g^*(\bar{x}). \quad (6.9)$$

Because it is a special case of Newton's method, it has local quadratic convergence to x_* if x_* is a local non-degenerate local minimum.

Another point of view for the local Newton method is that each iteration minimizes a *quadratic model* of the function $V(\bar{x} + z)$. The three term Taylor

series approximation to V about \bar{x} is

$$V(\bar{x} + z) \approx V^{(2)}(\bar{x}, z) = V(\bar{x}) + \nabla V(\bar{x})z + \frac{1}{2}z^*H(\bar{x})z. \quad (6.10)$$

If we minimize $V^{(2)}(\bar{x}, z)$ over z , the result is $z = -H(\bar{x})^{-1}\nabla V(\bar{x})^*$, which is the same as (6.9). As for Newton's method for nonlinear equations, the intuition is that $V^{(2)}(\bar{x}, z)$ will be close to $V(\bar{x} + z)$ for small z . This should make the minimum of $V^{(2)}(\bar{x}, z)$ close to the minimizer of V , which is x_* .

Unfortunately, this simple local method cannot distinguish between a local minimum, a local maximum, or even a saddle point. If x_* has $\nabla V(x_*) = 0$ (so x_* is a *stationary point*) and $H(x_*)$ is nonsingular, then the iterates $x_{k+1} = x_k - H(x_k)^{-1}g^*(x_k)$ will happily converge to x_* if $\|x_0 - x_*\|$ is small enough. This could be a local maximum or a saddle point. Moreover, if $\|x_0 - x_*\|$ is not small, we have no idea whether the iterates will converge to anything at all.

The main difference between the unsafeguarded Newton method optimization problem and general systems of nonlinear equations is that the Hessian is symmetric and (close enough to a non-degenerate local minimum) positive definite. The Jacobian f' need not be symmetric. The Cholesky decomposition requires storage for the roughly $\frac{1}{2}n^2$ distinct elements of H and takes about $\frac{1}{6}n^3$ floating points to compute L . This is about half the storage and work required for a general non-symmetric linear system using the LU factorization.

6.6 Safeguards and global optimization

The real difference between minimization and general systems of equations comes from the possibility of evaluating $V(x)$ and forcing it to decrease from iteration to iteration. It is remarkable that two simple *safeguards* turn the unreliable Newton's method into a much more robust (though not perfect) method that converges to a local minimum from almost any initial guess. These are (i) finding a *descent direction* by modifying $H(\bar{x})$ if necessary, and (ii) using a one dimensional *line search* to prevent wild steps. Both of the safeguards have the purpose of guaranteeing descent, that $V(x') < V(\bar{x})$.

In principle, this would allow the x_k to converge to a saddle point, but this is extremely unlikely in practice because saddle points are unstable for this process.

The safeguarded methods use the formulation of the *search directions*, p and the *step size*, $t > 0$. One iteration will take the form $x' = \bar{x} + z$, where the step is $z = tp$. We define the search direction to be a *descent direction* if

$$\left. \frac{d}{dt}V(\bar{x} + tp) \right|_{t=0} = g(\bar{x}) \cdot p < 0. \quad (6.11)$$

This guarantees that if $t > 0$ is small enough, then $V(\bar{x} + tp) < V(\bar{x})$. Then we find a step size, t , that actually achieves this property. If we prevent t from becoming too small, it will be impossible for the iterates to converge except to a stationary point.

We find the search direction by solving a modified Newton equation

$$\tilde{H}p = -g^*(\bar{x}) . \quad (6.12)$$

Putting this into (6.11) gives

$$\left. \frac{d}{dt} V(\bar{x} + tp) \right|_{t=0} = -g(\bar{x}) \tilde{H} g(\bar{x})^* .$$

This is negative if \tilde{H} is positive definite (the right hand side is a 1×1 matrix (a number) because g is a row vector). One algorithm for finding a descent direction would be to apply the Cholesky decomposition algorithm (see Section 5.4). If the algorithm finds L with $LL^* = H(\bar{x})$, use this L to solve the Newton equation (6.12) with $\tilde{H} = H(\bar{x}) = LL^*$. If the Cholesky algorithm fails to find L , then $H(\bar{x})$ is not positive definite. A possible substitute (but poor in practice, see below) is $\tilde{H} = I$, which turns Newton's method into *gradient descent*.

A better choice for \tilde{H} comes from the *modified Cholesky* algorithm. This simply replaces the equation

$$l_{kk} = (H_{kk} - l_{k1}^2 + \cdots + l_{k,k-1}^2)^{1/2}$$

with the modified equation using the absolute value

$$l_{kk} = |H_{kk} - l_{k1}^2 + \cdots + l_{k,k-1}^2|^{1/2} . \quad (6.13)$$

Here, H_{kk} is the (k, k) entry of $H(\bar{x})$. This modified Cholesky algorithm produces L with $LL^* = H(\bar{x})$ if and only if $H(\bar{x})$ is positive definite. In any case, we take $\tilde{H} = LL^*$, which is positive definite. Using these non-Cholesky factors, the Newton equations become:

$$LL^*p = -g(\bar{x})^* . \quad (6.14)$$

It is not entirely clear why the more complicated modified Cholesky algorithm is more effective than simply taking $\tilde{H} = I$ when $H(\bar{x})$ is not positive definite. One possible explanation has to do with units. Let us suppose that U_k represents the units of x_k , such as seconds, dollars, kilograms, etc. Let us also suppose that $V(x)$ is dimensionless. In this case the units of $H_{jk} = \partial_{x_j} \partial_{x_k} V$ are $[H_{jk}] = 1/U_j U_k$. We can verify by studying the Cholesky decomposition equations from Section ?? that the entries of L have units $[l_{jk}] = 1/U_j$, whether we use the actual equations or the modification (6.13). We solve (??) in two stages, first $Lq = -\nabla V^*$, then $L^*p = q$. Looking at units, it is clear that all the elements of q are dimensionless and that the elements of p have units $[p_k] = U_k$. Thus, the modified Cholesky algorithm produces a search direction that component by component has the same units as x . This allows the update formula $x' = \bar{x} + tp$ to make sense with a dimensionless t . The reader should check that the choice $\tilde{H} = I$ does not have this property in general, even if we allow t to have units, if the U_k are different.

The second safeguard is a limited *line search*. In general, line search means minimizing the function $\phi(t) = V(x + tp)$ over the single variable t . This could be done using golden section search, but a much more rudimentary binary search process suffices as a safeguard. In this binary search, we evaluate $\phi(0) = V(\bar{x})$ and $\phi(1) = V(\bar{x} + p)$. If $\phi(1) > \phi(0)$, the step size is too large. In that case, we keep reducing t by a factor of 2 ($t = t/2$;) until $\phi(t) < \phi(0)$, or we give up. If p is a search direction, we will have $\phi(t) < \phi(0)$ for small enough t and this bisection process will halt after finitely many reductions of t . If $\phi(1) < \phi(0)$, we enter a greedy process of increasing t by factors of 2 until $\phi(2t) > \phi(t)$. This process will halt after finitely many doublings if the set of x with $V(x) < V(\bar{x})$ is bounded.

A desirable feature is that the safeguarded algorithm gives the ordinary Newton step, and rapid (quadratic) local convergence, if \bar{x} is close enough to a nondegenerate local minimum. The modified Hessian will correspond to the actual Hessian if $H(x_*)$ is positive definite and \bar{x} is close enough to x_* . The step size will be the default $t = 1$ if \bar{x} is close enough to x_* because the quadratic model (6.10) will be accurate. The quadratic model has $V^{(2)}(\bar{x}, 2z) > V^{(2)}(\bar{x}, z)$, because z is the minimizer of $V^{(2)}$.

6.7 Determining convergence

There are two reasonable ways to judge when an iterative method is “close to” a right answer: a small *residual error* or a small error in the solution. For example, suppose we seek solutions to the equation $f(x) = 0$. The iterate x_k has a small residual error if $\|f(x_k)\|$ is small for some appropriate norm. There is a small error in the solution if $\|x_k - x_*\|$ is small, where $\|x_*\|$ is the true solution.

The residual error is easy to determine, but what about the error $\|x_k - x_*\|$? For Newton’s method, each successive iterate is much more accurate than the previous one, so that as long as x_* is a regular root (i.e. the Jacobian is nonsingular), then

$$\|x_k - x_*\| = \|x_k - x_{k+1}\| + O(\|x_k - x_{k+1}\|^2).$$

Therefore, it is natural to use the length of the Newton step from x_k to x_{k+1} as an estimate for the error in x_k , assuming that we are close to the solution and that we are taking full Newton steps. Of course, we would then usually return the approximate root x_{k+1} , even though the error estimate is for x_k . That is, we test convergence based on the difference between a less accurate and a more accurate formula, and return the more accurate result even though the error estimate is only realistic for the less accurate formula. This is exactly the same strategy we used in our discussion of integration, and we will see it again when we discuss step size selection in the integration of ordinary differential equations.

6.8 Gradient descent and iterative methods

The *gradient descent* optimization algorithm uses the identity matrix as the approximate Hessian, $H = I$, so the (negative of the) gradient becomes the search direction: $p = -\nabla V(\bar{x})^*$. This seems to make sense geometrically, as the negative gradient is the steepest downhill direction (leading to the name *method of steepest descent*). With a proper line search, gradient descent has the theoretical global robustness properties of the more sophisticated Newton method with the modified Cholesky approximate Hessian. But much of the research in sophisticated optimization methods is motivated by the fact that simple gradient descent converges slowly in many applications.

One indication of trouble with gradient descent is that the formula,

$$x' = \bar{x} - t\nabla V(\bar{x})^* , \quad (6.15)$$

does not make dimensional sense in general, see Section 6.6. Written in components, (6.15) is $x'_k = \bar{x}_k - t\partial_{x_k} V(\bar{x})$. Applied for $k = 1$, this makes dimensional sense if the units of t satisfy $[t] = [x_1^2]/[V]$. If the units of x_2 are different from those of x_1 , the x_2 equation forces units of t inconsistent with those from the x_1 equation.

We can understand the slow convergence of gradient descent by studying how it works on the model problem $V(x) = \frac{1}{2}x^*Hx$, with a symmetric and positive definite H . This is the same as assuming that the local minimum is nondegenerate and the local approximation (6.10) is exact¹¹. In this case the gradient satisfies $g(x)^* = \nabla V(x)^* = Hx$, so solving the Newton equations (6.9) gives the exact solution in one iteration. We study the gradient method with a fixed step size¹², t , which implies $x_{k+1} = x_k - tHx_k$. We write this as

$$x_{k+1} = Mx_k , \quad (6.16)$$

where

$$M = I - tH . \quad (6.17)$$

The convergence rate of gradient descent in this case is the rate at which $x_k \rightarrow 0$ in the iteration (6.16).

This, in turn, is related to the eigenvalues of M . Since H and M are symmetric, we may choose an orthonormal basis in which both are diagonal: $H = \text{diag}(\lambda_1, \dots, \lambda_n)$, and $M = \text{diag}(\mu_1, \dots, \mu_n)$. The λ_j are positive, so we may assume that $0 < \lambda_{\min} = \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n = \lambda_{\max}$. The formula (6.17) implies that

$$\mu_j = 1 - t\lambda_j . \quad (6.18)$$

After k iterations of (6.16), we have $x_{kj} = \mu_j^k x_{0j}$, where x_{kj} component j of the iterate x_k . Clearly, the rate at which $x_k \rightarrow 0$ depends on the *spectral gap*,

$$\rho = 1 - \max_j |\mu_j| ,$$

¹¹We simplified the problem but have not lost generality by taking $x_* = 0$ here.

¹²See Exercise 7 for an example showing that line search does not improve the situation very much.

in the sense that the estimate

$$\|x_k\| \leq (1 - \rho)^k \|x_0\|$$

is sharp (take $x_0 = e_1$ or $x_0 = e_n$). The optimal step size, t is the one that maximizes ρ , which leads to (see (6.18))

$$\begin{aligned} 1 - \rho &= \mu_{\max} = 1 - t\lambda_{\min} \\ \rho - 1 &= \mu_{\min} = 1 - t\lambda_{\max} . \end{aligned}$$

Solving these gives the optimizing value $t = 2/(\lambda_{\min} + \lambda_{\max})$ and

$$\rho = 2 \cdot \frac{\lambda_{\min}}{\lambda_{\max}} \approx \frac{2}{\kappa(H)} . \quad (6.19)$$

If we take $k = 2/\rho \approx \kappa(H)$ iterations, and H is ill conditioned so that k is large, the error is reduced roughly by a factor of

$$(1 - \rho)^k = \left(1 - \frac{2}{k}\right)^k \approx e^{-2} .$$

This justifies what we said in the Introduction, that it takes $k = \kappa(H)$ iterations to reduce the error by a fixed factor.

6.8.1 Gauss Seidel iteration

The Gauss-Seidel iteration strategy makes use of the fact that optimizing over one variable is easier than optimizing over several. It goes from \bar{x} to x' in n steps. Step j optimizes $V(x)$ over the single component x_j with all other components fixed. We write the n intermediate stages as $x^{(j)}$ with components $x^{(j)} = (x_1^{(j)}, \dots, x_n^{(j)})$. Starting with $x^{(0)} = \bar{x}$, we go from $x^{(j-1)}$ to $x^{(j)}$ by optimizing over component j . That is $x_m^{(j-1)} = x_m^{(j)}$ if $m \neq j$, and we get $x_j^{(j)}$ by solving

$$\min_{\xi} V((x_1^{(j-1)}, \dots, x_{j-1}^{(j-1)}, \xi, x_{j+1}^{(j-1)}, \dots, x_n^{(j-1)})) .$$

6.9 Resources and further reading

The book by Ortega and Rheinboldt has a more detailed discussion of Newton's method for solving systems of nonlinear equations [17]. The book *Practical Optimization* by Phillip Gill, Walter Murray, and my colleague Margaret Wright, has much more on nonlinear optimization including methods for constrained optimization problems [5]. There is much public domain software for smooth optimization problems, but I don't think much of it is useful.

The set of initial guesses x_0 so that $k_k \rightarrow x_*$ as $t \rightarrow \infty$ is the *basin of attraction* of x_* . If the method is locally convergent, the basin of attraction contains a ball of radius R about x_* . The boundary of the basin of attraction can be a beautiful fractal set. The picture book *Fractals* by Benoit Mandelbrot, some of the most attractive fractals arise in this way.

6.10 Exercises

1. Study the convergence of Newton's method applied to solving the equation $f(x) = x^2 = 0$. Show that the root $x_* = 0$ is *degenerate* in that $f'(x_*) = 0$. The Newton iterates are x_k satisfying $x_{k+1} = x_k - f(x_k)/f'(x_k)$. Show that the local convergence in this case is *linear*, which means that there is an $\alpha < 1$ with $|x_{k+1} - x_*| \approx \alpha |x_k - x_*|$. Note that so called linear convergence still implies that $x_k - x_* \rightarrow 0$ exponentially. Nevertheless, contrast this linear local convergence with the quadratic local convergence for a nondegenerate problem.
2. Use the Taylor expansion to second order to derive the approximation

$$f(x') \approx C(\bar{x})f(\bar{x})^2 = \frac{1}{2} \frac{f''(\bar{x})}{f'(\bar{x})^2} \cdot f(\bar{x})^2. \quad (6.20)$$

Derive a similar expression that shows that $(x' - x_*)$ is approximately proportional to $(\bar{x} - x_*)^2$. Use (6.20) to predict that applying Newton's method to finding solving the equation $\sin(x) = 0$ will have superquadratic convergence. What makes $C(x)$ large, and the convergence slow, is (i) small $f'(x)$ (a nearly degenerate problem), and (ii) large $f''(x)$ (a highly nonlinear problem).

3. The function $f(x) = x/\sqrt{1+x^2}$ has a unique root: $f(x) = 0$ only for $x = 0$. Show that the unsafeguarded Newton method gives $x_{k+1} = x_k^3$. Conclude that the method succeeds if and only if $|x_0| < 1$. Draw graphs to illustrate the first few iterates when $x_0 = .5$ and $x_0 = 1.5$. Note that Newton's method for this problem has local cubic convergence, which is even faster than the more typical local quadratic convergence. The formula (6.20) explains why.
4. Suppose $n = 2$ and x_1 has units of (electric) charge, x_2 has units of mass, $f_1(x_1, x_2)$ has units of length, and $f_2(x_1, x_2)$ has units of time. Find the units of each of the four entries of $(f')^{-1}$. Verify the claims about the units of the step, z , at the end of Section 6.3.
5. Suppose x_* satisfies $f(x_*) = 0$. The *basin of attraction* of x_* is the set of x so that if $x_0 = x$ then $x_k \rightarrow x_*$ as $k \rightarrow \infty$. If $f'(x_*)$ is non-singular, the basin of attraction of x_* under unsafeguarded Newton's method includes at least a neighborhood of x_* , because Newton's method is locally convergent. Exercise 3 has an example in one dimension where the basin of attraction of $x_* = 0$ is the open interval (endpoints not included) $(-1, 1)$. Now consider the two dimensional problem of finding roots of $f(z) = z^2 - 1$, where $z = x + iy$. Written out in its real components, $f(x, y) = (x^2 - y^2 - 1, 2xy)$. The basin of attraction of the solution $z_* = 1$ ($(x_*, y_*) = (1, 0)$) includes a neighborhood of $z = 1$ but surprisingly many many other points in the complex plane. This *Mandelbrot set* is one of the most beautiful examples of a two dimensional fractal. The purpose of this exercise is to make a pretty picture, not to learn about scientific computing.

- (a) Show that Newton iteration is $z_{k+1} = z_k - \frac{z_k^2 - 1}{2z_k}$.
 - (b) Set $z_k = 1 + w_k$ and show that $w_{k+1} = \frac{3}{2}w_k^2/(1 + w_k)$.
 - (c) Use this to show that if $|w_k| < \frac{1}{4}$, then $|w_{k+1}| < \frac{1}{2}|w_k|$. Hint: Show $|1 + w_k| > \frac{3}{4}$. Argue that this implies that the basin of attraction of $z_* = 1$ includes at least a disk of radius $\frac{1}{4}$ about z_* , which is a quantitative form of local convergence.
 - (d) Show that if $|z_k - 1| < \frac{1}{4}$ for some k , then z_0 is in the basin of attraction of $z_* = 1$. (This is the point of parts (b) and (c).)
 - (e) Use part (d) to make a picture of the Mandelbrot set. Hint: Divide the rectangle $|x| < R_x$, $0 \leq y \leq R_y$ into a regular grid of small cells of size $\Delta x \times \Delta y$. Start Newton's method from the center of each cell. Color the cell if $|z_k - 1| < \frac{1}{4}$ for some $k \leq N$. See how the picture depends on the parameters Δx , Δy , R_x , R_y , and N .
6. A *saddle point*¹³ is an x so that $\nabla V(x) = 0$ and the Hessian, $H(x)$, is nonsingular and has at least one negative eigenvalue. We do not want the iterates to converge to a saddle point, but most Newton type optimization algorithms seem to have that potential. All the safeguarded optimization methods we discussed have $\Phi(x) = x$ if x is a saddle point because they all find the search direction by solving $\tilde{H}p = -\nabla V(x)$.
- (a) Let $V(x) = x_1^2 - x_2^2$ and suppose \bar{x} is on the x_1 axis. Show that with the modified Cholesky, x' also is on the x_1 axis, so the iterates converge to the saddle point, $x = 0$. Hint: \tilde{H} has a simple form in this case.
 - (b) Show that if $\bar{x}_2 \neq 0$, and $t > 0$ is the step size, and we use the bisection search that increases the step size until $\phi(t) = V(\bar{x} + tp)$ satisfies $\phi(2t) > \phi(t)$, then one of the following occurs:
 - i. The bisection search does not terminate, $t \rightarrow \infty$, and $\phi(t) \rightarrow -\infty$. This would be considered good, since the minimum of V is $-\infty$.
 - ii. The line search terminates with t satisfying $\phi(t) = V(x') < 0$. In this case, subsequent iterates cannot converge to $x = 0$ because that would force V to converge to zero, while our modified Newton strategy guarantees that V decreases at each iteration.
 - (c) *Nonlinear least squares* means finding $x \in R^m$ to minimize $V(x) = \|f(x) - b\|_2^2$, where $f(x) = (f_1(x), \dots, f_n(x))^*$ is a column vector of n nonlinear functions of the m unknowns, and $b \in R^n$ is a vector we are trying to approximate. If $f(x)$ is linear (there is an $n \times m$ matrix A with $f(x) = Ax$), then minimizing $V(x)$ is a linear least squares problem. The *Marquart Levenberg* iterative algorithm solves a linear

¹³The usual definition of saddle point is that H should have at least one positive and one negative eigenvalue and no zero eigenvalues. The simpler criterion here suffices for this application.

least squares problem at each iteration. If the current iterate is \bar{x} , let the linearization of f be the $n \times m$ Jacobian matrix A with entries $a_{ij} = \partial_{x_j} f_i(\bar{x})$. Calculate the step, p , by solving

$$\min_p \|Ap - (b - f(\bar{x}))\|_{l^2} . \quad (6.21)$$

Then take the next iterate to be $x' = \bar{x} + p$.

- i. Show that this algorithm has local quadratic convergence if the residual at the solution has zero residual: $r(x_*) = f(x_*) - b = 0$, but not otherwise (in general).
 - ii. Show that p is a descent direction.
 - iii. Describe a safeguarded algorithm that probably will converge to at least a local minimum or diverge.
7. This exercise shows a connection between the slowness of gradient descent and the condition number of H in one very special case. Consider minimizing the model quadratic function in two dimensions $V(x) = \frac{1}{2} (\lambda_1 x_1^2 + \lambda_2 x_2^2)$ using gradient descent. Suppose the line search is done exactly, choosing t to minimize $\phi(t) = V(\bar{x} + tp)$, where $p = \nabla V(\bar{x})$. In general it is hard to describe the effect of many minimization steps because the iteration $\bar{x} \rightarrow x'$ is nonlinear. Nevertheless, there is one case we can understand.
- (a) Show that for any V , gradient descent with exact line search has p_{k+1} orthogonal to p_k . Hint: otherwise, the step size t_k was not optimal.
 - (b) In the two dimensional quadratic optimization problem at hand, show that if p_k is in the direction of $(-1, -1)^*$, then p_{k+1} is in the direction of $(-1, 1)^*$.
 - (c) Show that p_k is in the direction of $(-1, -1)^*$ if and only if $(x_1, x_2) = r(\lambda_2, \lambda_1)$, for some r ,
 - (d) Since the optimum is $x_* = (0, 0)^*$, the error is $\|(x_1, x_2)\|$. Show that if p_0 is in the direction of $(-1, -1)$, then the error decreases exactly by a factor of $\rho = (\lambda_1 - \lambda_2)/(\lambda_1 + \lambda_2)$ if $\lambda_1 \geq \lambda_2$ (including the case $\lambda_1 = \lambda_2$).
 - (e) Show that if $\lambda_1 \gg \lambda_2$, then $\rho \approx 1 - 2\lambda_2/\lambda_1 = 1 - 2/\kappa(H)$, where $\kappa(H)$ is the linear systems condition number of H .
 - (f) Still supposing $\lambda_1 \gg \lambda_2$, show that it takes roughly $n = 1/\kappa(H)$ iterations to reduce the error by a factor of e^2 .
8. This exercise walks you through construction of a robust optimizer. It is as much an exercise in constructing scientific software as in optimization techniques. You will apply it to finding the minimum of the two variable function

$$V(x, y) = \frac{\psi(x, y)}{\sqrt{1 + \psi(x, y)^2}} , \quad \psi(x, y) = \psi_0 + wx^2 + (y - a \sin(x))^2 .$$

Hand in output documenting what you do for each of the parts below.

- (a) Write procedures that evaluate $V(x, y)$, $g(x, y)$, and $H(x, y)$ analytically. Write a tester that uses finite differences to verify that g and H are correct.
- (b) Implement a local Newton's method without safeguards as in Section 6.5. Use the Cholesky decomposition code from Exercise 3. Report failure if H is not positive definite. Include a stopping criterion and a maximum iteration count, neither hard wired. Verify local quadratic convergence starting from initial guess $(x_0, y_0) = (.3, .3)$ with parameters $\psi_0 = .5$, $w = .5$, and $a = .5$. Find an initial condition from which the unsafeguarded method fails.
- (c) Modify the Cholesky decomposition code Exercise 5.4 to do the modified Cholesky decomposition described in Section 6.6. This should require you to change a single line of code.
- (d) Write a procedure that implements the limited line search strategy described in Section 6.6. This also should have a maximum iteration count that is not hard wired. Write the procedure so that it sees only a scalar function $\phi(t)$. Test on:
 - i. $\phi(t) = (t - .9)^2$ (should succeed with $t = 1$).
 - ii. $\phi(t) = (t - .01)^2$ (should succeed after several step size reductions).
 - iii. $\phi(t) = (t - 100)^2$ (should succeed after several step size doublings).
 - iv. $\phi(t) = t$ (should fail after too many step size reductions).
 - v. $\phi(t) = -t$ (should fail after too many doublings).
- (e) Combine the procedures from parts (c) and (d) to create a robust global optimization code. Try the code on our test problem with $(x_0, y_0) = (10, 10)$ and parameters $\psi_0 = .5$, $w = .02$, and $a = 1$. Make plot that shows contour lines of V and all the iterates.

Chapter 7

Approximating Functions

Scientific computations often call for the approximate representation of a function $f(x)$. That is, we seek a \tilde{f} in a specified class of functions, such as polynomials, piecewise polynomials, rational functions, or trigonometric polynomials, so that¹ $\tilde{f} \approx f$ in some sense. For example, we might know the values $f_k = f(x_k)$ for some set of points x_0, x_1, \dots and wish to find an *interpolating function* \tilde{f} so that $\tilde{f}(x_k) = f_k$.

Interpolation by low order polynomials is simple and ubiquitous in scientific computing. We will begin by describing some simple tricks for constructing polynomial interpolants, and by describing Horner's method for evaluating polynomials. Ideas from Chapters 3 and 4 will let us understand the accuracy of the interpolation. The order of accuracy of polynomial interpolation is $p + 1$, where p is the degree of polynomials used. This suggests that we could get very accurate approximations using high degree polynomials. Unfortunately, high degree polynomial interpolation on uniformly spaced points leads to linear systems of equations that are exponentially ill conditioned, $\kappa \sim e^{cp}$, where p is the degree and κ is the condition number. The condition number grows moderately as $p \rightarrow \infty$ only if the interpolation points cluster at the ends of the interval in a very specific way.

High accuracy interpolation on uniformly spaced points can be done using trigonometric polynomial interpolation, also called Fourier interpolation. More generally, Fourier analysis for functions defined at n uniformly spaced points can be done using the discrete Fourier transform, or DFT. The fast Fourier transform, or FFT, is an algorithm that computes the DFT of n values in $O(n \log(n))$ time. Besides trigonometric interpolation, the FFT gives highly accurate solutions to certain linear partial differential equations and allows us to compute large discrete convolutions, including the convolution that defines the time lag covariance function for a time series.

7.1 Polynomial evaluation

Consider the degree d polynomial

$$p(x) = p_0 + p_1x + \dots + p_dx^d,$$

7.2 Polynomial interpolation

Given points x_0, \dots, x_d and values f_0, \dots, f_d , there is a unique *interpolating* polynomial of degree d ,

$$p(x) = p_0 + p_1x + \dots + p_dx^d, \tag{7.1}$$

so that

$$p(x_k) = f_k \quad \text{for } k = 0, 1, \dots, d. \tag{7.2}$$

¹ In Chapters 2 and 3 we used the *hat* or *caret* notation of statisticians to denote approximation: $\hat{Q} \approx Q$. In this chapter, the hat refers to Fourier coefficients and the tilde represents approximation.

We give three proofs of this, each one illustrating a different aspect of polynomial interpolation.

We distinguish between *low order* or *local* interpolation and *high order* or *global* interpolation. In low order interpolation, we have a small number (at least two and probably not much more than five) of nearby points and we seek an interpolating polynomial that should be valid near these points. This leads to approximations of order $d + 1$ for degree d interpolation. For example, local linear interpolation (degree $d = 1$) is second order accurate. See Exercise ??.

7.2.1 Vandermonde theory

The most direct approach to interpolation uses the Vandermonde matrix. The equations (7.22) and (7.2) form a set of linear equations that determine the $d + 1$ unknown coefficients, p_j , from the $d + 1$ given function values, f_k . The k^{th} equation is

$$p_0 + x_k p_1 + x_k^2 p_2 + \cdots + x_k^d p_d = f_k ,$$

which we write abstractly as

$$Vp = f , \quad (7.3)$$

where

$$V = \begin{pmatrix} 1 & x_0 & \cdot & \cdot & \cdot & x_0^d \\ 1 & x_1 & \cdot & \cdot & \cdot & x_1^d \\ \vdots & \vdots & & & & \vdots \\ 1 & x_d & \cdot & \cdot & \cdot & x_d^d \end{pmatrix} , \quad (7.4)$$

$p = (p_0, \dots, p_d)^*$, and $f = (f_0, \dots, f_d)^*$. The equations (7.3) have a unique solution if and only if $\det(V) \neq 0$. We show $\det(V) \neq 0$ using the following famous formula:

Theorem 2 Define $D(x_0, \dots, x_d) = \det(V)$ as in (7.4). Then

$$D(x_0, \dots, x_d) = \prod_{j < k} (x_k - x_j) . \quad (7.5)$$

The reader should verify directly that $D(x_0, x_1, x_2) = (x_2 - x_0)(x_1 - x_0)(x_2 - x_1)$. It is clear that $D = 0$ whenever $x_j = x_k$ for some $j \neq k$ because $x_j = x_k$ makes row j and row k equal to each other. The formula (7.5) says that D is a product of factors coming from these facts.

Proof: The proof uses three basic properties of determinants. The first is that the determinant does not change if we perform an elimination operation on rows or columns. If we subtract a multiple of row j from row k or of column j from column k , the determinant does not change. The second is that if row k or column k has a common factor, we can pull that factor out of the determinant. The third is that if the first column is $(1, 0, \dots, 0)^*$, then the determinant is the determinant of the $d \times d$ matrix got by deleting the top row and first column.

We work by induction on the number of points. For $d = 1$ (7.5) is $D(x_0, x_1) = x_1 - x_0$, which is easy to verify. The induction step is the formula

$$D(x_0, \dots, x_d) = \left(\prod_{k=1}^d (x_k - x_0) \right) \cdot D(x_1, \dots, x_d). \quad (7.6)$$

We use the easily checked formula

$$x^k - y^k = (x - y)(x^{k-1} + x^{k-2}y + \dots + y^{k-1}). \quad (7.7)$$

To compute the determinant of V in (7.4), we use Gauss elimination to set all but the top entry of the first column of V to zero. This means that we replace row j by row j minus row 1. Next we find common factors in the columns. Finally we perform column operations to put the $d \times d$ matrix back into the form of a Vandermonde matrix for x_1, \dots, x_d , which will prove (7.6).

Rather than giving the argument in general, we give it for $d = 2$ and $d = 3$. The general case will be clear from this. For $d = 2$ we have

$$\begin{aligned} \det \begin{pmatrix} 1 & x_0 & x_0^2 \\ 1 & x_1 & x_1^2 \\ 1 & x_2 & x_2^2 \end{pmatrix} &= \det \begin{pmatrix} 1 & x_0 & x_0^2 \\ 0 & x_1 - x_0 & x_1^2 - x_0^2 \\ 0 & x_2 - x_0 & x_2^2 - x_0^2 \end{pmatrix} \\ &= \det \begin{pmatrix} x_1 - x_0 & x_1^2 - x_0^2 \\ x_2 - x_0 & x_2^2 - x_0^2 \end{pmatrix}. \end{aligned}$$

The formula (7.7) with $k = 2$ gives $x_1^2 - x_0^2 = (x_1 - x_0)(x_1 + x_0)$, so $(x_1 - x_0)$ is a common factor in the top row. Similarly, $(x_2 - x_0)$ is a common factor of the bottom row. Thus:

$$\begin{aligned} \det \begin{pmatrix} x_1 - x_0 & x_1^2 - x_0^2 \\ x_2 - x_0 & x_2^2 - x_0^2 \end{pmatrix} &= \det \begin{pmatrix} x_1 - x_0 & (x_1 - x_0)(x_1 + x_0) \\ x_2 - x_0 & x_2^2 - x_0^2 \end{pmatrix} \\ &= (x_1 - x_0) \det \begin{pmatrix} 1 & (x_1 + x_0) \\ x_2 - x_0 & x_2^2 - x_0^2 \end{pmatrix} \\ &= (x_1 - x_0)(x_2 - x_0) \det \begin{pmatrix} 1 & x_1 + x_0 \\ 1 & x_2 + x_0 \end{pmatrix}. \end{aligned}$$

The final step is to subtract x_0 times the first column from the second column, which does not change the determinant:

$$\begin{aligned} \det \begin{pmatrix} 1 & x_1 + x_0 \\ 1 & x_2 + x_0 \end{pmatrix} &= \det \begin{pmatrix} 1 & x_1 + x_0 - x_0 * 1 \\ 1 & x_2 + x_0 - x_0 * 1 \end{pmatrix} \\ &= \det \begin{pmatrix} 1 & x_1 \\ 1 & x_2 \end{pmatrix} \\ &= D(x_1, x_2). \end{aligned}$$

This proves (7.6) for $d = 2$.

For $d = 3$ there is one more step. If we subtract row 1 from row k for $k > 1$ and do the factoring using (7.7) for $k = 2$ and $x = 3$, we get

$$\det \begin{pmatrix} 1 & x_0 & x_0^2 & x_0^3 \\ 1 & x_1 & x_1^2 & x_1^3 \\ 1 & x_2 & x_2^2 & x_2^3 \\ 1 & x_3 & x_3^2 & x_3^3 \end{pmatrix} = (x_1 - x_0)(x_2 - x_0)(x_3 - x_0) \det \begin{pmatrix} 1 & x_1 + x_0 & x_1^2 + x_1x_0 + x_0^2 \\ 1 & x_2 + x_0 & x_2^2 + x_2x_0 + x_0^2 \\ 1 & x_3 + x_0 & x_3^2 + x_3x_0 + x_0^2 \end{pmatrix}.$$

We complete the proof of (7.6) in this case by showing that

$$\det \begin{pmatrix} 1 & x_1 + x_0 & x_1^2 + x_1x_0 + x_0^2 \\ 1 & x_2 + x_0 & x_2^2 + x_2x_0 + x_0^2 \\ 1 & x_3 + x_0 & x_3^2 + x_3x_0 + x_0^2 \end{pmatrix} = \det \begin{pmatrix} 1 & x_1 & x_1^2 \\ 1 & x_2 & x_2^2 \\ 1 & x_3 & x_3^2 \end{pmatrix}.$$

For this we first subtract x_0 times the first column from the second column, then subtract x_0^2 times the first column from the third column, then subtract x_0 times the second column from the third column. This completes the proof of Theorem 2 and shows that one can find the coefficients of the interpolating polynomial by solving a linear system of equations involving the Vandermonde matrix.

7.2.2 Newton interpolation formula

The Newton interpolation formula is a simple and insightful way to express the interpolating polynomial. It is based on repeated divided differences, done in a way to expose the leading terms of polynomials. These are combined with a specific basis for the vector space of polynomials of degree k so that in the end the interpolation property is obvious. In some sense, the Newton interpolation formula provides a formula for the inverse of the Vandermonde matrix.

We begin with the problem of estimating derivatives of $f(x)$ using a number of function values. Given nearby points, x_1 and x_0 , we have

$$f' \approx \frac{f(x_1) - f(x_0)}{x_1 - x_0}.$$

We know that the divided difference is particularly close to the derivative at the center of the interval, so we write

$$f' \left(\frac{x_1 + x_0}{2} \right) \approx \frac{f(x_1) - f(x_0)}{x_1 - x_0}, \quad (7.8)$$

with an error that is $O(|x_1 - x_0|^2)$. If we have three points that might not be

uniformly spaced, the second derivative estimate using (7.8) could be

$$\begin{aligned} f'' &\approx \frac{\frac{f'(\frac{x_2+x_1}{2}) - f'(\frac{x_1+x_0}{2})}{\frac{x_2+x_1}{2} - \frac{x_1+x_0}{2}}}{\frac{x_2-x_1}{2} - \frac{x_1-x_0}{2}} \\ f'' &\approx \frac{\frac{f(x_2) - f(x_1)}{x_2 - x_1} - \frac{f(x_1) - f(x_0)}{x_1 - x_0}}{\frac{1}{2}(x_2 - x_0)} . \end{aligned} \quad (7.9)$$

As we saw in Chapter 3, the approximation (7.9) is consistent (converges to the exact answer as $x_2 \rightarrow x$, $x_1 \rightarrow x$, and $x_0 \rightarrow x$) if it is exact for quadratics, $f(x) = ax^2 + bx + c$. Some algebra shows that both sides of (7.9) are equal to $2a$.

The formula (7.9) suggests the right way to do repeated divided differences. Suppose we have $d+1$ points² x_0, \dots, x_d , we define $f[x_k] = f(x_k)$ (exchange round parentheses for square brackets), and the first order divided difference is:

$$f[x_k, x_{k+1}] = \frac{f[x_{k+1}] - f[x_k]}{x_{k+1} - x_k} .$$

More generally, the *Newton divided difference* of order $k+1$ is a divided difference of divided differences of order k :

$$f[x_j, \dots, x_{k+1}] = \frac{f[x_{j+1}, \dots, x_{k+1}] - f[x_j, \dots, x_k]}{x_{k+1} - x_j} . \quad (7.10)$$

The denominator in (7.10) is the difference between the extremal x values, as (7.9) suggests it should be. If instead of a function $f(x)$ we just have values f_0, \dots, f_d , we define

$$[f_j, \dots, f_{k+1}] = \frac{[f_{j+1}, \dots, f_{k+1}] - [f_j, \dots, f_k]}{x_{k+1} - x_j} . \quad (7.11)$$

It may be convenient to use the alternative notation

$$D_k(f) = f[x_0, \dots, x_k] .$$

If $r(x) = r_k x^k + \dots + r_0$ is a polynomial of degree k , we will see that $D_k r = k! \cdot r_k$. We verified this already for $k=1$ and $k=2$.

The interpolation problem is to find a polynomial of degree d that satisfies the interpolation conditions (7.2). The formula (7.1) expresses the interpolating polynomial as a linear combination of pure monomials x^k . Using the monomials as a basis for the vector space of polynomials of degree d leads to the Vandermonde matrix (7.4). Here we use a different basis, which might be called the *Newton monomials* of degree k (although they strictly speaking are not monomials), $q_0(x) = 1$, $q_1(x) = x - x_0$, $q_2(x) = (x - x_1)(x - x_0)$, and generally,

$$q_k(x) = (x - x_{k-1}) \cdots (x - x_0) . \quad (7.12)$$

²The x_k must be distinct but they need not be in order. Nevertheless, it helps the intuition to think that $x_0 < x_1 < \dots < x_d$.

It is easy to see that $q_k(x)$ is a polynomial of degree k in x with leading coefficient equal to one:

$$q_k(x) = x^k + a_{k-1}x^{k-1} + \cdots .$$

Since this also holds for q_{k-1} , we may subtract to get:

$$q_k(x) - a_{k-1}q_{k-1}(x) = x^k + b_{k-2}x^{k-2} + \cdots .$$

Continuing in this way, we express x^k in terms of Newton monomials:

$$x^k = q_k(x) - a_{k,k-1}q_{k-1}(x) - b_{k,k-2} - \cdots . \quad (7.13)$$

This shows that the $q_k(x)$ are linearly independent and span the same space as the monomial basis.

The connection between repeated divided differences (7.10) and Newton monomials (7.12) is

$$D_k q_j = \delta_{kj} . \quad (7.14)$$

The intuition is that $D_k f$ plays the role $\frac{1}{k!} \partial_x^k f(0)$ and $q_j(x)$ plays the role of x^j . For $k > j$, $\partial_k x^j = 0$ because differentiation lowers the order of a monomial. For $k < j$, $\partial_x^k x^j = 0$ when evaluated at $x = 0$ because monomials vanish when $x = 0$. The remaining case is the interesting one, $\frac{1}{k!} \partial_x^k x^k = 1$.

We verify (7.14) by induction on k . We suppose that (7.14) holds for all $k < d$ and all j and use that to prove it for $k = d$ and all j , treating the cases $j = d$, $j < d$, and $j > d$ separately. The base case $k = 1$ explains the ideas. For $j = 1$ we have

$$D_1 q_1(x) = \frac{q_1(x_1) - q_1(x_0)}{x_1 - x_0} = \frac{(x_1 - x_0) - (x_0 - x_0)}{x_1 - x_0} = 1 , \quad (7.15)$$

as claimed. More generally, any first order divided difference of q_1 is equal to one,

$$q_1[x_{k+1}, x_k] = \frac{q_1(x_{k+1}) - q_1(x_k)}{x_{k+1} - x_k} = 1 ,$$

which implies that higher order divided differences of q_1 are zero. For example,

$$q_1[x_2, x_3, x_4] = \frac{q_1[x_3, x_4] - q_1[x_2, x_3]}{x_4 - x_2} = \frac{1 - 1}{x_4 - x_2} = 0 .$$

This proves the base case, $k = 1$ and all j .

The induction step has the same three cases. For $j > d$ it is clear that $D_d q_j = q_j[x_0, \dots, x_d] = 0$ because $q_j(x_k) = 0$ for all the x_k that are used in $q_j[x_0, \dots, x_d]$. The interesting case is $q_k[x_0, \dots, x_k] = 1$. From (7.10) we have that

$$q_k[x_0, \dots, x_k] = \frac{q_k[x_1, \dots, x_k] - q_k[x_0, \dots, x_{k-1}]}{x_k - x_0} = \frac{q_k[x_1, \dots, x_k]}{x_k - x_0} ,$$

because $q_k[x_0, \dots, x_{k-1}] = 0$ (it involves all zeros). The same reasoning gives $q_k[x_1, \dots, x_{k-1}] = 0$ and

$$q_k[x_1, \dots, x_k] = \frac{q_k[x_2, \dots, x_k] - q_k[x_1, \dots, x_{k-1}]}{x_k - x_1} = \frac{q_k[x_2, \dots, x_k]}{x_k - x_1}.$$

Combining these gives

$$q_k[x_0, \dots, x_k] = \frac{q_k[x_1, \dots, x_k]}{x_k - x_0} = \frac{q_k[x_2, \dots, x_k]}{(x_k - x_0)(x_k - x_1)},$$

and eventually, using the definition (7.12), to

$$q_k[x_0, \dots, x_k] = \frac{q_k(x_k)}{(x_k - x_0) \cdots (x_k - x_{k-1})} = \frac{(x_k - x_0) \cdots (x_k - x_{k-1})}{(x_k - x_0) \cdots (x_k - x_{k-1})} = 1,$$

as claimed. Now, using (7.13) we see that

$$D_k x^k = D_k(q_k - a_{k,k-1}q_{k-1} - \cdots) = D_k q_k = 1,$$

for any collection of distinct points x_j . This, in turn, implies that

$$q_k[x_{m+k}, \dots, x_m] = 1.$$

which, as in (7.15), implies that $D_{k+1}q_k = 0$. This completes the induction step.

The formula (7.14) allows us to verify the *Newton interpolation formula*, which states that

$$p(x) = \sum_{k=0}^d [f_0, \dots, f_k] q_k(x), \quad (7.16)$$

satisfies the interpolation conditions (7.2). We see that $p(x_0) = f_0$ because each term on the right $k = 0$ vanishes when $x = x_0$. The formula (7.14) also implies that $D_1 p = D_1 f$. This involves the values $p(x_1)$ and $p(x_0)$. Since we already know $p(x_0)$ is correct, this implies that $p(x_1)$ also is correct. Continuing in this way verifies all the interpolation conditions.

7.2.3 Lagrange interpolation formula

The Lagrange approach to polynomial interpolation is simpler mathematically but less useful than the others. For each k , define the polynomial³ of degree d

$$l_k(x) = \frac{\prod_{j \neq k} (x - x_j)}{\prod_{j \neq k} (x_k - x_j)}. \quad (7.17)$$

³We do not call these *Lagrange polynomials* because that term means something else.

For example, for $d = 2$, $x_0 = 0$, $x_1 = 2$, $x_2 = 3$ we have

$$\begin{aligned} l_0(x) &= \frac{(x - x_1)(x - x_2)}{(x_0 - x_1)(x_0 - x_2)} = \frac{(x - 2)(x - 3)}{(-2)(-3)} = \frac{1}{6}(x^2 - 5x + 6) \\ l_1(x) &= \frac{(x - x_0)(x - x_2)}{(x_1 - x_0)(x_1 - x_2)} = \frac{(x - 0)(x - 3)}{(1)(-1)} = x^2 - 3x \\ l_2(x) &= \frac{(x - x_0)(x - x_1)}{(x_2 - x_0)(x_2 - x_1)} = \frac{(x - 0)(x - 2)}{(3)(1)} = \frac{1}{3}(x^2 - 2x) , \end{aligned}$$

If $j = k$, the numerator and denominator in (7.17) are equal. If $j \neq k$, then $l_k(x_j) = 0$ because $(x_j - x_j) = 0$ is one of the factors in the numerator. Therefore

$$l_k(x_j) = \delta_{jk} = \begin{cases} 1 & \text{if } j = k \\ 0 & \text{if } j \neq k \end{cases} \quad (7.18)$$

The *Lagrange interpolation formula* is

$$p(x) = \sum_{k=0}^d f_k l_k(x) . \quad (7.19)$$

The right side is a polynomial of degree d . This satisfies the interpolation conditions (7.2) because of (7.18).

7.3 Discrete Fourier transform

The Fourier transform is one of the most powerful methods of applied mathematics. Its finite dimensional analogue, the *discrete Fourier transform*, or *DFT*, is just as useful in scientific computing. The DFT allows direct algebraic solution of certain differential and integral equations. It is the basis of computations with time series data and for digital signal processing and control. It leads to computational methods that have an infinite order of accuracy (which is not the same as being exact).

The drawback of DFT based methods is their geometric inflexibility. They can be hard to apply to data that are not sampled at uniformly spaced points. The multidimensional DFT is essentially a product of one dimensional DFTs. Therefore is it hard to apply DFT methods to problems in more than one dimension unless the computational domain has a simple shape. Even in one dimension, applying the DFT depends on boundary conditions.

7.3.1 Fourier modes

The simplest Fourier analysis is for periodic functions of one variable. We say $f(x)$ is *periodic* with *period* p if $f(x + p) = f(x)$ for all x . If α is an integer, then the *Fourier mode*

$$w_\alpha(x) = e^{2\pi i \alpha x / p} \quad (7.20)$$

is such a periodic function. Fourier analysis starts with the fact that Fourier modes form a basis for the vector space of periodic functions. That means that if f has period p , then there are *Fourier coefficients*, \hat{f}_α , so that

$$f(x) = \sum_{\alpha=-\infty}^{\infty} \hat{f}_\alpha e^{2\pi i \alpha x/p} . \quad (7.21)$$

More precisely, let V_p be the vector space of complex valued periodic functions so that

$$\|f\|_{L^2}^2 = \int_0^p |f(x)|^2 dx < \infty .$$

This vector space has an *inner product* that has the properties discussed in Section 4.2.2. The definition is

$$\langle f, g \rangle = \int_0^p \bar{f}(x) g(x) dx . \quad (7.22)$$

Clearly, $\|f\|_{L^2}^2 = \langle f, f \rangle > 0$ unless $f \equiv 0$. The inner product is linear in the g variable:

$$\langle f, ag_1 + bg_2 \rangle = a\langle f, g_1 \rangle + b\langle f, g_2 \rangle ,$$

and *antilinear* in the f variable:

$$\langle af_1 + bf_2, g \rangle = \bar{a}\langle f_1, g \rangle + \bar{b}\langle f_2, g \rangle .$$

Functions f and g are *orthogonal* if $\langle f, g \rangle = 0$. A set of functions is orthogonal if any two of them are orthogonal and none of them is zero. The Fourier modes (7.20) have this property: if $\alpha \neq \beta$ are two integers, then $\langle w_\alpha, w_\beta \rangle = 0$ (check this).

Any orthogonal system of functions is linearly independent. Linear independence means that if

$$f(x) = \sum_{\alpha} \hat{f}_\alpha w_\alpha(x) = \sum_{\alpha=-\infty}^{\infty} \hat{f}_\alpha e^{2\pi i \alpha x/p} , \quad (7.23)$$

then the \hat{f}_α are uniquely determined. For orthogonal functions, we show this by taking the inner product of w_β with f and use the linearity of the inner product to get

$$\langle w_\beta, f \rangle = \sum_{\alpha} \hat{f}_\alpha \langle w_\beta, w_\alpha \rangle = \hat{f}_\beta \langle w_\beta, w_\beta \rangle ,$$

so

$$\hat{f}_\beta = \frac{\langle w_\beta, f \rangle}{\|w_\beta\|^2} . \quad (7.24)$$

This formula shows that the coefficients are uniquely determined. Written more explicitly for the Fourier modes, (7.24) is

$$\hat{f}_\alpha = \frac{1}{p} \int_{x=0}^p e^{-2\pi i \alpha x/p} f(x) dx . \quad (7.25)$$

An orthogonal family, w_α , is *complete* if any f has a representation in terms of them as a linear combination (7.23). An orthogonal family need not be complete. Fourier conjectured that the Fourier modes (7.20) are complete, but this was first proven several decades later.

Much of the usefulness of Fourier series comes from the relationship between the series for f and for derivatives f' , f'' , etc. When we differentiate (7.23) with respect to x and differentiate under the summation on the right side, we get

$$f'(x) = \frac{2\pi i}{p} \sum_{\alpha} \alpha \hat{f}_\alpha e^{2\pi i \alpha x/p}.$$

This shows that the α Fourier coefficient of f' is

$$\hat{f}'_\alpha = \frac{2\pi i \alpha}{p} \hat{f}_\alpha. \quad (7.26)$$

Formulas like these allow us to express the solutions to certain ordinary and partial differential equations in terms of Fourier series. See Exercise 4 for one example.

We will see that the differentiation formula (7.26) also contains important information about the Fourier coefficients of smooth functions, that they are very small for large α . This implies that approximations

$$f(x) \approx \sum_{|\alpha| \leq R} \hat{f}_\alpha w_\alpha(x)$$

are very accurate if f is smooth. It also implies that the DFT coefficients (see Section 7.3.2) are very accurate approximations of \hat{f}_α . Both of these make DFT based computational methods very attractive (when they apply) for application to problems with smooth solutions.

We start to see this decay by rewriting (7.26) as

$$\hat{f}_\alpha = \frac{p \hat{f}'_\alpha}{2\pi i} \cdot \frac{1}{\alpha}. \quad (7.27)$$

The integral formula

$$\hat{f}'_\alpha = \frac{1}{p} \int_0^p \overline{w_\alpha(x)} f'(x) dx$$

shows that the Fourier coefficients \hat{f}'_α are bounded if f' is bounded, since (for some real θ) $|w_\alpha(x)| = |e^{i\theta}| = 1$. This, and (7.27) shows that

$$|\hat{f}_\alpha| \leq C \cdot \frac{1}{|\alpha|}.$$

We can go further by applying (7.27) to f' and f'' to get

$$\hat{f}_\alpha = \frac{p^2 \hat{f}''_\alpha}{-4\pi^2} \cdot \frac{1}{\alpha^2},$$

so that if f'' is bounded, then

$$|\widehat{f}_\alpha| \leq C \cdot \frac{1}{|\alpha^2|},$$

which is faster decay ($1/\alpha^2 \ll 1/\alpha$ for large α). Continuing in this way, we can see that if f has N bounded derivatives then

$$|\widehat{f}_\alpha| \leq C_N \cdot \frac{1}{|\alpha^N|}. \quad (7.28)$$

This shows, as we said, that the Fourier coefficients of smooth functions decay rapidly.

It is helpful in real applications to use real Fourier modes, particularly when $f(x)$ is real. The real modes are sines and cosines:

$$u_\alpha(x) = \cos(2\pi\alpha x/p), \quad v_\alpha(x) = \sin(2\pi\alpha x/p). \quad (7.29)$$

The u_α are defined for $\alpha \geq 0$ and the v_α for $\alpha \geq 1$. The special value $\alpha = 0$ corresponds to $u_0(x) = 1$ (and $v_0(x) = 0$). Exercise 3 shows that the u_α for $\alpha = 0, 1, \dots$ and v_α for $\alpha = 1, 2, \dots$ form an orthogonal family. The real Fourier series representation expresses f as a superposition of these functions:

$$f(x) = \sum_{\alpha=0}^{\infty} a_\alpha \cos(2\pi\alpha x/p) + \sum_{\alpha=1}^{\infty} b_\alpha \sin(2\pi\alpha x/p). \quad (7.30)$$

The reasoning that led to (7.25), and the normalization formulas of Exercise ?? below, (7.41), gives

$$\left. \begin{aligned} a_\alpha &= \frac{2}{p} \int_0^p \cos(2\pi\alpha x/p) f(x) dx \quad (\alpha \geq 1), \\ b_\alpha &= \frac{2}{p} \int_0^p \sin(2\pi\alpha x/p) f(x) dx \quad (\alpha \geq 1), \\ a_0 &= \frac{1}{p} \int_0^p f(x) dx. \end{aligned} \right\} \quad (7.31)$$

This real Fourier series (7.30) is basically the same as the complex Fourier series (7.23) when $f(x)$ is real. If f is real, then (7.25) shows that⁴ $\widehat{f}_{-\alpha} = \overline{\widehat{f}_\alpha}$. This determines the $\alpha < 0$ Fourier coefficients from the $\alpha \geq 0$ coefficients. If $\alpha > 0$ and $\widehat{f}_\alpha = g_\alpha + ih_\alpha$ (g_α and h_α being real and imaginary parts), then (using $e^{i\theta} = \cos(\theta) + i\sin(\theta)$), we have

$$\begin{aligned} & \widehat{f}_\alpha e^{2\pi i\alpha x/p} + \widehat{f}_{-\alpha} e^{-2\pi i\alpha x/p} \\ &= (g_\alpha + ih_\alpha)(\cos(\cdot) + i\sin(\cdot)) + (g_\alpha - ih_\alpha)(\cos(\cdot) - i\sin(\cdot)) \\ &= 2g_\alpha \cos(\cdot) - 2h_\alpha \sin(\cdot). \end{aligned}$$

⁴This also shows that the full Fourier series sum over positive and negative α is somewhat redundant for real f . This is another motivation for using the version with cosines and sines.

Moreover, when f is real,

$$\begin{aligned} g_\alpha &= \frac{1}{p} \operatorname{Re} \left[\int_0^p e^{-2\pi i \alpha x/p} f(x) dx \right] \\ &= \frac{1}{p} \int_0^p \cos(2\pi i \alpha x/p) f(x) dx \\ &= \frac{1}{2} a_\alpha . \end{aligned}$$

Similarly, $h_\alpha = \frac{-1}{2} b_\alpha$. This shows that the real Fourier series relations (7.23) and (7.25) directly follow from the complex Fourier series relations (7.30) and (7.31) without using Exercise 3.

7.3.2 The DFT

The DFT is a discrete analogue of Fourier analysis. The vector space V_p is replaced by sequences with period n : $f_{j+n} = f_j$. A periodic sequence is determined by the n entries⁵: $f = (f_0, \dots, f_{n-1})^*$, and the vector space of such sequences is C^n . *Sampling* a periodic function of x is one way to create an element of C^n . Choose $\Delta x = p/n$, take *sample points*, $x_j = j\Delta x$, then the samples are $f_j = f(x_j)$. If the continuous f has period p , then the discrete sampled f has period n , because $f_{j+n} = f(x_{j+n})$, and $x_{j+n} = (j+n)\Delta x$, and $n\Delta x = p$.

The DFT modes come from sampling the continuous Fourier modes (7.20) in this way. That is

$$w_{\alpha,j} = w_\alpha(x_j) = \exp(2\pi i \alpha x_j/p) .$$

Since $x_j = jp/n$, this gives

$$w_{\alpha,j} = \exp(2\pi i \alpha j/n) = w^{\alpha j} , \quad (7.32)$$

where w is a *primitive root of unity*⁶

$$w = e^{2\pi i/n} . \quad (7.33)$$

Aliasing is a difference between continuous and discrete Fourier analysis. If $\beta = \alpha + n$ then the samplings of w_β and w_α are the same, $w_\beta(x_j) = w_\alpha(x_j)$ for all j , even though w_β and w_α are different functions of x . Figure 7.1 illustrates this with sine functions instead of complex exponentials. Aliasing implies that the discrete sampled vectors $w_\alpha \in C^n$ with components given by (7.32) are not all different. In fact, the n vectors w_α for $0 \leq \alpha < n$ are distinct. After that they repeat: $w_{\alpha+n} = w_\alpha$ (this being an equality between two vectors in C^n).

These n discrete sampled modes form a basis, the DFT basis, for C^n . If f and g are two elements of C^n , the discrete inner product is

$$\langle f, g \rangle = \sum_{j=0}^{n-1} \bar{f}_j g_j .$$

⁵The $*$ in $(f_0, \dots, f_{n-1})^*$ indicates that we think of f as a column vector in C^n .

⁶*Unity* means the number 1. An n^{th} root of x is a y with $y^n = x$. An n^{th} root of unity is *primitive* if $w^n = 1$ but $w^k \neq 1$ for $0 \leq k < n$.

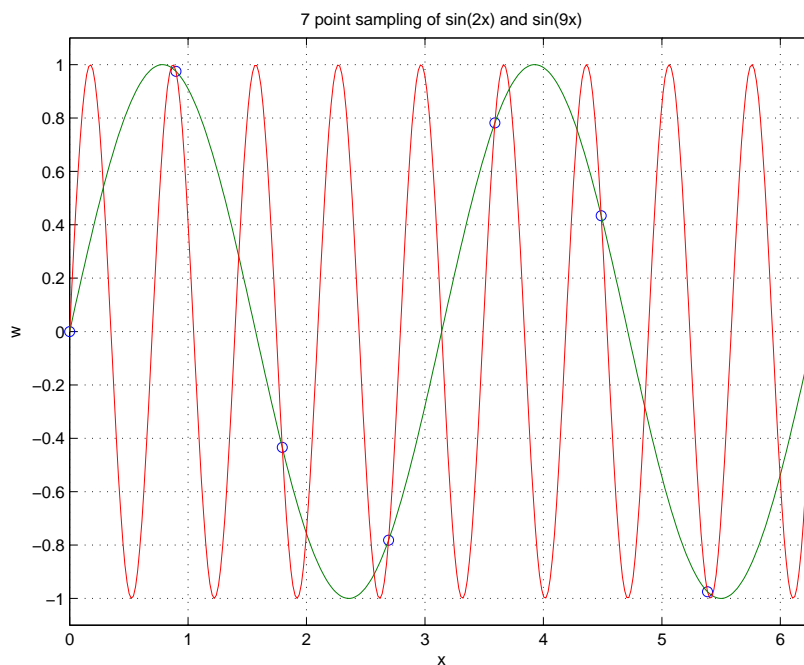


Figure 7.1: An illustration of aliasing with $n = 7$ sampling points and period $p = 2\pi$, with modes $\alpha = 2$ and $\beta = \alpha + n = 9$. The continuous curves are $\sin(2x)$ and $\sin(9x)$. The circles are these curves sampled at the sample points $x_j = 2\pi j/n$. The sampled values are identical even though the functions $\sin(2x)$ and $\sin(9x)$ are not.

This is the usual inner product on C^n and leads to the usual l^2 norm $\langle f, f \rangle = \|f\|_{l^2}^2$. We show that the discrete modes form an orthogonal family, but only as far as is allowed by aliasing. That is, if $0 \leq \alpha < n$ and $0 \leq \beta < n$, and $\alpha \neq \beta$, then $\langle w_\alpha, w_\beta \rangle = 0$.

Recall that for any complex number, z , $S(z) = \sum_{j=0}^{n-1} z^j$ has $S = n$ if $z = 1$ and $S = (z^n - 1)/(z - 1)$ if $z \neq 1$. Also,

$$\overline{w_{\alpha,j}} = \overline{e^{2\pi i \alpha j / n}} = e^{-2\pi i \alpha j / n} = w^{-\alpha j} ,$$

so we can calculate

$$\begin{aligned} \langle w_\alpha, w_\beta \rangle &= \sum_{j=0}^{n-1} w^{-\alpha j} w^{\beta j} \\ &= \sum_{j=0}^{n-1} w^{(\beta-\alpha)j} \\ &= \sum_{j=0}^{n-1} (w^{\beta-\alpha})^j \end{aligned}$$

Under our assumptions ($\alpha \neq \beta$ but $0 \leq \alpha < n$, $0 \leq \beta < n$) we have $0 < |\beta - \alpha| < n$, and $z = w^{\beta-\alpha} \neq 1$ (using the fact that w is a *primitive* root of unity). This gives

$$\langle w_\alpha, w_\beta \rangle = \frac{w^{n(\beta-\alpha)} - 1}{w^{\beta-\alpha} - 1} .$$

Also,

$$w^{n(\beta-\alpha)} = (w^n)^{\beta-\alpha} = 1^{\beta-\alpha} = 1 ,$$

because w is an n^{th} root of unity. This shows $\langle w_\alpha, w_\beta \rangle = 0$. We also can calculate that $\|w_\alpha\|^2 = \sum_{j=1}^n |w_{\alpha,j}|^2 = n$.

Since the n vectors w_α for $0 \leq \alpha < n$ are linearly independent, they form a basis of C^n . That is, any $f \in C^n$ may be written as a linear combination of the w_α :

$$f = \sum_{\alpha=0}^{n-1} \hat{f}_\alpha w_\alpha .$$

By the arguments we gave for Fourier series above, the DFT coefficients are

$$\hat{f}_\alpha = \frac{1}{n} \langle w_\alpha, f \rangle .$$

Expressed explicitly in terms of sums, these relations are

$$\hat{f}_\alpha = \frac{1}{n} \sum_{j=0}^{n-1} w^{-\alpha j} f_j , \quad (7.34)$$

and

$$f_j = \sum_{\alpha=0}^{n-1} \hat{f}_\alpha w^{\alpha j} . \quad (7.35)$$

These are the (*forward*) DFT and *inverse* DFT respectively. Either formula implies the other. If we start with the f_j and calculate the \hat{f}_α using (7.34), then we can use (7.35) to recover the f_j from the \hat{f}_α , and the other way around. These formulas are more similar to each other than are the corresponding formulas (7.25) and (7.21). Both are sums rather than one being a sum and the other an integral.

There are many other slightly different definitions of the DFT relations. Some people define the discrete Fourier coefficients using a variant of (7.34), such as

$$\hat{f}_\alpha = \sum_{j=0}^{n-1} w^{\alpha j} f_j .$$

This particular version changes (7.35) to

$$f_j = \frac{1}{n} \sum_{\alpha=0}^{n-1} w^{-\alpha j} \hat{f}_\alpha .$$

Still another way to express these relations is to use the *DFT matrix*, W , which is an orthogonal matrix whose (α, j) entry is

$$w_{\alpha,j} = \frac{1}{\sqrt{n}} w^{-\alpha j} .$$

The adjoint of W has entries

$$w_{j,\alpha}^* = \overline{w}_{\alpha,j} = \frac{1}{\sqrt{n}} w^{\alpha j} .$$

The DFT relations are equivalent to $W^*W = I$. In vector notation, this $\hat{f} = Wf$ transform differs from (7.34) by a factor of \sqrt{n} :

$$\hat{f}_\alpha = \frac{1}{\sqrt{n}} \sum_{j=0}^{n-1} w^{-\alpha j} f_j .$$

It has the advantage of making the direct and inverse DFT as similar as possible, with a factor $1/\sqrt{n}$ in both.

As for continuous Fourier series, there is a real discrete cosine and sine transform for real f . The complex coefficients \hat{f}_α , determine the real coefficients a_α and b_α as before. Aliasing is more complicated for the discrete sine and cosine transform, and depends on whether n is even or odd. The cosine and sine sums corresponding to (7.30) run from $\alpha = 0$ or $\alpha = 1$ roughly to $\alpha = n/2$.

We can estimate the Fourier coefficients of a continuous function by taking the DFT of its sampled values. We call the vector of samples $f^{(n)}$. It has

components $f_j^{(n)} = f(x_j)$, where $x_j = j\Delta x$ and $\Delta x = p/n$. The DFT coefficients $\widehat{f}_\alpha^{(n)}$ defined by (7.34) are rectangle rule approximations to the Fourier series coefficients (7.25). This follows from (7.32) and $\frac{1}{p}\Delta x = \frac{1}{n}$, since $\alpha j/n = \alpha x_j/p$ and

$$\begin{aligned}\widehat{f}_\alpha^{(n)} &= \frac{1}{n} \sum_{j=0}^{n-1} w^{-\alpha j} f_j^{(n)} \\ &= \frac{1}{p} \Delta x \sum_{j=0}^{n-1} e^{-2\pi i \alpha x_j/n} f(x_j) .\end{aligned}\tag{7.36}$$

There is a simple *aliasing formula* for the Fourier coefficients of $f^{(n)}$ in terms of those of f . It depends on aliasing when we put the continuous Fourier representation (7.25) into the discrete Fourier coefficient formula (7.37). If we rewrite (7.25) with β instead of α as the summation index, substitute into (7.37), and change the order of summation, we find

$$\widehat{f}_\alpha^{(n)} = \sum_{\beta=-\infty}^{\infty} \widehat{f}_\beta \frac{\Delta x}{p} \sum_{j=0}^{n-1} \exp[2\pi i (\beta - \alpha) x_j/n] .$$

We have shown that the inner sum is equal to zero unless mode β aliases to mode α on the grid, which means $\beta = \alpha + kn$ for some integer k . We also showed that if mode β does alias to mode α on the grid, then each of the summands is equal to one. Since $\Delta x/p = 1/n$, this implies that

$$\widehat{f}_\alpha^{(n)} = \sum_{k=-\infty}^{\infty} \widehat{f}_{\alpha+kn} .\tag{7.37}$$

This shows that the discrete Fourier coefficient is equal to the continuous Fourier coefficient (the $k = 0$ term on the right) plus all the coefficients that alias to it (the terms $k \neq 0$).

You might worry that the continuous function has Fourier coefficients with both positive and negative α while the DFT computes coefficients for $\alpha = 0, 1, \dots, n-1$. The answer to this is aliasing. We find approximations to the negative α Fourier coefficients using $\widehat{f}_{-\alpha}^{(n)} = \widehat{f}_{n-\alpha}^{(n)}$. It may be more helpful to think of the DFT coefficients as being defined for $\alpha \approx -\frac{n}{2}$ to $\alpha \approx \frac{n}{2}$ (the exact range depending on whether n is even or odd) rather than from $\alpha = 0$ to $\alpha = n-1$. The aliasing formula shows that if f is smooth, then $\widehat{f}_\alpha^{(n)}$ is a very good approximation to \widehat{f}_α .

7.3.3 FFT algorithm

It takes n^2 multiplies to carry out all the sums in (7.34) directly (n terms in each of n sums). The *Fast Fourier Transform*, or *FFT*, is an algorithm that

calculates the n components of \hat{f} from the n components of f using $O(n \log(n))$ operations, which is much less for large n .

The idea behind FFT algorithms is clearest when $n = 2m$. A single DFT of size $n = 2m$ is reduced to two DFT calculations of size $m = n/2$ followed by $O(n)$ work. If $m = 2r$, this process can be continued to reduce the two size m DFT operations to four size r DFT operations followed by $2 \cdot O(m) = O(n)$ operations. If $n = 2^p$, this process can be continued p times, where we arrive at $2^p = n$ trivial DFT calculations of size one each. The total work for the p levels is $p \cdot O(n) = O(n \log_2(n))$.

There is a variety of related methods to handle cases where n is not a power of 2, and the $O(n \log(n))$ work count holds for all n . The algorithm is simpler and faster for $n = 2^p$. For example, an FFT with $n = 2^{20} = 1,048,576$ should be significantly faster than with $n = 1,048,573$, which is a prime number.

Let $W_{n \times n}$ be the complex $n \times n$ matrix⁷ whose (α, j) entry is $w_{\alpha,j} = w^{-\alpha j}$, where w is a primitive n^{th} root of unity. The DFT (7.34), but for the factor of $\frac{1}{n}$, is the matrix product $\tilde{f} = W_{n \times n} f$. If $n = 2m$, then w^2 is a primitive m^{th} root of unity and an $m \times m$ DFT involves the matrix product $\tilde{g} = W_{m \times m} g$, where the (α, k) entry of $W_{m \times m}$ is $(w^2)^{-\alpha k} = w^{-2\alpha k}$. The reduction splits $f \in C^n$ into $g \in C^m$ and $h \in C^m$, then computes $\tilde{g} = W_{m \times m} g$ and $\tilde{h} = W_{m \times m} h$, then combines \tilde{g} and \tilde{h} to form \tilde{f} .

The elements of \tilde{f} are given by the sums

$$\tilde{f}_\alpha = \sum_{j=0}^{n-1} w^{-\alpha j} f_j .$$

We split these into even and odd parts, with $j = 2k$ and $j = 2k+1$ respectively. Both of these have k ranging from 0 to $\frac{n}{2} - 1 = m - 1$. For these sums, $-\alpha j = -\alpha(2k) = -2\alpha k$ (even), and $-\alpha j = -\alpha(2k+1) = -2\alpha k - \alpha$ (odd) respectively. Thus

$$\tilde{f}_\alpha = \sum_{k=0}^{m-1} w^{-2\alpha k} f_{2k} + w^{-\alpha} \sum_{k=0}^{m-1} w^{-2\alpha k} f_{2k+1} . \quad (7.38)$$

Now define $g \in C^m$ and $h \in C^m$ to have the even and odd components of f respectively:

$$g_k = f_{2k} , \quad h_k = f_{2k+1} .$$

The $m \times m$ operations $\tilde{g} = W_{m \times m} g$ and $\tilde{h} = W_{m \times m} h$, written out, are

$$\tilde{g}_\alpha = \sum_{k=0}^{m-1} (w^2)^{-\alpha k} g_k ,$$

and

$$\tilde{h}_\alpha = \sum_{k=0}^{m-1} (w^2)^{-\alpha k} h_k .$$

⁷This definition of W differs from that of Section 7.3 by a factor of \sqrt{n} .

Then (7.38) may be written

$$\tilde{f}_\alpha = \tilde{g}_\alpha + w^{-\alpha} \tilde{h}_\alpha . \quad (7.39)$$

This is the last step, which reassembles \tilde{f} from \tilde{g} and \tilde{h} . We must apply (7.39) for n values of α ranging from $\alpha = 0$ to $\alpha = n - 1$. The computed \tilde{g} and \tilde{h} have period m ($\tilde{g}_{\alpha+m} = \tilde{g}_\alpha$, etc.), but the factor $w^{-\alpha}$ in front of \tilde{h} makes \tilde{f} have period $n = 2m$ instead.

To summarize, an order n FFT requires first order n copying to form g and h , then two order $n/2$ FFT operations, then order n copying, adding, and multiplying. Of course, the order $n/2$ FFT operations themselves boil down to copying and simple arithmetic. As explained in Section 5.6, the copying and memory accessing can take more computer time than the arithmetic. High performance FFT software needs to be chip specific to take full advantage of cache.

7.3.4 Trigonometric interpolation

Interpolating $f(x)$ at points x_j , $j = 1, \dots, n$, means finding another function $F(x)$ so that $F(x_j) = f(x_j)$ for $j = 1, \dots, n$. In polynomial interpolation, $F(x)$ is a polynomial of degree $n - 1$. In *trigonometric interpolation*, $F(x)$ is a *trigonometric polynomial* with n terms. Because we are interested in values of $F(x)$ off the grid, do should not take advantage of aliasing to simplify notation. Instead, we let Z_n be a set of integers as symmetric about $\alpha = 0$ as possible. This depends on whether n is even or odd

$$Z_n = \begin{cases} \{-m, -m+1, \dots, m\} & \text{if } n = 2m+1 \text{ (i.e. } n \text{ is odd)} \\ \{-m+1, \dots, m\} & \text{if } n = 2m \text{ (i.e. } n \text{ is even)} \end{cases} \quad (7.40)$$

With this notation, an n term trigonometric polynomial may be written

$$F(x) = \sum_{\alpha \in Z_n} c_\alpha e^{2\pi i \alpha x / p} .$$

The DFT provides coefficients c_α of the trigonometric interpolating polynomial at n uniformly spaced points.

The high accuracy of DFT based methods comes from the fact that trigonometric polynomials are very efficient approximations for smooth periodic functions. This, in turn, follows from the rapid decay of Fourier coefficients of smooth periodic functions. Suppose, for example, that $f(x)$ is a periodic function that has N bounded derivatives. Let be the n term trigonometric polynomial consisting of n terms of the Fourier sum:

$$F^{(n)}(x) = \sum_{\alpha \in Z_n} \hat{f}_\alpha e^{2\pi i \alpha x / p} .$$

Note that for large m ,

$$\sum_{\alpha \geq m} \frac{1}{\alpha^N} \approx \int_{\alpha=m}^{\infty} \frac{1}{\alpha^N} d\alpha = \frac{1}{N-1} \frac{1}{\alpha^{N-1}} .$$

The rapid decay inequality (7.28) gives a simple error bound

$$\begin{aligned} \left| f(x) - F^{(n)}(x) \right| &= \left| \sum_{\alpha \notin Z_n} \hat{f}_\alpha e^{2\pi i \alpha x/p} \right| \\ &\leq C \sum_{\alpha \geq n/2} \frac{1}{\alpha^N} + C \sum_{\alpha \leq -n/2} \frac{1}{|\alpha|^N} \\ &\leq C \cdot \frac{1}{n^{N-1}} \end{aligned}$$

Thus, the smoother f is, the more accurate is the partial Fourier sum approximation.

7.4 Software

Performance tools.

7.5 References and Resources

The classical books on numerical analysis (Dahlquist and Björck, Isaacson and Keller, etc.) discuss the various facts and forms of polynomial interpolation. There are many good books on Fourier analysis. One of my favorites is by Harry Dym and my colleague Henry McKean⁸ I learned the aliasing formula (7.37) from a paper by Heinz Kreiss and Joseph Oliger.

7.6 Exercises

1. Verify that both sides of (7.9) are equal to $2a$ when $f(x) = ax^2 + bx + c$.
2. One way to estimate the derivative or integral from function values is to differentiate or integrate the interpolating polynomial.
3. Show that the real Fourier modes (7.29) form an orthogonal family. This means that

- (a) $\langle u_\alpha, u_\beta \rangle = 0$ if $\alpha \neq \beta$.
- (b) $\langle v_\alpha, v_\beta \rangle = 0$ if $\alpha \neq \beta$.
- (c) $\langle u_\alpha, v_\beta \rangle = 0$ for any α and β .

In (a), α or β may be zero. In (b), α and β start with one. In (c), α starts with zero and β starts with one. It is easy to verify these relations using

⁸He pronounces McKean as one would Senator McCain.

complex exponentials and the formula $e^{i\theta} = \cos(\theta) + i\sin(\theta)$. For example, we can write $\cos(\theta) = \frac{1}{2}(e^{i\theta} + e^{-i\theta})$, so that $u_\alpha = \frac{1}{2}(w_\alpha + w_{-\alpha})$. Therefore

$$\langle u_\alpha, u_\beta \rangle = \frac{1}{4}(\langle w_\alpha, w_\beta \rangle + \langle w_{-\alpha}, w_\beta \rangle + \langle w_\alpha, w_{-\beta} \rangle + \langle w_{-\alpha}, w_{-\beta} \rangle).$$

You can check that if $\alpha \geq 0$ and $\beta \geq 0$ and $\alpha \neq \beta$, then all four terms on the right side are zero because the w_α are an orthogonal family. Also check this way that

$$\|u_\alpha\|_{L^2}^2 = \langle u_\alpha, u_\alpha \rangle = \|v_\alpha\|_{L^2}^2 = \frac{p}{2}, \quad (7.41)$$

if $\alpha \geq 1$, and $\|u_0\|_{L^2}^2 = p$.

4. We wish to use Fourier series to solve the boundary value problem from 4.11.

- (a) Show that the solution to (4.43) and (4.44) can be written as a Fourier sine series

$$u(x) = \sum_{\alpha=1}^{n-1} c_\alpha \sin(\pi\alpha x). \quad (7.42)$$

One way to do this is to define new functions also called u and f , that *extend* the given ones in a special way to satisfy the boundary conditions. For $1 \leq x \leq 2$, we define $f(x) = -f(x-1)$. This defines f in the interval $[0, 2]$ in a way that makes it antisymmetric about $x = 1$. Next if $x \notin [0, 2]$ there is a k so that $x - 2k \in [0, 2]$. Define $f(x) = f(x - 2k)$ in that case. This makes f a periodic function with period $p = 2$ that is antisymmetric about any of the integer points $x = 0, x = 1$, etc. Draw a picture to make this two step extension clear. If we express this extended f as a real cosine and sine series, the coefficients of all the cosine terms are zero (why?), so $f(x) = \sum_{\alpha>0} b_\alpha \sin(\pi\alpha x)$. (Why π instead of 2π ?) Now determine the c_α in terms of the b_α so that u satisfies (4.43). Use the sine series for u to show that $u = 0$ for $x = 0$ and $x = 1$.

- (b) Write a program in Matlab that uses the Matlab `fft` function to calculate the discrete sine coefficients b_α for a sampled function $f^{(n)}$. The simplest way to program this is to extend f to the interval $[0, 2]$ as described, sample this extended f at $2n+1$ uniformly spaced points in the interval $[0, 2]$, compute the complex DFT of these samples, then extract the sine coefficients as described in Section 7.3.1.
- (c) Write a Matlab program that takes the b_α and computes $\tilde{f}(x) = \sum_{\alpha=1}^{n-1} b_\alpha \sin(\pi\alpha x)$ for an arbitrary x value. On one graph, make a plot of $f(x)$ and $\tilde{f}(x)$ for $x \in [0, 1]$. On another plot the error $f(x) - \tilde{f}(x)$. Check that the error is zero at the sampling points, as it should be. For plotting, you will have to evaluate $f(x)$ and $\tilde{f}(x)$

at many more than the sampling points. Do this for $n = 5, 20, 100$ and $f(x) = x^2$ and $f(x) = \sin(3 \sin(x))$. Comment on the accuracy in the two cases.

- (d) Use the relation between b_α and c_α to calculate the solution to (4.43) and (4.44) for the two functions f in part (c). Comment on the difference in accuracy.
- (e) Show that the eigenvectors of the matrix A in Exercise 4.11 are discrete sine modes. Use this to describe an algorithm to express any vector, F , in terms as a linear combination of eigenvectors of A . This is more or less what part (a) does, only stated in a different way.
- (f) Use part (e) to develop a *fast algorithm* (i.e. as fast as the FFT rather than the direct method) to solve $Au = F$. Write a Matlab code to do this. Compare the accuracy of the second order method from Exercise 4.11 to the DFT based algorithm of part (c).

Chapter 8

Dynamics and Differential Equations

Many dynamical systems are modeled by first order systems of differential equations. An n component vector $x(t) = (x_1(t), \dots, x_n(t))$, models the state of the system at time t . The dynamics are modeled by

$$\frac{dx}{dt} = \dot{x}(t) = f(x(t), t) , \quad (8.1)$$

where $f(x, t)$ is an n component function, $f(x, t) = f_1(x, t), \dots, f_n(x, t)$. The system is *autonomous* if $f = f(x)$, i.e., if f is independent of t . A *trajectory* is a function, $x(t)$, that satisfies (8.1). The *initial value problem* is to find a trajectory that satisfies the *initial condition*¹

$$x(0) = x_0 , \quad (8.2)$$

with x_0 being the *initial data*. In practice we often want to specify initial data at some time other than $t_0 = 0$. We set $t_0 = 0$ for convenience of discussion. If $f(x, t)$ is a differentiable function of x and t , then the initial value problem has a solution trajectory defined at least for t close enough to t_0 . The solution is unique as long as it exists.²

Some problems can be reformulated into the form (8.1), (8.2). For example, suppose $F(r)$ is the force on an object of mass m if the position of the object is $r \in R^3$. Newton's equation of motion: $F = ma$ is

$$m \frac{d^2 r}{dt^2} = F(r) . \quad (8.3)$$

This is a system of three second order differential equations. The velocity at time t is $v(t) = \dot{r}(t)$. The trajectory, $r(t)$, is determined by the initial position, $r_0 = r(0)$, and the initial velocity, $v_0 = v(0)$.

We can reformulate this as a system of six first order differential equations for the position and velocity, $x(t) = (r(t), v(t))$. In components, this is

$$\begin{pmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \\ x_4(t) \\ x_5(t) \\ x_6(t) \end{pmatrix} = \begin{pmatrix} r_1(t) \\ r_2(t) \\ r_3(t) \\ v_1(t) \\ v_2(t) \\ v_3(t) \end{pmatrix} .$$

The dynamics are given by $\dot{r} = v$ and $\dot{v} = \frac{1}{m}F(r)$. This puts the equations (8.3)

¹There is a conflict of notation that we hope causes little confusion. Sometimes, as here, x_k refers to component k of the vector x . More often x_k refers to an approximation of the vector x at time t_k .

²This is the existence and uniqueness theorem for ordinary differential equations. See any good book on ordinary differential equations for details.

into the form (8.1) where

$$f = \begin{pmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 \\ f_5 \\ f_6 \end{pmatrix} = \begin{pmatrix} x_4 \\ x_5 \\ x_6 \\ \frac{1}{m} F_1(x_1, x_2, x_3) \\ \frac{1}{m} F_2(x_1, x_2, x_3) \\ \frac{1}{m} F_3(x_1, x_2, x_3) \end{pmatrix} .$$

There are variants of this scheme, such as taking $x_1 = r_1$, $x_2 = \dot{r}_1$, $x_3 = r_2$, etc., or using the momentum, $p = m\dot{r}$ rather than the velocity, $v = \dot{r}$. The initial values for the six components $x_0 = x(t_0)$ are given by the initial position and velocity components.

8.1 Time stepping and the forward Euler method

For simplicity this section supposes f does not depend on t , so that (8.1) is just $\dot{x} = f(x)$. *Time stepping*, or *marching*, means computing approximate values of $x(t)$ by advancing time in a large number of small steps. For example, if we know $x(t)$, then we can estimate $x(t + \Delta t)$ using

$$x(t + \Delta t) \approx x(t) + \Delta t \dot{x}(t) = x(t) + \Delta t f(x(t)) . \quad (8.4)$$

If we have an approximate value of $x(t)$, then we can use (8.4) to get an approximate value of $x(t + \Delta t)$.

This can be organized into a method for approximating the whole trajectory $x(t)$ for $0 \leq t \leq T$. Choose a *time step*, Δt , and define discrete times $t_k = k\Delta t$ for $k = 0, 1, 2, \dots$. We compute a sequence of approximate values $x_k \approx x(t_k)$. The approximation (8.4) gives

$$x_{k+1} \approx x(t_{k+1}) = x(t_k + \Delta t) \approx x(t_k) + \Delta t f(x_k) \approx x_k + \Delta t f(x_k) .$$

The *forward Euler* method takes this approximation as the definition of x_{k+1} :

$$x_{k+1} = x_k + \Delta t f(x_k) . \quad (8.5)$$

Starting with initial condition x_0 (8.5) allows us to calculate x_1 , then x_2 , and so on as far as we like.

Solving differential equations sometimes is called integration. This is because of the fact that if $f(x, t)$ is independent of x , then $\dot{x}(t) = f(t)$ and the solution of the initial value problem (8.1) (8.2) is given by

$$x(t) = x(0) + \int_0^t f(s) ds .$$

If we solve this using the rectangle rule with time step Δt , we get

$$x(t_k) \approx x_k = x(0) + \Delta t \sum_{j=0}^{k-1} f(t_j) .$$

We see from this that $x_{k+1} = x_k + \Delta t f(t_k)$, which is the forward Euler method in this case. We know that the rectangle rule for integration is first order accurate. This is a hint that the forward Euler method is first order accurate more generally.

We can estimate the accuracy of the forward Euler method using an informal *error propagation equation*. The error, as well as the solution, evolves (or propagates) from one time step to the next. We write the value of the exact solution at time t_k as $\tilde{x}_k = x(t_k)$. The error at time t_k is $e_k = x_k - \tilde{x}_k$. The *residual* is the amount by which \tilde{x}_k fails to satisfy the forward Euler equations³ (8.5):

$$\tilde{x}_{k+1} = \tilde{x}_k + \Delta t f(\tilde{x}_k) + \Delta t r_k , \quad (8.6)$$

which can be rewritten as

$$r_k = \frac{x(t_k + \Delta t) - x(t_k)}{\Delta t} - f(x(t_k)) . \quad (8.7)$$

In Section 3.2 we showed that

$$\frac{x(t_k + \Delta t) - x(t_k)}{\Delta t} = \dot{x}(t_k) + \frac{\Delta t}{2} \ddot{x}(t_k) + O(\Delta t^2) .$$

Together with $\dot{x} = f(x)$, this shows that

$$r_k = \frac{\Delta t}{2} \ddot{x}(t_k) + O(\Delta t^2) , \quad (8.8)$$

which shows that $r_k = O(\Delta t)$.

The error propagation equation, (8.10) below, estimates e in terms of the residual. We can estimate $e_k = x_k - \tilde{x}_k = x_k - x(t_k)$ by comparing (8.5) to (8.6)

$$e_{k+1} = e_k + \Delta t (f(x_k) - f(\tilde{x}_k)) - \Delta t r_k .$$

This resembles the forward Euler method applied to approximating some function $e(t)$. Being optimistic, we suppose that x_k and $x(t_k)$ are close enough to use the approximation (f' is the Jacobian matrix of f as in Section ??)

$$f(x_k) = f(\tilde{x}_k + e_k) \approx f(\tilde{x}_k) + f'(\tilde{x}_k)e_k ,$$

and then

$$e_{k+1} \approx e_k + \Delta t (f'(\tilde{x}_k)e_k - r_k) . \quad (8.9)$$

If this were an equality, it would imply that the e_k satisfy the forward Euler approximation to the differential equation

$$\dot{e} = f'(x(t))e - r(t) , \quad (8.10)$$

where $x(t)$ satisfies (8.1), e has initial condition $e(0) = 0$, and $r(t)$ is given by (8.8):

$$r(t) = \frac{\Delta t}{2} \ddot{x}(t) . \quad (8.11)$$

³We take out one factor of Δt so that the order of magnitude of r is the order of magnitude of the error e .

The *error propagation equation* (8.10) is linear, so $e(t)$ should be proportional to⁴ r , which is proportional to Δt . If the approximate $e(t)$ satisfies $\|e(t)\| = C(t)\Delta t$, then the exact $e(t)$ should satisfy $\|e(t)\| = O(\Delta t)$, which means there is a $C(t)$ with

$$\|e(t)\| \leq C(t)\Delta t. \quad (8.12)$$

This is the first order accuracy of the forward Euler method.

It is important to note that this argument does not prove that the forward Euler method converges to the correct answer as $\Delta t \rightarrow 0$. Instead, it assumes the convergence and uses it to get a quantitative estimate of the error. The formal proof of convergence may be found in any good book on numerical methods for ODEs, such as the book by Arieh Iserles.

If this analysis is done a little more carefully, it shows that there is an asymptotic error expansion

$$x_k \sim x(t_k) + \Delta t u_1(t_k) + \Delta t^2 u_2(t_k) + \cdots. \quad (8.13)$$

The leading error coefficient, $u_1(t)$, is the solution of (8.10). The higher order coefficients, $u_2(t)$, etc. are found by solving higher order error propagation equations.

The *modified equation* is a different approach to error analysis that allows us to determine the long time behavior of the error. The idea is to modify the differential equation (8.1) so that the solution is closer to the forward Euler sequence. We know that $x_k = x(t_k) + O(\Delta t)$. We seek a differential equation $\dot{y} = g(y)$ so that $x_k = y(t_k) + O(\Delta t^2)$. We construct an error expansion for the equation itself rather than the solution.

It is simpler to require $y(t)$ so satisfy the forward Euler equation at each t , not just the discrete times t_k :

$$y(t + \Delta t) = y(t) + \Delta t f(y(t)). \quad (8.14)$$

We seek

$$g(y, \Delta t) = g_0(y) + \Delta t g_1(y) + \cdots \quad (8.15)$$

so that the solution of (8.14) satisfies $\dot{y} = g(y) + O(\Delta t^2)$. We combine the expansion (8.15) with the Taylor series

$$y(t + \Delta t) = y(t) + \Delta t \dot{y}(t) + \frac{\Delta t^2}{2} \ddot{y}(t) + O(\Delta t^3),$$

to get (dividing both sides by Δt , $O(\Delta t^3)/\Delta t = O(\Delta t^2)$):

$$\begin{aligned} y(t) + \Delta t \dot{y}(t) + \frac{\Delta t^2}{2} \ddot{y}(t) + O(\Delta t^3) &= y(t) + \Delta t f(y(t)) \\ g_0(y(t)) + \Delta t g_1(y(t)) + \frac{\Delta t}{2} \ddot{y}(t) &= f(y(t)) + O(\Delta t^2) \end{aligned}$$

⁴The value of $e(t)$ depends on the values of $r(s)$ for $0 \leq s \leq t$. We can solve $\dot{u} = f'(x)u - w$, where $w = \frac{1}{2}\ddot{x}$, then solve (8.10) by setting $e = \Delta t u$. This shows that $\|e(t)\| = \Delta t \|u(t)\|$, which is what we want, with $C(t) = \|u(t)\|$.

Equating the leading order terms gives the unsurprising result

$$g_0(y) = f(y) ,$$

and leaves us with

$$g_1(y(t)) + \frac{1}{2}\ddot{y}(t) = O(\Delta t) . \quad (8.16)$$

We differentiate $\dot{y} = f(y) + O(\Delta t)$ and use the chain rule, giving

$$\begin{aligned} \ddot{y} = \frac{d}{dt}\dot{y} &= \frac{d}{dt}\left(f(y(t)) + O(\Delta t)\right) \\ &= f'(y)\dot{y}(t) + O(\Delta t) \\ \ddot{y} &= f'(y)f(y) + O(\Delta t) \end{aligned}$$

Substituting this into (8.16) gives

$$g_1(y) = -\frac{1}{2}f'(y)f(y) .$$

so the modified equation, with the first correction term, is

$$\dot{y} = f(y) - \frac{\Delta t}{2}f'(y)f(y) . \quad (8.17)$$

A simple example illustrates these points. The nondimensional harmonic oscillator equation is $\ddot{r} = -r$. The solution is $r(t) = a \sin(t) + b \cos(t)$, which oscillates but does not grow or decay. We write this in first order as $\dot{x}_1 = x_2$, $\dot{x}_2 = -x_1$, or

$$\frac{d}{dt} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} x_2 \\ -x_1 \end{pmatrix} . \quad (8.18)$$

Therefore, $f(x) = \begin{pmatrix} x_2 \\ -x_1 \end{pmatrix}$, $f' = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$, and $f'f = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} x_2 \\ -x_1 \end{pmatrix} = \begin{pmatrix} -x_1 \\ -x_2 \end{pmatrix}$, so (8.17) becomes

$$\frac{d}{dt} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} y_2 \\ -y_1 \end{pmatrix} + \frac{\Delta t}{2} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} \frac{\Delta t}{2} & 1 \\ -1 & \frac{\Delta t}{2} \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} .$$

We can solve this by finding eigenvalues and eigenvectors, but a simpler trick is to use a partial integrating factor and set $y(t) = e^{\frac{1}{2}\Delta t \cdot t} z(t)$, where $\dot{z} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} z$. Since $z_1(t) = a \sin(t) + b \cos(t)$, we have our approximate numerical solution $y_1(t) = e^{\frac{1}{2}\Delta t \cdot t} (a \sin(t) + b \cos(t))$. Therefore

$$\|e(t)\| \approx \left(e^{\frac{1}{2}\Delta t \cdot t} - 1 \right) . \quad (8.19)$$

This modified equation analysis confirms that forward Euler is first order accurate. For small Δt , we write $e^{\frac{1}{2}\Delta t \cdot t} - 1 \approx \frac{1}{2}\Delta t \cdot t$ so the error is about $\Delta t \cdot t (a \sin(t) + b \cos(t))$. Moreover, it shows that the error grows with t . For each fixed t , the error satisfies $\|e(t)\| = O(\Delta t)$ but the implied constant $C(t)$ (in $\|e(t)\| \leq C(t)\Delta t$) is a growing function of t , at least as large as $C(t) \geq \frac{t}{2}$.

8.2 Runge Kutta methods

Runge Kutta⁵ methods are a general way to achieve higher order approximate solutions of the initial value problem (8.1), (8.2). Each time step consists of m stages, each stage involving a single evaluation of f . The relatively simple four stage fourth order method is in wide use today. Like the forward Euler method, but unlike multistep methods, Runge Kutta time stepping computes x_{k+1} from x_k without using values x_j for $j < k$. This simplifies error estimation and adaptive time step selection.

The simplest Runge Kutta method is forward Euler (8.5). Among the second order methods is Heun's⁶

$$\xi_1 = \Delta t f(x_k, t_k) \quad (8.20)$$

$$\xi_2 = \Delta t f(x_k + \xi_1, t_k + \Delta t) \quad (8.21)$$

$$x_{k+1} = x_k + \frac{1}{2} (\xi_1 + \xi_2) . \quad (8.22)$$

The calculations of ξ_1 and ξ_2 are the two *stages* of Heun's method. Clearly they depend on k , though that is left out of the notation.

To calculate x_k from x_0 using a Runge Kutta method, we apply take k time steps. Each time step is a transformation that may be written

$$x_{k+1} = \widehat{S}(x_k, t_k, \Delta t) .$$

As in Chapter 6, we express the general time step as⁷ $x' = \widehat{S}(\bar{x}, t, \Delta t)$. This \widehat{S} approximates the exact *solution operator*, $S(\bar{x}, t, \Delta t)$. We say that $x' = S(\bar{x}, t, \Delta t)$ if there is a trajectory satisfying the differential equation (8.1) so that $x(t) = \bar{x}$ and $x' = x(t + \Delta t)$. In this notation, we would give Heun's method as $x' = \widehat{S}(\bar{x}, \Delta t) = \bar{x} + \frac{1}{2} (\xi_1 + \xi_2)$, where $\xi_1 = f(\bar{x}, t, \Delta t)$, and $\xi_2 = f(\bar{x} + \xi_1, t, \Delta t)$.

The best known and most used Runge Kutta method, which often is called *the* Runge Kutta method, has four stages and is fourth order accurate

$$\xi_1 = \Delta t f(\bar{x}, t) \quad (8.23)$$

$$\xi_2 = \Delta t f(\bar{x} + \frac{1}{2}\xi_1, t + \frac{1}{2}\Delta t) \quad (8.24)$$

$$\xi_3 = \Delta t f(\bar{x} + \frac{1}{2}\xi_2, t + \frac{1}{2}\Delta t) \quad (8.25)$$

$$\xi_4 = \Delta t f(\bar{x} + \xi_3, t + \Delta t) \quad (8.26)$$

$$x' = \bar{x} + \frac{1}{6} (\xi_1 + 2\xi_2 + 2\xi_3 + \xi_4) . \quad (8.27)$$

⁵Carl Runge was Professor of applied mathematics at the turn of the 20th century in Göttingen, Germany. Among his colleagues were David Hilbert (of Hilbert space) and Richard Courant. But Courant was forced to leave Germany and came to New York to found the Courant Institute. Kutta was a student of Runge.

⁶Heun, whose name rhymes with "coin", was another student of Runge.

⁷The notation x' here does not mean the derivative of x with respect to t (or any other variable) as it does in some books on differential equations.

Understanding the accuracy of Runge Kutta methods comes down to Taylor series. The reasoning of Section 8.1 suggests that the method has error $O(\Delta t^p)$ if

$$\widehat{S}(\bar{x}, t, \Delta t) = S(\bar{x}, t, \Delta t) + \Delta t r, \quad (8.28)$$

where $\|r\| = O(\Delta t^p)$. The reader should verify that this definition of the residual, r , agrees with the definition in Section 8.1. The analysis consists of expanding both $S(\bar{x}, t, \Delta t)$ and $\widehat{S}(\bar{x}, t, \Delta t)$ in powers of Δt . If the terms agree up to order Δt^p but disagree at order Δt^{p+1} , then p is the order of accuracy of the overall method.

We do this for Heun's method, allowing f to depend on t as well as x . The calculations resemble the derivation of the modified equation (8.17). To make the expansion of S , we have $x(t) = \bar{x}$, so

$$x(t + \Delta t) = \bar{x} + \Delta t \dot{x}(t) + \frac{\Delta t^2}{2} \ddot{x}(t) + O(\Delta t^3).$$

Differentiating with respect to t and using the chain rule gives:

$$\ddot{x} = \frac{d}{dt} \dot{x} = \frac{d}{dt} f(x(t), t) = f'(x(t), t) \dot{x}(t) + \partial_t f(x(t), t),$$

so

$$\ddot{x}(t) = f'(\bar{x}, t) f(\bar{x}, t) + \partial_t f(\bar{x}, t).$$

This gives

$$S(\bar{x}, t, \Delta t) = \bar{x} + \Delta t f(\bar{x}, t) + \frac{\Delta t^2}{2} (f'(\bar{x}, t) f(\bar{x}, t) + \partial_t f(\bar{x}, t)) + O(\Delta t^3). \quad (8.29)$$

To make the expansion of \widehat{S} for Heun's method, we first have $\xi_1 = \Delta t f(\bar{x}, t)$, which needs no expansion. Then

$$\begin{aligned} \xi_2 &= \Delta t f(\bar{x} + \xi_1, t + \Delta t) \\ &= \Delta t (f(\bar{x}, t) + f'(\bar{x}, t) \xi_1 + \partial_t f(\bar{x}, t) \Delta t + O(\Delta t^2)) \\ &= \Delta t f(\bar{x}, t) + \Delta t^2 (f'(\bar{x}, t) f(\bar{x}, t) + \partial_t f(\bar{x}, t)) + O(\Delta t^3). \end{aligned}$$

Finally, (8.22) gives

$$\begin{aligned} x' &= \bar{x} + \frac{1}{2} (\xi_1 + \xi_2) \\ &= \bar{x} + \frac{1}{2} \left\{ \Delta t f(\bar{x}, t) + \left[\Delta t f(\bar{x}, t) + \Delta t^2 (f'(\bar{x}, t) f(\bar{x}, t) + \partial_t f(\bar{x}, t)) \right] \right\} + O(\Delta t^3) \end{aligned}$$

Comparing this to (8.29) shows that

$$\widehat{S}(\bar{x}, t, \Delta t) = S(\bar{x}, t, \Delta t) + O(\Delta t^3).$$

which is the second order accuracy of Heun's method. The same kind of analysis shows that the four stage Runge Kutta method is fourth order accurate, but it would take a full time week. It was Kutta's thesis.

8.3 Linear systems and stiff equations

A good way to learn about the behavior of a numerical method is to ask what it would do on a properly chosen *model problem*. In particular, we can ask how an initial value problem solver would perform on a linear system of differential equations

$$\dot{x} = Ax. \quad (8.30)$$

We can do this using the eigenvalues and eigenvectors of A if the eigenvectors are not too ill conditioned. If⁸ $Ar_\alpha = \lambda_\alpha r_\alpha$ and $x(t) = \sum_{\alpha=1}^n u_\alpha(t)r_\alpha$, then the components u_α satisfy the scalar differential equations

$$\dot{u}_\alpha = \lambda_\alpha u_\alpha. \quad (8.31)$$

Suppose $x_k \approx x(t_k)$ is the approximate solution at time t_k . Write $x_k = \sum_{\alpha=1}^n u_{\alpha k} r_\alpha$. For a majority of methods, including Runge Kutta methods and linear multistep methods, the $u_{\alpha k}$ (as functions of k) are what you would get by applying the same time step approximation to the scalar equation (8.31). The eigenvector matrix, R , (see Section ??), that diagonalizes the differential equation (8.30) also diagonalizes the computational method. The reader should check that this is true of the Runge Kutta methods of Section 8.2.

One question this answers, at least for linear equations (8.30), is how small the time step should be. From (8.31) we see that the λ_α have units of 1/time, so the $1/|\lambda_\alpha|$ have units of time and may be called *time scales*. Since Δt has units of time, it does not make sense to say that Δt is small in an absolute sense, but only relative to other time scales in the problem. This leads to the following:

Possibility: A time stepping approximation to (8.30) will be accurate only if

$$\max_{\alpha} \Delta t |\lambda_\alpha| \ll 1. \quad (8.32)$$

Although this *possibility* is not true in every case, it is a dominant technical consideration in most practical computations involving differential equations. The *possibility* suggests that the time step should be considerably smaller than the smallest time scale in the problem, which is to say that Δt should *resolve* even the fastest time scales in the problem.

A problem is called *stiff* if it has two characteristics: (i) there is a wide range of time scales, and (ii) the fastest time scale modes have almost no energy. The second condition states that if $|\lambda_\alpha|$ is large (relative to other eigenvalues), then $|u_\alpha|$ is small. Most time stepping problems for partial differential equations are stiff in this sense. For a stiff problem, we would like to take larger time steps than (8.32):

$$\Delta t |\lambda_\alpha| \ll 1 \quad \left\{ \begin{array}{l} \text{for all } \alpha \text{ with } u_\alpha \text{ signifi-} \\ \text{cantly different from zero.} \end{array} \right. \quad (8.33)$$

What can go wrong if we ignore (8.32) and choose a time step using (8.33) is *numerical instability*. If mode u_α is one of the large $|\lambda_\alpha|$ small $|u_\alpha|$ modes,

⁸We call the eigenvalue index α to avoid conflict with k , which we use to denote the time step.

it is natural to assume that the real part satisfies $\operatorname{Re}(\lambda_\alpha) \leq 0$. In this case we say the mode is *stable* because $|u_\alpha(t)| = |u_\alpha(0)| e^{\lambda_\alpha t}$ does not increase as t increases. However, if $\Delta t \lambda_\alpha$ is not small, it can happen that the time step approximation to (8.31) is unstable. This can cause the $u_{\alpha k}$ to grow exponentially while the actual u_α decays or does not grow. Exercise 8 illustrates this. Time step restrictions arising from stability are called *CFL* conditions because the first systematic discussion of this possibility in the numerical solution of partial differential equations was given in 1929 by Richard Courant, Kurt Friedrichs, and Hans Levy.

8.4 Adaptive methods

Adaptive means that the computational steps are not fixed in advance but are determined as the computation proceeds. Section 3.6, discussed an integration algorithm that chooses the number of integration points adaptively to achieve a specified overall accuracy. More sophisticated adaptive strategies choose the distribution of points to maximize the accuracy from a given amount of work. For example, suppose we want an \hat{I} for $I = \int_0^2 f(x)dx$ so that $|\hat{I} - I| < .06$. It might be that we can calculate $I_1 = \int_0^1 f(x)dx$ to within .03 using $\Delta x = .1$ (10 points), but that calculating $I_2 = \int_1^2 f(x)dx$ to within .03 takes $\Delta x = .02$ (50 points). It would be better to use $\Delta x = .1$ for I_1 and $\Delta x = .02$ for I_2 (60 points total) rather than using $\Delta x = .02$ for all of I (100 points).

Adaptive methods can use *local error estimates* to concentrate computational resources where they are most needed. If we are solving a differential equation to compute $x(t)$, we can use a smaller time step in regions where x has large acceleration. There is an active community of researchers developing systematic ways to choose the time steps in a way that is close to optimal without having the overhead in choosing the time step become larger than the work in solving the problem. In many cases they conclude, and simple model problems show, that a good strategy is to *equidistribute* the *local truncation error*. That is, to choose time steps Δt_k so that the local truncation error $\rho_k = \Delta t_k r_k$ is nearly constant.

If we have a variable time step Δt_k , then the times $t_{k+1} = t_k + \Delta t_k$ form an irregular *adapted mesh* (or *adapted grid*). Informally, we want to choose a mesh that *resolves* the solution, $x(t)$ being calculated. This means that knowing the $x_k \approx x(t_k)$ allows you make an accurate reconstruction of the function $x(t)$, say, by interpolation. If the points t_k are too far apart then the solution is *underresolved*. If the t_k are so close that $x(t_k)$ is predicted accurately by a few neighboring values ($x(t_j)$ for $j = k \pm 1, k \pm 2$, etc.) then $x(t)$ is *overresolved*, we have computed it accurately but paid too high a price. An efficient adaptive mesh avoids both underresolution and overresolution.

Figure 8.1 illustrates an adapted mesh with equidistributed interpolation error. The top graph shows a curve we want to resolve and a set of points that concentrates where the curvature is high. Also shown is the piecewise linear

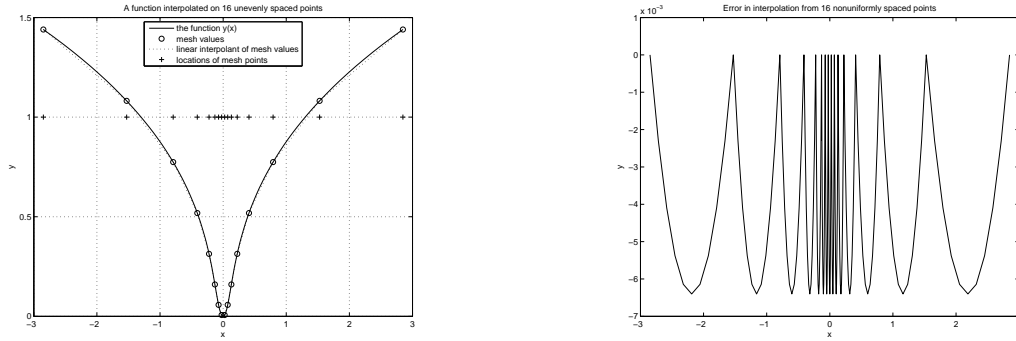


Figure 8.1: A nonuniform mesh for a function that needs different resolution in different places. The top graph shows the function and the mesh used to interpolate it. The bottom graph is the difference between the function and the piecewise linear approximation. Note that the interpolation error equidistributed though the mesh is much finer near $x = 0$.

curve that connects the interpolation points. On the graph it looks as though the piecewise linear graph is closer to the curve near the center than in the smoother regions at the ends, but the error graph in the lower frame shows this is not so. The reason probably is that what is plotted in the bottom frame is the vertical distance between the two curves, while what we see in the picture is the two dimensional distance, which is less if the curves are steep. The bottom frame illustrates equidistribution of error. The interpolation error is zero at the grid points and gets to be as large as about -6.3×10^{-3} in each interpolation interval. If the points were uniformly spaced, the interpolation error would be smaller near the ends and larger in the center. If the points were bunched even more than they are here, the interpolation error would be smaller in the center than near the ends. We would not expect such perfect equidistribution in real problems, but we might have errors of the same order of magnitude everywhere.

For a Runge Kutta method, the local truncation error is $\rho(x, t, \Delta t) = \widehat{S}(x, t, \Delta t) - S(x, t, \Delta t)$. We want a way to estimate ρ and to choose Δt so that $|\rho| = e$, where e is the value of the equidistributed local truncation error. We suggest a method related to Richardson extrapolation (see Section 3.3), comparing the result of

one time step of size Δt to two time steps of size $\Delta t/2$. The best adaptive Runge Kutta differential equation solvers do not use this generic method, but instead use ingenious schemes such as the Runge Kutta Fehlberg five stage scheme that simultaneously gives a fifth order \hat{S}_5 , but also gives an estimate of the difference between a fourth order and a fifth order method, $\hat{S}_5 - \hat{S}_4$. The method described here does the same thing with eight function evaluations instead of five.

The Taylor series analysis of Runge Kutta methods indicates that $\rho(x, t, \Delta t) = \Delta t^{p+1} \sigma(x, t) + O(\Delta t^{p+2})$. We will treat σ as a constant because all the x and t values we use are within $O(\Delta t)$ of each other, so variations in σ do not effect the principal error term we are estimating. With one time step, we get $x' = \hat{S}(\bar{x}, t, \Delta t)$. With two half size time steps we get first $\tilde{x}_1 = \hat{S}(\bar{x}, t, \Delta t/2)$, then $\tilde{x}_2 = \hat{S}(\tilde{x}_1, t + \Delta t/2, \Delta t/2)$.

We will show, using the Richardson method of Section 3.3, that

$$x' - \tilde{x}_2 = (1 - 2^{-p}) \rho(\bar{x}, t, \Delta t) + O(\Delta t^{p+1}). \quad (8.34)$$

We need to use the *semigroup* property of the solution operator: If we “run” the exact solution from \bar{x} for time $\Delta t/2$, then run it from there for another time $\Delta t/2$, the result is the same as running it from \bar{x} for time Δt . Letting x be the solution of (8.1) with $x(t) = \bar{x}$, the formula for this is

$$\begin{aligned} S(\bar{x}, t, \Delta t) &= S(x(t + \Delta t/2), t + \Delta t/2, \Delta t/2) \\ &= S(S(\bar{x}, t, \Delta t/2), t + \Delta t/2, \Delta t/2). \end{aligned}$$

We also need to know that $S(x, t, \Delta t) = x + O(\Delta t)$ is reflected in the Jacobian matrix S' (the matrix of first partials of S with respect to the x arguments with t and Δt fixed)⁹: $S'(x, t, \Delta t) = I + O(\Delta t)$.

The actual calculation starts with

$$\begin{aligned} \tilde{x}_1 &= \hat{S}(\bar{x}, t, \Delta t/2) \\ &= S(\bar{x}, t, \Delta t/2) + 2^{-(p+1)} \Delta t^{-(p+1)} \sigma + O(\Delta t^{-(p+2)}), \end{aligned}$$

and

$$\begin{aligned} \tilde{x}_2 &= \hat{S}(\tilde{x}_1, t + \Delta t/2, \Delta t/2) \\ &= S(\tilde{x}_1, t + \Delta t/2, \Delta t/2) + 2^{-(p+1)} \Delta t^{-(p+1)} \sigma + O(\Delta t^{-(p+2)}), \end{aligned}$$

We simplify the notation by writing $\tilde{x}_1 = x(t + \Delta t/2) + u$ with $u = 2^{-(p+1)} \Delta t^p \sigma + O(\Delta t^{-(p+2)})$. Then $\|u\| = O(\Delta t^{-(p+1)})$ and also (used below) $\Delta t \|u\| = O(\Delta t^{-(p+2)})$ and (since $p \geq 1$) $\|u\|^2 = O(\Delta t^{-(2p+2)}) = O(\Delta t^{-(p+2)})$. Then

$$\begin{aligned} S(\tilde{x}_1, t + \Delta t/2, \Delta t/2) &= S(x(t + \Delta t/2) + u, t + \Delta t/2, \Delta t/2) \\ &= S(x(t + \Delta t/2), t + \Delta t/2, \Delta t/2) + S' u + O(\|u\|^2) \\ &= S(x(t + \Delta t/2), t + \Delta t/2, \Delta t/2) + u + O(\|u\|^2) \\ &= S(\bar{x}, t, \Delta t) + 2^{-(p+1)} \Delta t^p \sigma + u + O(\Delta t^{p+2}). \end{aligned}$$

⁹This fact is a consequence of the fact that S is twice differentiable as a function of all its arguments, which can be found in more theoretical books on differential equations. The Jacobian of $f(x) = x$ is $f'(x) = I$.

Altogether, since $2 \cdot 2^{-(p+1)} = 2^{-p}$, this gives

$$\tilde{x}_2 = S(\bar{x}, t, \Delta t) + 2^{-p} \Delta t^{p+1} \sigma + O(\Delta t^{p+2}) .$$

Finally, a single size Δt time step has

$$x' = X(\bar{x}, \Delta t, t) + \Delta t^{p+1} \sigma + O(\Delta t^{p+2}) .$$

Combining these gives (8.34). It may seem like a mess but it has a simple underpinning. The whole step produces an error of order Δt^{p+1} . Each half step produces an error smaller by a factor of 2^{p+1} , which is the main idea of Richardson extrapolation. Two half steps produce almost exactly twice the error of one half step.

There is a simple adaptive strategy based on the local truncation error estimate (8.34). We arrive at the start of time step k with an estimated time step size Δt_k . Using that time step, we compute $x' = \hat{S}(x_k, t_k, \Delta t_k)$ and \tilde{x}_2 by taking two time steps from x_k with $\Delta t_k/2$. We then estimate ρ_k using (8.34):

$$\hat{\rho}_k = \frac{1}{1 - 2^{-p}} (x' - \tilde{x}_2) . \quad (8.35)$$

This suggests that if we adjust Δt_k by a factor of μ (taking a time step of size $\mu \Delta t_k$ instead of Δt_k), the error would have been $\mu^{p+1} \hat{\rho}_k$. If we choose μ to exactly equidistribute the error (according to our estimated ρ), we would get

$$e = \mu^{p+1} \|\hat{\rho}_k\| \implies \mu_k = (e / \|\hat{\rho}_k\|)^{1/(p+1)} . \quad (8.36)$$

We could use this estimate to adjust Δt_k and calculate again, but this may lead to an infinite loop. Instead, we use $\Delta t_{k+1} = \mu_k \Delta t_k$.

Chapter 3 already mentioned the paradox of error estimation. Once we have a quantitative error estimate, we should use it to make the solution more accurate. This means taking

$$x_{k+1} = \hat{S}(x_k, t_k, \Delta t_k) + \hat{\rho}_k ,$$

which has order of accuracy $p + 1$, instead of the order p time step \hat{S} . This increases the accuracy but leaves you without an error estimate. This gives an order $p + 1$ calculation with a mesh chosen to be nearly optimal for an order p calculation. Maybe the reader can find a way around this paradox. Some adaptive strategies reduce the overhead of error estimation by using the Richardson based time step adjustment, say, every fifth step.

One practical problem with this strategy is that we do not know the quantitative relationship between local truncation error and global error¹⁰. Therefore it is hard to know what e to give to achieve a given global error. One way to estimate global error would be to use a given e and get some time steps Δt_k , then

¹⁰ *Adjoint* based error control methods that address this problem are still in the research stage (2006).

redo the computation with each interval $[t_k, t_{k+1}]$ cut in half, taking exactly twice the number of time steps. If the method has order of accuracy p , then the global error should decrease very nearly by a factor of 2^p , which allows us to estimate that error. This is rarely done in practice. Another issue is that there can be isolated zeros of the leading order truncation error. This might happen, for example, if the local truncation error were proportional to a scalar function \ddot{x} . In (8.36), this could lead to an unrealistically large time step. One might avoid that, say, by replacing μ_k with $\min(\mu_k, 2)$, which would allow the time step to grow quickly, but not too much in one step. This is less of a problem for systems of equations.

8.5 Multistep methods

Linear multistep methods are the other class of methods in wide use. Rather than giving a general discussion, we focus on the two most popular kinds, methods based on difference approximations, and methods based on integrating $f(x(t))$, *Adams methods*. Hybrids of these are possible but often are unstable. For some reason, almost nothing is known about methods that both are multistage and multistep.

Multistep methods are characterized by using information from previous time steps to go from x_k to x_{k+1} . We describe them for a fixed Δt . A simple example would be to use the second order centered difference approximation $\dot{x}(t) \approx (x(t + \Delta t) - x(t - \Delta t)) / 2\Delta t$ to get

$$(x_{k+1} - x_{k-1}) / 2\Delta t = f(x_k) ,$$

or

$$x_{k+1} = x_{k-1} + 2\Delta t f(x_k) . \quad (8.37)$$

This is the *leapfrog*¹¹ method. We find that

$$\tilde{x}_{k+1} = \tilde{x}_{k-1} + 2\Delta t f(\tilde{x}_k) + \Delta t O(\Delta t^2) ,$$

so it is second order accurate. It achieves second order accuracy with a single evaluation of f per time step. Runge Kutta methods need at least two evaluations per time step to be second order. Leapfrog uses x_{k-1} and x_k to compute x_{k+1} , while Runge Kutta methods forget x_{k-1} when computing x_{k+1} from x_k .

The next method of this class illustrates the subtleties of multistep methods. It is based on the four point one sided difference approximation

$$\dot{x}(t) = \frac{1}{\Delta t} \left(\frac{1}{3}x(t + \Delta t) + \frac{1}{2}x(t) - x(t - \Delta t) + \frac{1}{6}x(t - 2\Delta t) \right) + O(\Delta t^3) .$$

This suggests the time stepping method

$$f(x_k) = \frac{1}{\Delta t} \left(\frac{1}{3}x_{k+1} + \frac{1}{2}x_k - x_{k-1} + \frac{1}{6}x_{k-2} \right) , \quad (8.38)$$

¹¹Leapfrog is a game in which two or more children move forward in a line by taking turns jumping over each other, as (8.37) jumps from x_{k-1} to x_{k+1} using only $f(x_k)$.

which leads to

$$x_{k+1} = 3\Delta t f(x_k) - \frac{3}{2}x_k + 3x_{k-1} - \frac{1}{2}x_{k-2} . \quad (8.39)$$

This method never is used in practice because it is *unstable* in a way that Runge Kutta methods cannot be. If we set $f \equiv 0$ (to solve the model problem $\dot{x} = 0$), (8.38) becomes the *recurrence relation*

$$x_{k+1} + \frac{3}{2}x_k - 3x_{k-1} + \frac{1}{2}x_{k-2} = 0 , \quad (8.40)$$

which has *characteristic polynomial*¹² $p(z) = z^3 + \frac{3}{2}z^2 - 3z + \frac{1}{2}$. Since one of the roots of this polynomial has $|z| > 1$, general solutions of (8.40) grow exponentially on a Δt time scale, which generally prevents approximate solutions from converging as $\Delta t \rightarrow 0$. This cannot happen for Runge Kutta methods because setting $f \equiv 0$ always gives $x_{k+1} = x_k$, which is the exact answer in this case.

Adams methods use old values of f but not old values of x . We can integrate (8.1) to get

$$x(t_{k+1}) = x(t_k) + \int_{t_k}^{t_{k+1}} f(x(t))dt . \quad (8.41)$$

An accurate estimate of the integral on the right leads to an accurate time step. *Adams Bashforth* methods estimate the integral using polynomial extrapolation from earlier f values. At its very simplest we could use $f(x(t)) \approx f(x(t_k))$, which gives

$$\int_{t_k}^{t_{k+1}} f(x(t))dt \approx (t_{k+1} - t_k)f(x(t_k)) .$$

Using this approximation on the right side of (8.41) gives forward Euler.

The next order comes from linear rather than constant extrapolation:

$$f(x(t)) \approx f(x(t_k)) + (t - t_k) \frac{f(x(t_k)) - f(x(t_{k-1}))}{t_k - t_{k-1}} .$$

With this, the integral is estimated as (the generalization to non constant Δt is Exercise ??):

$$\begin{aligned} \int_{t_k}^{t_{k+1}} f(x(t))dt &\approx \Delta t f(x(t_k)) + \frac{\Delta t^2}{2} \frac{f(x(t_k)) - f(x(t_{k-1}))}{\Delta t} \\ &= \Delta t \left[\frac{3}{2}f(x(t_k)) - \frac{1}{2}f(x(t_{k-1})) \right] . \end{aligned}$$

The second order Adams Bashforth method for constant Δt is

$$x_{k+1} = x_k + \Delta t \left[\frac{3}{2}f(x_k) - \frac{1}{2}f(x_{k-1}) \right] . \quad (8.42)$$

¹²If $p(z) = 0$ then $x_k = z^k$ is a solution of (8.40).

To program higher order Adams Bashforth methods we need to evaluate the integral of the interpolating polynomial. The techniques of polynomial interpolation from Chapter 7 make this simpler.

Adams Bashforth methods are attractive for high accuracy computations where stiffness is not an issue. They cannot be unstable in the way (8.39) is because setting $f \equiv 0$ results (in (8.42), for example) in $x_{k+1} = x_k$, as for Runge Kutta methods. Adams Bashforth methods of any order or accuracy require one evaluation of f per time step, as opposed to four per time step for the fourth order Runge Kutta method.

8.6 Implicit methods

Implicit methods use $f(x_{k+1})$ in the formula for x_{k+1} . They are used for stiff problems because they can be stable with large $\lambda\Delta t$ (see Section 8.3) in ways explicit methods, all the ones discussed up to now, cannot. An implicit method must solve a system of equations to compute x_{k+1} .

The simplest implicit method is *backward Euler*:

$$x_{k+1} = x_k + \Delta t f(x_{k+1}) . \quad (8.43)$$

This is only first order accurate, but it is stable for any λ and Δt if $\text{Re}(\lambda) \leq 0$. This makes it suitable for solving stiff problems. It is called implicit because x_{k+1} is determined implicitly by (8.43), which we rewrite as

$$F(x_{k+1}, \Delta t) = 0 \quad , \quad \text{where} \quad F(y, \Delta t) = y - \Delta t f(y) - x_k , \quad (8.44)$$

To find x_{k+1} , we have to solve this system of equations for y .

Applied to the linear scalar problem (8.31) (dropping the α index), the method (8.43) becomes $u_{k+1} = u_k + \Delta t \lambda u_{k+1}$, or

$$u_{k+1} = \frac{1}{1 - \Delta t \lambda} u_k .$$

This shows that $|u_{k+1}| < |u_k|$ if $\Delta t > 0$ and λ is any complex number with $\text{Re}(\lambda) \leq 0$. This is in partial agreement with the qualitative behavior of the exact solution of (8.31), $u(t) = e^{\lambda t} u(0)$. The exact solution decreases in time if $\text{Re}(\lambda) < 0$ but not if $\text{Re}(\lambda) = 0$. The backward Euler approximation decreases in time even when $\text{Re}(\lambda) = 0$. The backward Euler method artificially stabilizes a neutrally stable system, just as the forward Euler method artificially destabilizes it (see the modified equation discussion leading to (8.19)).

Most likely the equations (8.44) would be solved using an iterative method as discussed in Chapter 6. This leads to *inner iterations*, with the *outer* iteration being the time step. If we use the unsafeguarded local Newton method, and let j index the inner iteration, we get $F' = I - \Delta t f'$ and

$$y_{j+1} = y_j - \left(I - \Delta t f'(y_j) \right)^{-1} (y_j - \Delta t f(y_j) - x_k) , \quad (8.45)$$

hoping that $y_j \rightarrow x_{k+1}$ as $j \rightarrow \infty$. We can take initial guess $y_0 = x_k$, or, even better, an extrapolation such as $y_0 = x_k + \Delta t(x_k - x_{k-1})/\Delta t = 2x_k - x_{k-1}$. With a good initial guess, just one Newton iteration should give x_{k+1} accurately enough.

Can we use the approximation $J \approx I$ for small Δt ? If we could, the Newton iteration would become the simpler *functional iteration* (check this)

$$y_{j+1} = x_k + \Delta t f(y_j) . \quad (8.46)$$

The problem with this is that it does not work precisely for the stiff systems we use implicit methods for. For example, applied to $\dot{u} = \lambda u$, the functional iteration diverges ($|y_j| \rightarrow \infty$ as $j \rightarrow \infty$) for $\Delta t \lambda < -1$.

Most of the explicit methods above have implicit analogues. Among implicit Runge Kutta methods we mention the *trapezoid rule*

$$\frac{x_{k+1} - x_k}{\Delta t} = \frac{1}{2}(f(x_{k+1}) + f(x_k)) . \quad (8.47)$$

There are *backward differentiation formula*, or *BDF* methods based on higher order one sided approximations of $\dot{x}(t_{k+1})$. The second order BDF method uses (??):

$$\dot{x}(t) = \frac{1}{\Delta t} \left(\frac{3}{2}x(t) - 2x(t - \Delta t) + \frac{1}{2}x(t - 2\Delta t) \right) + O(\Delta t^2) ,$$

to get

$$f(x(t_{k+1})) = \dot{x}(t_{k+1}) = \left(\frac{3}{2}x(t_{k+1}) - 2x(t_k) + \frac{1}{2}x(t_{k-1}) \right) + O(\Delta t^2) ,$$

and, neglecting the $O(\Delta t^2)$ error term,

$$x_{k+1} - \frac{2\Delta t}{3}f(x_{k+1}) = \frac{4}{3}x_k - \frac{1}{3}x_{k-1} . \quad (8.48)$$

The *Adams Molton* methods estimate $\int_{t_k}^{t_{k+1}} f(x(t))dt$ using polynomial interpolation using the values $f(x_{k+1})$, $f(x_k)$, and possibly $f(x_{k-1})$, etc. The second order Adams Molton method uses $f(x_{k+1})$ and $f(x_k)$. It is the same as the trapezoid rule (8.47). The third order Adams Molton method also uses $f(x_{k-1})$. Both the trapezoid rule (8.47) and the second order BDF method (8.48) both have the property of being *A-stable*, which means being stable for (8.31) with any λ and Δt as long as $\text{Re}(\lambda) \leq 0$. The higher order implicit methods are more stable than their explicit counterparts but are not A stable, which is a constant frustration to people looking for high order solutions to stiff systems.

8.7 Computing chaos, can it be done?

In many applications, the solutions to the differential equation (8.1) are *chaotic*.¹³ The informal definition is that for large t (not so large in real applications) $x(t)$ is an unpredictable function of $x(0)$. In the terminology of Section 8.5, this means that the solution operator, $S(x_0, 0, t)$, is an ill conditioned function of x_0 .

The dogma of Section 2.7 is that a floating point computation cannot give an accurate approximation to S if the condition number of S is larger than $1/\epsilon_{\text{mach}} \sim 10^{16}$. But practical calculations ranging from weather forecasting to molecular dynamics violate this rule routinely. In the computations below, the condition number of $S(x, t)$ increases with t and crosses 10^{16} by $t = 3$ (see Figure 8.3). Still, a calculation up to time $t = 60$ (Figure 8.4, bottom), shows the beautiful butterfly shaped *Lorentz attractor*, which looks just as it should.

On the other hand, in this and other computations, it truly is impossible to calculate details correctly. This is illustrated in Figure 8.2. The top picture plots two trajectories, one computed with $\Delta t = 4 \times 10^{-4}$ (dashed line), and the other with the time step reduced by a factor of 2 (solid line). The difference between the trajectories is an estimate of the accuracy of the computations. The computation seems somewhat accurate (curves close) up to time $t \approx 5$, at which time the dashed line goes up to $x \approx 15$ and the solid line goes down to $x \approx -15$. At least one of these is completely wrong. Beyond $t \approx 5$, the two “approximate” solutions have similar qualitative behavior but seem to be independent of each other. The bottom picture shows the same experiment with Δt a hundred times smaller than in the top picture. With a hundred times more accuracy, the approximate solution loses accuracy at $t \approx 10$ instead of $t \approx 5$. If a factor of 100 increase in accuracy only extends the validity of the solution by 5 time units, it should be hopeless to compute the solution out to $t = 60$.

The present numerical experiments are on the *Lorentz equations*, which are a system of three nonlinear ordinary differential equations

$$\begin{aligned}\dot{x} &= \sigma(y - x) \\ \dot{y} &= x(\rho - z) - y \\ \dot{z} &= xy - \beta z\end{aligned}$$

with¹⁴ $\sigma = 10$, $\rho = 28$, and $\beta = \frac{8}{3}$. The C/C++ program outputs (x, y, z) for plotting every $t = .02$ units of time, though there many time steps in each plotting interval. The solution first finds its way to the butterfly shaped *Lorentz attractor* then stays on it, traveling around the right and left wings in a seemingly random (technically, *chaotic*) way. The initial data $x = y = z = 0$ are not close to the attractor, so we ran the differential equations for some time before time $t = 0$ in order that $(x(0), y(0), z(0))$ should be a typical point on the attractor.

¹³James Gleick has written a nice popular book on chaos. Steve Strogatz has a more technical introduction that does not avoid the beautiful parts.

¹⁴These can be found, for example, in <http://wikipedia.org> by searching on “Lorentz attractor”.

Figure 8.2 shows the chaotic sequence of wing choice. A trip around the left wing corresponds to a dip of $x(t)$ down to $x \approx -15$ and a trip around the right wing corresponds to x going up to $x \approx 15$.

Sections 2.7 and 4.3 explain that the condition number of the problem of calculating $S(x, t)$ depends on the Jacobian matrix $A(x, t) = \partial_x S(x, t)$. This represents the sensitivity of the solution at time t to small perturbations of the initial data. Adapting notation from (4.28), we find that the condition number of calculating $x(t)$ from initial conditions $x(0) = x$ is

$$\kappa(S(x, t)) = \|A(x(0), t)\| \frac{\|x(0)\|}{\|x(t)\|} . \quad (8.49)$$

We can calculate $A(x, t)$ using ideas of perturbation theory similar to those we used for linear algebra problems in Section 4.2.6. Since $S(x, t)$ is the value of a solution at time t , it satisfies the differential equation

$$f(S(x, t)) = \frac{d}{dt} S(x, t) .$$

We differentiate both sides with respect to x and interchange the order of differentiation,

$$\frac{\partial}{\partial x} f(S(x, t)) = \frac{\partial}{\partial x} \frac{d}{dt} S(x, t) = \frac{d}{dt} \frac{\partial}{\partial x} S(x, t) = \frac{d}{dt} A(x, t) ,$$

to get (with the chain rule)

$$\begin{aligned} \frac{d}{dt} A(x, t) &= \frac{\partial}{\partial x} f(S(x, t)) \\ &= f'(S(x, t)) \cdot \partial_x S \\ \dot{A} &= f'(S(x, t)) A(x, t) . \end{aligned} \quad (8.50)$$

Thus, if we have an initial value x and calculate the trajectory $S(x, t)$, then we can calculate the *first variation*, $A(x, t)$, by solving the linear initial value problem (8.50) with initial condition $A(x, 0) = I$ (why?). In the present experiment, we solved the Lorentz equations and the perturbation equation using forward Euler with the same time step.

In typical chaotic problems, the first variation grows exponentially in time. If $\sigma_1(t) \geq \sigma_2(t) \geq \dots \geq \sigma_n(t)$ are the singular values of $A(x, t)$, then there typically are *Lyapounov exponents*, μ_k , so that

$$\sigma_k(t) \sim e^{\mu_k t} ,$$

More precisely,

$$\lim_{t \rightarrow \infty} \frac{\ln(\sigma_k(t))}{t} = \mu_k .$$

In our problem, $\|A(x, t)\| = \sigma_1(t)$ seems to grow exponentially because $\mu_1 > 0$. Since $\|x = x(0)\|$ and $\|x(t)\|$ are both of the order of about ten, this, with

(8.49), implies that $\kappa(S(x, t))$ grows exponentially in time. That explains why it is impossible to compute the values of $x(t)$ with any precision at all beyond $t = 20$.

It is interesting to study the condition number of A itself. If $\mu_1 > \mu_n$, the l^2 this also grows exponentially,

$$\kappa_{l^2}(A(x, t)) = \frac{\sigma_1(t)}{\sigma_n(t)} \sim e^{(\mu_1 - \mu_n)t}.$$

Figure 8.3 gives some indication that our Lorentz system has differing Lyapunov exponents. The top figure shows computations of the three singular values for $A(x, t)$. For $0 \leq t < \approx 2$, it seems that σ_3 is decaying exponentially, making a downward sloping straight line on this log plot. When σ_3 gets to about 10^{-15} , the decay halts. This is because it is nearly impossible for a full matrix in double precision floating point to have a condition number larger than $1/\epsilon_{\text{mach}} \approx 10^{16}$. When σ_3 hits 10^{-15} , we have $\sigma_1 \sim 10^2$, so this limit has been reached. These trends are clearer in the top frame of Figure 8.4, which is the same calculation carried to a larger time. Here $\sigma_1(t)$ seems to be growing exponentially with a gap between σ_1 and σ_3 of about $1/\epsilon_{\text{mach}}$. Theory says that σ_2 should be close to one, and the computations are consistent with this until the condition number bound makes $\sigma_2 \sim 1$ impossible.

To summarize, some results are quantitatively right, including the butterfly shape of the attractor and the exponential growth rate of $\sigma_1(t)$. Some results are qualitatively right but quantitatively wrong, including the values of $x(t)$ for $t > \approx 10$. Convergence analyses (comparing Δt results to $\Delta t/2$ results) distinguishes right from wrong in these cases. Other features of the computed solution are consistent over a range of Δt and consistently wrong. There is no reason to think the condition number of $A(x, t)$ grows exponentially until $t \sim 2$ then levels off at about 10^{16} . Much more sophisticated computational methods using the semigroup property show this is not so.

8.8 Software: Scientific visualization

Visualization of data is indispensable in scientific computing and computational science. Anomalies in data that seem to jump off the page in a plot are easy to overlook in numerical data. It can be easier to interpret data by looking at pictures than by examining columns of numbers. For example, here are entries 500 to 535 from the time series that made the top curve in the top frame of Figure 8.4 (multiplied by 10^{-5}).

0.1028	0.1020	0.1000	0.0963	0.0914	0.0864	0.0820
0.0790	0.0775	0.0776	0.0790	0.0818	0.0857	0.0910
0.0978	0.1062	0.1165	0.1291	0.1443	0.1625	0.1844
0.2104	0.2414	0.2780	0.3213	0.3720	0.4313	0.4998
0.5778	0.6649	0.7595	0.8580	0.9542	1.0395	1.1034

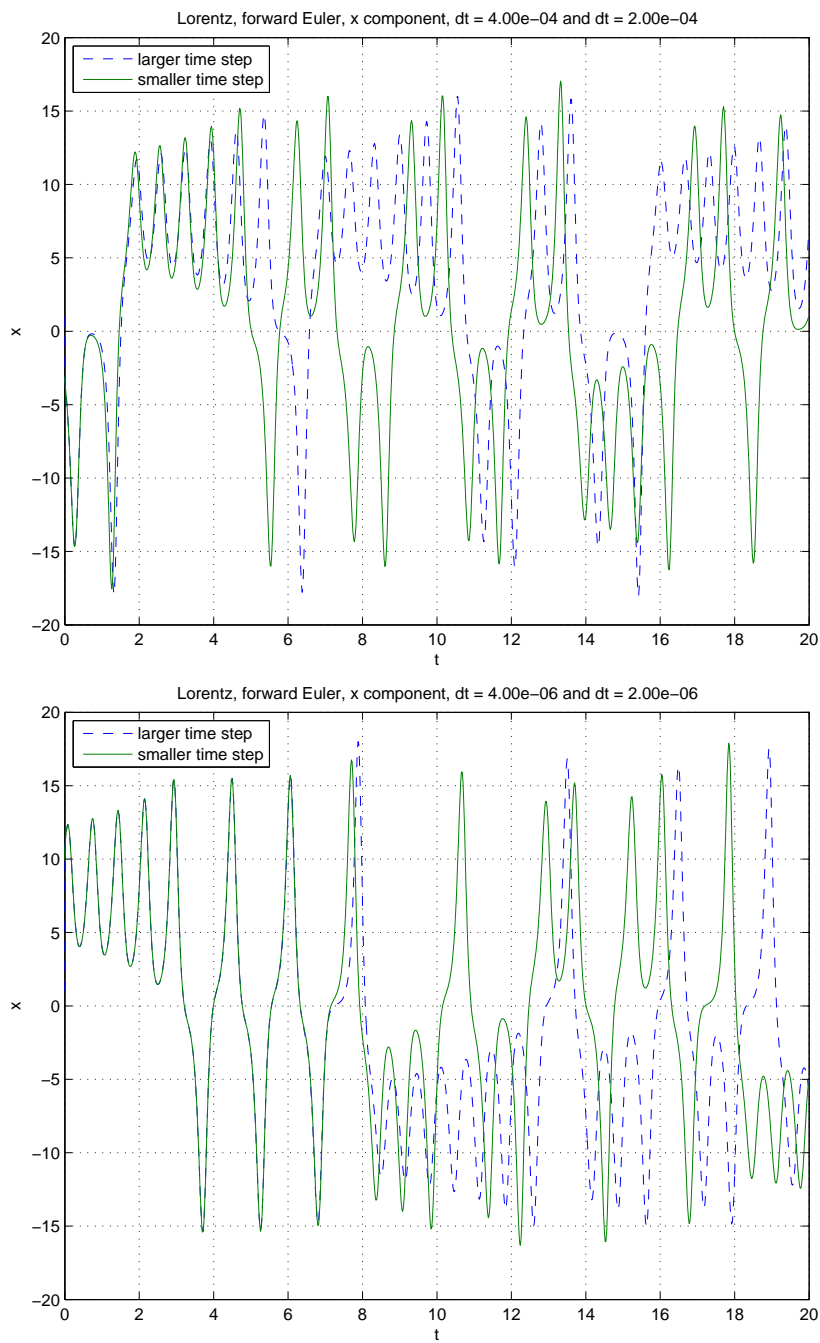


Figure 8.2: Two convergence studies for the Lorentz system. The time steps in the bottom figure are 100 times smaller than the time steps in the top figure. The more accurate calculation loses accuracy at $t \approx 10$, as opposed to $t \approx 5$ with a larger time step. The qualitative behavior is similar in all computations.

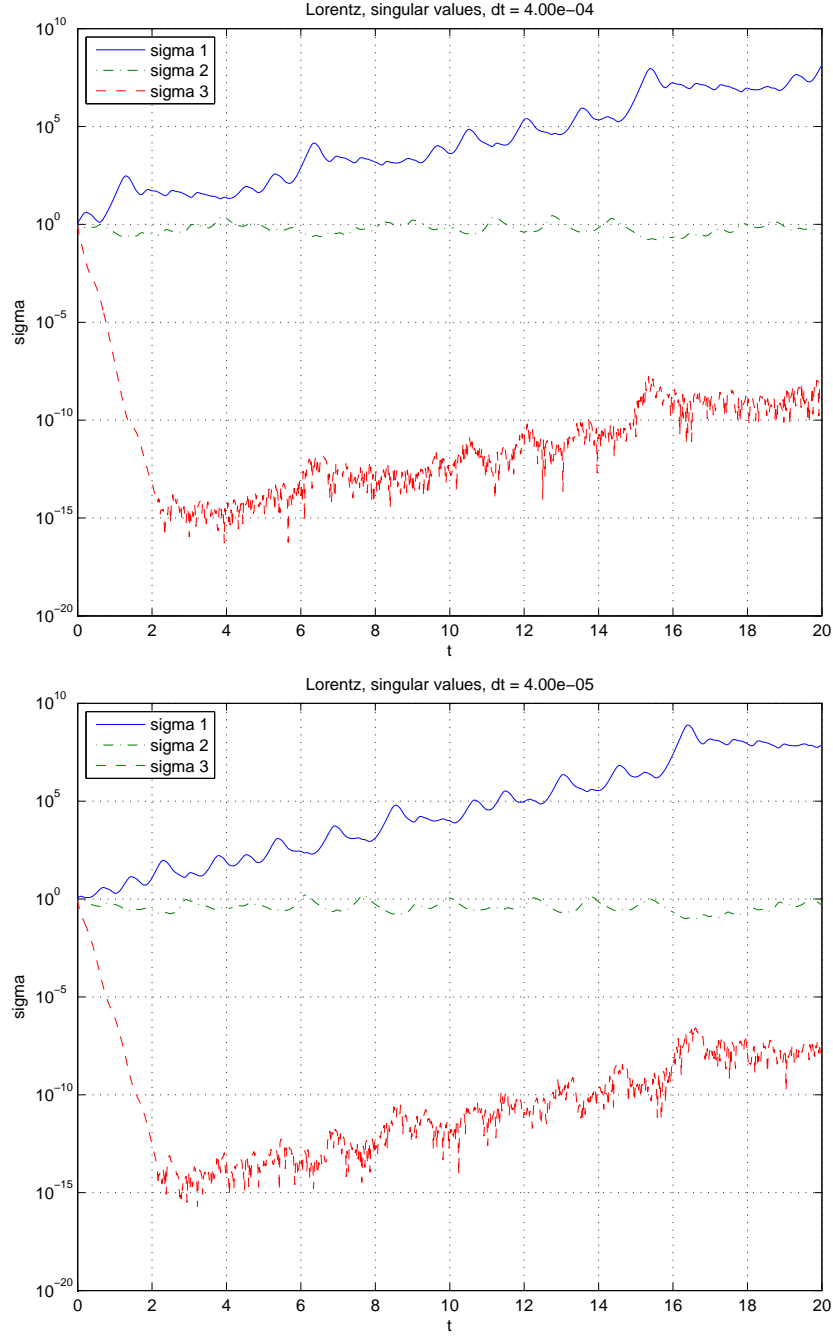


Figure 8.3: Computed singular values of the sensitivity matrix $A(x, t) = \partial_x S(x, t)$ with large time step (top) and ten times smaller time step (bottom). Top and bottom curves are similar qualitatively though the fine details differ. Theory predicts that middle singular value should be not grow or decay with t . The times from Figure 8.2 at which the numerical solution loses accuracy are not apparent here. In higher precision arithmetic, $\sigma_3(t)$ would have continued to decay exponentially. It is unlikely that computed singular values of any full matrix would differ by less than a factor of $1/\epsilon_{mach} \approx 10^{16}$.

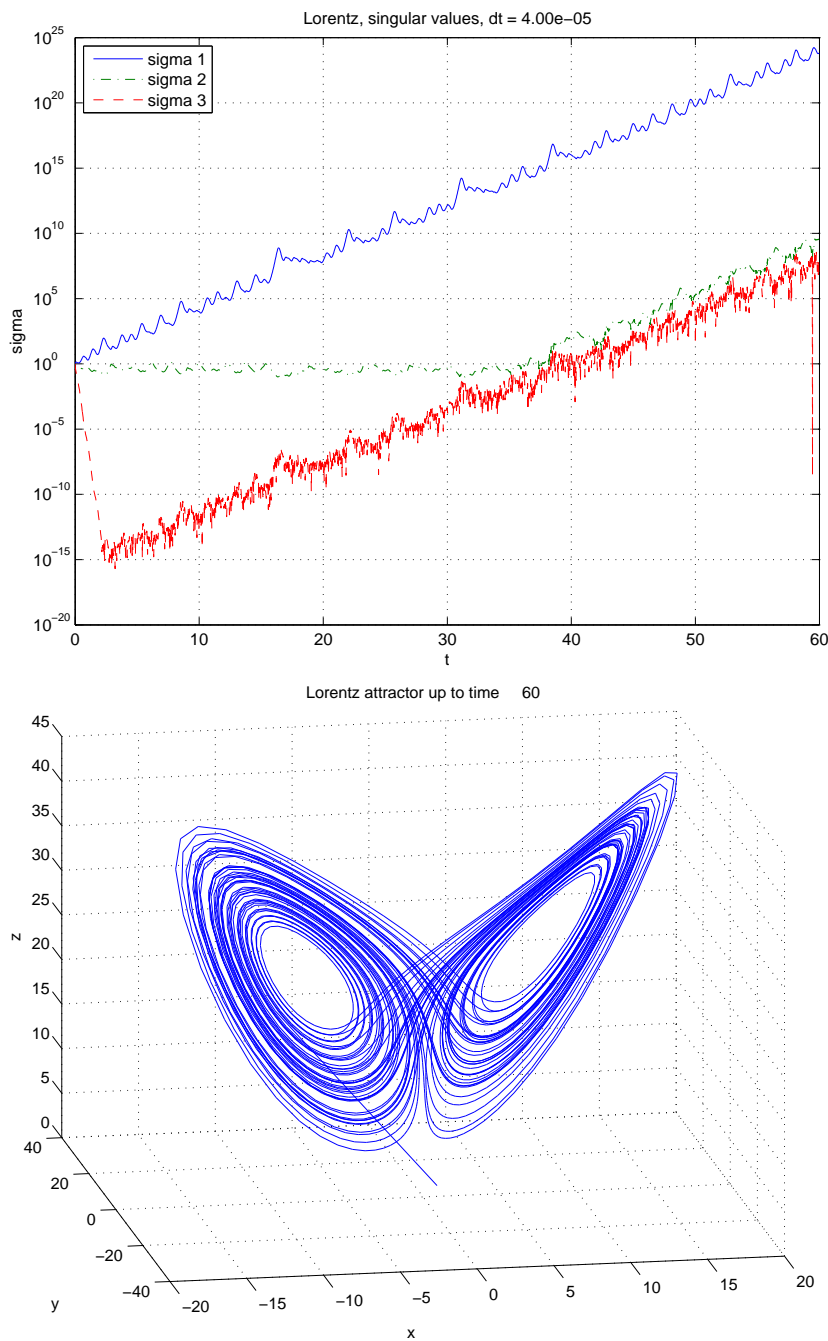


Figure 8.4: *Top: Singular values from Figure 8.3 computed for longer time. The $\sigma_1(t)$ grows exponentially, making a straight line on this log plot. The computed $\sigma_2(t)$ starts growing with the same exponential rate as σ_1 when roundoff takes over. A correct computation would show $\sigma_3(t)$ decreasing exponentially and $\sigma_2(t)$ neither growing nor decaying. Bottom: A beautiful picture of the butterfly shaped Lorenz attractor. It is just a three dimensional plot of the solution curve.*

Looking at the numbers, we get the overall impression that they are growing in an irregular way. The graph shows that the numbers have simple exponential growth with fine scale irregularities superposed. It could take hours to get that information directly from the numbers.

It can be a challenge to make visual sense of higher dimensional data. For example, we could make graphs of $x(t)$ (Figure 8.2), $y(t)$ and $z(t)$ as functions of t , but the single three dimensional plot in the lower frame of Figure 8.4 makes it clearer that the solution goes sometimes around the left wing and sometimes around the right. The three dimensional plot (`plot3` in Matlab) illuminates the structure of the solution better than three one dimensional plots.

There are several ways to visualize functions of two variables. A *contour plot* draws several *contour lines*, or *level lines*, of a function $u(x, y)$. A contour line for level u_k is the set of points (x, y) with $u(x, y) = u_k$. It is common to take about ten uniformly spaced values u_k , but most contour plot programs, including the Matlab program `contour`, allow the user to specify the u_k . Most contour lines are smooth curves or collections of curves. For example, if $u(x, y) = x^2 - y^2$, the contour line $u = u_k$ with $u_k \neq 0$ is a hyperbola with two components. An exception is $u_k = 0$, the contour line is an \times .

A *grid plot* represents a two dimensional rectangular array of numbers by colors. A *color map* assigns a color to each numerical value, that we call $c(u)$. In practice, usually we specify $c(u)$ by giving RGB values, $c(u) = (r(u), g(u), b(u))$, where r , g , and b are the intensities for red, green and blue respectively. These may be integers in a range (e.g. 0 to 255) or, as in Matlab, floating point numbers in the range from 0 to 1. Matlab uses the commands `colormap` and `image` to establish the color map and display the array of colors. The image is an array of boxes. Box (i, j) is filled with the color $c(u(i, j))$.

Surface plots visualize two dimensional surfaces in three dimensional space. The surface may be the graph of $u(x, y)$. The Matlab commands `surf` and `surfc` create surface plots of graphs. These look nice but often are harder to interpret than contour plots or grid plots. It also is possible to plot *contour surfaces* of a function of three variables. This is the set of (x, y, z) so that $u(x, y, z) = u_k$. Unfortunately, it is hard to plot more than one contour surface at a time.

Movies help visualize time dependent data. A movie is a sequence of frames, with each frame being one of the plots above. For example, we could visualize the Lorentz attractor with a movie that has the three dimensional butterfly together with a dot representing the position at time t .

The default in Matlab, and most other quality visualization packages, is to render the user's data as explicitly as possible. For example, the Matlab command `plot(u)` will create a piecewise linear "curve" that simply connects successive data points with straight lines. The plot will show the granularity of the data as well as the values. Similarly, the grid lines will be clearly visible in a color grid plot. This is good most of the time. For example, the bottom frame of Figure 8.4 clearly shows the granularity of the data in the wing tips. Since the curve is sampled at uniform time increments, this shows that the trajectory is moving faster at the wing tips than near the body where the wings meet.

Some plot packages offer the user the option of smoothing the data using

spline interpolation before plotting. This might make the picture less angular, but it can obscure features in the data and introduce artifacts, such as overshoots, not present in the actual data.

8.9 Resources and further reading

There is a beautiful discussion of computational methods for ordinary differential equations in *Numerical Methods* by Åke Björk and Germund Dahlquist. It was Dahlquist who created much of our modern understanding of the subject. A more recent book is *A First Course in Numerical Analysis of Differential Equations* by Arieh Iserles. The book *Numerical Solution of Ordinary Differential Equations* by Lawrence Shampine has a more practical orientation.

There is good public domain software for solving ordinary differential equations. A particularly good package is LSODE (google it).

The book by Sans-Serna explains symplectic integrators and their application to large scale problems such as the dynamics of large scale biological molecules. It is an active research area to understand the quantitative relationship between long time simulations of such systems and the long time behavior of the systems themselves. Andrew Stuart has written some thoughtful papers on the subject.

The numerical solution of partial differential equations is a vast subject with many deep specialized methods for different kinds of problems. For computing stresses in structures, the current method of choice seems to be *finite element* methods. For fluid flow and wave propagation (particularly nonlinear), the majority relies on finite difference and finite volume methods. For finite differences, the old book by Richtmeyer and Morton still merit though there are more up to date books by Randy LeVeque and by Bertil Gustavson, Heinz Kreiss, and Joe Oliger.

8.10 Exercises

1. We compute the second error correction $u_2(t)$ in (8.13). For simplicity, consider only the scalar equation ($n = 1$). Assuming the error expansion, we have

$$\begin{aligned} f(x_k) &= f(\tilde{x}_k + \Delta t u_1(t_k) + \Delta t^2 u_2(t_k) + O(\Delta t^3)) \\ &\approx f(\tilde{x}_k) + f'(\tilde{x}_k) (\Delta t u_1(t_k) + \Delta t^2 u_2(t_k)) \\ &\quad + \frac{1}{2} f''(\tilde{x}_k) \Delta t^2 u_1(t_k)^2 + O(\Delta t^3) . \end{aligned}$$

Also

$$\frac{x(t_k + \Delta t) - x(t_k)}{\Delta t} = \dot{x}(t_k) + \frac{\Delta t}{2} \ddot{x}(t_k) + \frac{\Delta t^2}{6} x^{(3)}(t_k) + O(\Delta t^3) ,$$

and

$$\Delta t \frac{u_1(t_k + \Delta t) - u_1(t_k)}{\Delta t} = \Delta t \dot{u}_1(t_k) + \frac{\Delta t^2}{2} \ddot{u}_1(t_k) + O(\Delta t^3) .$$

Now plug in (8.13) on both sides of (8.5) and collect terms proportional to Δt^2 to get

$$\dot{u}_2 = f'(x(t))u_2 + \frac{1}{6}x^{(3)}(t) + \frac{1}{2}f''(x(t))u_1(t)^2 + \dots .$$

2. This exercise confirms what was hinted at in Section 8.1, that (8.19) correctly predicts error growth even for t so large that the solution has lost all accuracy. Suppose $k = R/\Delta t^2$, so that $t_k = R/\Delta t$. The error equation (8.19) predicts that the forward Euler approximation x_k has grown by a factor of $e^R/2$ although the exact solution has not grown at all. We can confirm this by direct calculation. Write the forward Euler approximation to (8.18) in the form $x_{k+1} = Ax_k$, where A is a 2×2 matrix that depends on Δt . Calculate the eigenvalues of A up to second order in Δt : $\lambda_1 = 1 + i\Delta t + a\Delta t^2 + O(\Delta t^3)$, and $\lambda_2 = 1 - i\Delta t + b\Delta t^2 + O(\Delta t^3)$. Find the constants a and b . Calculate $\mu_1 = \ln(\lambda_1) = i\Delta t + c\Delta t^2 + O(\Delta t^3)$ so that $\lambda_1 = \exp(i\Delta t + c\Delta t^2 + O(\Delta t^3))$. Conclude that for $k = R/\Delta t^2$, $\lambda_1^k = \exp(k\mu_1) = e^{iR/\Delta t} e^{R/2 + O(\Delta t)}$, which shows that the solution has grown by a factor of nearly $e^{R/2}$ as predicted by (8.19). This s**t is good for something!
3. Another two stage second order Runge Kutta method sometimes is called the *modified Euler* method. The first stage uses forward Euler to predict the x value at the middle of the time step: $\xi_1 = \frac{\Delta t}{2} f(x_k, t_k)$ (so that $x(t_k + \Delta t/2) \approx x_k + \xi_1$). The second stage uses the midpoint rule with that estimate of $x(t_k + \Delta t/2)$ to step to time t_{k+1} : $x_{k+1} = x_k + \Delta t f(t_k + \frac{\Delta t}{2}, x_k + \xi_1)$. Show that this method is second order accurate.
4. Show that applying the four stage Runge Kutta method to the linear system (8.30) is equivalent to approximating the fundamental solution $S(\Delta t) = \exp(\Delta t A)$ by its Taylor series in Δt up to terms of order Δt^4 (see Exercise ??). Use this to verify that it is fourth order for linear problems.
5. Write a C/C++ program that solves the initial value problem for (8.1), with f independent of t , using a constant time step, Δt . The arguments to the initial value problem solver should be T (the final time), Δt (the time step), $f(x)$ (specifying the differential equations), n (the number of components of x), and x_0 (the initial condition). The output should be the approximation to $x(T)$. The code must do something to preserve the overall order of accuracy of the method in case T is not an integer multiple of Δt . The code should be able to switch between the three methods, forward Euler, second order Adams Bashforth, forth order four

state Runge Kutta, with a minimum of code editing. Hand in output for each of the parts below showing that the code meets the specifications.

- (a) The procedure should return an error flag or notify the calling routine in some way if the number of time steps called for is negative or impossibly large.
 - (b) For each of the three methods, verify that the coding is correct by testing that it gets the right answer, $x(.5) = 2$, for the scalar equation $\dot{x} = x^2$, $x(0) = 1$.
 - (c) For each of the three methods and the test problem of part b, do a convergence study to verify that the method achieves the theoretical order of accuracy. For this to work, it is necessary that T should be an integer multiple of Δt .
 - (d) Apply your code to problem (8.18) with initial data $x_0 = (1, 0)^*$. Repeat the convergence study with $T = 10$.
6. Verify that the recurrence relation (8.39) is unstable.
- (a) Let z be a complex number. Show that the sequence $x_k = z^k$ satisfies (8.39) if and only if z satisfies $0 = p(z) = z^3 + \frac{3}{2}z^2 - 3z + \frac{1}{2}$.
 - (b) Show that $x_k = 1$ for all k is a solution of the recurrence relation. Conclude that $z = 1$ satisfies $p(1) = 0$. Verify that this is true.
 - (c) Using polynomial division (or another method) to factor out the known root $z = 1$ from $p(z)$. That is, find a quadratic polynomial, $q(z)$, so that $p(z) = (z - 1)q(z)$.
 - (d) Use the quadratic formula and a calculator to find the roots of q as $z = \frac{-5}{4} \pm \sqrt{\frac{41}{16}} \approx -2.85, .351$.
 - (e) Show that the general formula $x_k = az_1^k + bz_2^k + cz_3^k$ is a solution to (8.39) if z_1, z_2 , and z_3 are three roots $z_1 = 1, z_2 \approx -2.85, z_3 \approx .351$, and, conversely, the general solution has this form. Hint: we can find a, b, c by solving a vanderMonde system (Section 7.4) using x_0, x_1 , and x_2 .
 - (f) Assume that $|x_0| \leq 1, |x_1| \leq 1$, and $|x_2| \leq 1$, and that b is on the order of double precision floating point roundoff (ϵ_{mach}) relative to a and c . Show that for $k > 80$, x_k is within ϵ_{mach} of bz_2^k . Conclude that for $k > 80$, the numerical solution has nothing in common with the actual solution $x(t)$.
7. Applying the implicit trapezoid rule (8.47) to the scalar model problem (8.31) results in $u_{k+1} = m(\lambda \Delta t)u_k$. Find the formula for m and show that $|m| \leq 1$ if $\text{Re}(\lambda) \leq 0$, so that $|u_{k+1}| \leq |u_k|$. What does this say about the applicability of the trapezoid rule to stiff problems?
8. Exercise violating time step constraint.

9. Write an adaptive code in C/C++ for the initial value problem (8.1) (8.2) using the method described in Section 8.4 and the four stage fourth order Runge Kutta method. The procedure that does the solving should have arguments describing the problem, as in Exercise 5, and also the local truncation error level, e . Apply the method to compute the trajectory of a comet. In nondimensionalized variables, the equations of motion are given by the inverse square law:

$$\frac{d^2}{dt^2} \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} = \frac{-1}{(r_1^2 + r_2^2)^{3/2}} \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} .$$

We always will suppose that the comet starts at $t = 0$ with $r_1 = 10$, $r_2 = 0$, $\dot{r}_1 = 0$, and $\dot{r}_2 = v_0$. If v_0 is not too large, the point $r(t)$ traces out an ellipse in the plane¹⁵. The shape of the ellipse depends on v_0 . The *period*, $P(v_0)$, is the first time for which $r(P) = r(0)$ and $\dot{r}(P) = \dot{r}(0)$. The solution $r(t)$ is *periodic* because it has a period.

- Verify the correctness of this code by comparing the results to those from the fixed time step code from Exercise 5 with $T = 30$ and $v_0 = .2$.
- Use this program, with a small modification to compute $P(v_0)$ in the range $.01 \leq v_0 \leq .5$. You will need a criterion for telling when the comet has completed one orbit since it will not happen that $r(P) = r(0)$ exactly. Make sure your code tests for and reports failure¹⁶.
- Choose a value of v_0 for which the orbit is rather but not extremely elliptical (width about ten times height). Choose a value of e for which the solution is rather but not extremely accurate – the error is small but shows up easily on a plot. If you set up your environment correctly, it should be quick and easy to find suitable parameters by trial and error using Matlab graphics. Make a plot of a single period with two curves on one graph. One should be a solid line representing a highly accurate solution (so accurate that the error is smaller than the line width – *plotting accuracy*), and the other being the modestly accurate solution, plotted with a little “o” for each time step. Comment on the distribution of the time step points.
- For the same parameters as part b, make a single plot of that contains three curves, an accurate computation of $r_1(t)$ as a function of t (solid line), a modestly accurate computation of r_1 as a function of t (“o” for each time step), and Δt as a function of t . You will need to use a different scale for Δt if for no other reason than that it has different units. Matlab can put one scale in the left and a different scale on

¹⁵Isaac Newton formulated these equations and found the explicit solution. Many aspects of planetary motion – elliptical orbits, sun at one focus, $|r|\dot{\theta} = \text{const}$ – had been discovered observationally by Johannes Kepler. Newton’s inverse square law *theory* fit Kepler’s *data*.

¹⁶This is not a drill.

the right. It may be necessary to plot Δt on a log scale if it varies over too wide a range.

- (e) Determine the number of adaptive time stages it takes to compute $P(.01)$ to .1% accuracy (error one part in a thousand) and how many fixed Δt time step stages it takes to do the same. The counting will be easier if you do it within the function f .
10. The vibrations of a two dimensional crystal lattice may be modelled in a crude way using the differential equations¹⁷

$$\ddot{r}_{jk} = r_{j-1,k} + r_{j+1,k} + r_{j,k-1} + r_{j,k+1} - 4r_{jk} . \quad (8.51)$$

Here $r_{jk}(t)$ represents the displacement (in the vertical direction) of an atom at the (j, k) location of a square crystal lattice of atoms. Each atom is *bonded* to its four neighbors and is pulled toward them with a linear force. A lattice of size L has $1 \leq j \leq L$ and $1 \leq k \leq L$. Apply *reflecting* boundary conditions along the four boundaries. For example, the equation for $r_{1,k}$ should use $r_{0,k} = r_{1,k}$. This is easy to implement using a *ghost cell* strategy. Create ghost cells along the boundary and copy appropriate values from the actual cells to the ghost cells before each evaluation of f . This allows you to use the formula (8.51) at every point in the lattice. Start with initial data $r_{jk}(0) = 0$ and $\dot{r}_{jk}(0) = 0$ for all j, k except $\dot{r}_{1,1} = 1$. Compute for $L = 100$ and $T = 300$. Use the fourth order Runge Kutta method with a fixed time step $\Delta t = .01$. Write the solution to a file every .5 time units then use Matlab to make a movie of the results, with a 2D color plot of the solution in each frame. The movie should be a circular wave moving out from the bottom left corner and bouncing off the top and right boundaries. There should be some beautiful wave patterns inside the circle that will be hard to see far beyond time $t = 100$. Hand in a few of your favorite stills from the movie. If you have a web site, post your movie for others to enjoy.

¹⁷This is one of Einstein's contributions to science.

Chapter 9

Monte Carlo methods

Monte Carlo means using random numbers in scientific computing. More precisely¹, it means using random numbers as a tool to compute something that is not random. We emphasize this point by distinguishing between Monte Carlo and *simulation*. Simulation means producing random variables with a certain distribution just to look at them. For example, we might have a model of a random process that produces clouds. We could simulate the model to generate cloud pictures, either out of scientific interest or for computer graphics. As soon as we start asking quantitative questions about, say, the average size of a cloud or the probability that it will rain, we move from pure simulation to Monte Carlo.

Many Monte Carlo computations have the following form. We express A , the quantity of interest, in the form

$$A = E_f [V(X)] . \quad (9.1)$$

This notation means² means that X is a random variable or collection of random variables whose joint distribution is given by f , and $V(x)$ is some scalar quantity determined by x . For example, x could be the collection of variables describing a random cloud and $V(x)$ could be the total moisture. The notation $X \sim f$ means that X has the distribution f , or is a *sample* or f .

A Monte Carlo code generates a large number of samples of f , written X_k , for $k = 1, \dots, N$. Sections 9.2 and 9.3 describe how this is done. The approximate answer is

$$A \approx \hat{A}_N = \frac{1}{N} \sum_{k=1}^N V(X_k) . \quad (9.2)$$

Section 9.4 discusses the error in (9.2). The error generally is of the order of $N^{-1/2}$, which means that it takes roughly four times the computation time to double the accuracy of a computation. By comparison, the least accurate integration method in Chapter 3 is the first order rectangle rule, with error roughly proportional to $1/N$.

The total error is the sum of the *bias* (error in the expected value of the estimator),

$$e_b = E [\hat{A}_N] - A ,$$

and the *statistical error*,

$$e_s = \hat{A}_N - E [\hat{A}_N] .$$

The equation (9.1) states that the estimator (9.2) is *unbiased*, which means that the error is purely statistical. If the samples X_k are independent of each other, the error is roughly proportional to the standard deviation, which is the square root of the variance of the random variable $V(X)$ when $X \sim f$:

$$\sigma = \sqrt{\text{var}_f(V(X))} . \quad (9.3)$$

¹This helpful definition is given in the book by Mal Kalos and Paula Whitlock.

²The notation used here is defined in Section 9.1 below.

More sophisticated Monte Carlo estimators have bias as well as statistical error, but the statistical error generally is much larger. For example, Exercise 1 discusses estimators that have statistical error of order $1/N^{1/2}$ and bias of order $1/N$.

There may be other ways to express A in terms of a random variable. That is, there may be a random variable, $Y \sim g$, with distribution g and a function W so that $A = E_g[W(Y)]$. This is *variance reduction* if $\text{var}_g[W(Y)] < \text{var}_f[V(X)]$. We distinguish between Monte Carlo and simulation to emphasize that variance reduction in Monte Carlo is possible. Even simple variance reduction strategies can make a big difference in computation time.

Given a choice between Monte Carlo and deterministic methods, the deterministic method usually is better. The large $O(1/\sqrt{N})$ statistical errors in Monte Carlo usually make it slower and less accurate than a deterministic method. For example, if X is a one dimensional random variable with probability density $f(x)$, then

$$A = E[V(X)] = \int V(x)f(x)dx .$$

Estimating the integral by deterministic panel methods as in Section ?? (error roughly proportional to $1/N$ for a first order method, $1/N^2$ for second order, etc.) usually estimates A to a given accuracy at a fraction of the cost of Monte Carlo. Deterministic methods are better than Monte Carlo in most situations where the deterministic methods are practical.

We use Monte Carlo because of the *curse of dimensionality*. The curse is that the work to solve a problem in many dimensions may grow exponentially with the dimension. For example, consider integrating over ten variables. If we use twenty points in each coordinate direction, the number of integration points is $20^{10} \approx 10^{13}$, which is on the edge of what a computer can do in a day. A Monte Carlo computation might reach the same accuracy with only, say, 10^6 points. There are many high dimensional problems that can be solved, as far as we know, only by Monte Carlo.

Another curse of high dimensions is that some common intuitions fail. This can lead to Monte Carlo algorithms that may work in principle but also are exponentially slow in high dimensions. An example is in Section 9.3.7. The probability that a point in a cube falls inside a ball is less than 10^{-20} in 40 dimensions, which is to say that it will never happen even on a teraflop computer.

One favorable feature of Monte Carlo is that it is possible to estimate the order of magnitude of statistical error, which is the dominant error in most Monte Carlo computations. These estimates are often called *error bars* because they are indicated as bars on plots. Error bars are statistical confidence intervals, which rely on elementary or sophisticated statistics depending on the situation. It is dangerous and unprofessional to do a Monte Carlo computation without calculating error bars. The practitioner should know how large the error is likely to be, even if he or she does not include that information in non-technical presentations.

In Monte Carlo, simple clever ideas can lead to enormous practical improvements in efficiency and accuracy. This is the main reason I emphasize so strongly that, while A is given, the algorithm for estimating it is not. The search for more accurate alternative algorithms is often called “variance reduction”. Common variance reduction techniques are importance sampling, antithetic variates, and control variates.

Many of the examples below are somewhat artificial because I have chosen not to explain specific applications. The techniques and issues raised here in the context of toy problems are the main technical points in many real applications. In many cases, we will start with the probabilistic definition of A , while in practice, finding this is part of the problem. There are some examples in later sections of choosing alternate definitions of A to improve the Monte Carlo calculation.

9.1 Quick review of probability

This is a quick review of the parts of probability needed for the Monte Carlo material discussed here. Please skim it to see what notation we are using and check Section 9.7 for references if something is unfamiliar.

9.1.1 Probabilities and events

Probability theory begins with the assumption that there are *probabilities* associated to *events*. Event B has probability $\Pr(B)$, which is a number between 0 and 1 describing what fraction of the time event B would happen if we could repeat the *experiment* many times. $\Pr(B) = 0$ and $\Pr(B) = 1$ are allowed. The exact meaning of probabilities is debated at length elsewhere.

An event is a set of possible outcomes. The set of all possible outcomes is called Ω and particular outcomes are called ω . Thus, an event is a subset of the set of all possible outcomes, $B \subseteq \Omega$. For example, suppose the experiment is to toss four coins (a penny, a nickel, a dime, and a quarter) on a table and record whether they were face up (*heads*) or face down (*tails*). There are 16 possible outcomes. The notation $THTT$ means that the penny was face down (tails), the nickel was up, and the dime and quarter were down. The event “all heads” consists of a single outcome, $HHHH$. The event $B =$ “more heads than tails” consists of the five outcomes: $B = \{HHHH, THHH, HTHH, HHTH, HHHT\}$.

The basic set operations apply to events. For example the intersection of events B and C is the set of outcomes both in B and in C : $\omega \in B \cap C$ means $\omega \in B$ and $\omega \in C$. For that reason, $B \cap C$ represents the event “ B and C ”. For example, if B is the event “more heads than tails” above and C is the event “the dime was heads”, then C has 8 outcomes in it, and $B \cap C = \{HHHH, THHH, HTHH, HHHT\}$. The set union, $B \cup C$ is $\omega \in B \cup C$ if $\omega \in B$ or $\omega \in C$, so we call it “ B or C ”. One of the axioms of probability is

$$\Pr(B \cup C) = \Pr(B) + \Pr(C) \quad , \quad \text{if } B \cap C \text{ is empty.}$$

This implies the intuitive fact that

$$B \subset C \implies \Pr(B) \leq \Pr(C) ,$$

for if D is the event “ C but not B ”, then $B \cap D$ is empty, so $\Pr(C) = \Pr(B) + \Pr(D) \geq \Pr(B)$. We also suppose that $\Pr(\Omega) = 1$, which means that Ω includes every possible outcome.

The *conditional* “probability of B given C ” is given by Bayes’ formula:

$$\Pr(B | C) = \frac{\Pr(B \cap C)}{\Pr(C)} = \frac{\Pr(B \text{ and } C)}{\Pr(C)} \quad (9.4)$$

Intuitively, this is the probability that a C outcome also is a B outcome. If we do an experiment and $\omega \in C$, $\Pr(B | C)$ is the probability that $\omega \in B$. The right side of (9.4) represents the fraction of C outcomes that also are in B . We often know $\Pr(C)$ and $\Pr(B | C)$, and use (9.4) to calculate $\Pr(B \cap C)$. Of course, $\Pr(C | C) = 1$. Events B and C are *independent* if $\Pr(B) = \Pr(B | C)$, which is the same as $\Pr(B \cap C) = \Pr(B)\Pr(C)$, so B being independent of C is the same as C being independent of B .

The *probability space* Ω is finite if it is possible to make a finite list³ of all the outcomes in Ω : $\Omega = \{\omega_1, \omega_2, \dots, \omega_n\}$. The space is *countable* if it is possible to make a possibly infinite list of all the elements of Ω : $\Omega = \{\omega_1, \omega_2, \dots, \omega_n, \dots\}$. We call Ω *discrete* in both cases. When Ω is discrete, we can specify the probabilities of each outcome

$$f_k = \Pr(\omega = \omega_k) .$$

Then an event B has probability

$$\Pr(B) = \sum_{\omega_k \in B} \Pr(\omega_k) = \sum_{\omega_k \in B} f_k .$$

9.1.2 Random variables and distributions

A discrete random variable⁴ is a number, $X(\omega)$, that depends on the random outcome, ω . In the coin tossing example, $X(\omega)$ could be the number of heads. The *expected value* is (defining $x_k = X(\omega_k)$)

$$E[X] = \sum_{\omega \in \Omega} X(\omega) \Pr(\omega) = \sum_{\omega_k} x_k f_k .$$

The *distribution* of a discrete random variable is determined by the values x_k and the probabilities f_k . Two discrete random variables have the same distribution if they have the same x_k and f_k . It often happens that the possible values x_k are understood and only the probabilities f_k are unknown, for example, if X is known to be an integer. We say that f (the numbers f_k) is the *distribution* of X if $\Pr(X = x_k) = f_k$. We write this as $X \sim f$, and say that X is a *sample*

³Warning: this list is impractically large even in common simple applications.

⁴Warning: sometimes ω is the random variable and X is a function of a random variable.

of f . If X' is another random variable, we write $X' \sim X$ and say that X' is a sample of X if X' and X have the same distribution.

Random variables X' and X are *independent* if any event determined by X' is independent of any event determined by X . If X' and X are independent, then they are *uncorrelated*, which means that $E[X'X] = E[X']E[X]$. The converse is not true (example below). Let X_1, \dots, X_N be a sequence of random variables. We say the X_j are independent if any collection of them is independent of any disjoint collection. This is not the same as saying that X_j is independent of X_i whenever $j \neq i$. We say the X_j are *i.i.d.* (for “independent, identically distributed”) samples of X if they are independent and $X_j \sim X$ for each j .

A complicated random variable or collection of random variables may be called an *experiment*. For example, we might consider the experiment of generating N independent samples of X and dividing the largest one by the smallest. More precisely, let X_j for $1 \leq j \leq N$ be i.i.d. samples of X , define $X_{\max} = \max_{1 \leq j \leq N} X_j$ and $X_{\min} = \min_{1 \leq j \leq N} X_j$, and $Y = X_{\max}/X_{\min}$. *Repeating* the experiment means generating a Y' that is an independent sample of Y .

A *bivariate* random variable is a two component vector, $X = (X_1, X_2)$. The distribution of X is determined by the probabilities $f_{jk} = \Pr(X_1 = x_j \text{ and } X_2 = x_k)$. A random variable with two or more components is *multicomponent*. Bivariate random variables X and X' have the same distribution if $f_{jk} = f'_{jk}$ for all j, k .

A *continuous* random variable may be described by a *probability density*, $f(x)$, written $X \sim f$. We say $X \sim f$ if $\Pr(X \in B) = \int_B f(x)dx$ for any measurable⁵ set B . If $X = (X_1, \dots, X_n)$ is a multivariate random variable with n components, then $B \subseteq R^n$. If $V(x)$ is a function of n variables, then $E[V(X)] = \int_{R^n} V(x)f(x)dx$.

If X is an n component random variable and Y is an m component random variable, then the pair (X, Y) is a random variable with $n + m$ components. The probability density $h(x, y)$ is the *joint density* of X and Y . The *marginal* densities f and g , with $X \sim f$ and $Y \sim g$, are given by $f(x) = \int h(x, y)dy$, and $g(y) = \int h(x, y)dx$. The random variables X and Y , are independent if and only if $h(x, y) = f(x)g(y)$. The scalar random variables $V = V(X)$ and $W = W(Y)$ are uncorrelated if $E[VW] = E[V]E[W]$. The random variables X and Y are independent if and only if $V(X)$ is independent of $W(Y)$ for any (measurable) functions V and W .

If X is a univariate random variable with probability density $f(x)$, then it has *CDF*, or *cumulative distribution function* (or just *distribution function*), $F(x) = \Pr(X \leq x) = \int_{x' \leq x} f(x')dx'$. The distribution function of a discrete random variable is $F(x) = \Pr(X \leq x) = \sum_{x_k \leq x} f_k$. Two random variables have the same distribution if they have the same distribution function. It may be convenient to calculate probability density $f(x)$ by first calculating the distribution

⁵It is an annoying technicality in probability theory that there may be sets or functions that are so complicated that they are not *measurable*. If the set B is not measurable, then we do not define $\Pr(B)$. None of the functions or sets encountered in applied probability fail to be measurable.

function, $F(x)$, then using

$$f(x)dx = \int_x^{x+dx} f(x')dx' = F(x+dx) - F(x) = \Pr(x \leq X \leq x+dx) .$$

If $V(X)$ is a multicomponent function of the possibly multicomponent random variable X , then $E[V(X)]$ is the vector whose components are the expected values of the components of V . We write this in vector notation simply as

$$\mu = E[V(X)] = \int_{x \in R^n} V(x)f(x)dx .$$

Here Vf is the product of the vector, V with the scalar, f .

For a scalar random variable, the *variance* is

$$\sigma^2 = \text{var}[X] = E[(X - \mu)^2] = \int_{-\infty}^{\infty} (x - \mu)^2 f(x)dx = E[X^2] - \mu^2 .$$

The variance is non-negative, and $\sigma^2 = 0$ only if X is not random. For a multicomponent random variable, the *covariance* matrix is the symmetric $n \times n$ matrix,

$$C = E[(X - \mu)(X - \mu)^*] = \int_{R^n} (x - \mu)(x - \mu)^* f(x)dx . \quad (9.5)$$

The diagonal entries of C are the variances

$$C_{jj} = \sigma_j^2 = \text{var}[X_j] = E[(X_j - \mu_j)^2] .$$

The off diagonal entries are the covariances

$$C_{jk} = E[(X_j - \mu_j)(X_k - \mu_k)] = \text{cov}[X_j, X_k] .$$

The covariance matrix is positive semidefinite in the sense that for any $y \in R^n$, $y^*Cy \geq 0$. To see this, for any y , we can define the scalar random variable $V = y^*X$. Up to a scaling, V is the component of X in the direction y . Then we have the formula

$$y^*Cy = \text{var}[V] \geq 0 .$$

If $y \neq 0$ and $y^*Cy = \text{var}[V] = 0$, then $V = y^*X$ is equal to $y^*\mu$ and is not random. This means that either C is positive definite or X always lies on the hyperplane $y^*X = C = y^*\mu$. The formula follows from (9.5)⁶:

$$\begin{aligned} y^*Cy &= y^* (E[(X - \mu)(X - \mu)^*]) y \\ &= E[y^*(X - \mu)((X - \mu)^*y)] \\ &= E[y^*XX^*\mu] - (y^*\mu\mu^*y) \\ &= E[V^2] - (E[V])^2 \\ &= \text{var}[V] . \end{aligned}$$

⁶We also use the fact that $y^*XX^*y = (y^*X)(X^*y) = (y^*X)^2$, by the associativity of matrix and vector multiplication and the fact that y^*X is a scalar.

9.1.3 Common random variables

There are several probability distributions that are particularly important in Monte Carlo. They arise in common applications, and they are useful technical tools.

Uniform

A *uniform* random variable is equally likely to anywhere inside its allowed range. In one dimension, a univariate random variable U is uniformly distributed in the interval $[a, b]$ if it has probability density $f(u)$ with $f(u) = 0$ for $u \notin [a, b]$ and $f(x) = 1/(b - a)$ if $a \leq x \leq b$. The *standard uniform* has $b = 1$ and $a = 0$. The standard uniform is important in Monte Carlo because that is what random number generators try to produce. In higher dimensions, let B be a set with area b (for $n = 2$) or volume b (for $n \geq 3$). The n component random variable X is uniformly distributed if its probability density has the value $f(x) = b$ for $x \in B$ and $f(x) = 0$ otherwise.

Exponential

The exponential random variable, T , with *rate constant* $\lambda > 0$, has probability density,

$$f(t) = \begin{cases} \lambda e^{-\lambda t} & \text{if } 0 \leq t \\ 0 & \text{if } t < 0. \end{cases} \quad (9.6)$$

The exponential is a model for the amount of time something until happens (e.g. how long a light bulb will function before it breaks). It is characterized by the *Markov property*, if it has not happened by time t , it is as good as new. To see this, let $\lambda = f(0)$, which means $\lambda dt = \Pr(0 \leq T \leq dt)$. The random time T is exponential if all the conditional probabilities are equal to this:

$$\Pr(t \leq T \leq t + dt \mid T \geq t) = \Pr(T \leq dt) = \lambda dt .$$

Using Bayes' rule (9.4), and the observation that $\Pr(T \leq t + dt \text{ and } T \geq t) = f(t)dt$, this becomes

$$\lambda dt = \frac{f(t)dt}{\Pr(T > t)} = \frac{f(t)dt}{1 - F(t)} ,$$

which implies $\lambda(1 - F(t)) = f(t)$. Differentiating this gives $-\lambda f(t) = f'(t)$, which implies that $f(t) = Ce^{-\lambda t}$ for $t > 0$. We find $C = \lambda$ using $\int_0^\infty f(t)dt = 1$.

Independent exponential *inter arrival* times generate the *Poisson arrival processes*. Let T_k , for $k = 1, 2, \dots$ be independent exponentials with rate λ . The k^{th} arrival time is

$$S_k = \sum_{j \leq k} T_j .$$

The expected number of arrivals in interval $[t_1, t_2]$ is $\lambda(t_2 - t_1)$ and all arrivals are independent. This is a fairly good model for the arrivals of telephone calls at a large phone bank.

Gaussian, or normal

We denote the *standard normal* by Z . The standard normal has PDF

$$f(z) = \frac{1}{\sqrt{2\pi}} e^{-z^2/2} . \quad (9.7)$$

The general normal with mean μ and variance σ^2 is given by $X = \sigma Z + \mu$ and has PDF

$$f(x) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-(x-\mu)^2/2\sigma^2} . \quad (9.8)$$

We write $X \sim \mathcal{N}(\mu, \sigma^2)$ in this case. A standard normal has distribution $\mathcal{N}(0, 1)$. An n component random variable, X , is a *multivariate normal* with mean μ and covariance C it has probability density⁷

$$f(x) = \frac{1}{Z} \exp(-(x - \mu)^* H (x - \mu)/2) , \quad (9.9)$$

where $H = C^{-1}$ and the *normalization constant* is (we don't use this) $Z = 1/\sqrt{(2\pi)^n \det(C)}$. The reader should check that this is the same as (9.8) when $n = 1$.

The class of multivariate normals has the *linear transformation property*. Suppose L is an $m \times n$ matrix with rank m (L is a linear transformation from R^n to R^m that is onto R^m). If X is an n dimensional multivariate normal, then Y is an m dimensional multivariate normal. The covariance matrix for Y is given by

$$C_Y = LC_X L^* . \quad (9.10)$$

We derive this, taking $\mu = 0$ without loss of generality (why?), as follows:

$$C_Y = E[YY^*] = E[(LX)((LX)^*)] = E[LXX^*L^*] = LE[XX^*]L^* = LC_X L^* .$$

9.1.4 Limit theorems

The two important theorems for simple Monte Carlo are the *law of large numbers* and the *central limit theorem*. Suppose $A = E[X]$, that the X_k for $k = 1, 2, \dots$ are iid samples of X , and \hat{A}_N is the basic estimator (9.2). The *law of large numbers* states⁸ that $\hat{A}_N \rightarrow A$ as $N \rightarrow \infty$.

More generally, let \hat{A}_N be a family of Monte Carlo estimators of A . (See Exercise 1 for other estimators.) As usual, the \hat{A}_N are random but A is not. The estimators are *consistent* if $\hat{A}_N \rightarrow A$ as $N \rightarrow \infty$. The law of large numbers states that the estimators (9.2) are consistent.

⁷The Z here is a normalization constant, not a standard normal random variable. The double use of Z is unfortunately standard.

⁸More precisely, the *Kolmogorov strong law* of large numbers states that $\lim_{N \rightarrow \infty} \hat{A}_N = A$ almost surely, i.e. that the probability of the limit not existing or being the wrong answer is zero. More useful for us is the *weak law*, which states that, for any $\epsilon > 0$, $P(|\hat{A}_N - A| > \epsilon) \rightarrow 0$ as $N \rightarrow \infty$. These are related but they are not the same.

Let V be a random variable with mean A and variance σ^2 . Let V_k be independent samples of V , and $\hat{A}_N = \frac{1}{N} \sum_{k=1}^N V_k$, and $R_N = \hat{A}_N - A$. A simple calculation shows that $E[R_N] = 0$ and $\text{var}[R_N] = \sigma^2/N$. The central limit theorem states that for large N , the distribution of R_N is approximately normal with that mean and variance. That is to say that the distribution of R_N is approximately the same as that of $\sigma Z/\sqrt{N}$ with $Z \sim \mathcal{N}(0, 1)$. The same relation may be expressed as saying that R_N has approximately the same distribution as $(\sigma/\sqrt{N}) Z$. In this form, it says that the error in a Monte Carlo computation is of the order of $1/\sqrt{N}$ with a prefactor of σ and a random, but Gaussian, contribution. See Section 9.4.

9.1.5 Markov chains

A discrete time Markov chain with m states is characterized by an $m \times m$ *transition matrix*, P . The *time* variable, t , is an integer. The *state* at time t is $X(t)$, which is one of the numbers $1, \dots, m$. The entries of P are the *transition probabilities*

$$p_{jk} = \Pr(X(t+1) = k \mid X(t) = j) . \quad (9.11)$$

Once the initial state, $X(0)$, is given, we simulate the Markov chain by choosing successively $X(1)$, then $X(2)$, etc. satisfying the transition probabilities (9.11).

A continuous time Markov chain with m states is given by an $m \times m$ matrix of *transition rates*, r_{jk} . The time variable, t , is continuous and $X(t)$ is the state at time t . If $j \neq k$ then

$$r_{jk} = \text{Rate}(j \rightarrow k) , \quad r_{jk} dt = \Pr((X(t+dt) = k \mid X(t) = j) . \quad (9.12)$$

We usually define the diagonal entries of the transition rate matrix to be $r_{jj} = -\sum_{k \neq j} r_{jk}$. With this convention, the *occupation probabilities*, $f_j(t) = \Pr(X(t) = j)$, satisfy

$$\frac{df_j}{dt} = \sum_{k=1}^m f_k(t) r_{kj} .$$

A common way to simulate a continuous time Markov chain is to use the embedded discrete time chain. Suppose the state at time t is $X(t) = j$. Then the time until the next transition is an exponential random variable with rate $\lambda_j = \sum_{k \neq j} r_{jk}$. The probability that this transition is to state k is

$$p_{jk} = \frac{r_{jk}}{\sum_{l \neq j} r_{jl}} . \quad (9.13)$$

To simulate the continuous chain, we start with the given $j_0 = X(0)$. We choose T_1 to be an exponential with rate parameter λ_{j_0} . We set $S_1 = T_1$ and choose⁹ $X(S_1) = k = j_1$ according to the probabilities (9.13). Suppose the time of

⁹If S_i is the time of the i^{th} transition, we set $X(S_i)$ to be the new state rather than the old one. This convention is called *CADLAG*, from the French *continue a droite, limite a gauche* (continuous from the right, limit from the left).

the n^{th} transition is S_n and the transition takes us to state j_n . The time to the next transition is T_{n1} , which is an exponential with rate parameter λ_{j_n} . The next state is $j_{n+1} = k$ with probability $p_{j_n,k}$. The next transition time is $S_{n+1} = S_n + T_n$.

9.2 Random number generators

The random variables used in Monte Carlo are generated by a (*pseudo*) *random number generator*. The procedure `double rng()` is a perfect random number generator if

```
for( k=0; k<n; k++ ) U[k] = rng();
```

produces an array of iid standard uniform random variables. The best available random number generators are nearly perfect in this sense for most Monte Carlo applications. The native C/C++ procedure `random()` is good enough for most Monte Carlo (I use it).

Bad ones, such as the native `rand()` in C/C++ and the procedure in *Numerical Recipes* give incorrect results in common simple cases. If there is a random number generator of unknown origin being passed from person to person in the office, do not use it (without a condom).

The computer itself is not random. A pseudo random number generator simulates randomness without actually being random. The *seed* is a collection of m integer variables: `int seed[m];`. Assuming standard C/C++ 32 bit integers, the number of bits in the seed is $32 \cdot m$. There is a *seed update* function $\Phi(s)$ and an *output* function $\Psi(s)$. The update function produces a new seed: $s' = \Phi(s)$. The output function produces a floating point number (check the precision) $u = \Psi(s) \in [0, 1]$. One call `u = rng();` has the effect

```
s ← Φ(s); return u = Ψ(s); .
```

The random number generator should come with procedures `s = getSeed();` and `setSeed(s);`, with obvious functions. Most random number generators set the initial seed to the value of the system clock as a default if the program has no `setSeed(s);` command. We use `setSeed(s)` and `getSeed()` for two things. If the program starts with `setSeed(s);`, then the sequence of seeds and “random” numbers will be the same each run. This is helpful in debugging and reproducing results across platforms. Most random number generators have a built-in way to get the initial seed if the program does not set one. One common way is to use the least significant bits from the system clock. In this way, if you run the program twice, the results will be different. The results of Figure ?? were obtained in that way. It sometimes is recommended to “warm up” the random number generator by throwing away the first few numbers:

```
setSeed(17); // Start with the seed equal to 17.
int RandWarmUp = 100;
for ( int i = 0; i < RandWarmUp; i++)
    rng(); // Throw away the first RandWarmUp random numbers.
```

This is harmless but with modern random number generators unnecessary. The other use is *checkpointing*. Some Monte Carlo runs take so long that there is a real chance the computer will crash during the run. We avoid losing too much work by storing the state of the computation to disk every so often. If the machine crashes, we restart from the most recent checkpoint. The random number generator seed is part of the checkpoint data.

The simplest random number generators use linear congruences. The seed represents an integer in the range $0 \leq s < c$ and Φ is the linear congruence (a and b positive integers) $s' = \Phi(s) = (as + b)_{\text{mod } c}$. If $c > 2^{32}$, then we need more than one 32 bit integer variable to store s . Both `rand()` and `random()` are of this type, but `rand()` has $m = 1$ and `random()` has $m = 4$. The output is $u = \Psi(s) = s/c$. The more sophisticated random number generators are of a similar computational complexity.

9.3 Sampling

Sampling means producing random variables with a specified distribution. More precisely, it means using iid standard uniform random variables produced by a random number generator to produce samples from other distributions. A *simple sampler* is a procedure that produces an independent sample of X each time it is called. Simple samplers are practical mainly for univariate or low dimensional random variables, or for multivariate normals. Complex high dimensional random variables often require more advanced techniques such as *Markov Chain Monte Carlo*. A large Monte Carlo computation may spend most of its time in the sampler, and it often is possible to improve the performance by paying attention to the details of the algorithm and coding. Monte Carlo practitioners are amply rewarded for time spent fine tuning their samplers.

9.3.1 Bernoulli coin tossing

A *Bernoulli* random variable with parameter p , or a *coin toss*, is a random variable, X , with $\Pr(X = 1) = p$ and $\Pr(X = 0) = 1 - p$. If U is a standard uniform, then $p = \Pr(U \leq p)$. Therefore we can sample X using the code fragment

```
X=0; if ( rng() <= p) X=1;
```

Similarly, we can sample a random variable with finitely many values $\Pr(X = x_k) = p_k$ (with $\sum_k p_k = 1$) by dividing the unit interval into disjoint sub intervals of length p_k . This is all you need, for example, to simulate a simple random walk or a finite state space Markov chain.

9.3.2 Exponential

If U is standard uniform, then

$$T = \frac{-1}{\lambda} \ln(U) \quad (9.14)$$

is an exponential with rate parameter λ . Before checking this, note first that $U > 0$ so $\ln(U)$ is defined, and $U < 1$ so $\ln(U)$ is negative and $T > 0$. Next, since λ is a rate, it has units of 1/Time, so (9.14) produces a positive number with the correct units. The code `T = -(1/lambda)*log(rng());` generates the exponential.

We verify (9.14) using the informal probability method. Let $f(t)$ be the PDF of the random variable of (9.14). We want to show $f(t) = \lambda e^{-\lambda t}$ for $t > 0$. Let B be the event $t \leq T \leq t + dt$. This is the same as

$$t \leq \frac{-1}{\lambda} \ln(U) \leq t + dt ,$$

which is the same as

$$-\lambda t - \lambda dt \leq \ln(U) \leq -\lambda t \quad (\text{all negative}),$$

and, because e^x is an increasing function of x ,

$$e^{-\lambda t - \lambda dt} \leq U \leq e^{-\lambda t} .$$

Now, $e^{-\lambda t - \lambda dt} = e^{-\lambda t} e^{-\lambda dt}$, and $e^{-\lambda dt} = 1 - \lambda dt$, so this is

$$e^{-\lambda t} - \lambda dt e^{-\lambda t} \leq U \leq e^{-\lambda t} .$$

But this is an interval within $[0, 1]$ of length $\lambda dt e^{-\lambda t}$, so

$$\begin{aligned} f(t)dt &= \Pr(t \leq T \leq t + dt) \\ &= \Pr(e^{-\lambda t} - \lambda dt e^{-\lambda t} \leq U \leq e^{-\lambda t}) \\ &= \lambda dt e^{-\lambda t} , \end{aligned}$$

which shows that (9.14) gives an exponential.

9.3.3 Markov chains

To simulate a discrete time Markov chain, we need to make a sequence of random choices from m possibilities. One way to do this uses the idea behind Bernoulli sampling. Suppose $X(t) = j$ and we want to choose $X(t+1)$ according to the probabilities (9.11). Define the partial sums

$$s_{jk} = \sum_{l < k} p_{jl} ,$$

so that $s_{j1} = 0$, $s_{j2} = p_{j1} = \Pr(X(t+1) = 1)$, $s_{j3} = p_{j1} + p_{j2} = \Pr(X(t+1) = 1 \text{ or } X(t+1) = 2)$, etc. The algorithm is to set $X(t+1) = k$ if $s_{jk} \leq U < s_{j,k+1}$. This insures that $\Pr(X(t+1) = k) = p_{jk}$, since $\Pr(s_{jk} \leq U < s_{j,k+1}) = s_{j,k+1} - s_{jk} = p_{jk}$. In specific applications, the description of the Markov chain often suggests different ways to sample $X(t+1)$.

We can simulate a continuous time Markov chain by simulating the embedded discrete time chain and the transition times. Let S_n be the times at which

the state changes, the transition times. Suppose the transition is from state $j = X(S_n - 0)$, which is known, to state $k = X(S_n + 0)$, which must be chosen. The probability of $j \rightarrow k$ is given by (9.13). The probability of $j \rightarrow j$ is zero, since S_n is a transition time. Once we know k , the *inter-transition* time is an exponential with rate r_{kk} , so we can take (see (9.14), X is the state and $\text{TRate}[j]$ is λ_j)

```
S = S - log( rng() ) / TRate[X]; // Update the time variable.
```

9.3.4 Using the distribution function

In principle, the CDF provides a simple sampler for any one dimensional probability distribution. If X is a one component random variable with probability density $f(x)$, the cumulative distribution function is $F(x) = \Pr(X \leq x) = \int_{x' \leq x} f(x') dx'$. Of course $0 \leq F(x) \leq 1$ for all x , and for any $u \in [0, 1]$, there is an x with $F(x) = u$. The reader can check that if $F(X) = U$ then $X \sim f$ if and only if U is a standard uniform. Therefore, if `double fSolve(double u)` is a procedure that returns x with $f(x) = u$, then `X = fSolve(rng());` produces an independent X sample each time. The procedure `fSolve` might involve a simple formula in simple cases. Harder problems might require Newton's method with a careful choice of initial guess. This procedure could be the main bottleneck of a big Monte Carlo code, so it would be worth the programmer's time to spend some time making it fast.

For the exponential random variable,

$$F(t) = \Pr(T \leq t) = \int_{t'=0}^t \lambda e^{-\lambda t'} dt' = 1 - e^{-\lambda t}.$$

Solving $F(t) = u$ gives $t = -\frac{1}{\lambda} \ln(1 - u)$. This is the same as the sampler we had before, since $1 - u$ also is a standard uniform.

For the standard normal we have

$$F(z) = \int_{z'=-\infty}^z \frac{1}{\sqrt{2\pi}} e^{-z'^2/2} dz' = N(z). \quad (9.15)$$

There is no *elementary*¹⁰ formula for the *cumulative normal*, $N(z)$, but there is good software to evaluate it to nearly double precision accuracy, both for $N(z)$ and for the inverse cumulative normal $z = N^{-1}(u)$. In many applications,¹¹ this is the best way to make standard normals. The general $X \sim \mathcal{N}(\mu, \sigma^2)$ may be found using $X = \sigma Z + \mu$.

¹⁰An elementary function is one that can be computed using sums and differences, products and quotients, exponentials and logs, trigonometric functions, and powers.

¹¹There are applications where the relationship between Z and U is important, not only the value of Z . These include sampling using normal copulas, and quasi (low discrepancy sequence) Monte Carlo.

9.3.5 The Box Muller method

The *Box Muller* algorithm generates two independent standard normals from two independent standard uniforms. The formulas are

$$\begin{aligned} R &= \sqrt{-2 \ln(U_1)} \\ \Theta &= 2\pi U_2 \\ Z_1 &= R \cos(\Theta) \\ Z_2 &= R \sin(\Theta) . \end{aligned}$$

We can make a thousand independent standard normals by making a thousand standard uniforms then using them in pairs to generate five hundred pairs of independent standard normals.

The idea behind the Box Muller method is related to the brilliant elementary derivation of the formula $\int_{-\infty}^{\infty} e^{-z^2/2} = \sqrt{2\pi}$. Let $Z = (Z_1, Z_2)$ be a bivariate normal whose components are independent univariate standard normals. Since Z_1 and Z_2 are independent, the joint PDF of Z is

$$f(z) = \frac{1}{\sqrt{2\pi}} e^{-z_1^2/2} \frac{1}{\sqrt{2\pi}} e^{-z_2^2/2} = \frac{1}{2\pi} e^{-(z_1^2 + z_2^2)/2} . \quad (9.16)$$

Let R and Θ be polar coordinates for Z , which means that $Z_1 = R \cos(\Theta)$ and $Z_2 = R \sin(\Theta)$. From (9.16) it is clear that the probability density is radially symmetric, so Θ is uniformly distributed in the interval $[0, 2\pi]$, and Θ is independent of R . The Distribution function of R is

$$\begin{aligned} F(r) &= \Pr(R \leq r) \\ &= \int_{\rho=0}^r \int_{\theta=0}^{2\pi} \frac{1}{2\pi} e^{-\rho^2/2} d\theta \rho d\rho \\ &= \int_{\rho=0}^r e^{-\rho^2/2} \rho d\rho . \end{aligned}$$

With the change of variables (the trick behind the integral) $t = \rho^2/2$, $dt = \rho d\rho$, this becomes

$$\int_{t=0}^{r^2/2} e^{-t} dt = 1 - e^{-r^2/2} .$$

To sample R , we solve $1 - U = F(R)$ (recall that $1 - U$ is a standard uniform is U is a standard uniform), which gives $R = \sqrt{-2 \ln(U)}$, as claimed.

9.3.6 Multivariate normals

Let $X \in R^n$ be a multivariate normal random variable with mean zero and covariance matrix C . We can sample X using the Choleski factorization $C = LL^T$, where L is lower triangular. Note that L exists because C is symmetric and positive definite. Let $Z \in R^n$ be a vector of n independent standard normals

generated using Box Muller or any other way. The covariance matrix of Z is I (check this). Therefore, if

$$X = LZ , \quad (9.17)$$

then X is multivariate normal (because it is a linear transformation of a multivariate normal) and has covariance matrix (see (9.10))

$$C_X = LIL^t = C .$$

If we want a multivariate normal with mean $\mu \in R^n$, we simply take $X = LZ + \mu$. In some applications we prefer to work with $H = C^{-1}$ than with C . Exercise 6 has an example.

9.3.7 Rejection

The *rejection* algorithm turns samples from one density into samples of another. Not only is it useful for sampling, but the idea is the basis of the *Metropolis* algorithm in Markov Chain Monte Carlo (See Exercise 7). One ingredient is a simple sampler for the *trial distribution*, or *proposal distribution*. Suppose `gSamp()` produces iid samples from the PDF $g(x)$. The other ingredient is an *acceptance probability*, $p(x)$, with $0 \leq p(x) \leq 1$ for all x . The algorithm generates a *trial*, $X \sim g$, and accepts this trial value with probability $p(X)$. The process is repeated until the first acceptance. All this happens in

$$\text{while (rng() > p(X = gSamp()))}; \quad (9.18)$$

We accept X if $U \leq p(X)$, so $U > p(X)$ means reject and try again. Each time we generate a new X , which must be independent of all previous ones.

The X returned by (9.18) has PDF

$$f(x) = \frac{1}{Z} p(x) g(x) , \quad (9.19)$$

where Z is a normalization constant that insures that $\int f(x) dx = 1$:

$$Z = \int_{x \in R^n} p(x) g(x) dx . \quad (9.20)$$

This shows that Z is the probability that any given trial will be an acceptance. The formula (9.19) shows that rejection goes from g to f by thinning out samples in regions where $f < g$ by rejecting some of them. We verify it informally in the one dimensional setting:

$$\begin{aligned} f(x)dx &= \Pr (\text{accepted } X \in (x, x + dx)) \\ &= \Pr (\text{generated } X \in (x, x + dx) \mid \text{acceptance}) \\ &= \frac{\Pr (\text{generated } X \in (x, x + dx) \text{ and accepted})}{\Pr (\text{accepted})} \\ &= \frac{g(x)dx p(x)}{Z} , \end{aligned}$$

where Z is the probability that a given trial generates an acceptance. An argument like this also shows the correctness of (9.19) also for multivariate random variables.

We can use rejection to generate normals from exponentials. Suppose $g(x) = e^{-x}$ for $x > 0$, corresponding to a standard exponential, and $f(x) = \frac{2}{\sqrt{2\pi}}e^{-x^2/2}$ for $x > 0$, corresponding to the positive half of the standard normal distribution. Then (9.19) becomes

$$\begin{aligned} p(x) &= Z \frac{f(x)}{g(x)} \\ &= Z \cdot \frac{2}{\sqrt{2\pi}} \cdot \frac{e^{-x^2/2}}{e^{-x}} \\ p(x) &= Z \cdot \frac{2}{\sqrt{2\pi}} e^{x-x^2/2}. \end{aligned} \quad (9.21)$$

This would be a formula for $p(x)$ if we know the constant, Z .

We maximize the efficiency of the algorithm by making Z , the overall probability of acceptance, as large as possible, subject to the constraint $p(x) \leq 1$ for all x . Therefore, we find the x that maximizes the right side:

$$e^{x-x^2/2} = \max \implies x - \frac{x^2}{2} = \max \implies x_{\max} = 1.$$

Choosing Z so that the maximum of $p(x)$ is one gives

$$1 = p_{\max} = Z \cdot \frac{2}{\sqrt{2\pi}} e^{x_{\max}-x_{\max}^2/2} = Z \frac{2}{\sqrt{2\pi}} e^{1/2},$$

so

$$p(x) = \frac{1}{\sqrt{e}} e^{x-x^2/2}. \quad (9.22)$$

It is impossible to go the other way. If we try to generate a standard exponential from a positive standard normal we get acceptance probability related to the reciprocal to (9.21):

$$p(x) = Z \frac{\sqrt{2\pi}}{2} e^{x^2/2-x}.$$

This gives $p(x) \rightarrow \infty$ as $x \rightarrow \infty$ for any $Z > 0$. The normal has thinner¹² tails than the exponential. It is possible to start with an exponential and thin the tails using rejection to get a Gaussian (Note: (9.21) has $p(x) \rightarrow 0$ as $x \rightarrow \infty$). However, rejection cannot fatten the tails by more than a factor of $\frac{1}{2}$. In particular, rejection cannot fatten a Gaussian tail to an exponential.

The *efficiency* of a rejection sampler is the expected number of trials needed to generate a sample. Let N be the number of samples to get a success. The efficiency is

$$E[N] = 1 \cdot \Pr(N = 1) + 2 \cdot \Pr(N = 2) + \dots$$

¹²The tails of a probability density are the parts for large x , where the graph of $f(x)$ gets thinner, like the tail of a mouse.

We already saw that $\Pr(N = 1) = Z$. To have $N = 2$, we need first a rejection then an acceptance, so $\Pr(N = 2) = (1 - Z)Z$. Similarly, $\Pr(N = k) = (1 - Z)^{k-1}Z$. Finally, we have the geometric series formulas for $0 < r < 1$:

$$\sum_{k=0}^{\infty} r^k = \frac{1}{1-r} \quad , \quad \sum_{k=1}^{\infty} k r^{k-1} = \sum_{k=0}^{\infty} k r^{k-1} = \frac{d}{dr} \sum_{k=0}^{\infty} r^k = \frac{1}{(1-r)^2} \quad .$$

Applying these to $r = 1 - Z$ gives $E[N] = \frac{1}{Z}$. In generating a standard normal from a standard exponential, we get

$$Z = \sqrt{\frac{\pi}{2e}} \approx .76 \quad .$$

The sampler is efficient in that more than 75% of the trials are successes.

Rejection samplers for other distributions, particularly in high dimensions, can be much worse. We give a rejection algorithm for finding a random point uniformly distributed inside the unit ball n dimensions. The algorithm is correct for any n in the sense that it produces at the end of the day a random point with the desired probability density. For small n , it even works in practice and is not such a bad idea. However, for large n the algorithm is very inefficient. In fact, Z is an exponentially decreasing function of n . It would take more than a century on any present computer to generate a point uniformly distributed inside the unit ball in $n = 100$ dimensions this way. Fortunately, there are better ways.

A point in n dimensions is $x = (x_1, \dots, x_n)$. The *unit ball* is the set of points with $\sum_{k=1}^n x_k^2 \leq 1$. We will use a trial density that is uniform inside the smallest (hyper)cube that contains the unit ball. This is the cube with $-1 \leq x_k \leq 1$ for each k . The uniform density in this cube is

$$g(x_1, \dots, x_n) = \begin{cases} 2^{-n} & \text{if } |x_k| \leq 1 \text{ for all } k = 1, \dots, n \\ 0 & \text{otherwise.} \end{cases}$$

This density is a product of the one dimensional uniform densities, so we can sample it by choosing n independent standard uniforms:

```
for( k = 0; k < n; k++) x[k] = 2*rng() - 1; // unif in [-1,1].
```

We get a random point inside the unit ball if we simply reject samples outside the ball:

```
while(1) {      // The rejection loop (possibly infinite!)

    for( k = 0; k < n; k++) x[k] = 2*rng() - 1; // Generate a trial vector
                                           // of independent uniforms
                                           // in [-1,1].

    ssq = 0;    // ssq means "sum of squares"
    for( k = 0; k < n; k++) ssq += x[k]*x[k];
    if ( ssq <= 1.) break ; // You accepted and are done with the loop.
                          // Otherwise go back and do it again.
}
```

Table 9.1: Acceptance fractions for producing a random point in the unit ball in n dimensions by rejection.

dimension	$vol(\text{ball})$	$vol(\text{cube})$	ratio	$-\ln(\text{ratio})/\text{dim}$
2	π	4	.79	.12
3	$4\pi/3$	8	.52	.22
4	$\pi^2/2$	16	.31	.29
10	$2\pi^{n/2}/(n\Gamma(n/2))$	2^n	.0025	.60
20	$2\pi^{n/2}/(n\Gamma(n/2))$	2^n	2.5×10^{-8}	.88
40	$2\pi^{n/2}/(n\Gamma(n/2))$	2^n	3.3×10^{-21}	1.2

The probability of accepting in a given trial is equal to the ratio of the volume (or area in 2D) of the ball to the cube that contains it. In 2D this is

$$\frac{\text{area}(\text{disk})}{\text{area}(\text{square})} = \frac{\pi}{4} \approx .79 ,$$

which is pretty healthy. In 3D it is

$$\frac{\text{vol}(\text{ball})}{\text{vol}(\text{cube})} = \frac{\frac{4\pi}{3}}{8} \approx .52 ,$$

Table 9.1 shows what happens as the dimension increases. By the time the dimension reaches $n = 10$, the expected number of trials to get a success is about $1/.0025 = 400$, which is slow but not entirely impractical. For dimension $n = 40$, the expected number of trials has grown to about 3×10^{20} , which is entirely impractical. Monte Carlo simulations in more than 40 dimensions are common. The last column of the table shows that the acceptance probability goes to zero faster than any exponential of the form e^{-cn} , because the numbers that would be c , listed in the table, increase with n .

9.3.8 Histograms and testing

Any piece of scientific software is presumed wrong until it proves itself correct in tests. We can test a one dimensional sampler using a *histogram*. Divide the x axis into neighboring *bins* of length Δx centered about bin centers $x_j = j\Delta x$. The corresponding bins are $B_j = [x_j - \frac{\Delta x}{2}, x_j + \frac{\Delta x}{2}]$. With n samples, the *bin counts* are¹³ $N_j = \#\{X_k \in B_j, 1 \leq k \leq n\}$. The probability that a given sample lands in B_j is $\Pr(B_j) = \int_{x \in B_j} f(x)dx \approx \Delta x f(x_j)$. The expected bin count is $E[N_j] \approx n\Delta x f(x_j)$, and the standard deviation (See Section 9.4) is $\sigma_{N_j} \approx \sqrt{n\Delta x f(x_j)}$. We generate the samples and plot the N_j and $E[N_j]$ on the same plot. If $E[N_j] \gg \sigma_{N_j}$, then the two curves should be relatively close. This condition is

$$\frac{1}{\sqrt{n\Delta x}} \ll \sqrt{f(x_j)} .$$

¹³Here $\#\{\dots\}$ means the number of elements in the set $\{\dots\}$.

In particular, if f is of order one, $\Delta x = .01$, and $n = 10^6$, we should have reasonable agreement if the sampler is correct. If Δx is too large, the approximation $\int_{x \in B_j} f(x) dx \approx \Delta x f(x_j)$ will not hold. If Δx is too small, the histogram will not be accurate.

It is harder to test higher dimensional random variables. We can test two and possibly three dimensional random variables using multidimensional histograms. We can test that various one dimensional functions of the random X have the right distributions. For example, the distributions of $R^2 = \sum X_k^2 = \|X\|_{l_2}^2$ and $Y = \sum a_k X_k = a \cdot X$ are easy to figure out if X is uniformly distributed in the ball.

9.4 Error bars

For large N , the error in (??) is governed by the central limit theorem. This is because the numbers $V_k = V(X_k)$ are independent random variables with mean A and variance σ^2 with σ given by (9.3). The theorem states that $\hat{A}_N - A$ is approximately normal with mean zero and variance¹⁴ σ^2/N . This may be expressed by writing ($\sim\sim$ indicates that the two sides have approximately the same distribution)

$$\hat{A}_N - A \sim\sim \frac{\sigma}{\sqrt{N}} Z, \quad (9.23)$$

where $Z \sim \mathcal{N}(0, 1)$ is a standard normal.

The normal distribution is tabulated in any statistics book or scientific software package. In particular, $\Pr(|Z| \leq 1) \approx .67$ and $\Pr(|Z| \leq 2) \approx .95$. This means that

$$\Pr\left(\left|\hat{A}_N - A\right| \leq \frac{\sigma}{\sqrt{N}}\right) \approx .67. \quad (9.24)$$

In the computation, A is unknown but \hat{A}_N is known, along with N and (approximately, see below) σ . The equation (9.24) then says that the *confidence interval*,

$$\left[\hat{A}_N - \frac{\sigma}{\sqrt{N}}, \hat{A}_N + \frac{\sigma}{\sqrt{N}}\right], \quad (9.25)$$

has a 67% chance of containing the actual answer, A . The interval (9.25) is called the one standard deviation *error bar* because it is indicated as a bar in plots. The center of the bar is \hat{A}_N . The interval itself is a line segment (a bar) that extends $\frac{\sigma}{\sqrt{N}}$ in either direction about its center.

There is a standard way to estimate σ from the Monte Carlo data, starting with

$$\sigma^2 = \text{var}[V(X)] = E\left[(V(X) - A)^2\right] \approx \frac{1}{N} \sum_{k=1}^N (V(X_k) - A)^2.$$

¹⁴The mean and variance are exact. Only the normality is approximate.

Although A is unknown, using \hat{A}_N is good enough. This gives the variance estimator:¹⁵

$$\sigma^2 \approx \hat{\sigma}_N^2 = \frac{1}{N} \sum_{k=1}^N (V(X_k) - A)^2. \quad (9.26)$$

The full Monte Carlo estimation procedure is as follows. Generate N independent samples $X_k \sim f$ and evaluate the numbers $V(X_k)$. Calculate the *sample mean* from (9.2), and the sample variance $\hat{\sigma}_N^2$ from (9.26). The one standard deviation error bar half width is

$$\frac{1}{\sqrt{N}} \sigma \approx \frac{1}{\sqrt{N}} \hat{\sigma}_N = \frac{1}{\sqrt{N}} \sqrt{\hat{\sigma}_N^2}. \quad (9.27)$$

The person doing the Monte Carlo computation should be aware of the error bar size. Error bars should be included in any presentation of Monte Carlo results that are intended for technically trained people. This includes reports in technical journals and homework assignments for Scientific Computing.

The error bar estimate (9.27) is not exact. It has statistical error of order $1/\sqrt{N}$ and bias on the order of $1/N$. But there is a sensible saying: “Do not put error bars on error bars.” Even if we would know the exact value of σ , we would have only a rough idea of the actual error $\hat{A}_N - A$. Note that the estimate (9.26) has a bias on the order of $1/N$. Even if we used the unbiased estimate of σ^2 , which has $N - 1$ instead of N , the estimate of σ still would be biased because the square root function is nonlinear. The bias in (9.27) would be on the order of $1/N$ in either case.

It is common to report a Monte Carlo result in the form $A = \hat{A}_N \pm \text{error bar}$ or $A = \hat{A}_N(\text{error bar})$. For example, writing $A = .342 \pm .005$ or $A = .342(.005)$ indicates that $\hat{A}_N = .342$ and $\hat{\sigma}/\sqrt{N} = .005$. Note that we do not report significant digits of \hat{A}_N beyond the error bar size because they have no information. Also, we report only one significant digit of the error bar itself. If our estimated error bar actually were $\hat{\sigma}/\sqrt{N} = .0048523$, reporting the extra digits would only be annoying to the reader.

It is the custom in Monte Carlo practice to plot and report one standard deviation error bars. This requires the consumer to understand that the exact answer is outside the error bar about a third of the time. Plotting two or three standard deviation error bars would be safer but would give an inaccurate picture of the probable error size.

9.5 Variance reduction

Variance reduction means reformulating a problem to lower the variance without changing the expected value.¹⁶ We have seen that the statistical error in a

¹⁵The estimator with $N - 1$ in place of N on the right of (9.26) is often called s^2 in statistics. If the difference between $N - 1$ and N is significant, N is too small for Monte Carlo estimation.

¹⁶There are some variance reduction methods that are not unbiased, for example, in approximate solution of stochastic differential equations.

Monte Carlo computation is of the order of σ/\sqrt{N} . We can reduce the error by increasing N or decreasing σ . Reducing σ allows us to reduce N without losing accuracy. For example, if we could reduce σ by a factor of 2, we would be able to reduce N by a factor of 4 and achieve the same accuracy.

There are several general variance reduction tricks, but whether a given one will help depends very much on the details of the specific application. We give just a few examples to illustrate the possibilities. The applications are artificial because explaining real applications would take more pages and time than the average reader (or writer) wants to spend.

9.5.1 Control variates

A *control variate* is a function $W(X)$ whose expected value $B = E[W(X)]$ is known. We can estimate A using

$$A = E[V(X) - \alpha(W(X) - B)] . \quad (9.28)$$

If $W(X)$ is correlated with $V(X)$, then we can reduce the variance of \hat{A}_N using (9.28) instead of (9.1). We generate the samples X_k as usual. For each sample, we evaluate both $V(X_k)$ and $W(X_k)$ and estimate A using

$$\hat{A}_N = \frac{1}{N} \sum_{k=1}^N (V(X_k) - \alpha(W(X_k) - B)) .$$

The optimal value of α , the value that minimized the variance, is (see Exercise 9) $\alpha^* = \rho_{VW} \sqrt{\text{var}[V]/\text{var}[W]}$. In practical applications, we would estimate α^* from the same Monte Carlo data. It is not likely that we would know $\text{cov}[V, W]$ without knowing $A = E[V]$.

Obviously, this method depends on having a good control variate.

9.5.2 Antithetic variates

We say that R is a *symmetry* of a probability distribution f if the random variable $Y = R(X)$ is distributed as f whenever X is. For example, if $f(x)$ is a probability density so that $f(-x) = f(x)$ (e.g. a standard normal), then $R(x) = -x$ is a symmetry of f . If U is a standard uniform then $1 - U = R(U)$ also is standard uniform. If X is a sample of f ($X \sim f$), then $Y = R(X)$ also is a sample of f . This is called the *antithetic* sample. The method of *antithetic variates* is to use antithetic pairs of samples. If $A = E_f[V(X)]$, then also $A = E_f[V(R(X))]$ and

$$A = E \left[\frac{1}{2} (V(X) + V(R(X))) \right] .$$

9.5.3 Importance sampling

Suppose X is a multivariate random variable with probability density $f(x)$ and that we want to estimate

$$A = E_f [V(X)] = \int_{R^n} V(x)f(x)dx . \quad (9.29)$$

Let $g(x)$ be any other probability density subject only to the condition that $g(x) > 0$ whenever $f(x) > 0$. The *likelihood ratio* is $L(x) = f(x)/g(x)$. *Importance sampling* is the strategy of rewriting (9.29) as

$$A = \int_{R^n} V(x)L(x)g(x)dx = E_g [V(X)L(X)] . \quad (9.30)$$

The variance is reduced if

$$\text{var}_g [V(X)L(X)] \leq \text{var}_f [V(X)] .$$

The term “importance sampling” comes from the idea that the values of x that are most important in the expected value of $V(X)$ are unlikely in the f distribution. The g distribution puts more weight in these important areas.

9.6 Software: performance issues

Monte Carlo methods raises many performance issues. Naive coding following the text can lead to poor performance. Two significant factors are frequent branching and frequent procedure calls.

9.7 Resources and further reading

There are many good books on the probability background for Monte Carlo, the book by Sheldon Ross at the basic level, and the book by Sam Karlin and Gregory Taylor for more the ambitious. Good books on Monte Carlo include the still surprisingly useful book by Hammersley and Handscomb, the physicists’ book (which gets the math right) by Mal Kalos and Paula Whitlock, and the broader book by George Fishman. I also am preparing a book on Monte Carlo, with many parts already posted.

The sampling methods of Section 9.3 all *simple samplers*. They produce independent samples of f . There are many applications for which there is no known practical simple sampler. Most of these can be treated by *dynamic* samplers that produce X_k that are not independent of each other. The best known such *Markov Chain Monte Carlo* method is the *Metropolis* algorithm. The disadvantage of MCMC samplers is that the standard deviation of the estimator (9.2) may be much larger than σ/\sqrt{N} with σ given by (9.3). This is discussed in the online lecture Monte Carlo notes of Alan Sokal and the book by Jun Liu.

The ability to sample essentially arbitrary distributions has led to an explosion of new applications and improvements in technique. Bayesian statistics and Bayesian methods in machine learning rely on sampling posterior distributions. Statisticians use Monte Carlo to calibrate hypothesis testing methods. In rare event simulation using importance sampling, we may consider essentially arbitrary importance functions, even when the original has a simple sampler.

Statistical error makes it hard to use the optimization algorithms of Chapter 6 to compute

$$\max_t \phi(t) ,$$

where

$$\phi(t) = E_t[V(X)] . \quad (9.31)$$

Here $E_t[\cdot]$ means that the distribution of X depends on the parameter t in some way. For example, suppose X is normal with mean zero and variance t and we want samples to have X^2 as large as possible without exceeding 1. This could be formulated as finding the maximum over t of

$$\phi(t) = E_t[X^2 \cdot \mathbf{1}(X^2 \leq 1)] = \frac{1}{\sqrt{2\pi t}} \int_{-1}^1 x^2 e^{-x^2/2t} dx .$$

If we estimate $\phi(t)$ by

$$\hat{\phi}(t) = \frac{1}{N} \sum_{k=1}^N V(X_k) ,$$

the statistical error will be different for each value of t (see exercise 8).

In computational chemistry, using the embedded discrete time Markov chain as in Section 9.3.3 is called SSA (for “stochastic simulation algorithm”) and is attributed to Gillespie. *Event driven simulation* is a more elaborate but often more efficient way to simulate complex Markov chains either in discrete or continuous time.

Choosing a good random number generator is important yet subtle. The native C/C++ function `rand()` is suitable only for the simplest applications because it cycles after only a billion samples. The function `random()` is much better. The random number generators in Matlab are good, which cannot be said for the generators in other scientific computing and visualization packages. Joseph Marsaglia has a web site with the latest and best random number generators.

9.8 Exercises

1. Suppose we wish to estimate $A = \phi(B)$, where $B = E_f[V(X)]$. For example, suppose $E[X] = 0$ and we want the standard deviation $\sigma = (E[X^2])^{1/2}$. Here $B = E[X^2]$, $V(x) = x^2$, and $\phi(B) = \sqrt{B}$. Let \hat{B}_N be an unbiased estimator of B whose variance is σ_B^2/N . If we estimate A using $\hat{A}_N = \phi(\hat{B}_N)$, and N is large, find approximate expressions for the

bias and variance of \hat{A}_N . Show that \hat{A}_N is biased in general, but that the bias is much smaller than the statistical error for large N . Hint: assume N is large enough so that $\widehat{B}_N \approx B$ and

$$\phi(\widehat{B}_N) - \phi(B) \approx \phi'(B)(\widehat{B}_N - B) + \frac{1}{2}\phi''(B)(\widehat{B}_N - B)^2.$$

Use this to estimate the bias, $E[\hat{A}_N - A]$, and the variance $\text{var}[\hat{A}_N]$. What are the powers of N in the bias and statistical error?

2. Let (X, Y) be a bivariate random variable that is uniformly distributed in the unit disk. This means that $h(x, y)$, the joint density of x and y , is equal to $1/\pi$ if $x^2 + y^2 \leq 1$, and $h(x, y) = 0$ otherwise. Show that X and Y are uncorrelated but not independent (hint: calculate the covariance of X^2 and Y^2).
3. What is wrong with the following piece of code?

```
for ( k = 0; k < n; k++ ) {
    setSeed(s);
    U[k] = rng();
}
```

4. Calculate the distribution function for an exponential random variable with rate constant λ . Show that the sampler using the distribution function given in Section 9.3.4 is equivalent to the one given in Section 9.3.2. Note that if U is a standard uniform, then $1 - U$ also is standard uniform.
5. If S_1 and S_2 are independent standard exponentials, then $T = S_1 + S_2$ has PDF $f(t) = te^{-t}$.
 - (a) Write a simple sampler of T that generates S_1 and S_2 then takes $T = S_1 + S_2$.
 - (b) Write a simpler sampler of T that uses rejection from an exponential trial. The trial density must have $\lambda < 1$. Why? Look for a value of λ that gives reasonable efficiency. Can you find the optimal λ ?
 - (c) For each sampler, use the histogram method to verify its correctness.
 - (d) Program the Box Muller algorithm and verify the results using the histogram method.
6. A common problem that arises in statistics and stochastic control is to sample the n component normal

$$f(x) = \frac{1}{Z} \exp \left(\frac{-D}{2} \sum_{j=0}^n (x_{j+1} - x_j)^2 \right). \quad (9.32)$$

with the convention that $x_0 = x_{n+1} = 0$. The exponent in (9.32) is a function of the n variables x_1, \dots, x_n .

- (a) Determine the entries of the symmetric $n \times n$ matrix H so that the exponent in (9.32) has the form $x^* H x$. Show that H is tridiagonal.
 - (b) Show that the Choleski factorization of H takes the form $H = M M^*$ where M has nonzero entries only on the main diagonal and the first subdiagonal. How many operations and how much memory are required to compute and store M , what power of n ?
 - (c) ¹⁷ Find expressions for the entries of M .
 - (d) Show that if Z is an n component standard normal and we use back substitution to find X satisfying $M^* X = Z$, then X is a multivariate normal with covariance $C = H^{-1}$. What is the work to do this, as a function of n ?
 - (e) Let $C = H^{-1}$ have Choleski factorization $C = L L^*$. Show that $L = M^{-1}$. Show that L , and C itself are full matrices with all entries non-zero. Use this to find the work and memory needed to generate a sample of (9.32) using the method of Section 9.3.6.
 - (f) Find an expression for the entries of C . Hint: Let e_j be the standard column vector whose only non-zero entry is a one in component j . Let v satisfy $H v = e_j$. Show that v is linear except at component j . That is, $v_k = a k + b$ for $k \leq j$, and $v_k = c k + d$ for $k \geq j$.
7. Suppose X is a discrete random variable whose possible values are $x = 1, \dots, n$, and that $f(x) = \Pr(X = x)$ is a desired probability distribution. Suppose p_{jk} are the transition probabilities for an n state Markov chain. Then f is an *invariant* or *steady state* distribution for p if (supposing $X(t)$ is the state after t transitions)

$$\begin{aligned} \Pr(X(t) = x) &= f(x) \text{ for all } x \\ \implies \Pr(X(t+1) = x) &= f(x) \text{ for all } x. \end{aligned}$$

This is the same as $fP = f$, where f is the row vector with components $f(x)$ and P is the transition matrix. The *ergodic theorem for Markov chains* states that if f is an invariant distribution and if the chain is non-degenerate, then the *Markov chain Monte Carlo* (MCMC) estimator,

$$\hat{A}_T = \sum_{t=1}^T V(X(t)),$$

converges to $A = E_f[V(X)]$, as $t \rightarrow \infty$. This holds for any transition probabilities P that leave f invariant, though some lead to more accurate estimators than others. The *Metropolis* method is a general way to construct a suitable P .

¹⁷Parts c and f are interesting facts but may be time consuming and are not directly related to Monte Carlo

- (a) The transition matrix P satisfies *detailed balance* for f if the steady state probability of an $x \rightarrow y$ transition is the same as the probability of the reverse, $y \rightarrow x$. That is, $f(x)p_{xy} = f(y)p_{yx}$ for all x, y . Show that $fP = f$ if P satisfies detailed balance with respect to x .
- (b) Suppose Q is any $n \times n$ transition matrix, which we call the *trial probabilities*, and R another $n \times n$ matrix whose entries satisfy $0 \leq r_{xy} \leq 1$ and are called *acceptance probabilities*. The *Metropolis Hastings* (more commonly, simply Metropolis) method generates $X(t+1)$ from $X(t)$ as follows. First generate $Y \sim Q_{X(t), \cdot}$, which means $\Pr(Y = y) = Q_{X(t), y}$. Next, *accept* Y (and set $X(t+1) = Y$) with probability $R_{X(t), y}$. If Y is rejected, then $X(t+1) = X(t)$. Show that the transition probabilities satisfy $p_{xy} = q_{xy}r_{xy}$ as long as $x \neq y$.
- (c) Show that P satisfies detailed balance with respect to f if

$$r_{xy} = \min \left(\frac{f(x)q_{xy}}{f(y)q_{yx}}, 1 \right). \quad (9.33)$$

8. Suppose $\phi(t)$ is given by (9.31) and we want to estimate $\phi'(t)$. We consider three ways to do this.

- (a) The most straightforward way would be to take N samples $X_k \sim f_t$ and another N independent samples $\tilde{X}_k \sim f_{t+\Delta t}$ then take

$$\hat{\phi}' = \frac{1}{\Delta t} \left(\frac{1}{N} V(\tilde{X}_k) - \frac{1}{N} V(X_k) \right).$$

9. Write an expression for the variance of the random variable $Y = V(X) - \alpha(W(X) - B)$ in terms of $\text{var}[V(X)]$, $\text{var}[W(X)]$, the covariance $\text{cov}[V(X), W(X)]$, and α . Find the value of α that minimized the variance of Y . Show that the variance of Y , with the optimal α is less than the variance of $V(X)$ by the factor $1 - \rho_{VW}^2$, where ρ_{VW} is the correlation between V and W .
10. A *Poisson* random walk has a position, $X(t)$ that starts with $X(0) = 0$. At each time T_k of a Poisson process with rate λ , the position moves (jumps) by a $\mathcal{N}(0, \sigma^2)$, which means that $X(T_k + 0) = X(T_k - 0) + \sigma Z_k$ with iid standard normal Z_k . Write a program to simulate the Poisson random walk and determine $A = \Pr(X(T) > B)$. Use (but do not hard wire) two parameter sets:
- (a) $T = 1$, $\lambda = 4$, $\sigma = .5$, and $B = 1$.
- (b) $T = 1$, $\lambda = 20$, $\sigma = .2$, and $B = 2$.

Use standard error bars. In each case, choose a sample size, n , so that you calculate the answer to 5% relative accuracy and report the sample size needed for this.

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