

Exercise 1:

• Dirac's notation: $\langle 00100 \rangle_{AB} = \langle 010 \rangle_A \langle 010 \rangle_B = 1$

Array form: $\langle 00100 \rangle_{AB} = [1\ 0\ 0\ 0] \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = 1 + 0 + 0 + 0 = 0$

• Dirac's notation $\langle 01100 \rangle_{AB} = \langle 010 \rangle_A \langle 110 \rangle_B = 0$

Array form: $\langle 01100 \rangle_{AB} = [0\ 1\ 0\ 0] \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = 0$

• Dirac's notation $\langle 10100 \rangle_{AB} = \langle 110 \rangle_A \langle 010 \rangle_B = 0$

Array form: $\langle 10100 \rangle_{AB} = [0\ 0\ 1\ 0] \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = 0$

• Dirac's notation $\langle 11100 \rangle_{AB} = \langle 110 \rangle_A \langle 110 \rangle_B = 0$

Array form: $\langle 11100 \rangle_{AB} = [0\ 0\ 0\ 1] \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = 0$

• Dirac's notation: $\langle 01101 \rangle_{AB} = \langle 010 \rangle_A \langle 111 \rangle_B = 1$

Array form: $\langle 01101 \rangle_{AB} = [0\ 1\ 0\ 0] \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = 0 + 1 + 0 + 0 = 0$

• Dirac's notation: $\langle 10101 \rangle_{AB} = \langle 110 \rangle_A \langle 011 \rangle_B = 0$

Array form: $\langle 10101 \rangle_{AB} = [0\ 0\ 1\ 0] \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = 0$

• Dirac's notation: $\langle 11101 \rangle_{AB} = \langle 110 \rangle_A \langle 111 \rangle_B = 0$

Array form: $\langle 11101 \rangle_{AB} = [0\ 0\ 0\ 1] \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = 0$

• Dirac's notation: $\langle 10110 \rangle_{AB} = \langle 111 \rangle_A \langle 010 \rangle_B = 1$

Array form: $\langle 10110 \rangle_{AB} = [0\ 0\ 1\ 0] \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = 1$

• Dirac's notation: $\langle 11110 \rangle_{AB} = \langle 111 \rangle_A \langle 110 \rangle_B = 0$

Array form: $\langle 11110 \rangle_{AB} = [0\ 0\ 0\ 1] \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = 0$

• Dirac's notation: $\langle 11111 \rangle_{AB} = \langle 111 \rangle_A \langle 111 \rangle_B = 1$

Array form: $\langle 11111 \rangle_{AB} = [0\ 0\ 0\ 1] \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = 1$

Exercise 2:

$(\hat{x}_A \otimes \mathbb{1}_B) |101\rangle :$

• Dirac's notation: $(\hat{x}_A \otimes \mathbb{1}_B) |101\rangle = \hat{x}_A |1\rangle_A \otimes \mathbb{1}_B |0\rangle_B = |1\rangle_A \otimes |0\rangle_B = |11\rangle_{AB}$

• Matrix form: $(\hat{x}_A \otimes \mathbb{1}_B) |101\rangle = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = |11\rangle_{AB}$

$(\hat{x}_A \otimes \mathbb{1}_B) |10\rangle :$

• Dirac's notation: $(\hat{x}_A \otimes \mathbb{1}_B) |10\rangle = \hat{x}_A |1\rangle_A \otimes \mathbb{1}_B |0\rangle_B = |10\rangle_A \otimes |0\rangle_B = |100\rangle_{AB}$

• Matrix form: $(\hat{X}_A \otimes \mathbb{1}_B) |10\rangle = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = |100\rangle_{AB}$

$(\hat{X}_A \otimes \mathbb{1}_B) |11\rangle$:

• Dirac's notation: $(\hat{X}_A \otimes \mathbb{1}_B) |11\rangle = \hat{X}_A |1\rangle_A \otimes \mathbb{1}_B |1\rangle_B = |10\rangle_A \otimes |1\rangle_B = |101\rangle_{AB}$

• Matrix form: $(\hat{X}_A \otimes \mathbb{1}_B) |11\rangle = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = |101\rangle_{AB}$

I clearly prefer Dirac's notation which is much more easier.

Exercise 3:

$(\mathbb{1}_A \otimes \hat{X}_B) |101\rangle$:

• Dirac's notation: $(\mathbb{1}_A \otimes \hat{X}_B) |101\rangle = \mathbb{1}_A |10\rangle_A \otimes \hat{X}_B |1\rangle_B = |10\rangle_A \otimes |0\rangle_B = |100\rangle_{AB}$

• Matrix form: $(\mathbb{1}_A \otimes \hat{X}_B) |101\rangle = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = |100\rangle_{AB}$

$(\mathbb{1}_A \otimes \hat{X}_B) |110\rangle$:

• Dirac's notation: $(\mathbb{1}_A \otimes \hat{X}_B) |110\rangle = \mathbb{1}_A |11\rangle_A \otimes \hat{X}_B |0\rangle_B = |11\rangle_A \otimes |0\rangle_B = |110\rangle_{AB}$

• Matrix form: $(\mathbb{1}_A \otimes \hat{X}_B) |110\rangle = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = |110\rangle_{AB}$

$(\mathbb{1}_A \otimes \hat{X}_B) |111\rangle$:

• Dirac's notation: $(\mathbb{1}_A \otimes \hat{X}_B) |111\rangle = \mathbb{1}_A |11\rangle_A \otimes \hat{X}_B |1\rangle_B = |11\rangle_A \otimes |1\rangle_B = |110\rangle_{AB}$

• Matrix form: $(\mathbb{1}_A \otimes \hat{X}_B) |111\rangle = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = |110\rangle_{AB}$

Exercise 4:

$$(\hat{X}_A \otimes \mathbb{1}_B)(\mathbb{1}_A \otimes \hat{X}_B) = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} = \hat{X}_A \otimes \hat{X}_B.$$

Exercise 5:

Dirac's notation:

• $\hat{X}_A \otimes \hat{X}_B |100\rangle_{AB} = \hat{X}_A |10\rangle_A \otimes \hat{X}_B |0\rangle_B = |1\rangle_A \otimes |0\rangle_B = |100\rangle_{AB}$

• $\hat{x}_A \otimes \hat{x}_B |10\rangle_{AB} = \hat{x}_A |10\rangle \otimes \hat{x}_B |0\rangle = |11\rangle_A \otimes |00\rangle_B = |110\rangle_{AB}$: the quanta of B has been transferred to A.

• $\hat{x}_A \otimes \hat{x}_B |10\rangle_{AB} = \hat{x}_A |1\rangle \otimes \hat{x}_B |0\rangle = |10\rangle_A \otimes |1\rangle_B = |101\rangle_{AB}$: the quanta of A has been transferred to B.

• $\hat{x}_A \otimes \hat{x}_B |11\rangle_{AB} = \hat{x}_A |11\rangle_A \otimes \hat{x}_B |1\rangle_B = |10\rangle_A \otimes |10\rangle_B = |100\rangle_{AB}$.

Matrix form:

$$\cdot \hat{x}_A \otimes \hat{x}_B |100\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = |11\rangle_{AB}$$

$$\cdot \hat{x}_A \otimes \hat{x}_B |10\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = |10\rangle_{AB} \text{ : the quanta of B has been transferred to A.}$$

$$\cdot \hat{x}_A \otimes \hat{x}_B |101\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = |101\rangle_{AB} \text{ : the quanta of A has been transferred to B.}$$

$$\cdot \hat{x}_A \otimes \hat{x}_B |111\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = |100\rangle_{AB}$$

Exercise 6:

$$\cdot |++\rangle = |+\rangle_A \otimes |+\rangle_B = \frac{1}{\sqrt{2}}(|0\rangle_A + |1\rangle_A) \otimes \frac{1}{\sqrt{2}}(|0\rangle_B + |1\rangle_B) = \frac{1}{2}(|100\rangle_{AB} + |101\rangle_{AB} + |110\rangle_{AB} + |111\rangle_{AB}) = \begin{bmatrix} 1/2 \\ 1/2 \\ 1/2 \\ 1/2 \end{bmatrix}$$

$$\cdot |+-\rangle_{AB} = |+\rangle_A \otimes |- \rangle_B = \frac{1}{\sqrt{2}}(|0\rangle_A + |1\rangle_A) \otimes \frac{1}{\sqrt{2}}(|0\rangle_B - |1\rangle_B) = \frac{1}{2}(|100\rangle_{AB} - |101\rangle_{AB} + |110\rangle_{AB} - |111\rangle_{AB}) = \begin{bmatrix} 1/2 \\ -1/2 \\ 1/2 \\ -1/2 \end{bmatrix}$$

$$\cdot |-+\rangle_{AB} = |- \rangle_A \otimes |+\rangle_B = \frac{1}{\sqrt{2}}(|0\rangle_A - |1\rangle_A) \otimes \frac{1}{\sqrt{2}}(|0\rangle_B + |1\rangle_B) = \frac{1}{2}(|100\rangle_{AB} + |101\rangle_{AB} - |110\rangle_{AB} - |111\rangle_{AB}) = \begin{bmatrix} 1/2 \\ 1/2 \\ -1/2 \\ -1/2 \end{bmatrix}$$

$$\cdot |--\rangle_{AB} = |- \rangle_A \otimes |- \rangle_B = \frac{1}{\sqrt{2}}(|0\rangle_A - |1\rangle_A) \otimes \frac{1}{\sqrt{2}}(|0\rangle_B - |1\rangle_B) = \frac{1}{2}(|100\rangle_{AB} - |101\rangle_{AB} - |110\rangle_{AB} + |111\rangle_{AB}) = \begin{bmatrix} 1/2 \\ -1/2 \\ -1/2 \\ 1/2 \end{bmatrix}$$

Exercise 7:

$$\cdot (\hat{x}_A \otimes \hat{x}_B) |++\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1/2 \\ 1/2 \\ 1/2 \\ 1/2 \end{bmatrix} = \begin{bmatrix} 1/2 \\ 1/2 \\ 1/2 \\ 1/2 \end{bmatrix} = |++\rangle_{AB} \text{ : } |++\rangle_{AB} \text{ is an eigenvector of } \hat{x}_A \otimes \hat{x}_B \text{ of eigenvalue } +1$$

$$\cdot (\hat{x}_A \otimes \hat{x}_B) |+-\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1/2 \\ -1/2 \\ 1/2 \\ -1/2 \end{bmatrix} = \begin{bmatrix} -1/2 \\ 1/2 \\ -1/2 \\ 1/2 \end{bmatrix} = |-+\rangle_{AB} \text{ : } |-+\rangle_{AB} \text{ is an eigenvector of } \hat{x}_A \otimes \hat{x}_B \text{ of eigenvalue } -1$$

$$\cdot (\hat{x}_A \otimes \hat{x}_B) |-\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1/2 \\ 1/2 \\ -1/2 \\ -1/2 \end{bmatrix} = \begin{bmatrix} -1/2 \\ -1/2 \\ 1/2 \\ 1/2 \end{bmatrix} = -|-\rangle_{AB} : |-\rangle_{AB} \text{ is an eigenvector of } \hat{x}_A \otimes \hat{x}_B \text{ of eigenvalue } -1$$

$$\cdot (\hat{x}_A \otimes \hat{x}_B) |--\rangle_{AB} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1/2 \\ -1/2 \\ -1/2 \\ 1/2 \end{bmatrix} = \begin{bmatrix} 1/2 \\ -1/2 \\ -1/2 \\ 1/2 \end{bmatrix} = |--\rangle_{AB} : |--\rangle_{AB} \text{ is an eigenvector of } \hat{x}_A \otimes \hat{x}_B \text{ of eigenvalue } +1$$

Exercise 8:

$$\hat{z}_A + \hat{z}_B \equiv \hat{z}_A \otimes \mathbb{1}_B + \mathbb{1}_A \otimes \hat{z}_B = \begin{bmatrix} 1[1 & 0] & 0[1 & 0] \\ 0[0 & 1] & -1[1 & 0] \end{bmatrix} + \begin{bmatrix} 1[1 & 0] & 0[1 & 0] \\ 0[0 & -1] & 1[1 & 0] \end{bmatrix}$$

$$= \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

$$\hat{z}_A + \hat{z}_B = \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 \end{bmatrix} = 2|00\rangle_{AB} - 2|11\rangle_{AB} \text{ has eigenvalue } -2. \text{ The two others can be chosen but must satisfy orthonormality so we choose } |01\rangle_{AB}, |10\rangle_{AB} \text{ both of eigenvalue } 0.$$

Explicit calculation:

$$(\hat{z}_A + \hat{z}_B) |00\rangle_{AB} = \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = 2|00\rangle_{AB} : |00\rangle_{AB} \text{ is an eigenvector of } \hat{z}_A + \hat{z}_B \text{ of eigenvalue } 2$$

$$(\hat{z}_A + \hat{z}_B) |01\rangle_{AB} = \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix} = 0|01\rangle_{AB} : |01\rangle_{AB} \text{ is an eigenvector of } \hat{z}_A + \hat{z}_B \text{ of eigenvalue } 0$$

$$(\hat{z}_A + \hat{z}_B) |10\rangle_{AB} = \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix} = 0|10\rangle_{AB} : |10\rangle_{AB} \text{ is an eigenvector of } \hat{z}_A + \hat{z}_B \text{ of eigenvalue } 0$$

$$(\hat{z}_A + \hat{z}_B) |11\rangle_{AB} = \begin{bmatrix} 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} = -2|11\rangle_{AB} : |11\rangle_{AB} \text{ is an eigenvector of } \hat{z}_A + \hat{z}_B \text{ of eigenvalue } -2$$

Exercise 9:

$$|\psi_A\rangle = \alpha_0|0\rangle_A + \alpha_1|1\rangle_A \quad |\psi_B\rangle = \beta_0|0\rangle_B + \beta_1|1\rangle_B \quad |\psi\rangle_{AB} = |\psi_A\rangle \otimes |\psi_B\rangle = \begin{bmatrix} \alpha_0\beta_0 \\ \alpha_0\beta_1 \\ \alpha_1\beta_0 \\ \alpha_1\beta_1 \end{bmatrix}$$

Dirac's notation:

$$\langle \psi | \psi \rangle_{AB} = (\langle \psi_A | \otimes \langle \psi_B |) (|\psi\rangle_{AB} \otimes |\psi\rangle_{AB}) = \underbrace{\langle \psi_A | \psi_A \rangle}_{=1} \cdot \underbrace{\langle \psi_B | \psi_B \rangle}_{=1} = 1 \quad \blacksquare$$

Matrix form:

$$\langle \psi | \psi \rangle_{AB} = [\alpha_0^* \beta_0^* \alpha_0^* \beta_1^* \alpha_1^* \beta_0^* \alpha_1^* \beta_1^*] \begin{bmatrix} \alpha_0\beta_0 \\ \alpha_0\beta_1 \\ \alpha_1\beta_0 \\ \alpha_1\beta_1 \end{bmatrix} = \alpha_0\alpha_0^* \beta_0\beta_0^* + \alpha_0\alpha_0^* \beta_1\beta_1^* + \alpha_1\alpha_1^* \beta_0\beta_0^* + \alpha_1\alpha_1^* \beta_1\beta_1^*$$

$$\langle \psi | \psi \rangle_{AB} = |\alpha_0|^2 |\beta_0|^2 + |\alpha_0|^2 |\beta_1|^2 + |\alpha_1|^2 |\beta_0|^2 + |\alpha_1|^2 |\beta_1|^2 = \underbrace{|\alpha_0|^2 + |\alpha_1|^2}_{=1} \underbrace{|\beta_0|^2 + |\beta_1|^2}_{=1} = 1 \quad \blacksquare$$

Exercise 10:

$$\hat{z}_A \hat{x}_B = \hat{z}_A \otimes \hat{x}_B - \hat{x}_A \otimes \hat{z}_B = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

$$(\hat{z}_A \otimes \hat{x}_B) (|\psi\rangle_A \otimes |\psi\rangle_B) = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \end{bmatrix} \begin{bmatrix} \alpha_0\beta_0 \\ \alpha_0\beta_1 \\ \alpha_1\beta_0 \\ \alpha_1\beta_1 \end{bmatrix} = \begin{bmatrix} \alpha_0\beta_1 \\ \alpha_0\beta_0 \\ -\alpha_1\beta_1 \\ -\alpha_1\beta_0 \end{bmatrix}$$

$$\hat{z}_A |\psi_A\rangle \otimes \hat{x}_B |\psi_B\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} \otimes \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} \beta_0 \\ \beta_1 \end{bmatrix} = \begin{bmatrix} \alpha_0 \\ -\alpha_1 \end{bmatrix} \otimes \begin{bmatrix} \beta_1 \\ \beta_0 \end{bmatrix} = \begin{bmatrix} \alpha_0\beta_1 \\ \alpha_0\beta_0 \\ -\alpha_1\beta_1 \\ -\alpha_1\beta_0 \end{bmatrix}$$

We see that $(\hat{z}_A \otimes \hat{x}_B) (|\psi\rangle_A \otimes |\psi\rangle_B) = \hat{z}_A |\psi_A\rangle \otimes \hat{x}_B |\psi_B\rangle$

Exercise 11:

The length of the resulting state vector is 2^N

Number of real parameters: each qubit only needs 2 real parameters so N qubits require $2 \cdot N$ real parameters.

Exercise 12:

N -Qubits so $2 \cdot 2^N$ real numbers, but the norm removes one real number and global phase removes also one. It turns out in this case we need $2 \cdot 2^N - 2 = 2(2^N - 1)$

Number of real numbers required to describe this state is $2(2^N - 1)$

Exercise 13:

$N = 256$ qubits

- Real number for a product state: $\# = 2 \cdot N = 2 \cdot 256 = 512$

- Real parameters for a general quantum state $\# = 2(2^N - 1) = 2(2^{256} - 1) \cong 2,3158 \cdot 10^{77}$
but 10^{80} is slightly greater than the number of real parameters required

We can use all atoms in the universe but this does not really make any sense.

Exercise 14:

- $|14E1\rangle = \frac{1}{\sqrt{2}}|100\rangle + \frac{1}{\sqrt{2}}|111\rangle$. If this is a product state, it could be written in the following form:

$$|14_A\rangle \otimes |14_B\rangle = \begin{bmatrix} \alpha_0\beta_0 \\ \alpha_0\beta_1 \\ \alpha_1\beta_0 \\ \alpha_1\beta_1 \end{bmatrix}^{\begin{matrix} |100\rangle \\ |101\rangle \\ |110\rangle \\ |111\rangle \end{matrix}}$$

so: $\alpha_0\beta_0 = \frac{1}{\sqrt{2}} \Rightarrow \alpha_0, \beta_0 \neq 0$

$\alpha_0\beta_1 = 0 \Rightarrow \beta_1 = 0$

$\alpha_1\beta_0 = 0 \Rightarrow \alpha_1 = 0$

$\alpha_1\beta_1 = \frac{1}{\sqrt{2}} \Rightarrow \alpha_1, \beta_1 \neq 0$

We have a contradiction since α_1 and β_1 should be in the mean time 0 and non-zero. $|14E1\rangle$ can not be factorized as a product state!

- $|14E2\rangle = \frac{1}{\sqrt{2}}|101\rangle + \frac{1}{\sqrt{2}}|110\rangle \stackrel{?}{=} |14_A\rangle \otimes |14_B\rangle$

$$\left. \begin{array}{l} \alpha_0\beta_0 = 0 \Rightarrow \alpha_0 \text{ or } \beta_0 = 0 \\ \alpha_0\beta_1 = \frac{1}{\sqrt{2}} \Rightarrow \alpha_0, \beta_1 \neq 0 \\ \alpha_1\beta_0 = \frac{1}{\sqrt{2}} \Rightarrow \alpha_1, \beta_0 \neq 0 \\ \alpha_1\beta_1 = 0 \Rightarrow \alpha_1 \text{ or } \beta_1 = 0 \end{array} \right\}$$

We have a contradiction because we need $\alpha_0, \alpha_1, \beta_0, \beta_1$ to be non zero but $\alpha_0\beta_0$ and $\alpha_1\beta_1$ should be equal to zero which is not possible in the mean time! Therefore $|14E2\rangle$ can't be factorized as a product state.

Exercise 15:

Here we have $a=1$ and measure correspond to a collapse of $|14\rangle$ onto state $|11\rangle$:

$$\text{Prob}(M=-2) = |\langle 11|14\rangle|^2 = |14_{11}|^2 = \left| \langle 11| \left(\frac{1}{2}|100\rangle + \frac{1}{2}|101\rangle + \frac{1}{2}|110\rangle + \frac{1}{2}|111\rangle \right) \right|^2 = \left| \frac{1}{2} \right|^2 = \frac{1}{4}.$$

$$\text{Prob}(M=-2) = \frac{1}{4}$$

Exercise 16:

$$\text{Prob}(M \neq 2, -2) = |\langle 01|14\rangle|^2 + |\langle 10|14\rangle|^2 = |14_{01}|^2 + |14_{10}|^2 = \left| \frac{1}{2} \right|^2 + \left| \frac{1}{2} \right|^2 = \frac{1}{4} + \frac{1}{4} = \frac{1}{2}$$

$$\text{or Prob}(M \neq 2, -2) = 1 - |14_{00}|^2 + |14_{11}|^2 = 1 - \frac{1}{4} - \frac{1}{4} = \frac{1}{2}.$$

$$\text{Prob}(M \neq 2, -2) = \frac{1}{2}$$

Exercise 17:

$$\text{Prob}(M=0) = |\langle \psi_0 | \hat{\gamma} \rangle|^2$$

$$\begin{aligned} |\langle \psi_0 | \hat{\gamma} \rangle|^2 &= \left| \frac{1}{\sqrt{|\psi_{01}|^2 + |\psi_{10}|^2}} (\psi_{01}^* \langle 01 | + \psi_{10}^* \langle 10 |) \hat{\gamma} \right|^2 \\ &= \left| \frac{1}{\sqrt{|\psi_{01}|^2 + |\psi_{10}|^2}} (|\psi_{01}|^2 + |\psi_{10}|^2) \right|^2 \\ &= \left| \sqrt{|\psi_{01}|^2 + |\psi_{10}|^2} \right|^2 \\ &= |\psi_{01}|^2 + |\psi_{10}|^2 \end{aligned}$$

$$\text{Prob}(M=0) = |\langle \psi_0 | \hat{\gamma} \rangle|^2 = \frac{1}{2} = \text{Prob}(M \neq -2)$$

Exercise 18:

We know that $P(M=0) = \frac{1}{2}$ so in order to obtain an entangled state, we have to try the measurement on a product state, on average $N = \frac{1}{P(M=0)} = \frac{1}{1/2} = 2$.

We need to try measurement on average $N = \frac{1}{P(M=0)} = 2$

Exercise 19:

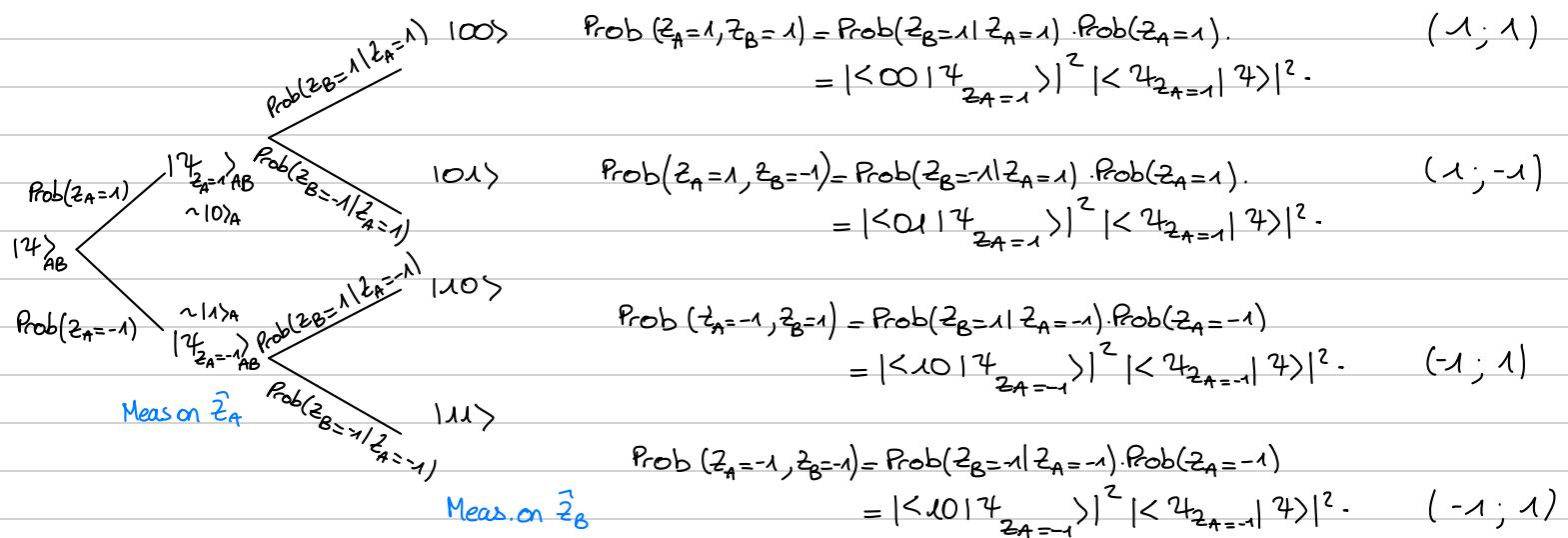
$$|\psi\rangle = \psi_{00} |00\rangle + \psi_{01} |01\rangle + \psi_{10} |10\rangle + \psi_{11} |11\rangle$$

$$|\psi_{z_A=1}\rangle = |\psi\rangle_A \otimes \frac{\psi_{00} |0\rangle_B + \psi_{01} |1\rangle_B}{\sqrt{|\psi_{00}|^2 + |\psi_{01}|^2}}$$

$$|\psi_{z_A=-1}\rangle = |\psi\rangle_A \otimes \frac{\psi_{10} |0\rangle_B + \psi_{11} |1\rangle_B}{\sqrt{|\psi_{10}|^2 + |\psi_{11}|^2}}$$

First we measure on \hat{z}_A . The output is either +1 or -1 and state $|\psi\rangle$ collapses on state $|\psi_{z_A=\pm 1}\rangle$ with proba $|\langle \psi_{z_A=\pm 1} | \psi \rangle|^2$ or on $|\psi_{z_A=-1}\rangle$ with proba $|\langle \psi_{z_A=-1} | \psi \rangle|^2$. If $z_A=1$, we have $|0\rangle_A$ and if $z_A=-1$ we collapse on Alice's lab on $|1\rangle_A$.

Then we measure on \hat{z}_B . The output is either +1 or -1 and state $|\psi_{z_A=\pm 1}\rangle$ (depending on outcome of z_A) collapses on $|0\rangle_B$ or $|1\rangle_B$ for Bob particle. We need to use conditional probabilities to see each probability of each outcome.



(1,1):

$$\text{Prob}((1,1)) = \text{Prob}(\hat{z}_B=1, \hat{z}_A=1) = \text{Prob}(\hat{z}_B=1 | \hat{z}_A=1) \cdot \text{Prob}(\hat{z}_A=1)$$

$$\begin{aligned} &= |\langle 4_{z_A=1} | \rangle|^2 |\langle 001 | \rangle|^2 = \left[14_{001}^2 + 14_{011}^2 \right] |\langle 001 \cdot 10 \rangle_A \otimes \frac{4_{0010} \rangle_B + 4_{0111} \rangle_B}{\sqrt{14_{001}^2 + 14_{011}^2}}|^2 \\ &= \left[14_{001}^2 + 14_{011}^2 \right] |\langle 010 \rangle_A \cdot \frac{4_{00} \langle 010 \rangle_B + 4_{01} \langle 011 \rangle_B}{\sqrt{14_{001}^2 + 14_{011}^2}}|^2 \\ &= \left(14_{001}^2 + 14_{011}^2 \right) \frac{14_{001}^2}{14_{001}^2 + 14_{011}^2} \end{aligned}$$

$$\text{Prob}((1,1)) = 14_{001}^2$$

(1,-1):

$$\text{Prob}((1,-1)) = \text{Prob}(\hat{z}_B=-1, \hat{z}_A=1) = \text{Prob}(\hat{z}_B=-1 | \hat{z}_A=1) \cdot \text{Prob}(\hat{z}_A=1)$$

$$\begin{aligned} &= |\langle 4_{z_A=1} | \rangle|^2 |\langle 011 | \rangle|^2 = \left[14_{001}^2 + 14_{011}^2 \right] |\langle 011 \cdot 10 \rangle_A \otimes \frac{4_{0010} \rangle_B + 4_{0111} \rangle_B}{\sqrt{14_{001}^2 + 14_{011}^2}}|^2 \\ &= \left[14_{001}^2 + 14_{011}^2 \right] |\langle 010 \rangle_A \cdot \frac{4_{00} \langle 110 \rangle_B + 4_{01} \langle 111 \rangle_B}{\sqrt{14_{001}^2 + 14_{011}^2}}|^2 \\ &= \left(14_{001}^2 + 14_{011}^2 \right) \frac{14_{011}^2}{14_{001}^2 + 14_{011}^2} \quad \text{Prob}((1,-1)) = 14_{011}^2 \end{aligned}$$

(-1,1):

$$\text{Prob}((-1,1)) = \text{Prob}(\hat{z}_B=1, \hat{z}_A=-1) = \text{Prob}(\hat{z}_B=1 | \hat{z}_A=-1) \cdot \text{Prob}(\hat{z}_A=-1)$$

$$\begin{aligned} &= |\langle 4_{z_A=-1} | \rangle|^2 |\langle 101 | \rangle|^2 = \left[14_{101}^2 + 14_{111}^2 \right] |\langle 101 \cdot 11 \rangle_A \otimes \frac{4_{1010} \rangle_B + 4_{1111} \rangle_B}{\sqrt{14_{101}^2 + 14_{111}^2}}|^2 \\ &= \left[14_{101}^2 + 14_{111}^2 \right] |\langle 111 \rangle_A \cdot \frac{4_{10} \langle 010 \rangle_B + 4_{11} \langle 011 \rangle_B}{\sqrt{14_{101}^2 + 14_{111}^2}}|^2 \\ &= \left(14_{101}^2 + 14_{111}^2 \right) \frac{14_{101}^2}{14_{101}^2 + 14_{111}^2} \end{aligned}$$

$$\text{Prob}((-1,1)) = 14_{101}^2$$

(-1,-1):

$$\text{Prob}((-1,-1)) = \text{Prob}(\hat{z}_B=-1, \hat{z}_A=-1) = \text{Prob}(\hat{z}_B=-1 | \hat{z}_A=-1) \cdot \text{Prob}(\hat{z}_A=-1)$$

$$\begin{aligned} &= |\langle 4_{z_A=-1} | \rangle|^2 |\langle 111 | \rangle|^2 = \left[14_{101}^2 + 14_{111}^2 \right] |\langle 111 \cdot 11 \rangle_A \otimes \frac{4_{1010} \rangle_B + 4_{1111} \rangle_B}{\sqrt{14_{101}^2 + 14_{111}^2}}|^2 \\ &= \left[14_{101}^2 + 14_{111}^2 \right] |\langle 111 \rangle_A \cdot \frac{4_{10} \langle 110 \rangle_B + 4_{11} \langle 111 \rangle_B}{\sqrt{14_{101}^2 + 14_{111}^2}}|^2 \\ &= \left(14_{101}^2 + 14_{111}^2 \right) \frac{14_{111}^2}{14_{101}^2 + 14_{111}^2} \end{aligned}$$

$$\text{Prob}((-1,-1)) = 14_{111}^2$$

Exercise 20:

Measurements are done locally, so doing before the measurement on Alice's lab or in Bob's lab, should not change the probability of the different output. In any case we compute $P(z_A=i, z_B=j)$ so nothing should change.

$$\hat{z}_B |00\rangle = +|100\rangle \quad \text{for a general 2 Qubit-state: } |1\rangle = |00100\rangle + |01101\rangle + |10110\rangle + |11111\rangle.$$

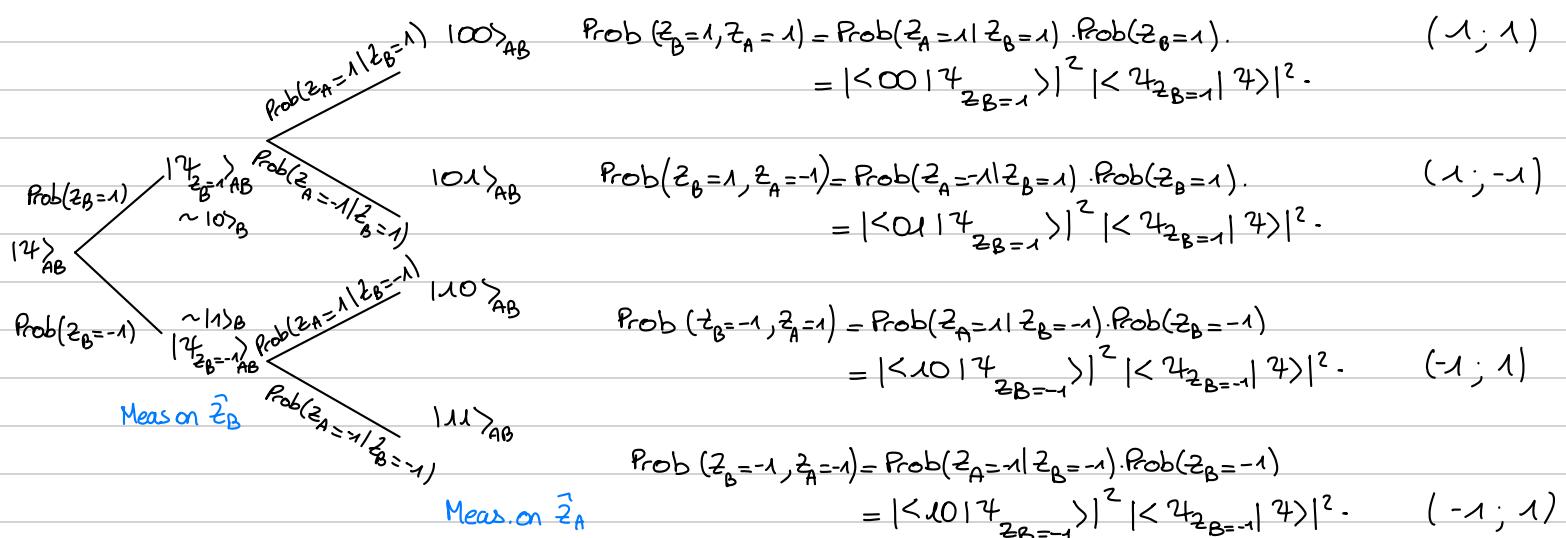
$$\hat{z}_B |01\rangle = -|01\rangle \quad \text{We have a partial collapse as previously seen.}$$

$$\hat{z}_B |10\rangle = +|10\rangle$$

$$\hat{z}_B |11\rangle = -|11\rangle \quad |1\rangle_{z_B=1} = \frac{|00100\rangle_A + |10110\rangle_A}{\sqrt{|1400|^2 + |1410|^2}} \otimes |10\rangle_B \quad |1\rangle_{z_B=-1} = \frac{|01101\rangle_A + |11111\rangle_A}{\sqrt{|1401|^2 + |1411|^2}} \times |11\rangle_B$$

First we measure on \hat{z}_B . The output is either +1 or -1 and state $|1\rangle$ collapses on state $|1\rangle_{z_B=1}$ with proba $|\langle 1\rangle_{z_B=1}|^2$ or on $|1\rangle_{z_B=-1}$ with proba $|\langle 1\rangle_{z_B=-1}|^2$. If $z_B=1$, we have $|10\rangle_B$ and if $z_B=-1$ we collapse on Bob's lab on $|11\rangle_B$.

Then we measure on \hat{z}_A . The output is either +1 or -1 and state $|1\rangle_{z_B=1}$ or $|1\rangle_{z_B=-1}$ (depending on outcome of z_B) collapses on $|10\rangle_A$ or $|11\rangle_A$ for Alice particle. We need to use conditional probabilities to see each probability of each outcome.



$$\text{Prob}(z_B=1) = |\langle 1 | 1 \rangle_{z_B=1}|^2 = |1400|^2 + |1410|^2.$$

$$\text{Prob}(z_B=-1) = |\langle 1 | 1 \rangle_{z_B=-1}|^2 = |1401|^2 + |1411|^2$$

(1, 1):

$$\text{Prob}((1, 1)) = \text{Prob}(\hat{z}_B=1, \hat{z}_A=1) = \text{Prob}(\hat{z}_A=1 | \hat{z}_B=1) \cdot \text{Prob}(\hat{z}_B=1)$$

$$= |\langle 00 |_{AB} \cdot \frac{|00100\rangle_A + |10110\rangle_A}{\sqrt{|1400|^2 + |1410|^2}} \otimes |10\rangle_B|^2 (|1400|^2 + |1410|^2)$$

$$= \left| \frac{400 \langle 010 \rangle_A}{\sqrt{|1400|^2 + |1410|^2}} + \frac{410 \langle 011 \rangle_A}{\sqrt{|1400|^2 + |1410|^2}} \right|^2 (|1400|^2 + |1410|^2)$$

$$= \frac{|1400|^2}{|1400|^2 + |1410|^2} \cdot (|1400|^2 + |1410|^2)$$

$$\text{Prob}((1, 1)) = |1400|^2$$

(1,-1):

$$\text{Prob}((1; -1)) = \text{Prob}(\hat{z}_B = -1, \hat{z}_A = 1) = \text{Prob}(\hat{z}_A = 1 | \hat{z}_B = -1) \cdot \text{Prob}(\hat{z}_B = -1)$$

$$= |\langle 4_{z_B=-1}| 4 \rangle|^2 |\langle 01 | 4_{z_B=-1} \rangle|^2 = \left[14_{01}^2 + 14_{11}^2 \right] \left| \langle 01 | \frac{4_{01} 10\rangle_A + 4_{11} 11\rangle_A}{\sqrt{14_{01}^2 + 14_{11}^2}} \otimes 10 \rangle_B \right|^2$$

$$= \left[14_{01}^2 + 14_{11}^2 \right] \left| \frac{4_{01} \langle 01 | 0 \rangle_A + 4_{11} \langle 01 | 1 \rangle_A}{\sqrt{14_{01}^2 + 14_{11}^2}} \cdot \langle 11 | 1 \rangle_B \right|^2$$

$$= \left(14_{01}^2 + 14_{11}^2 \right) \frac{14_{01}^2}{14_{01}^2 + 14_{11}^2}$$

$$\text{Prob}((1; -1)) = 14_{01}^2$$

(-1,1):

$$\text{Prob}((-1; 1)) = \text{Prob}(\hat{z}_B = 1, \hat{z}_A = -1) = \text{Prob}(\hat{z}_A = -1 | \hat{z}_B = 1) \cdot \text{Prob}(\hat{z}_B = 1)$$

$$= \left| \langle 10 | \frac{4_{00} 10\rangle_A + 4_{10} 11\rangle_A}{\sqrt{14_{00}^2 + 14_{10}^2}} \otimes 10 \rangle_B \right|^2 (14_{00}^2 + 14_{10}^2)$$

$$= \left| \frac{4_{00} \langle 1 | 0 \rangle_A}{\sqrt{14_{00}^2 + 14_{10}^2}} + \frac{4_{10} \langle 1 | 1 \rangle_A}{\sqrt{14_{00}^2 + 14_{10}^2}} \right|^2 (14_{00}^2 + 14_{10}^2)$$

$$= \frac{14_{10}^2}{14_{00}^2 + 14_{10}^2} (14_{00}^2 + 14_{10}^2)$$

$$\text{Prob}((1; 1)) = 14_{10}^2$$

(-1,-1):

$$\text{Prob}((-1; -1)) = \text{Prob}(\hat{z}_B = -1, \hat{z}_A = -1) = \text{Prob}(\hat{z}_A = -1 | \hat{z}_B = -1) \cdot \text{Prob}(\hat{z}_B = -1)$$

$$= |\langle 4_{z_B=-1}| 4 \rangle|^2 |\langle 11 | 4_{z_B=-1} \rangle|^2 = \left[14_{01}^2 + 14_{11}^2 \right] \left| \langle 11 | \frac{4_{01} 10\rangle_A + 4_{11} 11\rangle_A}{\sqrt{14_{01}^2 + 14_{11}^2}} \otimes 11 \rangle_B \right|^2$$

$$= \left[14_{01}^2 + 14_{11}^2 \right] \left| \frac{4_{01} \langle 1 | 0 \rangle_A + 4_{11} \langle 1 | 1 \rangle_A}{\sqrt{14_{01}^2 + 14_{11}^2}} \cdot \langle 11 | 1 \rangle_B \right|^2$$

$$= \left(14_{01}^2 + 14_{11}^2 \right) \frac{14_{11}^2}{14_{01}^2 + 14_{11}^2}$$

$$\text{Prob}((1; -1)) = 14_{11}^2$$

We indeed obtain same probabilities.

Exercise 21:

$$\begin{aligned} \hat{z}_A \otimes \hat{z}_B |100\rangle_{AB} &= \hat{z}_A |10\rangle_A \otimes \hat{z}_B |10\rangle_B = (+1) |10\rangle_A \otimes (+1) |10\rangle_B = + |100\rangle_{AB} && \text{eigenvector: } |100\rangle ; \text{ eigenvalue: } +1 \\ \hat{z}_A \otimes \hat{z}_B |101\rangle_{AB} &= \hat{z}_A |10\rangle_A \otimes \hat{z}_B |11\rangle_B = (+1) |10\rangle_A \otimes (-1) |10\rangle_B = - |101\rangle_{AB} && \text{eigenvector: } |101\rangle ; \text{ eigenvalue: } -1 \\ \hat{z}_A \otimes \hat{z}_B |110\rangle_{AB} &= \hat{z}_A |11\rangle_A \otimes \hat{z}_B |10\rangle_B = (-1) |11\rangle_A \otimes (+1) |10\rangle_B = - |110\rangle_{AB} && \text{eigenvector: } |110\rangle ; \text{ eigenvalue: } -1 \\ \hat{z}_A \otimes \hat{z}_B |111\rangle_{AB} &= \hat{z}_A |11\rangle_A \otimes \hat{z}_B |11\rangle_B = (-1) |11\rangle_A \otimes (-1) |11\rangle_B = + |111\rangle_{AB} && \text{eigenvector: } |111\rangle ; \text{ eigenvalue: } +1 \end{aligned}$$

Since we cannot make the difference between $|100\rangle$ and $|111\rangle$ ($M=+1$) but also $|101\rangle_{AB}$ and $|110\rangle_{AB}$ ($M=-1$)
One can define the partial collapse states as: $-|14_{M=1}\rangle = \frac{1}{\sqrt{14_{00}^2 + 14_{11}^2}} (4_{00} |100\rangle_{AB} + 4_{11} |111\rangle_{AB})$

These two vectors are the outcomes

$$-|14_{M=-1}\rangle = \frac{1}{\sqrt{14_{01}^2 + 14_{10}^2}} (4_{01} |101\rangle_{AB} + 4_{10} |110\rangle_{AB})$$

The probabilities are for a general 2-Qubit state $| \psi \rangle$:

$$\begin{aligned} \text{Prob}(M=1) &= |\langle \psi_{M=1} | \psi \rangle|^2 = \left| \frac{1}{\sqrt{|\psi_{00}|^2 + |\psi_{11}|^2}} (|\psi_{00}\rangle_{AB} + |\psi_{11}\rangle_{AB}) (|\psi_{00}\rangle_{AB} + |\psi_{11}\rangle_{AB} + |\psi_{01}\rangle_{AB} + |\psi_{10}\rangle_{AB}) \right|^2 \\ &= \left| \frac{|\psi_{00}|^2 + |\psi_{11}|^2}{\sqrt{|\psi_{00}|^2 + |\psi_{11}|^2}} \right|^2 \end{aligned}$$

$$\text{Prob}(M=1) = |\psi_{00}|^2 + |\psi_{11}|^2$$

$$\begin{aligned} \text{Prob}(M=-1) &= |\langle \psi_{M=-1} | \psi \rangle|^2 = \left| \frac{1}{\sqrt{|\psi_{01}|^2 + |\psi_{10}|^2}} (|\psi_{01}\rangle_{AB} + |\psi_{10}\rangle_{AB}) (|\psi_{00}\rangle_{AB} + |\psi_{11}\rangle_{AB} + |\psi_{01}\rangle_{AB} + |\psi_{10}\rangle_{AB}) \right|^2 \\ &= \left| \frac{|\psi_{01}|^2 + |\psi_{10}|^2}{\sqrt{|\psi_{01}|^2 + |\psi_{10}|^2}} \right|^2 \end{aligned}$$

$$\text{Prob}(M=-1) = |\psi_{01}|^2 + |\psi_{10}|^2$$

Finally for $| \psi \rangle = |+\rangle_A \otimes |+\rangle$ we get:

$$\text{Prob}(M=1) = |\langle \psi_{M=1} | \psi \rangle|^2 = |\psi_{00}|^2 + |\psi_{11}|^2 = \frac{1}{4} + \frac{1}{4} = \frac{1}{2}.$$

\hookrightarrow outcome $|\psi_{M=1}\rangle = \frac{1}{\sqrt{2}} (|00\rangle_{AB} + |11\rangle_{AB})$ with Proba: $\text{Prob}(M=1) = \frac{1}{2}$ (Bell state $|B00\rangle$)

$$\text{Prob}(M=-1) = |\langle \psi_{M=-1} | \psi \rangle|^2 = |\psi_{01}|^2 + |\psi_{10}|^2 = \frac{1}{4} + \frac{1}{4} = \frac{1}{2}.$$

\hookrightarrow outcome: $|\psi_{M=-1}\rangle = \frac{1}{\sqrt{2}} (|01\rangle_{AB} + |10\rangle_{AB})$ with Proba: $\text{Prob}(M=-1) = \frac{1}{2}$ (Bell state $|B01\rangle$.)

Exercise 22:

$$|\psi_p\rangle = |\psi_A\rangle \otimes |\psi_B\rangle \quad |\psi_A\rangle = \alpha_0 |0\rangle_A + \alpha_1 |1\rangle_A. \quad \alpha_0 = \cos\left(\frac{\theta}{2}\right) \quad \alpha_1 = \sin\left(\frac{\theta}{2}\right) e^{i\phi}. \quad \theta \in [0; \pi] \quad \phi \in [0; 2\pi]$$

$$\langle \psi_A | \hat{x} | \psi_A \rangle = \sin(\theta) \cos(\phi) \Rightarrow |\langle \psi_A | \hat{x} | \psi_A \rangle|^2 = \sin^2(\theta) \cos^2(\phi)$$

$$\langle \psi_A | \hat{y} | \psi_A \rangle = \sin(\theta) \sin(\phi) \Rightarrow |\langle \psi_A | \hat{y} | \psi_A \rangle|^2 = \sin^2(\theta) \sin^2(\phi)$$

$$\langle \psi_A | \hat{z} | \psi_A \rangle = \cos(\theta) \Rightarrow |\langle \psi_A | \hat{z} | \psi_A \rangle|^2 = \cos^2(\theta)$$

With Bloch representation:

$$\begin{aligned} |\langle \psi_A | \hat{x} | \psi_A \rangle|^2 + |\langle \psi_A | \hat{y} | \psi_A \rangle|^2 + |\langle \psi_A | \hat{z} | \psi_A \rangle|^2 &= \sin^2(\theta) \cos^2(\phi) + \sin^2(\theta) \sin^2(\phi) + \cos^2(\theta) \\ &= \sin^2(\theta) (\underbrace{\cos^2(\phi) + \sin^2(\phi)}_{=1}) + \cos^2(\theta) \\ &= \sin^2(\theta) + \cos^2(\theta) \\ &= 1 \end{aligned}$$

without Bloch representation:

$$\langle \psi_A | \hat{x} | \psi_A \rangle = (\alpha_0^* \langle 0 | + \alpha_1^* \langle 1 |) (\alpha_0 |0\rangle_A + \alpha_1 |1\rangle) = \alpha_0^* \alpha_1 + \alpha_1^* \alpha_0.$$

$$\langle \psi_A | \hat{y} | \psi_A \rangle = [\alpha_0^* \alpha_1] \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix} = [\alpha_0^* \alpha_1^*] \begin{bmatrix} -i\alpha_1 \\ i\alpha_0 \end{bmatrix} = -i\alpha_0^* \alpha_1 + i\alpha_1^* \alpha_0 \quad \langle \psi_A | \hat{z} | \psi_A \rangle = \frac{1}{\sqrt{2}} \quad \langle \psi_A | \hat{z} | \psi_A \rangle = \frac{-i}{\sqrt{2}}$$

$$\langle \psi_A | \hat{z} | \psi_A \rangle = (\alpha_0^* \langle 0 | + \alpha_1^* \langle 1 |) (\alpha_0 |0\rangle - \alpha_1 |1\rangle) = |\alpha_0|^2 - |\alpha_1|^2$$

$$\begin{aligned}
|\langle \Psi_A | \vec{x} | \Psi_A \rangle|^2 &= |(a_0 - ib_0)(a_1 + ib_1) + (a_0 + ib_0)(a_1 - ib_1)|^2 \\
&= |a_0 a_1 + ia_0 b_1 - ia_0 b_1 + a_0 b_0 + a_0 a_1 - ia_0 b_1 + ib_0 a_1 + b_0 b_1|^2 \\
&= |2a_0 a_1 + 2b_0 b_1|^2 = 4a_0^2 a_1^2 + 4b_0^2 b_1^2 + 8a_0 a_1 b_0 b_1
\end{aligned}$$

$$\begin{aligned}
|\langle \Psi_A | \vec{y} | \Psi_A \rangle|^2 &= |-i(a_0 - ib_0)(a_1 + ib_1) + i(a_0 + ib_0)(a_1 - ib_1)|^2 \\
&= |i(a_0 a_1 - ia_0 b_1 + ia_1 b_0 - b_0 b_1) + i(a_0 a_1 - ia_0 b_1 + ib_0 a_1 + b_0 b_1)|^2 \\
&= |-2ia_0 b_1 + 2ia_1 b_0|^2 \\
&= |2a_0 b_1 - 2a_1 b_0|^2 = 4a_0^2 b_1^2 + 4a_1^2 b_0^2 - 8a_0 a_1 b_0 b_1
\end{aligned}$$

$$|\langle \Psi_A | \vec{z} | \Psi_A \rangle|^2 = |\alpha_0|^2 - |\alpha_1|^2 = |\alpha_0|^4 + |\alpha_1|^4 - 2|\alpha_0|^2 |\alpha_1|^2$$

$$\begin{aligned}
|\langle \Psi_A | \vec{x} | \Psi_A \rangle|^2 + |\langle \Psi_A | \vec{y} | \Psi_A \rangle|^2 + |\langle \Psi_A | \vec{z} | \Psi_A \rangle|^2 &= 4a_0^2 a_1^2 + 4b_0^2 b_1^2 + 8a_0 a_1 b_0 b_1 \\
&\quad + 4a_0^2 b_1^2 + 4a_1^2 b_0^2 - 8a_0 a_1 b_0 b_1 \\
&\quad + |\alpha_0|^4 + |\alpha_1|^4 - 2|\alpha_0|^2 |\alpha_1|^2 \\
&= 4|\alpha_0|^2 (a_1^2 + b_1^2) + 4b_0^2 (a_1^2 + b_1^2) + |\alpha_0|^4 + |\alpha_1|^4 - 2|\alpha_0|^2 |\alpha_1|^2 \\
&= 4|\alpha_0|^2 |\alpha_1|^2 + |\alpha_0|^4 + |\alpha_1|^4 - 2|\alpha_0|^2 |\alpha_1|^2 \\
&= |\alpha_0|^4 + |\alpha_1|^4 + 2|\alpha_0|^2 |\alpha_1|^2 \\
&= (\alpha_0^2 + \alpha_1^2)^2 \\
&= 1
\end{aligned}$$

$$\begin{aligned}
|\langle \Psi_P | \vec{x}_A | \Psi_P \rangle|^2 &= |\langle \Psi_A | \otimes \langle \Psi_B | (\vec{x}_A \otimes \vec{x}_B) | \Psi_A \rangle \otimes | \Psi_B \rangle|^2 \\
&= |\langle \Psi_A | \vec{x}_A | \Psi_A \rangle \langle \Psi_B | \vec{x}_B | \Psi_B \rangle|^2 \\
&= |\langle \Psi_A | \vec{x}_A | \Psi_A \rangle|^2 = 1
\end{aligned}$$

$$\begin{aligned}
|\langle \Psi_P | \vec{y}_A | \Psi_P \rangle|^2 &= |\langle \Psi_A | \otimes \langle \Psi_B | (\vec{y}_A \otimes \vec{y}_B) | \Psi_A \rangle \otimes | \Psi_B \rangle|^2 \\
&= |\langle \Psi_A | \vec{y}_A | \Psi_A \rangle \langle \Psi_B | \vec{y}_B | \Psi_B \rangle|^2 \\
&= |\langle \Psi_A | \vec{y}_A | \Psi_A \rangle|^2 = 1
\end{aligned}$$

$$\begin{aligned}
|\langle \Psi_P | \vec{z}_A | \Psi_P \rangle|^2 &= |\langle \Psi_A | \otimes \langle \Psi_B | (\vec{z}_A \otimes \vec{z}_B) | \Psi_A \rangle \otimes | \Psi_B \rangle|^2 \\
&= |\langle \Psi_A | \vec{z}_A | \Psi_A \rangle \langle \Psi_B | \vec{z}_B | \Psi_B \rangle|^2 \\
&= |\langle \Psi_A | \vec{z}_A | \Psi_A \rangle|^2 = 1
\end{aligned}$$

Therefore $|\langle \Psi_P | \vec{x}_A | \Psi_P \rangle|^2 + |\langle \Psi_P | \vec{y}_A | \Psi_P \rangle|^2 + |\langle \Psi_P | \vec{z}_A | \Psi_P \rangle|^2 = |\langle \Psi_A | \vec{x} | \Psi_A \rangle|^2 + |\langle \Psi_A | \vec{y} | \Psi_A \rangle|^2 + |\langle \Psi_A | \vec{z} | \Psi_A \rangle|^2 = 1$ (see above)

Exercise 23:

$$\begin{aligned} \langle \Psi_p | \vec{x}_A \vec{x}_B | \Psi_p \rangle &= \langle \Psi_A | \otimes \langle \Psi_B | (\vec{x}_A \otimes \vec{x}_B) | \Psi_A \rangle \otimes | \Psi_B \rangle \\ &= \langle \Psi_A | \vec{x}_A | \Psi_A \rangle \cdot \langle \Psi_B | \vec{x}_B | \Psi_B \rangle \\ &= \sin(\theta_A) \cos(\phi_A) \cdot \sin(\theta_B) \cos(\phi_B) \end{aligned}$$

$$\begin{aligned} \langle \Psi_p | \vec{y}_A \vec{y}_B | \Psi_p \rangle &= \langle \Psi_A | \otimes \langle \Psi_B | (\vec{y}_A \otimes \vec{y}_B) | \Psi_A \rangle \otimes | \Psi_B \rangle \\ &= \langle \Psi_A | \vec{y}_A | \Psi_A \rangle \cdot \langle \Psi_B | \vec{y}_B | \Psi_B \rangle \\ &= \sin(\theta_A) \sin(\phi_A) \cdot \sin(\theta_B) \sin(\phi_B) \end{aligned}$$

$$\begin{aligned} \langle \Psi_p | \vec{z}_A \vec{z}_B | \Psi_p \rangle &= \langle \Psi_A | \otimes \langle \Psi_B | (\vec{z}_A \otimes \vec{z}_B) | \Psi_A \rangle \otimes | \Psi_B \rangle \\ &= \langle \Psi_A | \vec{z}_A | \Psi_A \rangle \cdot \langle \Psi_B | \vec{z}_B | \Psi_B \rangle \\ &= \cos(\theta_A) \cos(\theta_B) \end{aligned}$$

$$\langle \Psi_p | \vec{x}_A \vec{x}_B | \Psi_p \rangle + \langle \Psi_p | \vec{y}_A \vec{y}_B | \Psi_p \rangle + \langle \Psi_p | \vec{z}_A \vec{z}_B | \Psi_p \rangle = \sin(\theta_A) \cos(\phi_A) \cdot \sin(\theta_B) \cos(\phi_B) + \sin(\theta_A) \sin(\phi_A) \cdot \sin(\theta_B) \sin(\phi_B) + \cos(\theta_A) \cos(\theta_B)$$

$$= \begin{bmatrix} \sin(\theta_A) \cos(\phi_A) \\ \sin(\theta_A) \sin(\phi_A) \\ \cos(\theta_A) \end{bmatrix} \cdot \begin{bmatrix} \sin(\theta_B) \cos(\phi_B) \\ \sin(\theta_B) \sin(\phi_B) \\ \sin(\theta_B) \end{bmatrix} = \underbrace{\|\vec{r}_A\|}_{=1} \cdot \underbrace{\|\vec{r}_B\|}_{=1} \cdot \underbrace{\cos(\theta)}_{\in [-1, 1]} \quad \theta = \text{angle between } \vec{r}_A \text{ and } \vec{r}_B$$

These are spherical coordinates so the length of these two vector is 1. The scalar product is maximized when \vec{r}_A and \vec{r}_B are aligned in the same direction giving a scalar product of 1 (because $\|\vec{r}_A\| = \|\vec{r}_B\| = 1$). The minimum value is obtained when \vec{r}_A and \vec{r}_B are aligned in the opposite direction giving a scalar product of -1.

Therefore $-1 \leq \langle \Psi_p | \vec{x}_A \vec{x}_B | \Psi_p \rangle + \langle \Psi_p | \vec{y}_A \vec{y}_B | \Psi_p \rangle + \langle \Psi_p | \vec{z}_A \vec{z}_B | \Psi_p \rangle \leq 1$

Exercise 24:

$|B_0\rangle$:

$$\langle B_0 | B_0 \rangle = \frac{1}{\sqrt{2}} (\langle 011 - 101 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle - \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 01101 \rangle}_{=0} - \underbrace{\langle 01110 \rangle}_{=0} - \underbrace{\langle 10101 \rangle}_{=0} + \underbrace{\langle 10110 \rangle}_{=1}) = \frac{1}{2} (2) = 1$$

$$\langle B_1 | B_0 \rangle = \frac{1}{\sqrt{2}} (\langle 011 + 101 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle - \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 01101 \rangle}_{=1} - \underbrace{\langle 01110 \rangle}_{=0} + \underbrace{\langle 10101 \rangle}_{=0} - \underbrace{\langle 10110 \rangle}_{=1}) = \frac{1}{2} (1 - 1) = 0 = \langle B_0 | B_1 \rangle$$

$$\langle B_2 | B_0 \rangle = \frac{1}{\sqrt{2}} (\langle 001 - 111 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle - \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 00101 \rangle}_{=0} - \underbrace{\langle 00110 \rangle}_{=0} - \underbrace{\langle 11101 \rangle}_{=0} + \underbrace{\langle 11110 \rangle}_{=0}) = 0 = \langle B_0 | B_2 \rangle$$

$$\langle B_3 | B_0 \rangle = \frac{1}{\sqrt{2}} (\langle 001 + 111 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle - \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 00101 \rangle}_{=0} - \underbrace{\langle 00110 \rangle}_{=0} + \underbrace{\langle 11101 \rangle}_{=0} - \underbrace{\langle 11110 \rangle}_{=0}) = 0 = \langle B_0 | B_3 \rangle$$

$|B_1\rangle$:

$$\langle B_1 | B_1 \rangle = \frac{1}{\sqrt{2}} (\langle 011 + 101 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle + \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 01101 \rangle}_{=1} + \underbrace{\langle 01110 \rangle}_{=0} + \underbrace{\langle 10101 \rangle}_{=0} + \underbrace{\langle 10110 \rangle}_{=1}) = \frac{1}{2} (1 + 1) = 1$$

$$\langle B_2 | B_1 \rangle = \frac{1}{\sqrt{2}} (\langle 001 - 111 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle + \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 00101 \rangle}_{=0} + \underbrace{\langle 00110 \rangle}_{=0} - \underbrace{\langle 11101 \rangle}_{=0} - \underbrace{\langle 11110 \rangle}_{=0}) = 0 = \langle B_1 | B_2 \rangle$$

$$\langle B_3 | B_1 \rangle = \frac{1}{\sqrt{2}} (\langle 001 + 111 \rangle) \frac{1}{\sqrt{2}} (\langle 101 \rangle + \langle 110 \rangle) = \frac{1}{2} (\underbrace{\langle 00101 \rangle}_{=0} + \underbrace{\langle 00110 \rangle}_{=0} + \underbrace{\langle 11101 \rangle}_{=0} + \underbrace{\langle 11110 \rangle}_{=0}) = 0 = \langle B_1 | B_3 \rangle$$

$|B_2\rangle$:

- $\langle B_2 | B_2 \rangle = \frac{1}{\sqrt{2}} (\langle 001 - \langle 111 \rangle) \frac{1}{\sqrt{2}} (\langle 100 \rangle - \langle 111 \rangle) = \frac{1}{2} (\underbrace{\langle 00100 \rangle}_{=1} - \underbrace{\langle 00111 \rangle}_{=0} - \underbrace{\langle 11100 \rangle}_{=0} + \underbrace{\langle 11111 \rangle}_{=1}) = \frac{1}{2} (1+1) = 1$
- $\langle B_3 | B_2 \rangle = \frac{1}{\sqrt{2}} (\langle 001 + \langle 111 \rangle) \frac{1}{\sqrt{2}} (\langle 100 \rangle - \langle 111 \rangle) = \frac{1}{2} (\underbrace{\langle 00100 \rangle}_{=1} - \underbrace{\langle 00111 \rangle}_{=0} + \underbrace{\langle 11100 \rangle}_{=0} - \underbrace{\langle 11111 \rangle}_{=1}) = \frac{1}{2} (1-1) = 0 = \langle B_2 | B_3 \rangle$

$|B_3\rangle$:

- $\langle B_3 | B_3 \rangle = \frac{1}{\sqrt{2}} (\langle 001 + \langle 111 \rangle) \frac{1}{\sqrt{2}} (\langle 100 \rangle + \langle 111 \rangle) = \frac{1}{2} (\underbrace{\langle 00100 \rangle}_{=1} + \underbrace{\langle 00111 \rangle}_{=0} + \underbrace{\langle 11100 \rangle}_{=0} + \underbrace{\langle 11111 \rangle}_{=1}) = \frac{1}{2} (1+1) = 1$

Exercise 25:

- $|B_1\rangle = \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle)$
- $\hat{Z}_A |B_1\rangle = (\hat{Z}_A \otimes \mathbb{1}_B) \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle) = \frac{1}{\sqrt{2}} (\hat{Z}_A |0\rangle_A \otimes \mathbb{1}_B |1\rangle_B + \hat{Z}_A |1\rangle_A \otimes \mathbb{1}_B |0\rangle_B) = \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle) = |B_0\rangle$
 - $(\hat{X}_A \hat{Z}_A) |B_1\rangle = (\hat{X}_A \hat{Z}_A \otimes \mathbb{1}_B) \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle) = \frac{1}{\sqrt{2}} (\hat{X}_A \hat{Z}_A |0\rangle_A \otimes \mathbb{1}_B |1\rangle_B + \hat{X}_A \hat{Z}_A |1\rangle_A \otimes \mathbb{1}_B |0\rangle_B) = \frac{1}{\sqrt{2}} (\hat{X}_A |0\rangle_A \otimes \mathbb{1}_B - \hat{X}_A |1\rangle_A \otimes \mathbb{1}_B) = \frac{1}{\sqrt{2}} (|11\rangle - |00\rangle)$
 - $(\hat{X}_A \hat{Z}_A) |B_1\rangle = e^{i\pi} |B_2\rangle \equiv |B_2\rangle$ (global phase is irrelevant).
 - $\hat{X}_A |B_1\rangle = (\hat{X}_A \otimes \mathbb{1}_B) \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle) = \frac{1}{\sqrt{2}} (\hat{X}_A |0\rangle_A \otimes \mathbb{1}_B |1\rangle_B + \hat{X}_A |1\rangle_A \otimes \mathbb{1}_B |0\rangle_B) = \frac{1}{\sqrt{2}} (|11\rangle + |00\rangle) = |B_3\rangle$

Exercise 26:

$$\hat{U} = |B_0\rangle \langle 111 + |B_1\rangle \langle 011 + |B_2\rangle \langle 101 + |B_3\rangle \langle 001$$

$$\hat{U} = \begin{bmatrix} \langle 001 \hat{U} | 000 \rangle & \langle 001 \hat{U} | 01 \rangle & \langle 001 \hat{U} | 10 \rangle & \langle 001 \hat{U} | 11 \rangle \\ \langle 011 \hat{U} | 00 \rangle & \langle 011 \hat{U} | 01 \rangle & \langle 011 \hat{U} | 10 \rangle & \langle 011 \hat{U} | 11 \rangle \\ \langle 101 \hat{U} | 00 \rangle & \langle 101 \hat{U} | 01 \rangle & \langle 101 \hat{U} | 10 \rangle & \langle 101 \hat{U} | 11 \rangle \\ \langle 111 \hat{U} | 00 \rangle & \langle 111 \hat{U} | 01 \rangle & \langle 111 \hat{U} | 10 \rangle & \langle 111 \hat{U} | 11 \rangle \end{bmatrix}$$

$$\hat{U} = \begin{bmatrix} \langle 001 B_3 \rangle & \langle 001 B_1 \rangle & \langle 001 B_2 \rangle & \langle 001 B_0 \rangle \\ \langle 011 B_3 \rangle & \langle 011 B_1 \rangle & \langle 011 B_2 \rangle & \langle 011 B_0 \rangle \\ \langle 101 B_3 \rangle & \langle 101 B_1 \rangle & \langle 101 B_2 \rangle & \langle 101 B_0 \rangle \\ \langle 111 B_3 \rangle & \langle 111 B_1 \rangle & \langle 111 B_2 \rangle & \langle 111 B_0 \rangle \end{bmatrix}$$

$$\hat{U} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 0 & 1 & 0 & -1 \\ 1 & 0 & -1 & 0 \end{bmatrix}$$

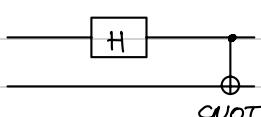
otherwise, if we consider a change of basis from $|000\rangle, |011\rangle, |101\rangle, |111\rangle$ to $|B_0\rangle, |B_1\rangle, |B_2\rangle, |B_3\rangle$ (with respect to this order):

$$\hat{U} = |B_0\rangle \langle 001 + |B_1\rangle \langle 011 + |B_2\rangle \langle 101 + |B_3\rangle \langle 111$$

$$\hat{U} = \begin{bmatrix} \langle 001 B_0 \rangle & \langle 001 B_1 \rangle & \langle 001 B_2 \rangle & \langle 001 B_3 \rangle \\ \langle 011 B_0 \rangle & \langle 011 B_1 \rangle & \langle 011 B_2 \rangle & \langle 011 B_3 \rangle \\ \langle 101 B_0 \rangle & \langle 101 B_1 \rangle & \langle 101 B_2 \rangle & \langle 101 B_3 \rangle \\ \langle 111 B_0 \rangle & \langle 111 B_1 \rangle & \langle 111 B_2 \rangle & \langle 111 B_3 \rangle \end{bmatrix}$$

$$\hat{U} = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 & 0 & 1 & 1 \\ 1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & -1 & 1 \end{bmatrix}$$

This is



Exercise 27:

$$\begin{aligned} \langle B_0 | \hat{x}_A | B_0 \rangle &= \frac{1}{\sqrt{2}} (\langle 011 \rangle - \langle 101 \rangle) (\hat{x}_A \otimes \mathbb{1}_B) \frac{1}{\sqrt{2}} (\langle 101 \rangle - \langle 110 \rangle) \\ &= \frac{1}{2} (\langle 011 \hat{x}_A \otimes \mathbb{1}_B | 101 \rangle - \langle 011 \hat{x}_A \otimes \mathbb{1}_B | 110 \rangle - \langle 101 \hat{x}_A \otimes \mathbb{1}_B | 101 \rangle + \langle 101 \hat{x}_A \otimes \mathbb{1}_B | 110 \rangle) \\ &= \frac{1}{2} (\underbrace{\langle 01111 \rangle}_{=0} - \underbrace{\langle 01100 \rangle}_{=0} - \underbrace{\langle 10111 \rangle}_{=0} + \underbrace{\langle 10100 \rangle}_{=0}) \end{aligned}$$

$$\langle B_0 | \hat{x}_A | B_0 \rangle = 0$$

$$\begin{aligned} \langle B_0 | \hat{y}_A | B_0 \rangle &= \langle B_0 | \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} | B_0 \rangle \\ &= \frac{1}{2} [0 \ 1 \ -1 \ 0] \begin{bmatrix} 0 & 0 & -i & 0 \\ 0 & 0 & 0 & -i \\ i & 0 & 0 & 0 \\ 0 & i & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ -1 \\ 0 \end{bmatrix} \\ &= \frac{i}{2} [0 \ 1 \ -1 \ 0] \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix} \end{aligned}$$

$$\langle B_0 | \hat{y}_A | B_0 \rangle = 0$$

$$\begin{aligned} \langle B_0 | \hat{z}_A | B_0 \rangle &= \langle B_0 | \hat{z}_A \otimes \mathbb{1}_B | B_0 \rangle = \frac{1}{2} (\langle 011 \rangle - \langle 101 \rangle) (\hat{z}_A \otimes \mathbb{1}_B) (\langle 101 \rangle - \langle 110 \rangle) \\ &= \frac{1}{2} (\langle 011 \hat{z}_A \otimes \mathbb{1}_B | 101 \rangle - \langle 011 \hat{z}_A \otimes \mathbb{1}_B | 110 \rangle - \langle 101 \hat{z}_A \otimes \mathbb{1}_B | 101 \rangle + \langle 101 \hat{z}_A \otimes \mathbb{1}_B | 110 \rangle) \\ &= \frac{1}{2} (\underbrace{\langle 01101 \rangle}_{=1} + \underbrace{\langle 01110 \rangle}_{=0} - \underbrace{\langle 10101 \rangle}_{=0} - \underbrace{\langle 10110 \rangle}_{=1}) \\ &= \frac{1}{2} (1 + 0 - 0 - 1) \end{aligned}$$

$$\langle B_0 | \hat{z}_A | B_0 \rangle = 0$$

This would be a point located at the origin $(0,0,0)$ of the "Bloch Ball", representing the completely mixed state (most entangled here). Since all spin projections are 0, this state corresponds to maximum uncertainty. It is no longer a valid pure state on the Bloch sphere, but instead a mixed state at the center of the ball.

Exercise 28:

$$\begin{aligned} \langle B_0 | \hat{x}_B | B_0 \rangle &= \frac{1}{2} (\langle 011 \rangle - \langle 101 \rangle) (\mathbb{1}_A \otimes \hat{x}_B) (\langle 101 \rangle - \langle 110 \rangle) \\ &= \frac{1}{2} (\langle 011 \mathbb{1}_A \otimes \hat{x}_B | 101 \rangle - \langle 011 \mathbb{1}_A \otimes \hat{x}_B | 110 \rangle - \langle 101 \mathbb{1}_A \otimes \hat{x}_B | 101 \rangle + \langle 101 \mathbb{1}_A \otimes \hat{x}_B | 110 \rangle) \\ &= \frac{1}{2} (\underbrace{\langle 01100 \rangle}_{=0} - \underbrace{\langle 01111 \rangle}_{=0} - \underbrace{\langle 10100 \rangle}_{=0} + \underbrace{\langle 10111 \rangle}_{=0}) \end{aligned}$$

$$\langle B_0 | \hat{x}_B | B_0 \rangle = 0$$

$$\langle B_0 | \vec{Y}_B | B_0 \rangle = \langle B_0 | \vec{y}_A \otimes \vec{y}_A | B_0 \rangle = \frac{1}{2} [0 \ 1 -1 \ 0] \begin{bmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ -1 \\ 0 \end{bmatrix}$$

$$= \frac{i}{2} [0 \ 1 -1 \ 0] \begin{bmatrix} -1 \\ 0 \\ 0 \\ -1 \end{bmatrix}$$

$$\langle B_0 | \vec{Y}_B | B_0 \rangle = 0$$

$$\langle B_0 | \vec{z}_B | B_0 \rangle = \frac{1}{2} (\langle 011 - \langle 101 \rangle) (\vec{z}_A \otimes \vec{z}_B) (|01\rangle - |10\rangle)$$

$$= \frac{1}{2} (\langle 011 \vec{z}_A \otimes \vec{z}_B |01\rangle - \langle 011 \vec{z}_A \otimes \vec{z}_B |10\rangle - \langle 101 \vec{z}_A \otimes \vec{z}_B |01\rangle + \langle 101 \vec{z}_A \otimes \vec{z}_B |10\rangle)$$

$$= \frac{1}{2} \left(\underbrace{-\langle 01101 \rangle}_{=1} - \underbrace{\langle 01110 \rangle}_{=0} + \underbrace{\langle 10101 \rangle}_{=0} + \underbrace{\langle 10110 \rangle}_{=1} \right)$$

$$= \frac{1}{2} (-1 + 1)$$

$$\langle B_0 | \vec{z}_B | B_0 \rangle = 0$$

Exercise 29:

$$\langle B_0 | \vec{x}_A \otimes \vec{x}_B | B_0 \rangle = \frac{1}{2} (\langle 011 - \langle 101 \rangle) (\vec{x}_A \otimes \vec{x}_B) (|01\rangle - |10\rangle)$$

$$= \frac{1}{2} (\langle 011 \vec{x}_A \otimes \vec{x}_B |01\rangle - \langle 011 \vec{x}_A \otimes \vec{x}_B |10\rangle - \langle 101 \vec{x}_A \otimes \vec{x}_B |01\rangle + \langle 101 \vec{x}_A \otimes \vec{x}_B |10\rangle)$$

$$= \frac{1}{2} \left(\underbrace{\langle 01110 \rangle}_{=0} - \underbrace{\langle 01101 \rangle}_{=1} - \underbrace{\langle 10110 \rangle}_{=1} + \underbrace{\langle 10101 \rangle}_{=0} \right)$$

$$= \frac{1}{2} (-1 - 1)$$

$$\langle B_0 | \vec{x}_A \otimes \vec{x}_B | B_0 \rangle = -1$$

$$\langle B_0 | \vec{y}_A \otimes \vec{y}_A | B_0 \rangle = \frac{1}{2} [0 \ 1 -1 \ 0] \begin{bmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ -1 \\ 0 \end{bmatrix}$$

$$= \frac{1}{2} [0 \ 1 -1 \ 0] \begin{bmatrix} 0 \\ -1 \\ 1 \\ 0 \end{bmatrix}$$

$$= \frac{1}{2} (-1 - 1)$$

$$\langle B_0 | \vec{y}_A \otimes \vec{y}_A | B_0 \rangle = -1$$

$$\langle B_0 | \vec{z}_A \otimes \vec{z}_B | B_0 \rangle = \frac{1}{2} (\langle 011 - \langle 101 \rangle) (\vec{z}_A \otimes \vec{z}_B) (|01\rangle - |10\rangle)$$

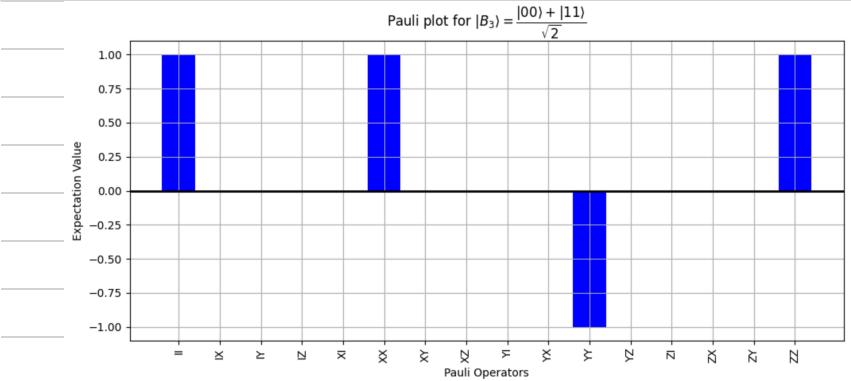
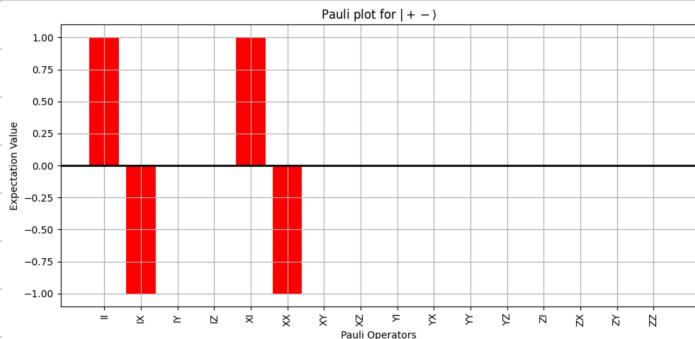
$$= \frac{1}{2} (\langle 011 \vec{z}_A \otimes \vec{z}_B |01\rangle - \langle 011 \vec{z}_A \otimes \vec{z}_B |10\rangle - \langle 101 \vec{z}_A \otimes \vec{z}_B |01\rangle + \langle 101 \vec{z}_A \otimes \vec{z}_B |10\rangle)$$

$$= \frac{1}{2} (-\underbrace{\langle 01|01 \rangle}_{=1} + \underbrace{\langle 01|10 \rangle}_{=0} + \underbrace{\langle 10|01 \rangle}_{=0} - \underbrace{\langle 10|10 \rangle}_{=1})$$

$$= \frac{1}{2} (-1 - 1)$$

$$\langle B_0 | \hat{z}_A \otimes \hat{z}_B | B_0 \rangle = -1$$

Exercise 30:



Axis not displayed represent a mean of 0.

Exercise 31:

(i) Alice measures +1 in her lab:

If A measures $\hat{z}_A = 1$ and B does nothing, the state after measurement is:

$$|14\rangle = \frac{(10)\langle 01_A \otimes 11_B)}{\sqrt{\langle B_3 | (10)\langle 01_A \otimes 11_B | B_3 \rangle}} |B_3\rangle$$

normalization

$$\text{Prob}_A : \langle B_3 | (10)\langle 01_A \otimes 11_B | B_3 \rangle = \frac{1}{2} (\langle 001 + 111 \rangle) (10)\langle 01_A \otimes 11_B | (100 + 111) = \frac{1}{2} (\langle 001 + 111 \rangle) (\underbrace{\langle 10 \times 010_A \otimes 11_B | 10 \rangle + \langle 10 \times 011_A \otimes 11_B | 11 \rangle}_{=0})$$

$$= \frac{1}{2} (\langle 001 + 111 \rangle) (100).$$

$$\text{Prob}(z_A=1) = \langle B_3 | (10)\langle 01_A \otimes 11_B | B_3 \rangle = \frac{1}{2}$$

$$\text{normalization: } \sqrt{\langle B_3 | (10)\langle 01_A \otimes 11_B | B_3 \rangle} = \frac{1}{\sqrt{2}}$$

$$\text{state after measurement in Alice's lab: } |14\rangle = \frac{(10)\langle 01_A \otimes 11_B)}{\sqrt{1/2}} |B_3\rangle$$

$$|14\rangle = \frac{1/\sqrt{2}}{1/\sqrt{2}} (10)\langle 01_A \otimes 11_B | (100 + 111)$$

$$|14\rangle = 10\rangle \underbrace{\langle 010}_{=1} \otimes 11_B 10\rangle_B + 10\rangle \underbrace{\langle 011}_{=0} \otimes 11_B 11\rangle_B$$

$$|14\rangle = 10\rangle_A \otimes 10\rangle_B$$

Now Bob measure in his lab without knowing the state of his qubit. We need to use state $|14\rangle$ since Alice's measurement has projected initial state. Since Bob measures \hat{z}_B , the two possible

outcomes are $|0\rangle_B$ or $|1\rangle_B$ for his partite. We again use projector:

$$\text{Bob measures } -1: |1\rangle \otimes (\mathbb{1}_A \otimes |1\rangle \langle 1|_B) |1\rangle = (\mathbb{1}_A \otimes |1\rangle \langle 1|) |100\rangle_{AB} = \cancel{\mathbb{1}_A} |10\rangle \otimes |1\rangle \langle 1|_B = 0$$

$$\text{Also } \text{Prob}(|z_B = -1| z_A = +1) = \langle 4|(\mathbb{1}_A \otimes |1\rangle \langle 1|_B) |1\rangle = \langle 00|(\mathbb{1}_A |1\rangle \otimes |1\rangle \langle 1|_B) = 0$$

We see that the probability that Bob measures -1 in his lab is 0 if Alice measured $+1$.

$$\text{• Bob measure } +1: |1\rangle \otimes (\mathbb{1}_A \otimes |0\rangle \langle 0|_B) |1\rangle = (\mathbb{1}_A \otimes |0\rangle \langle 0|) |100\rangle = \mathbb{1}_A |10\rangle \otimes |0\rangle \cancel{\langle 0|} = 1$$

$$\propto |10\rangle \otimes |0\rangle \text{ "already normalized"}$$

$$|1\rangle = |10\rangle \otimes |0\rangle = |100\rangle$$

$$\text{Prob}(z_B = +1 | z_A = +1) = \langle 4|(\mathbb{1}_A \otimes |0\rangle \langle 0|) |1\rangle = \langle 00|00\rangle = 1.$$

\Rightarrow normalization = 1

State after measurement in Bob's lab is still $|1\rangle = |100\rangle$.

We see that state is unchanged and Bob measures $+1$ with probability 1 if Alice measured $+1$ previously \blacksquare not asked?

(ii) Alice measures -1 in her lab:

If A measures $\frac{z}{A} = -1$ and B does nothing, the state after measurement is:

$$|1\rangle = \frac{(\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B)}{\sqrt{|B_3|} (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |B_3\rangle} |B_3\rangle$$

normalization

$$\begin{aligned} \text{Prob}_A : & \langle B_3 | (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |B_3\rangle = \frac{1}{2} (\langle 001 \rangle + \langle 111 \rangle) (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) (\langle 100 \rangle + \langle 111 \rangle) = \frac{1}{2} (\langle 001 \rangle + \langle 111 \rangle) (\cancel{\mathbb{1}\rangle \langle 1|_A} \cancel{\otimes \mathbb{1}_B}) = \cancel{\frac{1}{2}} \cancel{\langle 001 \rangle} \cancel{\langle 111 \rangle} = 1 \\ & = \frac{1}{2} (\langle 001 \rangle + \langle 111 \rangle) (\langle 111 \rangle). \end{aligned}$$

$$\text{Prob}(z_A = -1) = \langle B_3 | (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |B_3\rangle = \frac{1}{2}$$

$$\text{normalization: } \sqrt{\langle B_3 | (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |B_3\rangle} = \frac{1}{\sqrt{2}}$$

$$\text{state after measurement in Alice's lab: } |1\rangle = \frac{(\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |B_3\rangle}{1/\sqrt{2}}$$

$$|1\rangle = \frac{1/\sqrt{2}}{1/\sqrt{2}} (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) (\langle 100 \rangle + \langle 111 \rangle)$$

$$|1\rangle = \cancel{\frac{1}{\sqrt{2}}} (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |10\rangle_B + \cancel{\frac{1}{\sqrt{2}}} (\mathbb{1}\rangle \langle 1|_A \otimes \mathbb{1}_B) |11\rangle_B$$

$$|1\rangle = |1\rangle_A \otimes |1\rangle_B$$

Now Bob measures in his lab without knowing the state of his qubit. We need to use state $|1\rangle$ since Alice's measurement has projected initial state. Since Bob measures \hat{Z}_B , the two possible outcomes are $|0\rangle_B$ or $|1\rangle_B$ for his particle. We again use projector:

- Bob measures +1: $|1\rangle \propto (\mathbb{1}_A \otimes |0\rangle \langle 0|_B) |1\rangle = (\mathbb{1}_A \otimes |0\rangle \langle 0|) |11\rangle_{AB} = \frac{1}{\sqrt{2}} |\mathbb{1}\rangle_A \otimes |0\rangle \langle 0|_B = \frac{1}{\sqrt{2}} = 0$

Also $\text{Prob}(\hat{Z}_B = +1 | Z_A = -1) = \langle +1 | (\mathbb{1}_A \otimes |0\rangle \langle 0|_B) |1\rangle = \langle +1 | (\mathbb{1}_A |1\rangle \otimes |0\rangle \langle 0|_B) = \frac{1}{\sqrt{2}} = 0$

We see that the probability that Bob measures +1 in his lab is 0 if Alice measured -1.

- Bob measure -1: $|1\rangle \propto (\mathbb{1}_A \otimes |1\rangle \langle 1|_B) |1\rangle = (\mathbb{1}_A \otimes |1\rangle \langle 1|) |11\rangle = \frac{1}{\sqrt{2}} |1\rangle \otimes |1\rangle \langle 1|_B = \frac{1}{\sqrt{2}} = 1$

$\propto |1\rangle \otimes |1\rangle_B$ "already normalized"

$$|1\rangle = \frac{1}{\sqrt{2}} |\mathbb{1}\rangle_A \otimes |1\rangle_B = |\mathbb{1}\rangle$$

$\text{Prob}(\hat{Z}_B = -1 | Z_A = -1) = \langle +1 | (\mathbb{1}_A \otimes |1\rangle \langle 1|) |1\rangle = \langle +1 | +1\rangle = 1.$

\Rightarrow normalization = 1

state after measurement in Bob's lab $|1\rangle = |\mathbb{1}\rangle$.

We see that state is unchanged and Bob measures +1 with probability 1 if Alice measured +1 previously.

Conclusion:

- If Alice measured $Z_A = +1$, Bob measures $Z_B = +1$ with probability 1 so always.
- If Alice measured $Z_A = -1$, Bob measures $Z_B = -1$ with probability 1 so always. \blacksquare

Exercise 32:

(i) Alice measures $\hat{X}_A = +1$ and then Bob measures $Z_B = +1$:

Alice measures +1 so we project state on $|+\rangle \otimes |+\rangle_A$ but Bob does nothing:

state after measurement: $|1\rangle = \frac{(|+\rangle \otimes \mathbb{1}_B)}{\sqrt{\langle B_3 | (|+\rangle \otimes |+\rangle_A) |B_3 \rangle}} |B_3\rangle$

normalization:

$$\text{Prob}(\hat{X}_A = 1) = \langle B_3 | (|+\rangle \otimes |+\rangle_A) |B_3\rangle = \frac{1}{2} (\langle 001 | + \langle 111 |) (|+\rangle \otimes |+\rangle_B) (|00\rangle + |11\rangle)$$

$$= \frac{1}{2} (\langle 001 | + \langle 111 |) (|+\rangle \otimes |0\rangle_A \otimes |0\rangle_B + |+\rangle \otimes |1\rangle_A \otimes |1\rangle_B)$$

$$= \frac{1}{2} (\langle 001 | + \langle 111 |) \left(\frac{1}{\sqrt{2}} |+\rangle_A \otimes |0\rangle_B + \frac{1}{\sqrt{2}} |+\rangle_A \otimes |1\rangle_B \right)$$

$$= \frac{1}{2\sqrt{2}} \left(\underbrace{\langle 01+ |}_{=1/\sqrt{2}} \underbrace{\langle 010 |}_B + \underbrace{\langle 01+ |}_{=0} \underbrace{\langle 011 |}_B + \underbrace{\langle 11+ |}_{=0} \underbrace{\langle 110 |}_B + \underbrace{\langle 11+ |}_{=1/\sqrt{2}} \underbrace{\langle 111 |}_B \right) = \frac{1}{2\sqrt{2}} = 0$$

$$= \frac{1}{2} \cdot \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} + \frac{1}{\sqrt{2}} \right)$$

$$\text{Prob}(\hat{x}_A = +1) = \frac{1}{2}$$

$$\text{normalization: } \sqrt{\langle B_3 | (|+\rangle \otimes |+\rangle_A \otimes |+\rangle_B) |B_3 \rangle} = \frac{1}{\sqrt{2}}$$

$$\text{state after Alice's measurement: } |14\rangle = \frac{|+\rangle \otimes |+\rangle_A \otimes |+\rangle_B}{\sqrt{2}} \frac{1}{\sqrt{2}} (|100\rangle + |111\rangle)$$

$$|14\rangle = \frac{1}{\sqrt{2}} \left(|+\rangle_A \otimes |0\rangle_B + |+\rangle_A \otimes |1\rangle_B \right) = |+\rangle_A \otimes \frac{|0\rangle_B + |1\rangle_B}{\sqrt{2}}$$

Now Bob measures in his lab without knowing the state of his qubit. We need to use state $|14\rangle$ since Alice's measurement has projected initial state. Since Bob measures \hat{z}_B , the two possible outcomes are $|0\rangle_B$ or $|1\rangle_B$ for his partite. We again use projector:

- Bob measures $\hat{z}_B = +1$:

This projects Bob's partite onto $|0\rangle$, we use projector $|0\rangle \otimes |0\rangle_B$:

$$|14'\rangle = \frac{|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B}{\sqrt{\langle 2|(|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B)|14\rangle}} |14\rangle$$

$$\begin{aligned} \text{Prob}(\hat{z}_B = +1 | x_A = +1) &= \langle 2|(|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B)|14\rangle = \left(\langle +|_A \otimes \frac{|0\rangle_B + |1\rangle_B}{\sqrt{2}} \right) (|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B) \left(|+\rangle_A \otimes \frac{|0\rangle + |1\rangle_B}{\sqrt{2}} \right) \\ &= \left(\langle +|_A \otimes \frac{|0\rangle_B + |1\rangle_B}{\sqrt{2}} \right) \left(|+\rangle_A \otimes \frac{|0\rangle \langle 0|0\rangle_B + |0\rangle \langle 0|1\rangle_B}{\sqrt{2}} \right) \\ &= \left(\langle +|_A \otimes \frac{|0\rangle_B + |1\rangle_B}{\sqrt{2}} \right) \left(|+\rangle_A \otimes \frac{1}{\sqrt{2}} |0\rangle_B \right) \\ &= \underbrace{\langle +|_A}_{=1} \cdot \frac{1}{2} \left(\underbrace{\langle 0|0\rangle_B}_{=1} + \underbrace{\langle 1|0\rangle_B}_{=0} \right) \end{aligned}$$

$$\text{Prob}(\hat{z}_B = +1 | x_A = +1) = \frac{1}{2}$$

$$\text{normalization: } \sqrt{\langle 2|(|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B)|14\rangle} = \frac{1}{\sqrt{2}}$$

$$\text{state after Bob's measurement: } |14'\rangle = \frac{|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B}{\sqrt{\langle 2|(|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B)|14\rangle}} |14\rangle = \frac{|+\rangle_A \otimes |0\rangle \otimes |0\rangle_B}{\sqrt{2}} \left(|+\rangle_A \otimes \frac{|0\rangle + |1\rangle_B}{\sqrt{2}} \right)$$

$$|14'\rangle = |+\rangle_A \otimes \left(|0\rangle \underbrace{\langle 0|0\rangle_B}_{=1} + |0\rangle \underbrace{\langle 0|1\rangle_B}_{=0} \right)$$

$$|14'\rangle = |+\rangle_A \otimes |0\rangle_B = \frac{1}{\sqrt{2}} (|100\rangle + |101\rangle)$$

\Rightarrow If A measures $\hat{x}_A = +1$ then if B measures on \hat{z}_A , he will record +1 with Prob $(\hat{z}_B = +1 | \hat{x}_A = +1) = \frac{1}{2}$ and final state after all measurements is $|14'\rangle = |+\rangle_A \otimes |0\rangle_B$

(iii) Alice measures $\hat{X}_A = +1$ and then Bob measures $\hat{Z}_B = -1$:

- Bob measures $\hat{Z}_B = -1$:

This projects Bob's partite onto $|1\rangle$, we use projector $|1\rangle\langle 1|_B$:

$$|14'\rangle = \frac{|1\rangle_A \otimes |1\rangle\langle 1|_B}{\sqrt{\langle 1|(|1\rangle_A \otimes |1\rangle\langle 1|_B)|14\rangle}} |14\rangle$$

$$\begin{aligned} \text{Prob}(\hat{Z}_B = -1 | \hat{X}_A = +1) &= \langle 24 | \left(\frac{|1\rangle_A \otimes |1\rangle\langle 1|_B}{\sqrt{\langle 1|(|1\rangle_A \otimes |1\rangle\langle 1|_B)|14\rangle}} \right) |14\rangle = \left(\langle +|_A \otimes \frac{\langle 0|_B + \langle 1|_B}{\sqrt{2}} \right) \left(\frac{|1\rangle_A \otimes |1\rangle\langle 1|_B}{\sqrt{\langle 1|(|1\rangle_A \otimes |1\rangle\langle 1|_B)|14\rangle}} \right) \left(|+\rangle_A \otimes \frac{|0\rangle + |1\rangle_B}{\sqrt{2}} \right) \\ &= \left(\langle +|_A \otimes \frac{\langle 0|_B + \langle 1|_B}{\sqrt{2}} \right) \left(|+\rangle_A \otimes \frac{|1\rangle\langle 1|_B + |1\rangle\langle 1|_B}{\sqrt{2}} \right) \\ &= \left(\langle +|_A \otimes \frac{\langle 0|_B + \langle 1|_B}{\sqrt{2}} \right) \left(|+\rangle_A \otimes \frac{1}{\sqrt{2}} |1\rangle_B \right) \\ &= \underbrace{\langle +|_A}_{=1} \cdot \frac{1}{2} \left(\underbrace{\langle 0|_B}_{=0} + \underbrace{\langle 1|_B}_{=1} \right) \end{aligned}$$

$$\text{Prob}(\hat{Z}_B = -1 | \hat{X}_A = +1) = \frac{1}{2}$$

$$\text{normalization: } \sqrt{\langle 1|(|1\rangle_A \otimes |1\rangle\langle 1|_B)|14\rangle} = \frac{1}{\sqrt{2}}$$

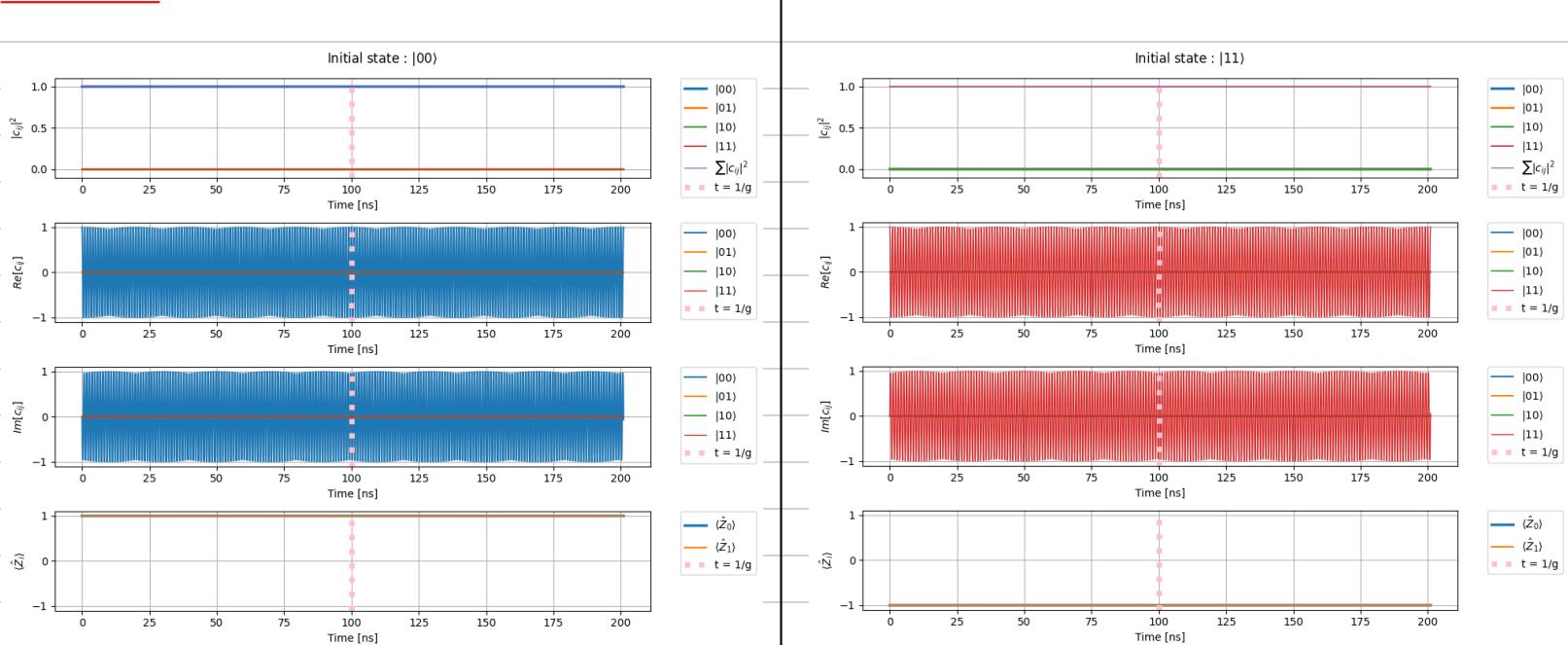
$$\text{state after Bob's measurement: } |14'\rangle = \frac{|1\rangle_A \otimes |1\rangle\langle 1|_B}{\sqrt{\langle 1|(|1\rangle_A \otimes |1\rangle\langle 1|_B)|14\rangle}} |14\rangle = \frac{|1\rangle_A \otimes |1\rangle\langle 1|_B}{1/\sqrt{2}} \left(|+\rangle_A \otimes \frac{|0\rangle + |1\rangle_B}{\sqrt{2}} \right)$$

$$|14'\rangle = |+\rangle_A \otimes \left(|1\rangle\langle 1|_B \underbrace{+ |1\rangle\langle 1|_B}_{=0} \right)$$

$$|14'\rangle = |+\rangle_A \otimes |1\rangle_B = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

\Rightarrow If A measures $\hat{X}_A = +1$ then if B measures on \hat{Z}_B , he will record -1 with $\text{Prob}(\hat{Z}_B = -1 | \hat{X}_A = +1) = \frac{1}{2}$ and final state after all measurements is $|14'\rangle = |+\rangle_A \otimes |1\rangle_B$

Exercise 33:



coeff c_{ij} : we see that only $|c_{00}|=1$ and all the others are 0, meaning that initial state $|100\rangle$ stays $|100\rangle$ during the whole time evolution.

Averages over observables: One can see that averages over z_1 (z-spin component of qubit 1) stays at +1 and average over z_2 (z-spin component of qubit 2) stays also at +1. This again confirm that initial state $|100\rangle$ stays at $|100\rangle$ since each qubit stays at $|10\rangle$ in its partite.

coeff c_{ij} : we see that only $|c_{11}|=1$ and all the others are 0, meaning that initial state $|11\rangle$ stays $|11\rangle$ during the whole time evolution.

Averages over observables: One can see that averages over z_1 (z-spin component of qubit 1) stays at -1 and average over z_2 (z-spin component of qubit 2) stays also at -1. This again confirm that initial state $|11\rangle$ stays at $|11\rangle$ since each qubit stays at $|11\rangle$ in its partite.

Conclusion: We can conclude that in both case, the initial state stays the same during time evolution

Why?:

$$\text{Hamiltonian: } H = \begin{bmatrix} -\frac{\omega_1}{2} - \frac{\omega_2}{2} & 0 & 0 & g \\ 0 & -\frac{\omega_1}{2} + \frac{\omega_2}{2} & g & 0 \\ 0 & g & \frac{\omega_1}{2} - \frac{\omega_2}{2} & 0 \\ g & 0 & 0 & \frac{\omega_1}{2} + \frac{\omega_2}{2} \end{bmatrix} \quad H = \begin{bmatrix} -2\pi & 0 & 0 & 0.02\pi \\ 0 & 0 & 0.02\pi & 0 \\ 0 & 0.02\pi & 0 & 0 \\ 0.02\pi & 0 & 0 & 2\pi \end{bmatrix}$$

eigenvalues: eigenvalues = $\left\{ -\frac{\sqrt{4g^2 + \omega_1^2 - 2\omega_1\omega_2 + \omega_2^2}}{2}, 1, \frac{\sqrt{4g^2 + \omega_1^2 - 2\omega_1\omega_2 + \omega_2^2}}{2}, -\frac{\sqrt{4g^2 + \omega_1^2 + 2\omega_1\omega_2 + \omega_2^2}}{2}, 1, \frac{\sqrt{4g^2 + \omega_1^2 + 2\omega_1\omega_2 + \omega_2^2}}{2}, 1 \right\}$
eigenvalues = $\{-6.28349945859136 : 1, -0.0628318530717959 : 1, 0.0628318530717959 : 1, 6.28349945859136 : 1\}$

$$(g \ll \omega_1 = \omega_2)$$

$$(g = \frac{2\pi}{100}; \omega_1 = 2\pi)$$

eigenvectors:

$$\text{eigenvectors} = \left[\left(\begin{array}{c} \frac{-\sqrt{4g^2 + \omega_1^2 - 2\omega_1\omega_2 + \omega_2^2}}{2}, 1, \begin{bmatrix} 0 \\ -\frac{\omega_1 + \omega_2}{2g} - \frac{\sqrt{4g^2 + \omega_1^2 - 2\omega_1\omega_2 + \omega_2^2}}{2g} \\ 1 \\ 0 \end{bmatrix} \end{array} \right), \left(\begin{array}{c} \frac{\sqrt{4g^2 + \omega_1^2 - 2\omega_1\omega_2 + \omega_2^2}}{2}, 1, \begin{bmatrix} 0 \\ -\frac{\omega_1 + \omega_2}{2g} + \frac{\sqrt{4g^2 + \omega_1^2 - 2\omega_1\omega_2 + \omega_2^2}}{2g} \\ 1 \\ 0 \end{bmatrix} \end{array} \right), \left(\begin{array}{c} -\frac{\sqrt{4g^2 + \omega_1^2 + 2\omega_1\omega_2 + \omega_2^2}}{2}, 1, \begin{bmatrix} \frac{-\omega_1 - \omega_2}{2g} - \frac{\sqrt{4g^2 + \omega_1^2 + 2\omega_1\omega_2 + \omega_2^2}}{2g} \\ 0 \\ 0 \\ 1 \end{bmatrix} \end{array} \right), \left(\begin{array}{c} \frac{\sqrt{4g^2 + \omega_1^2 + 2\omega_1\omega_2 + \omega_2^2}}{2}, 1, \begin{bmatrix} \frac{-\omega_1 - \omega_2}{2g} + \frac{\sqrt{4g^2 + \omega_1^2 + 2\omega_1\omega_2 + \omega_2^2}}{2g} \\ 0 \\ 0 \\ 1 \end{bmatrix} \end{array} \right) \right]$$

$$\text{eigenvectors} = \left[\left(-0.0628318530717959, 1, \begin{bmatrix} 0 \\ -0.707106781186548 \\ 0.707106781186548 \\ 0 \end{bmatrix} \right), \left(0.0628318530717959, 1, \begin{bmatrix} 0 \\ -0.707106781186548 \\ -0.707106781186548 \\ 0 \end{bmatrix} \right), \left(-6.28349945859136, 1, \begin{bmatrix} -0.999987500859308 \\ 0 \\ 0 \\ 0.00499981251210846 \end{bmatrix} \right), \left(6.28349945859136, 1, \begin{bmatrix} 0.00499981251210846 \\ 0 \\ 0 \\ 0.999987500859308 \end{bmatrix} \right) \right]$$

$$\hookrightarrow |\lambda_3\rangle \cong |100\rangle$$

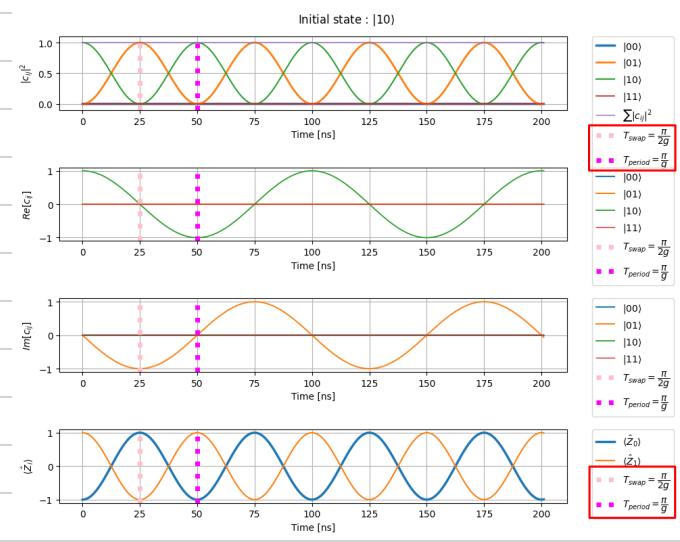
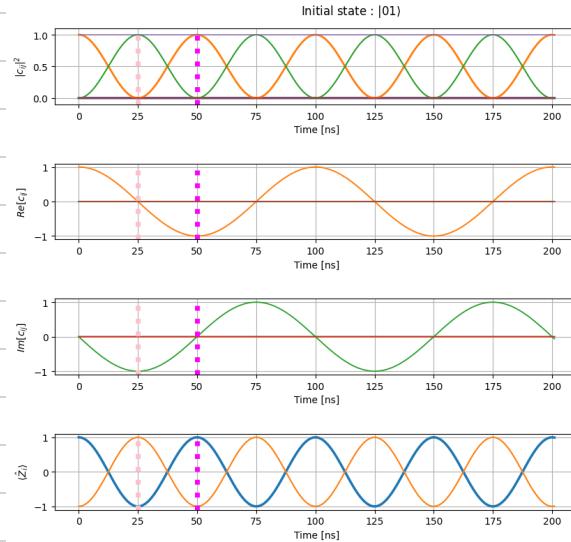
$$\hookrightarrow |\lambda_4\rangle \cong |11\rangle.$$

$$H|00\rangle = \begin{bmatrix} -2\pi \\ 0 \\ 0 \\ 0.02\pi \end{bmatrix} \cong -2\pi|100\rangle \quad H|11\rangle = \begin{bmatrix} 0.02\pi \\ 0 \\ 0 \\ 2\pi \end{bmatrix} \cong 2\pi|11\rangle.$$

We see that for this config. of $g = \frac{2\pi}{100}$ and $\omega_1 = \omega_2 = 2\pi$, $|100\rangle$ and $|11\rangle$ are eigenstates of this Hamiltonian. Therefore, they stay the same under time evolution with a global phase that is irrelevant. (Hamiltonian is time independant, we can solve Schrödinger equation such that $|\psi(t)\rangle = e^{-iH/\hbar t}|\psi_0\rangle = \sum_i e^{-iE_i\hbar t}|\lambda_i\rangle\langle\lambda_i|\psi_0\rangle$. where $|11\rangle$ and $|100\rangle$ are two of the 4 eigenstates).

$$H = \begin{bmatrix} -\frac{\omega_1}{2} - \frac{\omega_2}{2} & 0 & 0 & g \\ 0 & -\frac{\omega_1}{2} + \frac{\omega_2}{2} & g & 0 \\ 0 & g & \frac{\omega_1}{2} - \frac{\omega_2}{2} & 0 \\ g & 0 & 0 & \frac{\omega_1}{2} + \frac{\omega_2}{2} \end{bmatrix}$$

Exercise 34:



$$H = \begin{bmatrix} -2\pi & 0 & 0 & 0,02\pi \\ 0 & 0 & 0,02\pi & 0 \\ 0 & 0,02\pi & 0 & 0 \\ 0,02\pi & 0 & 0 & 2\pi \end{bmatrix} = -2\pi|00><00| + 0,02\pi|100><11| + 0,02\pi|101><10| + 0,02\pi|10><01| + 0,02\pi|11><00| + 2\pi|111><111|.$$

$g = 0,02\pi$

$$\frac{H}{\hbar}(-\frac{1}{\sqrt{2}}|101> + \frac{1}{\sqrt{2}}|110>) = -\frac{0,02\pi}{\sqrt{2}}|110> + \frac{0,02\pi}{\sqrt{2}}|101> = -0,02\pi\left(-\frac{1}{\sqrt{2}}|101> + \frac{1}{\sqrt{2}}|110>\right). \text{ eigenvector: } -\frac{1}{\sqrt{2}}|101> + \frac{1}{\sqrt{2}}|110>$$

$$\frac{H}{\hbar}\left(-\frac{1}{\sqrt{2}}|101> - \frac{1}{\sqrt{2}}|110>\right) = -\frac{0,02\pi}{\sqrt{2}}|110> - \frac{0,02\pi}{\sqrt{2}}|101> = 0,02\pi\left(-\frac{1}{\sqrt{2}}|101> + \frac{1}{\sqrt{2}}|110>\right). \text{ eigenvector: } -\frac{1}{\sqrt{2}}|101> - \frac{1}{\sqrt{2}}|110>$$

eigenvalue

$$U = e^{-i\frac{H}{\hbar}t} \cong e^{+i2\pi t}[-100>(-<001|) + e^{-i2\pi t}|111><111| + e^{+i0,02\pi t}\left(-\frac{1}{\sqrt{2}}|101> + \frac{1}{\sqrt{2}}|110>\right)\left(\frac{1}{\sqrt{2}}<011| + \frac{1}{\sqrt{2}}<101|\right) + e^{-i0,02\pi t}\left(-\frac{1}{\sqrt{2}}|101> - \frac{1}{\sqrt{2}}|110>\right)\left(\frac{1}{\sqrt{2}}<011| - \frac{1}{\sqrt{2}}<101|\right)]$$

Initial state |01>:

$$\bullet |14_{01}(t)> = U|01> = e^{+i0,02\pi t}\left(-\frac{1}{\sqrt{2}}|101> + \frac{1}{\sqrt{2}}|110>\right)\left(\frac{1}{\sqrt{2}}\right) + e^{-i0,02\pi t}\left(-\frac{1}{\sqrt{2}}|101> - \frac{1}{\sqrt{2}}|110>\right)\left(-\frac{1}{\sqrt{2}}\right)$$

$$= \frac{1}{2} \left(e^{+i0,02\pi t} + e^{-i0,02\pi t} \right) |01> - \frac{1}{2} \left(e^{+i0,02\pi t} - e^{-i0,02\pi t} \right) |110>$$

$$|14_{01}(t)> = \cos(0,02\pi t)|01> - i \sin(0,02\pi t)|110>$$

$\stackrel{=g}{=} \stackrel{=g}{=}$

$$\bullet P(|14_{01}(t)> = |101>) = |\cos|^2 = \cos^2(0,02\pi t) \quad T = \frac{\pi}{\omega} = \frac{\pi}{g} = \frac{\pi}{2\pi/100} = \frac{100}{2} = 50 \quad (\text{see graph } |101> \text{ above})$$

$$\bullet P(|14_{01}(t)> = |110>) = |\cos|^2 = \sin^2(0,02\pi t). \quad T = \frac{\pi}{\omega} = \frac{\pi}{g} = 50 \quad (\text{see graph } |110> \text{ above})$$

One can see that the proba of being in state |101> is 1 (cosine) every 50 ns and proba of being |110> is also 1 every 50 ns also (but sine). The period of alternating between the two states is therefore $T_{swap} = \frac{T}{2} = \frac{\pi}{2g} = 25$ ns. As we can see in graph |101> above, $|c_{101}|^2$ (Proba of being in state |101>) is 1 but $|c_{110}|^2$ is 0 and 25 ns later $|c_{101}|^2$ is 0 and $|c_{110}|^2$ is 1 (Proba of being in |110>).

$$T_{swap} = \frac{\pi}{2g} = 25 \text{ ns}$$

Initial state $|10\rangle$:

$$|\psi_{10}(t)\rangle = \mathcal{U}|10\rangle = e^{+i\frac{\omega}{2}\sin(\omega t)} \left(-\frac{1}{\sqrt{2}}|10\rangle + \frac{1}{\sqrt{2}}|10\rangle \right) \left(+\frac{1}{\sqrt{2}} \right) + e^{-i\frac{\omega}{2}\sin(\omega t)} \left(-\frac{1}{\sqrt{2}}|10\rangle - \frac{1}{\sqrt{2}}|10\rangle \right) \left(-\frac{1}{\sqrt{2}} \right).$$

$$= -\frac{1}{2} \left(e^{+i\frac{\omega}{2}\sin(\omega t)} - e^{-i\frac{\omega}{2}\sin(\omega t)} \right) |10\rangle + \frac{1}{2} \left(e^{+i\frac{\omega}{2}\sin(\omega t)} + e^{-i\frac{\omega}{2}\sin(\omega t)} \right) |10\rangle$$

$$|\psi_{10}(t)\rangle = -i\sin(\omega t)|10\rangle + \cos(\omega t)|10\rangle$$

$$\cdot P(|\psi_{10}(t)\rangle = |10\rangle) = |\cos|^2 = \sin^2(\omega t)$$

$$T = \frac{\pi}{\omega} = \frac{\pi}{g} = \frac{\pi}{2\pi/100} = \frac{100}{2} = 50 \text{ (see graph } |10\rangle \text{ above)}$$

$$\cdot P(|\psi_{10}(t)\rangle = |10\rangle) = |\cos|^2 = \cos^2(\omega t)$$

$$T = \frac{\pi}{\omega} = \frac{\pi}{g} = 50 \text{ (see graph } |10\rangle \text{ above)}$$

One can see that the proba of being in state $|10\rangle$ is 1 (sine) every 50 ns and proba of being $|10\rangle$ is also 1 every 50 ns also (but cosine). The period of alternating between the two states is therefore $T_{\text{swap}} = \frac{T}{2} = \frac{\pi}{2g} = 25 \text{ ns}$. As we can see in graph $|10\rangle$ above, $|\cos|^2$ (Proba of being in state $|10\rangle$) is 1 but $|\cos|^2$ is 0 and 25 ns later $|\cos|^2$ is 0 and $|\cos|^2$ is 1 (Proba of being in $|10\rangle$).

$$T_{\text{swap}} = \frac{\pi}{2g} = 25 \text{ ns}$$

Exercise 35:

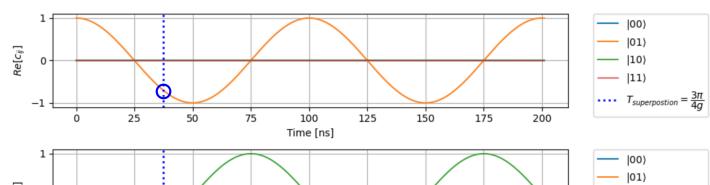
Taking what we used in previous question we can rewrite:

Initial state $|10\rangle$:

$$|\psi_{10}(t)\rangle = \mathcal{U}|10\rangle = \cos(gt)|10\rangle - i\sin(gt)|10\rangle \quad \tau = \frac{3}{8} \cdot \frac{2\pi}{g}$$

* (1) what we want but global phase so irrelevant

$$|\psi_{10}(t = \frac{3}{8} \cdot \frac{2\pi}{g})\rangle = \cos(\frac{3\pi}{4g}g)|10\rangle - i\sin(\frac{3\pi}{4g}g)|10\rangle$$



$$|\psi_{10}(\tau)\rangle = \cos(\frac{3\pi}{4})|10\rangle - i\sin(\frac{3\pi}{4})|10\rangle.$$

$$|\psi_{10}(\tau)\rangle = \frac{e^{i\frac{3\pi}{4}}}{\sqrt{2}} (|10\rangle + i|10\rangle)$$

(the global phase is irrelevant)

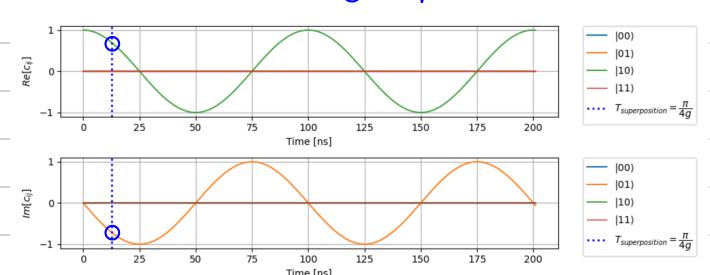
$$\Rightarrow |\psi_{10}(\tau)\rangle = \frac{1}{\sqrt{2}} (|10\rangle + i|10\rangle) \quad \text{for } \tau = \frac{3}{8} \cdot \frac{2\pi}{g} = \frac{3\pi}{4g} = \frac{3}{4} \text{ Tperiod.}$$

Initial state $|10\rangle$:

$$|\psi_{10}(t)\rangle = \mathcal{U}|10\rangle = -i\sin(gt)|10\rangle + \cos(gt)|10\rangle \quad \tau = \frac{1}{4} \cdot \frac{\pi}{g}$$

* (-i) what we want but global phase so irrelevant

$$|\psi_{10}(t = \frac{1}{8} \cdot \frac{2\pi}{g})\rangle = -i\sin(\frac{\pi}{4g}g)|10\rangle + \cos(\frac{\pi}{4g}g)|10\rangle$$



$$|\psi_{10}(\tau)\rangle = -i\sin(\frac{\pi}{4})|10\rangle + \cos(\frac{\pi}{4})|10\rangle.$$

$$|\psi_{10}(\tau)\rangle = -i\frac{1}{\sqrt{2}}|10\rangle + \frac{1}{\sqrt{2}}|10\rangle.$$

$$|1^A 10(\tau)\rangle = \frac{e^{i\pi}}{\sqrt{2}} e^{i\pi/2} (|10\rangle - i|10\rangle) = \frac{e^{i3\pi/2}}{\sqrt{2}} (|10\rangle - i|10\rangle) \quad (\text{the global phase is irrelevant})$$

$$\Leftrightarrow |1^A 10(\tau)\rangle = \frac{1}{\sqrt{2}} (|10\rangle - i|10\rangle) \quad \text{for } \tau = \frac{\pi}{4g} = \frac{1}{4} \text{ period}$$

Exercise 36:

$$S = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/2} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \quad S|0\rangle = |0\rangle \quad S|1\rangle = i|1\rangle$$

$$(S_A \otimes \mathbb{1}_B) (|10\rangle + i|10\rangle) = S_A|10\rangle \otimes \mathbb{1}_B|1\rangle_B + iS_A|1\rangle \otimes \mathbb{1}_B|0\rangle_B$$

$$\begin{aligned} &= |10\rangle + i(i)|10\rangle \\ &= |10\rangle - |10\rangle \end{aligned}$$

$$(S_A \otimes \mathbb{1}_B) (|10\rangle + i|10\rangle) = |10\rangle - |10\rangle = |B\rangle$$

We need to apply a phase gate to Qubit A.

Exercise 37:

$$\text{We recall that: } \frac{\hat{H}}{\hbar} = -\omega|00\rangle\langle 00| + g|00\rangle\langle 11| + g|10\rangle\langle 10| + g|11\rangle\langle 11| + \omega|11\rangle\langle 11|.$$

Considering that $|00\rangle$ and $|11\rangle$ don't evolve, we can truncate our Hamiltonian by considering the evolution of the two state $|10\rangle$ and $|10\rangle$, and we are reduced to a 2-level qubit system. Doing so, we map $|10\rangle \equiv |1\rangle \equiv |0\rangle$ (north pole of the Bloch sphere) and $|10\rangle \equiv |1\rangle \equiv |1\rangle$ (south pole of the Bloch sphere).

$\frac{\hat{H}}{\hbar} \stackrel{4 \times 4}{=} g(|10\rangle\langle 10| + |10\rangle\langle 01|) \Rightarrow$ Plays the role of an \hat{X} -gate flipping $|10\rangle$ to $|10\rangle$ and vice versa.

$$\Rightarrow \frac{\hat{H}}{\hbar} \stackrel{2 \times 2}{=} g(|10\rangle\langle 11| + |11\rangle\langle 01|) = (g) \hat{X} = (2g) \hat{X}/2, \quad \text{with } g = 0, \omega\pi \text{ and } \omega_1 = \omega_2 = \omega = \omega\pi$$

$$U_t = e^{-i\frac{(2g)}{2}\hat{X}t} = \cos(gt)\mathbb{I} - i\sin(gt)\hat{X} \quad \exp\left(i\frac{a}{2}\mathbf{n} \cdot \vec{\sigma}\right) = \left(\cos\frac{a}{2}\right)\mathbb{I} + i\left(\sin\frac{a}{2}\right)\mathbf{n} \cdot \vec{\sigma}$$

$$|1^A(t)\rangle_{01} = U_t |10\rangle = \cos(gt)|10\rangle - i\sin(gt)|11\rangle \quad \text{in order to be in state } |11\rangle \text{ we need } gT_{\text{swap}} = \frac{\pi}{2}$$

$$\Rightarrow T_{\text{swap}} = \frac{\pi}{2g}.$$

Under time evolution, this will make the qubit $|10\rangle \equiv |10\rangle$ evolve by doing a rotation around \hat{X} axis and flipping to $|11\rangle \equiv |10\rangle$ with a period of $T = \frac{\pi}{2g}$ as expected in our previous 4×4 Hamiltonian.

The mathematical way of seeing this is the following:

we consider the original Hamiltonian:

$$H = \begin{bmatrix} -\frac{\omega_1}{2} - \frac{\omega_2}{2} & 0 & 0 & g \\ 0 & -\frac{\omega_1}{2} + \frac{\omega_2}{2} & g & 0 \\ 0 & g & \frac{\omega_1}{2} - \frac{\omega_2}{2} & 0 \\ g & 0 & 0 & \frac{\omega_1}{2} + \frac{\omega_2}{2} \end{bmatrix} = \left(-\frac{\omega_1}{2} - \frac{\omega_2}{2} \right) |100\rangle\langle 100| + g |100\rangle\langle 111| + \left(-\frac{\omega_1}{2} + \frac{\omega_2}{2} \right) |101\rangle\langle 101| + g |101\rangle\langle 101| + \dots + g |110\rangle\langle 101| + \left(\frac{\omega_1}{2} - \frac{\omega_2}{2} \right) |110\rangle\langle 101| + g |110\rangle\langle 111| + \left(\frac{\omega_1}{2} + \frac{\omega_2}{2} \right) |111\rangle\langle 111|.$$

$$\frac{\hat{H}^{2x2}}{\hbar} = \begin{bmatrix} \langle 011 | \hat{H} | 101 \rangle & \langle 011 | \hat{H} | 110 \rangle \\ \langle 101 | \hat{H} | 101 \rangle & \langle 101 | \hat{H} | 110 \rangle \end{bmatrix}$$

$$\frac{\hat{H}^{2x2}}{\hbar} = \begin{bmatrix} -\frac{\omega_1}{2} + \frac{\omega_2}{2} & g \\ g & \frac{\omega_1}{2} - \frac{\omega_2}{2} \end{bmatrix} = \left(-\frac{\omega_1}{2} + \frac{\omega_2}{2} \right) \hat{z} + g \hat{x}$$

If we go back to $\omega_1 = \omega_2$ we find back $\frac{\hat{H}^{2x2}}{\hbar} = g \hat{x}$ with $g = 0, 0.2\pi$ as explained before. ■