BIPOLAR TRANSISTOR

Bibliography:

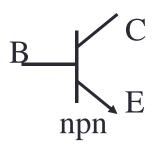
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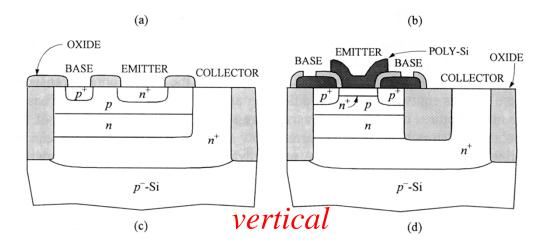
Plan

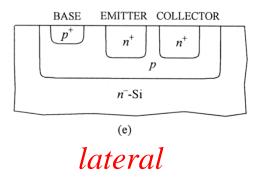
- Geometry
- Principle of operating
- Static characteristics
- Ebers-Moll Equations
- Static parameters (gain...)
- Second order effects
- Switching performances
- Transistor in HF domain
- Hetero-junction Bipolar Transistor (HBT or TBH)

Geometry

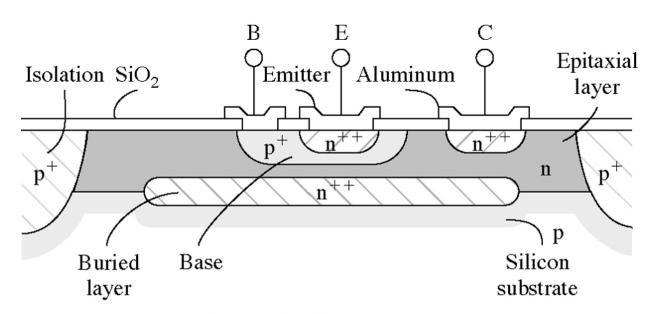
- Geometry:
 - Lateral
 - Vertical
- For digital circuits,
 vertical design







Geometry for IC

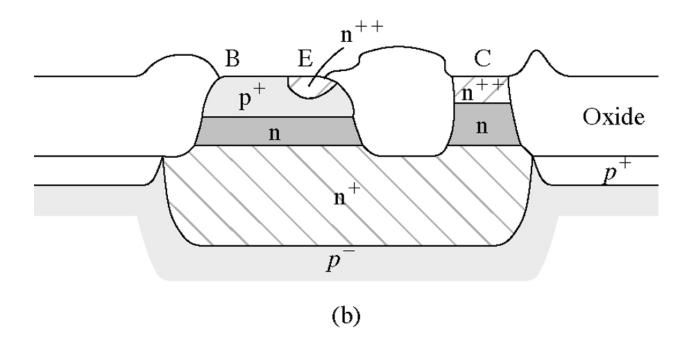


Conventional npn transistor

(a)

Muller et Kamins, « device electronics for IC »,2nd Ed., Wiley, 1986

Geometry with insulating oxide



Muller et Kamins, « device electronics for IC »,2nd Ed., Wiley, 1986

BJT always present: why?

- speed
- Low noise
- High gain
- Low output resistance

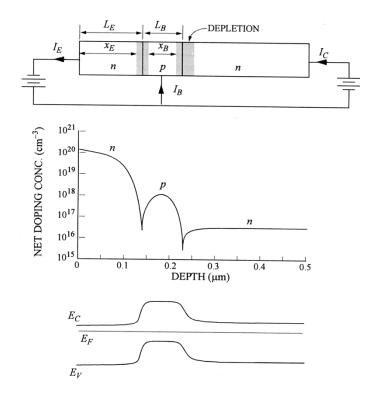


Analogue amplifier

- Still present in mobile phone (analog part)
- Low density, mainly in power stage
- BiCMOS

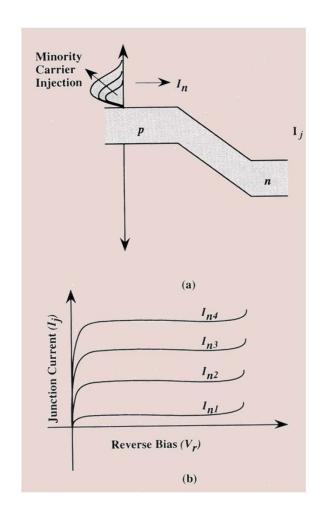
Working principle

- 2 PN Junctions with one common region (base)
 - The first Junction (EB) will inject carriers
 - The second one (CB) will collect carriers
- The base must be thin (lower than diffusion)length



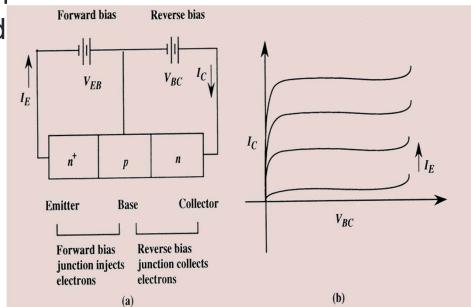
Working principle

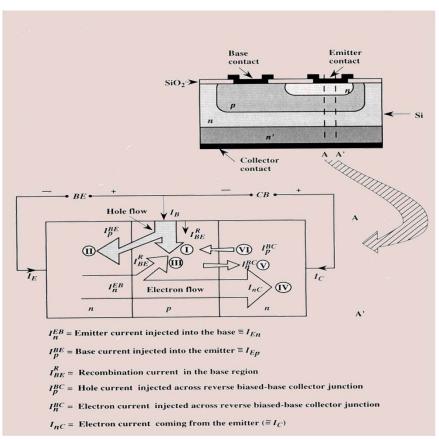
- Reverse biased Junction:
 - Low current due to empty « tank »
 - En modulant le remplissage du réservoir, modulation du courant inverse collecté (collecteur)
 - On remplit le réservoir (la base) en polarisant en direct la jonction EB

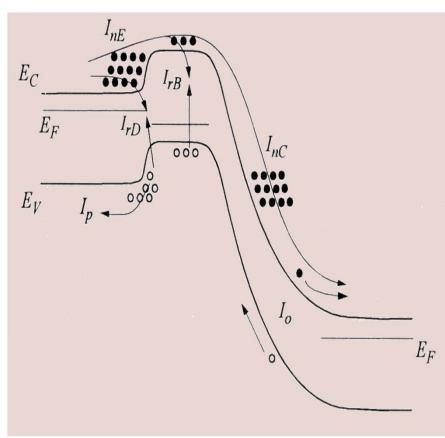


Working principle

- The reverve biasing CB junction creates a favorable electric field for the collect
- Conditions:
 - Thin base:
 - Avoid recombination effects
 - Ligthly doped base compared to emitter
 - favors one type of injected carriers (better injection efficiency)



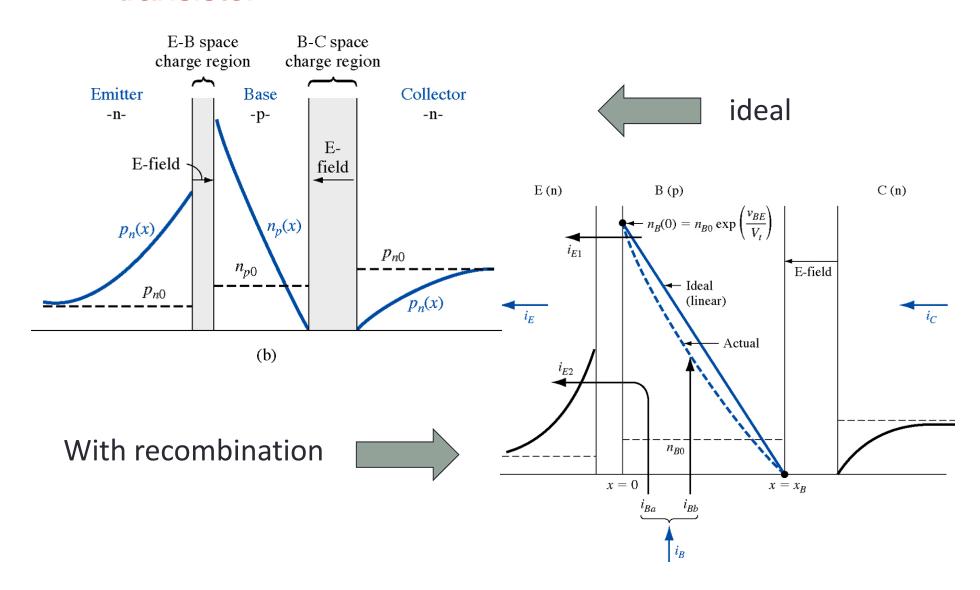




NPN Transistor

NPN Transistor

Minority carrier distribution in a forward biased npn transistor



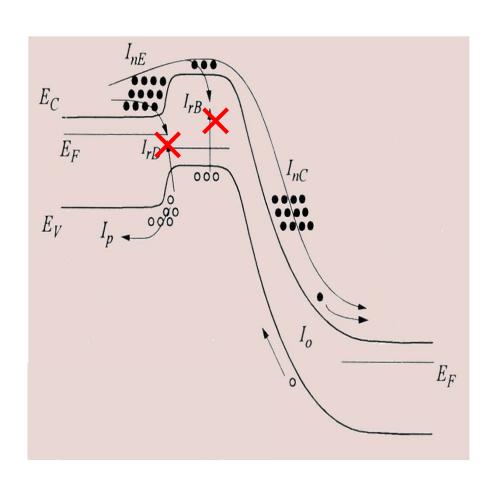
Statics characteristics + simplifying hypotheses

No recombination in the Base!(X)

1D Approximation

Homogeneous doping in the Base

Low Injection



PNP Transistor

- In the neutral Base region:
 - Continuity equation

$$p\frac{J_n}{eD_n} - n\frac{J_p}{eD_p} = \frac{d(p.n)}{dx}$$

• and $J_n >> J_p$, n << p



$$p\frac{J_n}{eD_n} = \frac{d(p.n)}{dx}$$

Integrating from E-B to C-B:

$$J_{n} = -eD_{nb}n_{iB}^{2} \frac{\left[\exp(\frac{eV_{BE}}{kT}) - \exp(\frac{eV_{BC}}{kT})\right]}{\int_{E}^{C} p(x)dx}$$

• finally:

$$J_{n} = -eD_{nb}n_{iB}^{2} \frac{\left[e^{\frac{eV_{BE}}{kT}} - 1 - (e^{\frac{eV_{BC}}{kT}} - 1) \right]}{\int_{E}^{C} p(x)dx}$$

In forward active regime,
 J_n is negative (e⁻ vers
 x<0)

$$J_{n} = -eD_{nb}n_{iB}^{2} \frac{\left[(e^{\frac{eV_{BE}}{kT}} - 1) - (e^{\frac{eV_{BC}}{kT}} - 1) \right]}{\int_{E'}^{C} p(x)dx} = -\left[\frac{eD_{nb}n_{iB}^{2}}{\int_{E'}^{C} p(x)dx} (e^{\frac{eV_{BE}}{kT}} - 1) - \frac{eD_{nb}n_{iB}^{2}}{\int_{E'}^{C} p(x)dx} (e^{\frac{eV_{BC}}{kT}} - 1) \right]$$

and:
$$\int_{E'}^{C'} p(x)dx = N_{A_B} \times W_{Beff}$$

So:
$$I_{n} = -A_{E} \left[\frac{e n_{iB}^{2} D_{nb}}{N_{A_{B}} W_{Beff}} (e^{\frac{e V_{BE}}{kT}} - 1) - \frac{e n_{iB}^{2} D_{nb}}{N_{A_{B}} W_{Beff}} (e^{\frac{e V_{BC}}{kT}} - 1) \right] = -I_{Sn} \left[(e^{\frac{e V_{BE}}{kT}} - 1) - (e^{\frac{e V_{BC}}{kT}} - 1) \right]$$

with :
$$I_{Sn} = A_E \frac{e n_{iB}^2 D_{nb}}{N_{A_B} W_{Beff}} = A_E \frac{e n_i^2}{G_B}$$
 Saturation current of electrons in the PN narrow junction (or without recombination in the Base).

Gummel number in the Base (s/cm⁴)

$$G_{B} = \frac{n_{i}^{2}}{n_{iB}^{2}} \frac{N_{A_{B}}}{D_{nb}} W_{Beff} = \frac{n_{i}^{2}}{n_{iB}^{2}} \frac{p}{D_{nb}} W_{Beff}$$

Emitter current

Collector current

$$J_{pE} = -J_{spE} \left(\exp(\frac{eV_{BE}}{kT}) - 1 \right)$$

$$J_{pC} = J_{spC} \left(\exp(\frac{eV_{BC}}{kT}) - 1 \right)$$

$$J_{pC} = J_{spC} \left(\exp(\frac{eV_{BC}}{kT}) - 1 \right)$$

finally

$$\begin{split} \boldsymbol{I}_E &= \boldsymbol{I}_{pE} + \boldsymbol{I}_n \\ \boldsymbol{I}_C &= + \boldsymbol{I}_{pC} - \boldsymbol{I}_n \\ \boldsymbol{I}_B &= - \boldsymbol{I}_C - \boldsymbol{I}_E = - \boldsymbol{I}_{pE} - \boldsymbol{I}_{pC} \end{split}$$

$$I_B$$
 I_E
 I_{pE}
 I_n
 I_{pC}
 I_C
 I_C
 I_C
 I_C

• And (!):

$$I_{E} = I_{n} + I_{pE} = -\left[\frac{Aen_{i}^{2}D_{nb}}{\int_{E^{'}}^{C}p(x)dx} + I_{spE}\right](\exp\frac{eV_{BE}}{kT} - 1) + \frac{Aen_{i}^{2}D_{nb}}{\int_{E^{'}}^{C}p(x)dx}(\exp\frac{eV_{BC}}{kT} - 1)$$

$$I_{sn}$$

$$I_{C} = -I_{n} - I_{pC} = +\frac{Aen_{i}^{2}D_{nb}}{\int_{E'}^{C}p(x)dx} (\exp\frac{eV_{BE}}{kT} - 1) - \left[\frac{Aen_{i}^{2}D_{nb}}{\int_{E'}^{C}p(x)dx} + I_{spC}\right] (\exp\frac{eV_{BC}}{kT} - 1)$$

$$I_{B} = -I_{E} - I_{C} = -I_{pE} - I_{pC} = +I_{spE} \left(\exp(\frac{eV_{BE}}{kT}) - 1 \right) - I_{spC} \left(\exp(\frac{eV_{BC}}{kT}) - 1 \right)$$

- Forward active region:
 - E-B (forward) and C-B (reverse)

$$I_E = -(\frac{Ae^2n_i^2D_{nB}}{Q_B + Q_S} + I_{SpE})\exp\frac{eV_{BE}}{kT} = -(I_{Sn} + I_{SpE})\exp\frac{eV_{BE}}{kT}$$

$$I_C = +\left(\frac{Ae^2n_i^2D_{nB}}{Q_B + Q_S}\right)\exp\frac{eV_{BE}}{kT} = I_{Sn}\exp\frac{eV_{BE}}{kT}$$

$$I_{B}^{*} = -I_{E} - I_{C} = -I_{SpE} \exp \frac{eV_{BE}}{kT}$$

Emitter injection efficiency:

$$\gamma_E = \frac{I_n}{I_{Ep}}$$

DC common base current gain:

$$\alpha = \frac{I_C}{I_E} = \frac{1}{1 + \frac{J_{Sp}.(Q_B + Q_S)}{e^2 n_i^2 D_{nB}}} = \frac{I_{Sn}}{I_{Sn} + I_{Sp}}$$

DC common emitter current gain :

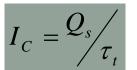
$$\beta = \frac{I_C}{I_B} = \frac{\alpha}{1 - \alpha}$$

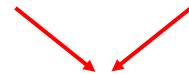
 $\beta = \frac{I_C}{I_R} = \frac{\alpha}{1 - \alpha}$ NB:if we neglect recombination process β and γ_E are equivalent

Base transport factor:

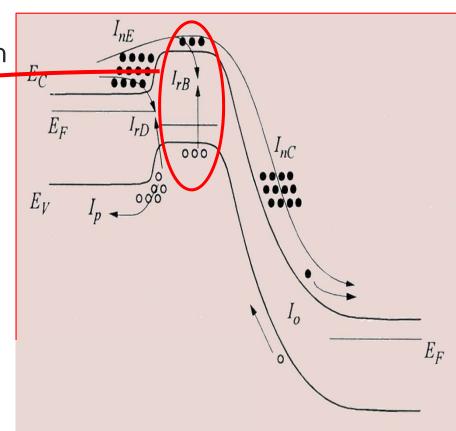
 Take into account recombination in neutral base region

$$I_{rB} = \frac{Q_s}{\tau_n} \approx \frac{AeX_{B_{eff}}(n_p(0) - n_p)/2}{\tau_n}$$





$$\delta = \frac{I_C}{I_{rB}} = \frac{\tau_n}{\tau_t} = \frac{2L_n^2}{X_{B_{eff}}^2} > 1$$



Recombination factor:

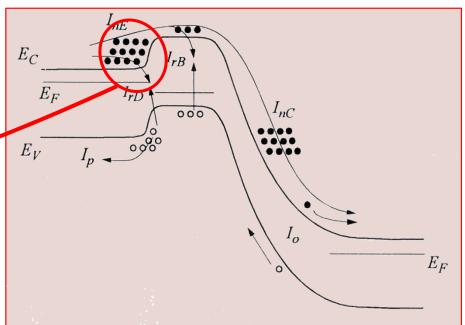
 Take into account recombination in depleted base region

$$I_{rD} = \frac{Aen_i}{2\tau} W_T \exp\left(\frac{eV_{BE}}{2kT}\right)$$

with W_T , width of E-B space charges.

When we add up all the contribution to the base current we get:

$$I_B = I_B^* + I_{rB} + I_{rD}$$



The common emitter gain can be expressed as:

$$\frac{1}{\beta} = \frac{I_B}{I_C} = \frac{I_B^* + I_{rB} + I_{rD}}{I_C} = \frac{1}{\gamma_E} + \frac{1}{\delta} + \frac{I_{rD}}{I_C}$$

- I_B^* intriseque base current (no recombination)
- I_{rB} recombination current in the neutral base region
- $I_{\rm rD}$ recombination current in the depletion zone of EB junction

Other modes of operation

- Saturation mode (regime):
 - The two junctions are forward biased.

$$I_{E} \cong -\left(\frac{Ae^{2}n_{i}^{2}D_{nB}}{Q_{S} + Q_{B}} + I_{spE}\right)e^{\frac{eV_{BE}}{kT}} + \frac{Ae^{2}n_{i}^{2}D_{nB}}{Q_{S} + Q_{B}}e^{\frac{eV_{BC}}{kT}}$$

$$Base$$

$$I_{C} \cong \frac{Ae^{2}n_{i}^{2}D_{nB}}{Q_{S} + Q_{B}}e^{\frac{eV_{BE}}{kT}} - \left(\frac{Ae^{2}n_{i}^{2}D_{nB}}{Q_{S} + Q_{B}} + I_{spC}\right)e^{\frac{eV_{BC}}{kT}}$$

$$Q_{SI}$$

$$Q_{SI}$$

$$Q_{SI}$$

$$Q_{SI}$$

Saturation mode

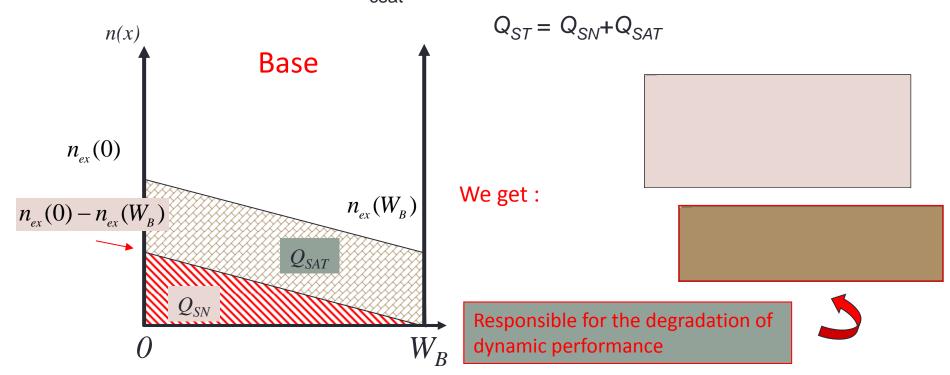
- Low injection : $(Q_S << Q_B)$:
 - Current is due to saturation charge injected into the base, ie $Q_{ST} = Q_{S1} + Q_{S2}$
 - If we deal with « narrow » junction, this charge is simply given by the surface of ½ trapeze (linear decay)

$$Q_{S1} = -\frac{1}{2}W_B en(x = 0) = -\frac{1}{2}W_B \frac{n_i^2}{N_A} e^{\frac{eV_{BE}}{kT}} = -\tau_t J_{sn} e^{\frac{eV_{BE}}{kT}}$$

$$Q_{S2} = -\frac{1}{2}W_B en(x = W_B) = -\frac{1}{2}W_B \frac{n_i^2}{N_A} e^{\frac{eV_{BC}}{kT}} = -\tau_t J_{sn} e^{\frac{eV_{BC}}{kT}}$$

Saturation mode

- Low injection: $(Q_S << Q_B)$:
 - An other « view » of saturation charge (from Ablard):
 - We consider transistor in active mode with a charge Q_{SN} and a charge Q_{SAT} (we have to determine) which supply the same saturation current I_{csat} .



Saturation mode

- High injection level
 - In this case, injected electron density reaches holes density into the base ($n \approx p$
 - Equivalent study leads to the following result :

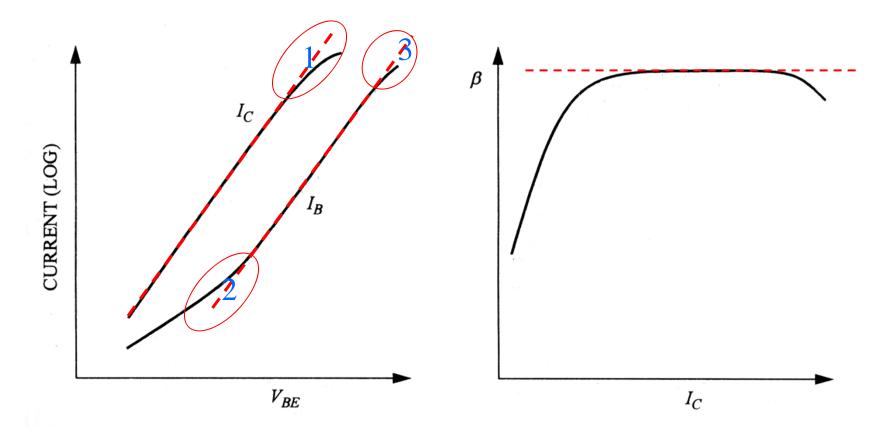
$$Q_{S} = -\frac{1}{2}en_{i}\left(e^{\frac{eV_{BE}}{2kT}} + e^{\frac{eV_{BC}}{2kT}}\right)W_{B}$$

$$J_{n} = -\frac{2eD_{nB}}{W_{B}}n_{i}e^{\frac{eV_{BE}}{2kT}} - \frac{2eD_{nB}}{W_{B}}n_{i}e^{\frac{eV_{BC}}{2kT}}$$

•

Nonideal Effects

- « Gummel plot »:
 - Graph of I_C and I_B versus V_{BE}



Nonideal Effects

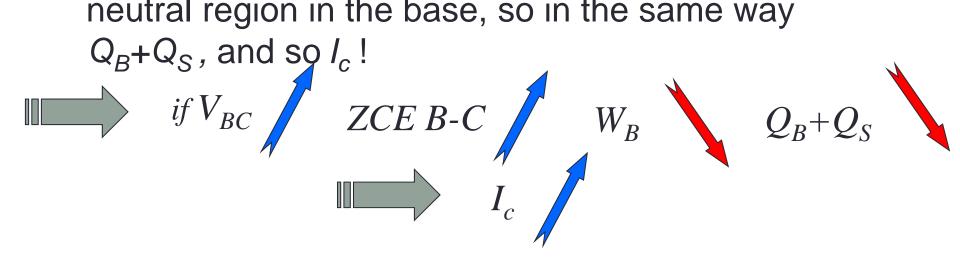
- Early effect / collector punchthrough
- Base collector junction breakdown
- Emitter and Base serie resistance
- I_c « collapse » fort high current
- « crowding effect »

Early effect / collector punchthrough

At "first glance" Ic is independent of V_{CR}

$$I_C = -\left(\frac{Ae^2n_i^2D_{nB}}{Q_B + Q_S}\right)\exp\frac{eV_{BE}}{kT}$$

 In fact we have the modulation of the width of the neutral region in the base, so in the same way

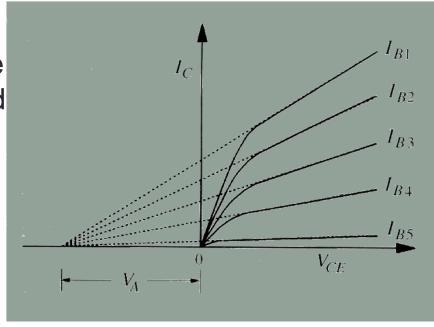


Early effect / collector punchthrough

- At the limit:
 - The space charge BC « depletes » totaly the neutral base
 - collector injects directly current into the emitter.
 - Current only limited by serie resistance R_{serie} form E and C

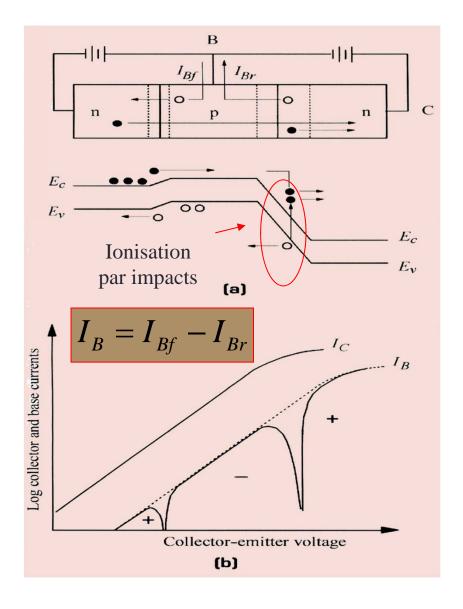
$$V_{pt} = \frac{eW_{B}^{2}N_{A_{B}}(N_{A_{B}} + N_{D_{C}})}{2\varepsilon_{SC}N_{D_{C}}}$$

$$|V_A| \approx \frac{Q_{pB}}{C_{dBC}} = \frac{eN_BW_B}{\varepsilon_{SC}/W_{ZCE_{BC}}}$$



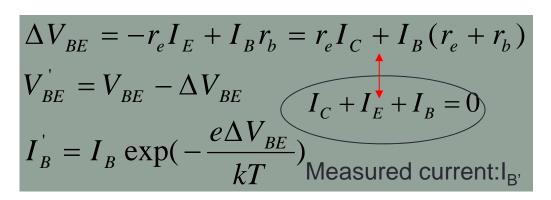
Base - Collector junction breakdown

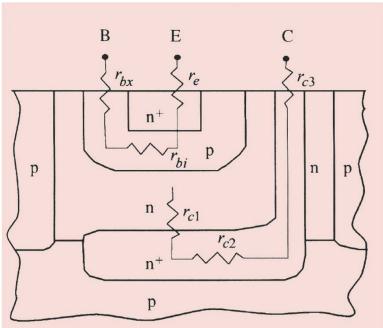
- Avalanche of BC:
 - Often occurs before punchthrough
 - How to prevent it?
 - Lowering the electric field:
 - Reduce the doping gradient in the collector
 - Lightly doped layer between base and collector



Emitter and Base resistance (effet 3)

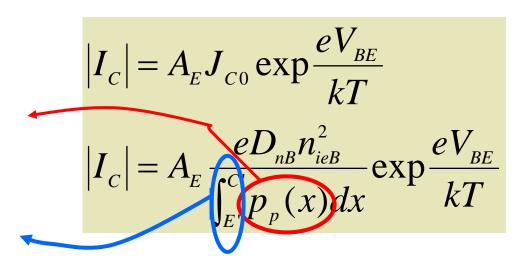
- At low current negligeable effects
- In high speed circuits, B-C allways reverse biased (r_{c2} and r_{c3} as low as possible)
 - Resistances r_c have no effects on current flow
- only r_e and r_b play a major role.
 - IR drop Voltage

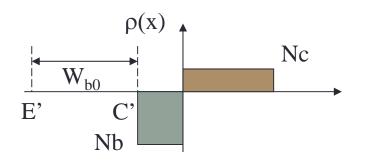


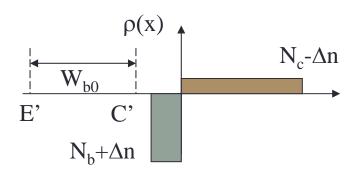


Collector « collapse » for high injection level (effect 1)

- Numbers of physical mechanisms can cause this falloff of I_{CO:}
 - Increase of the charge (holes) into the base (maintain neutrality)
 - Increase the width of the quasineutral base layer (push of the space charge into the collector): Kirk effect



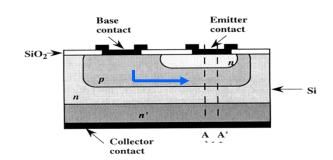


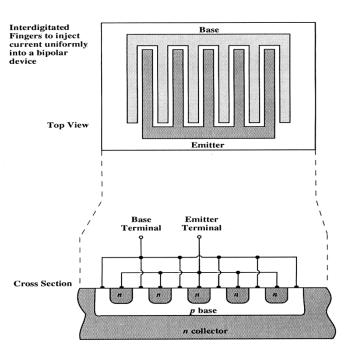


And the effect #2 ???????????????

« crowding effect »

- The view of a 1D devices is an approximation
- Drop voltage IR in the base.
- The edge of the emitter contact is more biased than the core
- Favour a high density of current essentially along the edges
- Not a good thing for high power devices
- Solutions: interdigitated approach

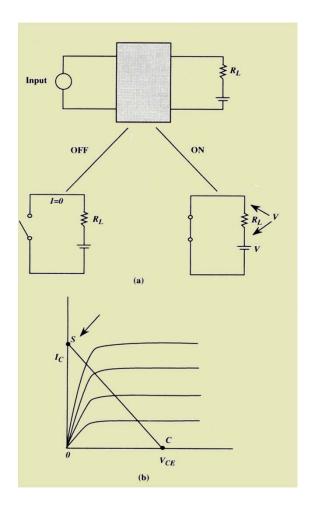




Bipolar Transistor: a switch?

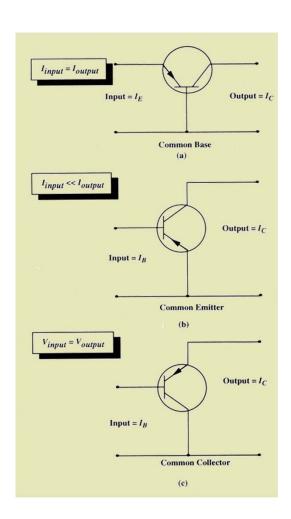
ON state: the switch is closed (saturation mode)

OFF state: the switch is open (cutoff mode)



Input signal as low as possible input power as low as possible

common emitter



Switching velocity? Limiting factors?

- Turn on transient:
 - Continuity equation fir the charge:

$$I_n = \frac{dQ_B}{dt} + \frac{Q_B}{\tau_n}$$

Base charge can be written:

$$Q_B(t) = I_B \tau_n [1 - \exp(-\frac{t}{\tau_n})]$$

Collector current is given by:

$$I_C(t) = \frac{Q_B(t)}{\tau_t}$$
 transit time

in the base (narrow)

$$\frac{Q_B(t)}{\tau_t} = I_c = I_B \frac{\tau_n}{\tau_t} [1 - \exp(-\frac{t}{\tau_n})]$$

- Turn on:
 - I_C increases until the saturation regime is reached:

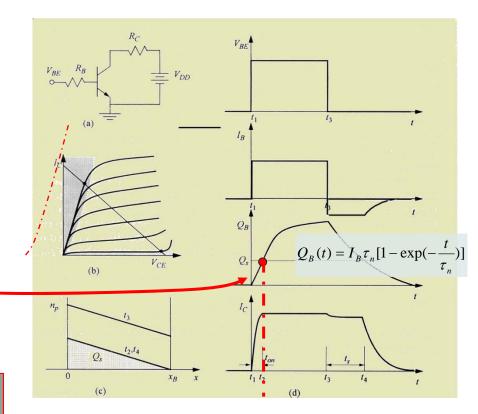
$$I_{C_{sat}} \approx \frac{V_{DD}}{R_C}$$
 (we neglect V_{CEsat})

• The limit charge $Q_B(t_{on})$ to saturate the transistor is given by :

$$Q_S = \frac{I_{C_{sat}} d_{pB}^2}{2D_{nB}}$$

• The on state time is given by:

$$t_{ON} = \tau_n \ln \left[\frac{1}{1 - (Q_S/I_B \tau_n)} \right] = \tau_n \ln \left[\frac{1}{1 - (I_c/I_B \delta)} \right]$$



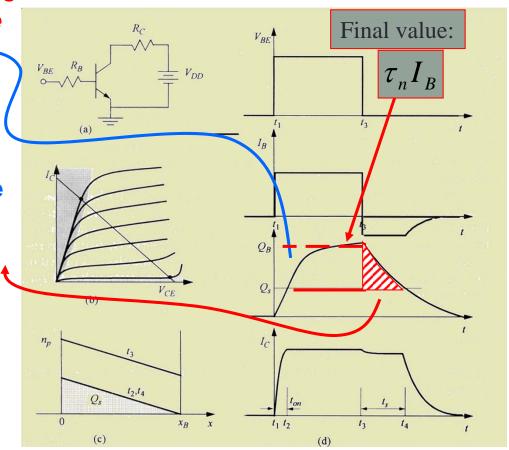
 Remark: the charge can increases over QB(ON) to over saturate the transistor

- Turn off time: input with a « 0 »:
 - Evacuation of the stored charge

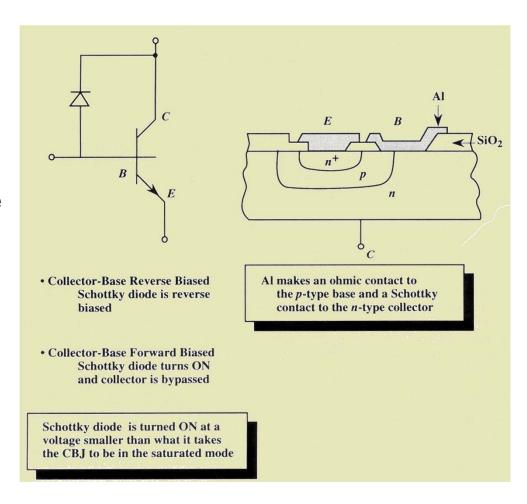
This is the storage time t_s

$$t_S = \tau_n \ln \left(\frac{I_B \tau_n}{Q_S} \right)$$

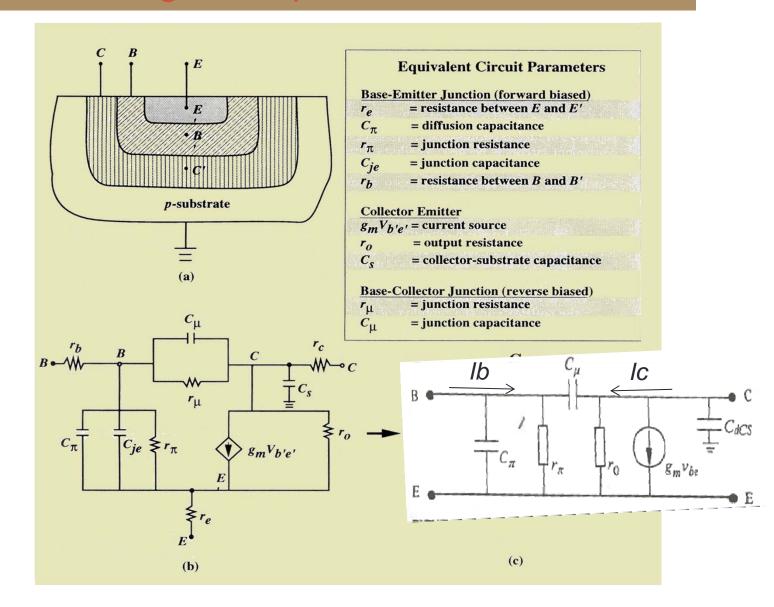
 after, this is the same mechanism than for PN Junction



- Storage time (desaturation) limits the switching time
- 2 ways to reduce it:
 - Add impurities which decrease strongly the lifetime into the base
 - Schottky Diode in // with CB junction: avoid saturation of the transistor



AC signal: equivalent circuit



 Transconductance :links the variation of the collector current to the base – emiter voltage:

$$g_{m} = \frac{\partial I_{C}}{\partial V_{BE}} = \frac{eI_{C}}{kT}$$

Input resistance: links the variation of Base current to the base – emiter voltage:

$$r_{\pi} = \left(\frac{\partial I_B}{\partial V_{BE}}\right)^{-1} = \frac{kT}{eI_B} = \frac{\beta}{g_m} = h_{11}$$

Output resistance:

$$r_o = \left(\frac{\partial I_C}{\partial V_{CE}}\right)^{-1} \approx \frac{V_A}{I_C} = \frac{1}{h_{22}}$$

• Capacitance C_π :

$$C_{\pi} = C_{SE} + C_{T_{EB}}$$

Storage capacitance

$$C_{SE} = \tau_F g_m$$

transit time

$$\tau_{F} = t_{E} + t_{t_{B}} + t_{t_{BE}} + t_{t_{BC}}$$

• Capacitance C_u : CB junction capacitance reverse biased

$$C_{\mu} = C_{T_{CB}}$$

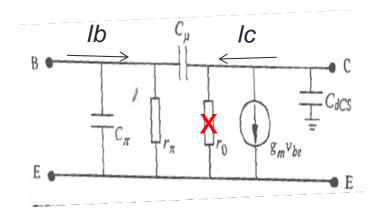
• Collector –substrate (depletion layer) capacitance:



- Cut off frequency (current gain =1)
 - We neglect r0

$$i_{c} = g_{m} v_{be} - j\omega C_{\mu} v_{be}$$

$$i_{b} = \left(\frac{1}{r_{\pi}} + j\omega C_{\pi} + j\omega C_{\mu}\right) v_{be}$$



Current gain is given by:

$$\beta(\omega) = \frac{i_c}{i_b} = \frac{g_m - j\omega C_{\mu}}{(1/r_{\pi}) + j\omega (C_{\pi} + C_{\mu})} = h_{21}$$

- At low frequency:
 - In modern devices, in general,

$$\omega C_{\mu} \ll g_{m} \qquad \beta(\omega) \approx \frac{i_{c}}{i_{b}} = \frac{g_{m} r_{\pi}}{1 + j \omega r_{\pi} (C_{\pi} + C_{\mu})}$$

At high frequecy, the imaginary term dominates and :

$$\beta(\omega) \approx \frac{g_m}{j\omega(C_{\pi} + C_{\mu})}$$

Transistor en ac: schéma équivalent

On obtient alors la fréquence de coupure (« cutoff frequency ») en faisant i_C/i_B=1

Soit encore

$$2\pi f_{T} = \frac{g_{m}}{C_{\pi} + C_{\mu}}$$

$$\frac{1}{2\pi f_{T}} = (\tau_{F}) + \frac{kT}{eI_{C}}(C_{SE} + C_{T_{BC}}) + C_{T_{BC}}(r_{e} + r_{c})$$

Temps de transit en direct

Transistor en ac: schéma équivalent

- Maximun oscillation frequency ⇔Power gain=1
 - Laborious and tedious calculation/ we have to take into account the base resistance

$$f_{\text{max}} = \sqrt{\frac{f_T}{8\pi r_b C_{dBC}}}$$

Current gain:

$$\alpha = \frac{\gamma_E}{1 - \gamma_E} \frac{\delta - 1}{\delta} = \left[1 - \frac{n_{i_E}^2}{N_{D_E}} \frac{D_{p_E}}{D_{n_B}} \frac{N_{A_B}}{n_{i_B}^2} \frac{W_{Beff}}{L_{p_E}} \right] \left[1 - \frac{W_{Beff}^2}{2L_{n_B}^2} \right]$$

In the case of narrow base:

$$\alpha = \frac{\gamma_E}{1 - \gamma_E} \frac{\delta - 1}{\delta} = \left[1 - \frac{n_{i_E}^2}{N_{D_E}} \frac{D_{p_E}}{D_{n_B}} \frac{N_{A_B}}{n_{i_B}^2} \frac{W_{Beff}}{L_{p_E}} \right]$$

$$\alpha = \frac{\gamma_E}{1 - \gamma_E} \frac{\delta - 1}{\delta} = \left[1 - \frac{n_{i_E}^2}{N_{D_E}} \frac{D_{p_E}}{D_{n_B}} \frac{N_{A_B}}{n_{i_B}^2} \frac{W_{Beff}}{L_{p_E}} \right]$$

- If we want a gain β with a high value, we need α close to 1. It means:
 - Lowering base doping
 - Lowering base width (careful with punchthrough!)

Increase of the base resistance $\Leftrightarrow f_{max}$ decreases

$$\alpha = \frac{\gamma_E}{1 - \gamma_E} \frac{\delta - 1}{\delta} = \left[1 - \frac{n_{i_E}^2}{N_{D_E}} \frac{D_{p_E}}{D_{n_B}} \frac{N_{A_B}}{n_{i_B}^2} \frac{W_{Beff}}{L_{p_E}} \right]$$

- Other solution:
 - Increase emitter doping
 - Improve emitter efficiency
 - Problem: « gap shrinking » \Leftrightarrow $\Delta E_g = E_{g(base)} E_{g(\acute{e}metteur)} > 0$.

$$n_i^2(\acute{e}metteur) = n_i^2(Base) \exp(\frac{\Delta E_g}{kT})$$

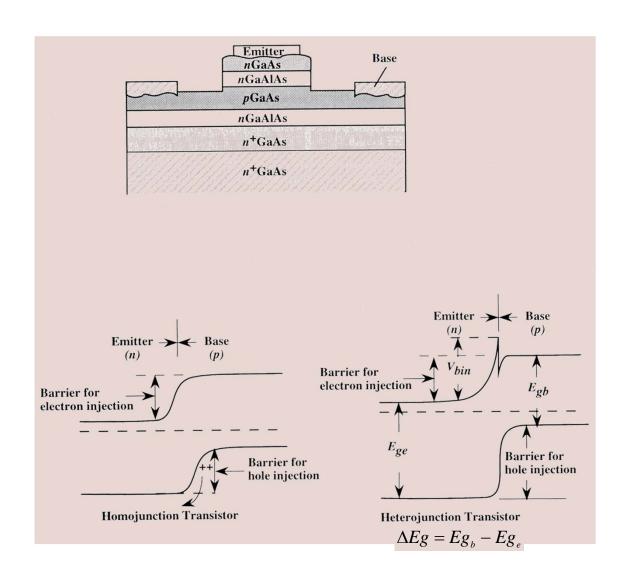
This « problem » can be transform in opportunity !! (next two slides)

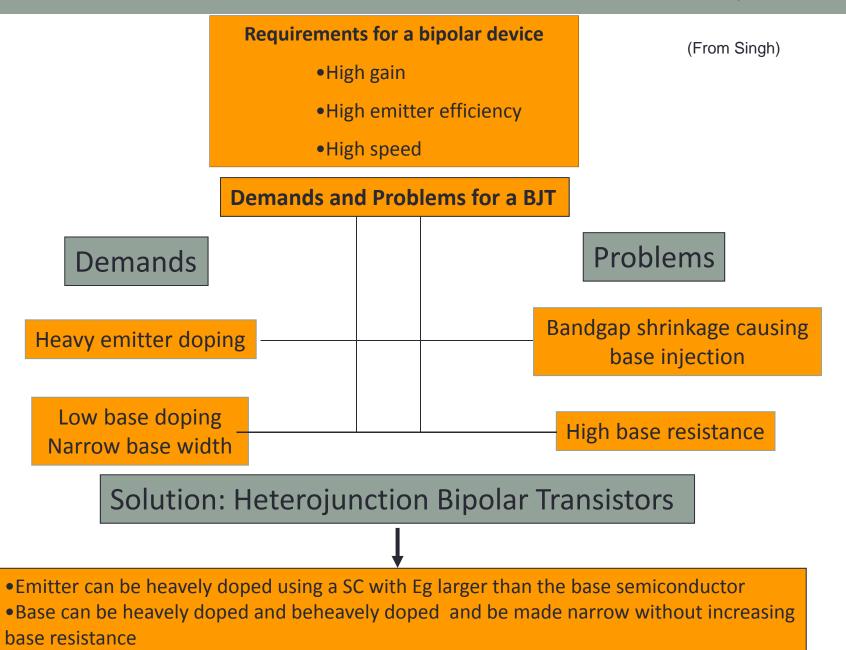
$$\alpha = \left[1 - \frac{n_{i_e}^2}{N_{D_E}} \frac{D_{p_E}}{D_{n_B}} \frac{N_{A_B}}{n_{i_b}^2} \frac{W_{Beff}}{L_{p_E}} \exp \frac{\Delta E_g}{kT}\right]$$

$$\beta = \frac{\alpha}{1 - \alpha} = \left[\frac{N_{D_E}}{N_{A_B}} \frac{D_{n_B}}{D_{p_E}} \frac{L_{p_E}}{W_{Beff}} \exp(-\frac{\Delta E_g}{kT}) - 1 \right]$$

We see that it is difficult to have at the same time a large emitter doping, a lightly doped and thin base and a large gain

 We design a structure with a negative energy gap difference: this is an HBT





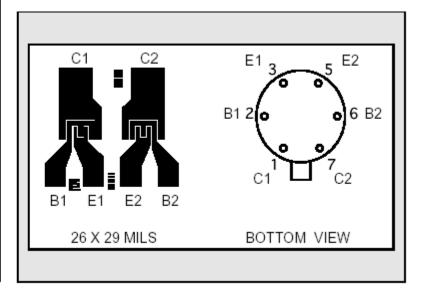
• Collector can be chosen from a material to increase breakdown voltage



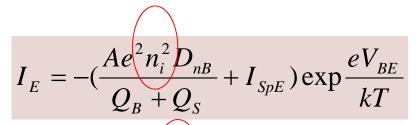
Linear Integrated Systems

FEATURES		
VERY HIGH GAIN	h _{FE} ≥ 2000 @ 1.0μA TYP.	
LOW OUTPUT CAPACITANCE	C _{OBO} ≤2.0pF	
TIGHT V _{BE} MATCHING	$ V_{BE1}-V_{BE2} = 0.2 \text{mV TYP}.$	
HIGH f _T	100MHz	
ABSOLUTE MAXIMUM RATINGS <u>NOTE 1</u> @ 25°C (unless otherwise noted)		
I _C Collector Current	5mA	
Maximum Temperatures		
Storage Temperature	-65° to +200°C	
Operating Junction Temperature	+150°C	
Maximum Power Dissipation	ONE SIDE	BOTH SIDES
Device Dissipation @ Free Air	250mW	500mW
Linear Derating Factor	2.3mW/°C	4.3mVV/°C

LS301 LS302 LS303 HIGH VOLTAGE SUPER-BETA MONOLITHIC DUAL NPN TRANSISTORS



Transistor with a Silicon – Germanium base



$$I_{C} = +\left(\frac{Ae^{2}n_{i}^{2}D_{nB}}{Q_{B}+Q_{S}}\right)\exp\frac{eV_{BE}}{kT}$$

Electrons density

