

Manuscript Number:

Title: Discontinuous Diffusion Synthetic Acceleration for Sn Transport on 2D Arbitrary Polygonal Meshes

Article Type: Regular Article

Keywords: Diffusion Synthetic Acceleration  
DSA  
Sn transport equation  
piece-wise linear discontinuous finite element  
discontinuous finite element method  
interior penalty method.

Corresponding Author: Dr. Jean C Ragusa, PhD

Corresponding Author's Institution: Texas A&M University

First Author: Bruno Turcksin, PhD

Order of Authors: Bruno Turcksin, PhD; Jean C Ragusa, PhD

Abstract: In this paper, a Diffusion Synthetic Acceleration (DSA) technique applied to the Sn radiation transport equation is developed using Piece-Wise Linear Discontinuous (PWLD) finite elements on arbitrary polygonal grids.

The discretization of the DSA equations employs an Interior Penalty technique, as is classically done for the stabilization of the diffusion equation using discontinuous finite element approximations. The penalty method yields a system of linear equations that is Symmetric Positive Definite (SPD). Thus, solution techniques such as Preconditioned Conjugate Gradient (PCG) can be effectively employed. Algebraic MultiGrid (AMG) and Symmetric Gauss-Seidel (SGS) are employed as conjugate gradient preconditioners for the MIP system. AMG is shown to be significantly more efficient than SGS.

Fourier analyses are carried out and we show that this scheme is always stable and effective at reducing the spectral radius for iterative transport solves, even for grids with high-aspect ratio cells. Numerical results are presented for different grid types: quadrilateral, hexagonal, and polygonal grids as well as grids with local mesh adaptivity.

Suggested Reviewers: Ed Larsen PhD  
Professor, Nuclear Engineering and Radiological Sciences, University of Michigan  
edlarsen@umich.edu  
expert in synthetic acceleration schemes for radiation transport

Teresa Bailey PhD  
research scientist, LLNL  
bailey42@llnl.gov

expert in radiation transport; has developed a piece-wise linear DFEM solution technique for radiation transport on polygonal/polyhedral grids

Dmitriy Anistratov PhD  
Associate Professor, Nuclear engineering, NCSU  
anistratov@ncsu.edu  
expert in radiation transport and acceleration schemes

Jim Warsa PhD  
research scientist, LANL  
warsa@lanl.gov  
expert in radiation transport and synthetic acceleration schemes; has developed a code using piece-wise linear DFEM

Jean Ragusa  
Department of Nuclear Engineering  
Texas A&M University  
College Station, TX 77843-3133, USA  
phone: (979) 862 2033  
e-mail: jean.ragusa@tamu.edu

December 24, 2013

Professor William Martin  
Editor,  
Journal of Computational Physics

Dear Professor Martin,

Please find attached a copy of our manuscript titled “Discontinuous Diffusion Synthetic Acceleration for  $S_n$  Transport on Arbitrary Polygonal Meshes” for submission to the *Journal of Computational Physics*.

Some years ago (starting circa 2003), a piecewise linear *discontinuous* (PWLD) finite element discretization of the  $S_n$  transport equation has been proposed by Marv Adams et al. Their PWLD discontinuous finite element discretization can be advantageously employed with arbitrary polygonal and polyhedral meshes and satisfies the thick diffusion limit. However, in thick and diffusive configurations, the standard iterative techniques for  $S_n$  transport can be slowly converging and need to be preconditioned, usually via a Diffusion Synthetic Acceleration (DSA) scheme.

This manuscript presents a DSA scheme that is fully compatible with a PWLD discretization for arbitrary polygonal meshes. The DSA diffusion matrix is SPD, thus amenable to efficient solution techniques such as PCG and we present results with algebraic multigrid (AMG) as the preconditioner for CG. Our DSA scheme is based on a common interior penalty technique for the discontinuous finite element discretization of the diffusion equation. Numerical results are provided for grids with arbitrary quadrangles, arbitrary polygons, a mesh with a juxtaposition of triangles, rectangles and hexagons, and a locally refined mesh as obtained in adaptive mesh refinement. Cells with high-aspect ratios are also considered and the proposed DSA scheme performed quite well, especially with CG preconditioned with AMG.

This work follows closely prior works on DSA schemes for DFEM discretizations of the transport equation, for instance, works by Jim Morel (LANL, now TAMU), Jim Warsa (LANL), Marv Adams (TAMU), Teresa Bailey (LLNL), Marvin Adams (TAMU), Ed Larsen (U Mich), and Todd Palmer (OSU), among (many) others.

Thank you for considering this manuscript for publication in JCP.

Best regards,

Bruno Turcksin and Jean Ragusa

# Discontinuous Diffusion Synthetic Acceleration for $S_n$ Transport on 2D Arbitrary Polygonal Meshes

Bruno Turcksin<sup>\*,a</sup>, Jean C. Ragusa<sup>\*\*,a</sup>

<sup>a</sup>*Department of Nuclear Engineering, Texas A&M University College Station, TX 77843, USA*

---

## Abstract

In this paper, a Diffusion Synthetic Acceleration (DSA) technique applied to the  $S_n$  radiation transport equation is developed using Piece-Wise Linear Discontinuous (PWLD) finite elements on arbitrary polygonal grids. The discretization of the DSA equations employs an Interior Penalty technique, as is classically done for the stabilization of the diffusion equation using discontinuous finite element approximations. The penalty method yields a system of linear equations that is Symmetric Positive Definite (SPD). Thus, solution techniques such as Preconditioned Conjugate Gradient (PCG) can be effectively employed. Algebraic MultiGrid (AMG) and Symmetric Gauss-Seidel (SGS) are employed as conjugate gradient preconditioners for the MIP system. AMG is shown to be significantly more efficient than SGS. Fourier analyses are carried out and we show that this scheme is always stable and effective at reducing the spectral radius for iterative transport solves, even for grids with high-aspect ratio cells. Numerical results are presented for different grid types: quadrilateral, hexagonal, and polygonal grids as well as grids with local mesh adaptivity.

*Key words:* Diffusion Synthetic Acceleration, DSA,  $S_n$  transport equation, piece-wise linear finite element, discontinuous finite element method, interior penalty method.

---

## 1. Introduction

In this paper, we present a Diffusion Synthetic Acceleration (DSA) scheme that is fully compatible with the Piece-Wise Linear Discontinuous (PWLD) finite element discretization of the transport equation on arbitrary polygonal cells.

Arbitrary polygonal (polyhedral in 3D) cells can advantageously be employed, especially in the context of spatial discretizations based on discontinuous finite elements (DFE), for the following reasons: polygonal grids (i) allow for a reduced numbers of unknowns and (ii) provide a natural transition for locally adapted meshes.

To illustrate these two points, first consider a hexagonal cell. Employing a PWLD discretization, such a cell possesses six unknowns. Using other DFE that perform well in the thick diffusive limit [1], such as linear discontinuous on triangles and bilinear discontinuous on quadrangles, the same hexagonal cell could be

---

\*Current address: Department of Mathematics, Texas A&M University College Station, TX 77843, USA

\*\*Corresponding author

*Email addresses:* `bruno.turcksin@math.tamu.edu` (Bruno Turcksin ), `jean.ragusa@tamu.edu` (Jean C. Ragusa)

12 split into: two quadrangles (for a total of 8 unknowns), two triangles and one quadrangle (10 unknowns), or  
 13 four triangles (12 unknowns). By inserting an extra point inside the cell, the hexagon could also be divided  
 14 into: three quadrangles (12 unknowns), four quadrangles (16 unknowns) or six triangles (18 unknowns).  
 15 A similar reasoning can be applied to any  $n$ -polygon. Arbitrary polygonal grids can also handle locally  
 16 refined meshes in a natural manner. The example given in Figure 1 is typical of configurations obtained  
 17 with Adaptive Mesh Refinement (AMR), where neighboring cells possess different refinement level depths.  
 18 Solvers based on arbitrary polygonal cells can easily handle cells with various numbers of edges as follows:  
 19 on Figure 1, the left cell is actually interpreted as a (degenerate) pentagon whereas the two cells on the right  
 20 are quadrilaterals. The PWLD spatial discretization can handle locally adapted meshes without any special  
 21 treatment or further approximation of the coupling between cells. Moreover, the finite element representation  
 22 along the refined side is piece-wise linear. Had the finite element representation been, for instance, bilinear  
 23 discontinuous (BLD), the unrefined cell in Figure 1 would remain rectangular with a linear (not piece-wise  
 24 linear) representation of the solution along the refined edge.

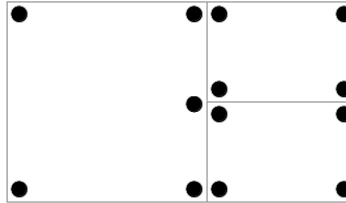


Figure 1: Adaptive mesh refinement grid with PWLD finite elements: left cell is interpreted as a pentagon while right cells are quadrangles.

25 Next, we recall the rationale for solving transport problems iteratively using diffusion acceleration schemes  
 26 (preconditioners). Because analytical solutions are unavailable for most radiation transport problems of  
 27 practical interest, one typically employs iterative techniques to solve the large system of equations that  
 28 results from the spatial and angular discretizations of the transport equation. Standard iterative techniques  
 29 for the first-order form of the discrete-ordinate ( $S_n$ ) transport equation include Source Iteration (SI) and  
 30 Krylov subspace techniques (usually GMRes [2]). For highly diffusive materials (i.e., with scattering ratios  
 31  $c = \Sigma_s/\Sigma_t$  close to 1) and optically thick configurations (i.e., problems that are not leakage-dominated), these  
 32 iterative techniques can become quite ineffective, requiring high iteration counts and possibly leading to false  
 33 convergence. To mitigate these issues, SI and GMRes-based transport solves can be effectively accelerated  
 34 (preconditioned) using Diffusion Synthetic Acceleration (DSA) [3, 4, 5, 6, 7, 8].

The spatial discretization of the DSA equations must be somewhat “consistent” with the one used for the  $S_n$  transport equations in order to yield unconditionally stable and efficient DSA schemes [3, 4, 5, 6, 7, 8]. However, the search for full consistency between the discretized transport equations and the discretized diffusion may not be computationally practical (especially for unstructured arbitrary meshes, [3]). For instance, Warsa, Wareing, and Morel [5] derived a fully consistent DSA scheme for linear discontinuous finite elements on unstructured tetrahedral meshes; their DSA scheme yields a  $P_1$  system of equations that is computationally more expensive than partially consistent DSA schemes that are based upon discretizations of a standard diffusion equation. Several partially consistent schemes have been analyzed for discontinuous finite element discretizations of the transport equation on unstructured meshes, for example, the modified-four-step (M4S) scheme [6], the Wareing-Larsen-Adams (WLA) scheme [7], and the Modified Interior Penalty (MIP) scheme [8].

To the authors’ knowledge, no work is currently ongoing to adapt the M4S technique to polygonal meshes. This is very likely due to the fact that the M4S scheme does not yield a Symmetric Positive Definite (SPD) matrix and was found to be divergent for 3D tetrahedral meshes with linear discontinuous elements [5]. Recent work to develop a DSA scheme for polygonal cells has mainly focused on adapting the WLA scheme to polygonal meshes [9, 10]. The WLA scheme is a two-stage process, where first a diffusion solution is obtained using a *continuous* finite element discretization and then a *discontinuous* update is performed cell-by-cell in order to provide an appropriate discontinuous scalar flux correction to the discontinuous finite element transport solver. In [5], the WLA scheme was found to be a stable and effective DSA technique, though its efficiency degraded as the problem became optically thick and highly diffusive. In this paper, we extend the MIP technique to the PWLD discretization technique for arbitrary polygonal meshes. The MIP scheme is based on the standard Interior Penalty (IP) method for the discontinuous finite element discretization of diffusion equations. MIP was first derived in [8], where it was applied to triangular unstructured meshes (with locally adapted cells). MIP does not suffer from the same issues as the WLA technique when the problem becomes optically thick and highly diffusive and, therefore, can be a useful alternate DSA for to accelerate DFE transport solves. In [8], a Preconditioned Conjugate Gradient (PCG) technique (with Symmetric Gauss-Seidel, SGS, as preconditioner) was used to solve the MIP-DSA equation. In this paper, we also analyze the effectiveness of algebraic multigrid methods (AMG) [11, 12] as a preconditioner.

The remainder of this paper is organized as follows. In Section 2, we briefly review the PWLD discontinuous finite element discretization and the iterative solution techniques applied to the  $S_n$  transport

equation. The MIP-DSA scheme is extended to the PWLD discretization for arbitrary polygons in Section 3. In Section 4, we introduce the Algebraic MultiGrid (AMG) approaches used here, which are based on the ML package of Trilinos [13] and the AGMG (AGgregation-based algebraic MultiGrid) technique of [14, 15, 16, 17]. In Section 5, we present a Fourier analysis for the MIP-DSA scheme discretized with PWLD, and we compare the different preconditioned CG approaches. Conclusions are given in Section 6.

## 2. Discretization and Solution Techniques for the $S_n$ Transport Equation

In this Section, we review the  $S_n$  transport equation and the iterative solution techniques typically employed to solve it. We then describe the PWLD discontinuous spatial discretization for the transport equation with an emphasis on arbitrary polygonal grids.

### 2.1. The $S_n$ Transport Equations

Given an angular quadrature set  $\{\mathbf{\Omega}_m, w_m\}_{1 \leq m \leq M}$ , the one-group  $S_n$  transport equation with isotropic source and isotropic scattering is:

$$(\mathbf{\Omega}_m \cdot \nabla + \Sigma_t(\mathbf{r})) \psi_m(\mathbf{r}) = \frac{1}{4\pi} (\Sigma_s(\mathbf{r})\phi(\mathbf{r}) + S(\mathbf{r})), \text{ for } \mathbf{r} \in \mathcal{D}, 1 \leq m \leq M, \quad (1)$$

with  $\psi_m(\mathbf{r}) = \psi(\mathbf{r}, \mathbf{\Omega}_m)$  the angular flux at position  $\mathbf{r}$  in direction  $\mathbf{\Omega}_m$ ,  $\Sigma_t$  and  $\Sigma_s$  the total and scattering cross sections, respectively, and  $\mathcal{D}$  the spatial domain. The scalar flux is defined and numerically evaluated as follows:

$$\phi(\mathbf{r}) \equiv \int_{4\pi} \psi(\mathbf{r}, \mathbf{\Omega}) d\mathbf{\Omega} \approx \sum_{m=1}^M w_m \psi_m(\mathbf{r}). \quad (2)$$

For brevity, we assume only incoming boundary conditions, that is,  $\psi_m(\mathbf{r}_b) = \psi_m^{inc}(\mathbf{r}_b)$  for any point on the inflow boundary:  $\mathbf{r}_b \in \partial\mathcal{D}_m^- = \{\partial\mathcal{D} \text{ such that } \mathbf{\Omega}_m \cdot \mathbf{n}_b < 0\}$ , with  $\mathbf{n}_b = \mathbf{n}(\mathbf{r}_b)$  the outward unit normal vector at  $\mathbf{r}_b$ . Equation (1) can be written in a compact form using operators:

$$\mathbf{L}\Psi = \mathbf{M}\mathbf{\Sigma}\Phi + S \equiv q, \quad (3)$$

$$\Phi = \mathbf{D}\Psi, \quad (4)$$

where  $\Psi = (\psi_1, \dots, \psi_M)^T$  is the vector of angular fluxes,  $\Phi$  the scalar flux,  $q$  is the total (scattering+external) source,  $\mathbf{L}$  is the streaming operator,  $\mathbf{\Sigma}$  is the scattering matrix,  $\mathbf{M}$  is the moment-to-direction operator, and

$\mathbf{D}$  is the direction-to-moment operator.  $\mathbf{L} = \text{diag}(\mathbf{L}_1, \dots, \mathbf{L}_m, \dots, \mathbf{L}_M)$  is a block-diagonal operator; given a total source  $q$ , one can solve independently for the resulting angular fluxes in all directions. The action of  $\mathbf{L}^{-1}$  is often referred to as a transport sweep when discontinuous spatial approximations are employed: for any direction  $\boldsymbol{\Omega}_m$ , the action of  $\mathbf{L}_m^{-1}$  can be obtained by traversing the mesh (i.e., sweeping) in a specific ordering of the cells, thus requiring that a small linear system of equations be solved cell by cell. The order in which the elements are solved constitutes the graph of the sweep; for brevity and since this is not the focus of this article, we do not expand on situations where the graph presents dependencies (cycles). In such a case, these dependencies can either be lagged within the iterative procedure [18] or the solution vector consisting of the scalar flux is augmented by the angular flux unknowns that cause the cycle [8].

## 2.2. Solution Techniques

Equations (3) and (4) can be solved using the Source Iteration (SI) method (a stationary iterative technique also known as Richardson iteration). The  $\ell^{\text{th}}$  iteration of SI is given by :

$$\Phi^{(\ell+1)} = \mathbf{D}\mathbf{L}^{-1} \left( \mathbf{M}\boldsymbol{\Sigma}\Phi^{(\ell)} + S \right). \quad (5)$$

Alternatively, a subspace Krylov method (usually GMRes) can be employed to solve the linear system of equations, recast as follows:

$$(\mathbf{I} - \mathbf{D}\mathbf{L}^{-1}\mathbf{M}\boldsymbol{\Sigma}) \Phi = \mathbf{D}\mathbf{L}^{-1}S. \quad (6)$$

SI and GMRes-based transport solves require transport sweeps (the action of  $\mathbf{L}^{-1}$ ). When the scattering ratio  $c = \frac{\Sigma_s}{\Sigma_t}$  tends to one in optically thick domains, the number of SI and GMRes iterations can become large. To speed up convergence, a DSA preconditioner is needed; this is further discussed in Section 3.

## 2.3. Discontinuous Finite Element Discretization on Arbitrary Grids

To seek a DFE solution for the angular flux, the domain  $\mathcal{D}$  is partitioned into arbitrary polygonal cells. For a given streaming direction  $\boldsymbol{\Omega}_m$ , the discontinuous finite element formulation, on a given spatial cell  $K$ , is given by:

$$-\int_K (\psi_m \boldsymbol{\Omega}_m \cdot \nabla b_i + \Sigma_t \Psi_m b_i) \, d\mathbf{r} + \int_{\partial K^+} \boldsymbol{\Omega}_m \cdot \mathbf{n} \Psi_m b_i \, d\mathbf{r} = \int_K q b_i \, d\mathbf{r} + \int_{\partial K^-} |\boldsymbol{\Omega}_m \cdot \mathbf{n}| \psi_m^\uparrow b_i \, d\mathbf{r}, \quad (7)$$



where  $b_i$  represents a generic discontinuous finite element basis function,  $\partial K^-$  is the inflow face of element  $K$ , and  $\partial K^+$  is the outflow face of element  $K$ . The angular flux values on an inflow face, denoted by  $\psi_m^\uparrow$  in Eq. (7), are taken from the upwind neighbor element of that face.

Next, we define the  $b_i$  basis functions for the PWLD discretization. First, we introduce the within-cell point  $c$  for any 2D polygonal cell. The coordinates of  $c$  are the weighted averages of the polygon's vertices:

$$u_c = \sum_{i=1}^{N_V} \alpha_i u_i, \quad (8)$$

where  $u = x$  or  $y$ ,  $N_V$  is the number of vertices for the cell under consideration, and the (positive) weights  $\alpha_i$  are such that  $\sum_{i=1}^{N_V} \alpha_i = 1$ . The basis function at vertex  $i$  is defined by (see also [19]):

$$b_i(x, y) = t_i(x, y) + \alpha_i t_c(x, y), \quad (9)$$

where  $t_i(x, y)$  is a linear function such that  $t_i(x, y)$  is unity at vertex  $i$  and zero at vertices  $i - 1$ ,  $i + 1$ , and  $c$ . The  $t_c(x, y)$  function is a “tent” function in the interior of the cell; it is unity at the within-cell point  $c$  and zero at all of the vertices of the cell. The number of PWLD basis functions is equal to the number of vertices in the polygon (i.e.,  $1 \leq i \leq N_V$ ). In this paper, the arbitrary positive weights  $\alpha_i$  are chosen to be equal to  $\frac{1}{N_V}$ . For example, on a quadrilateral cell, we employ  $\alpha_i = \frac{1}{4}$ . Finally, note that on triangular cells, the PWLD basis functions reduce to the standard Linear Discontinuous (LD) basis functions if one chooses  $\alpha_i = \frac{1}{3}$ . Given the definition of the PWLD finite elements, it may seem complicated to build the elementary transport matrices on an arbitrary polygonal cell, but the construction of such matrices can be greatly simplified using the notion of “sub-cells”, where a “sub-cell”, in 2D, is defined as the triangular cell linking an edge of the polygonal cell to its within-cell point. Note that these “sub-cells” are never created nor stored as part of the polygonal grid. However, from a code implementation perspective, the computation of the elementary matrices on a polygonal cell is greatly simplified by looping over its sub-cells.

### 3. Diffusion Synthetic Acceleration (DSA) on Arbitrary Polygonal Cells

#### 3.1. DSA Solution Principle

As noted earlier, standard iterative techniques applied to transport solves can be slowly converging in thick diffusive configurations. A DSA scheme must be employed to accelerate their convergence. The idea

126 behind synthetic acceleration is that the error between the (yet unknown) transport solution and the current  
 127 iterate can be estimated from a computationally less expensive process, yielding a corrective term to be added  
 128 to the current iterate in order to improve the next iterate. In DSA, a corrective scalar flux contribution,  $\delta\Phi$ ,  
 129 is sought through the following diffusion equation [3],

$$\mathbf{A} \delta\Phi = \Sigma \left( \Phi^{(\ell+1/2)} - \Phi^{(\ell)} \right). \quad (10)$$

130 where the source term is a scattering term due to the difference between the previous iterate scalar flux  $\Phi^{(\ell)}$   
 131 and the newest scalar flux,  $\Phi^{(\ell+1/2)}$ , obtained after a single transport sweep. The next scalar flux iterate is  
 132 obtained by adding the scalar flux correction to the latest scalar flux, yielding:

$$\Phi^{(\ell+1)} = \Phi^{(\ell+1/2)} + \delta\Phi. \quad (11)$$

133  $\mathbf{A}$  is the diffusion matrix of the DSA scheme. Ideally,  $\mathbf{A}$  should be SPD (such that efficient iterative techniques  
 134 can be employed in its linear solve) and easy to form (even on arbitrary grids).

### 135 3.2. Modified Interior Penalty DSA Scheme for Polygons

Here, we discuss the Modified Interior Penalty (MIP) method for the diffusion equation for arbitrary polygons as a DSA scheme. In its discretization, this DSA scheme employs the same discontinuous finite elements used in the discretization of the transport operator. MIP as a DSA scheme has been shown to be always stable on triangular cells for isotropic scattering [8] (for highly forward-peaked scattering, DSA is not an efficient preconditioner and one needs to employ an angular multigrid technique [20], whose coarsest angular level can be a diffusion solve; see [21] for instance). The MIP diffusion solve is based on the standard Interior Penalty (IP) method [22]. Consider the following diffusion update equation:

$$(-\nabla \cdot \mathbf{D} \nabla + \Sigma_a) \delta\phi = Q_0 \quad \text{for } \mathbf{r} \in \mathcal{D} \quad (12)$$

$$\frac{1}{4} \delta\phi - \frac{1}{2} \mathbf{D} \partial_n \delta\phi = 0 \quad \text{for } \mathbf{r} \in \partial\mathcal{D} \quad (13)$$

136 where  $\mathbf{D}$  is the diffusion coefficient and  $\Sigma_a$  is the absorption cross section. For DSA, the volumetric source term  
 137 is given by the scattering due to the difference in scalar flux between two iterations,  $Q_0 = \Sigma_s (\Phi^{(\ell+1/2)} - \Phi^{(\ell)})$ ,  
 138 and the known incoming angular flux boundary condition from the transport problem translates into the

139 Robin-type vacuum boundary condition given in Equation (13), where  $\partial_n = \mathbf{n} \cdot \nabla$  and  $\mathbf{n}$  is the outer normal  
 140 unit vector on the domain's boundary,  $\partial\mathcal{D}$ . The MIP weak form for the DSA update equation is given by:

$$a(\delta\phi, \phi^*) = l(\phi^*) \quad (14)$$

141 with the bilinear (matrix) form:

$$\begin{aligned} a(\delta\phi, \phi^*) = & (\Sigma_a \delta\phi, \phi^*)_{\mathcal{D}} + (\mathbf{D} \nabla \delta\phi, \nabla \phi^*)_{\mathcal{D}} + (\kappa_e^{MIP} \llbracket \delta\phi \rrbracket, \llbracket \phi^* \rrbracket)_{E_h^i} + (\llbracket \delta\phi \rrbracket, \{\!\!\{ \mathbf{D} \partial_n \phi^* \}\!\!\})_{E_h^i} \\ & + (\{\!\!\{ \mathbf{D} \partial_n \delta\phi \}\!\!\}, \llbracket \phi^* \rrbracket)_{E_h^i} + (\kappa_e^{MIP} \delta\phi, \phi^*)_{\partial\mathcal{D}^d} - \frac{1}{2} (\delta\phi, \mathbf{D} \partial_n \phi^*)_{\partial\mathcal{D}^d} - \frac{1}{2} (\mathbf{D} \partial_n \delta\phi, \phi^*)_{\partial\mathcal{D}^d} \end{aligned} \quad (15)$$

142 and the linear (right-hand-side) form

$$l(\phi^*) = (Q_0, \phi^*)_{\mathcal{D}}. \quad (16)$$

143 The notations used are as follows: the domain integral is split onto the element integrals,  $(f, g)_{\mathcal{D}} =$   
 144  $\sum_{K \in \mathbb{T}_h} (f, g)_K$ , with the element integrals defined as  $(f, g)_K = \int_K f g \, d\mathbf{r}$ ; and the integral over all inte-  
 145 rior edges is  $(f, g)_{E_h^i} = \sum_{e \in E_h^i} (f, g)_e$ , with the edge integrals defined as  $(f, g)_e = \int_e f g \, ds$ . In the above,  
 146  $\mathbb{T}_h$  denotes the partition of domain  $\mathcal{D}$  into non-overlapping elements  $K$ ,  $E_h^i$  is the set of interior edges,  $\delta\phi$   
 147 is the numerical solution for the diffusion correction, and  $\phi^*$  is any test function. Any entry  $A_{ij}$  of matrix  
 148  $\mathbf{A}$  is simply given by  $A_{ij} = a(b_j, b_i)$ . The resulting matrix  $\mathbf{A}$  is SPD [23, 24] and thus can be solved using  
 149 a preconditioned conjugate gradient technique. The jump and mean value of a variable  $u$  at the interface  
 150 between two elements is defined as follows:

$$\llbracket u \rrbracket = u^+ - u^- \text{ and } \{\!\!\{ u \}\!\!\} = \frac{u^+ + u^-}{2}, \quad (17)$$

151 respectively. The definition of the left/right values along a given edge  $e$  is

$$u^{\pm} = \lim_{s \rightarrow 0^{\pm}} u(\mathbf{r} + s \mathbf{n}_e), \quad (18)$$

152 where  $\mathbf{n}_e$  is the normal unit vector associated with the given edge  $e$  (on the boundary  $\mathbf{n}_e$  is the external  
 153 normal). The MIP penalty coefficient is given by

$$\kappa_e^{MIP} = \max \left( \kappa_e^{IP}, \frac{1}{4} \right) \quad (19)$$

154 with:

$$\kappa_e^{IP} = \begin{cases} \frac{c}{2} \left( \frac{D^+}{h_\perp^+} + \frac{D^-}{h_\perp^-} \right) & \text{on interior edges, i.e., for } e \in E_h^i \\ c \frac{D}{h_\perp} & \text{on boundary edges, i.e., for } e \in \partial\mathcal{D} \end{cases} \quad (20)$$

155 where  $c$  is a constant (chosen equal to 4 here) and  $h_\perp$  is the length of the cell in the direction orthogonal to the  
 156 edge  $e$ . When polygonal cells are employed, there are no simple ways in which to compute  $h_\perp$ . To estimate  
 157  $h_\perp$ , we assume that the polygonal cells do not deviate significantly from regular polygons. In such cases, if  
 158 the regular polygon has an even number of edges, the orthogonal length equals two times the apothem. The  
 159 apothem is the line segment between the polygon's center and the midpoint of the polygon's side and its length  
 160 is equal to twice the polygon's area divided by the polygon's perimeter,  $\text{apothem} = 2 \times \frac{\text{area}}{\text{perimeter}}$ . If polygon  
 161 has an odd number of sides, its orthogonal length is given by the sum of the apothem and the circumradius,  
 162 the latter being the radius of the circle circumscribed to the polygon  $\left( \text{circumradius} = \sqrt{\frac{2 \times \text{area}}{N_V \sin\left(\frac{2\pi}{N_V}\right)}} \right)$ . A  
 163 summary of the definitions used for  $h_\perp$  for any polygon as a function of the number of vertices is given in  
 164 Table 1.

Table 1: Orthogonal length  $h_\perp$  for different polygonal types : area is the polygon's area; perimeter is the polygon's perimeter;  $N_V$  is the number of vertices;  $L_e$  is the edge's length.

Number of vertices	3	4	> 4 and even	> 4 and odd
$h_\perp$	$2 \times \frac{\text{area}}{L_e}$	$\frac{\text{area}}{L_e}$	$4 \times \frac{\text{area}}{\text{perimeter}}$	$2 \times \frac{\text{area}}{\text{perimeter}} + \sqrt{\frac{2 \times \text{area}}{N_V \sin\left(\frac{2\pi}{N_V}\right)}}$

## 165 4. MIP-DSA Solves Based on Algebraic Multigrid Techniques

### 166 4.1. AMG Principles

167 A common way to solve an SPD system of equations is to use a Conjugate Gradient (CG) technique  
 168 preconditioned with Symmetric Gauss-Seidel (PCG-SGS) or SSOR (PCG-SSOR); SGS is simply SSOR with  
 169 a damping factor of one. Here, we compare CG preconditioned with Symmetric Gauss-Seidel (PCG-SGS) with  
 170 CG preconditioned with an algebraic multigrid method. PCG-SGS was chosen because little difference was  
 171 noted when employing other damping factors for MIP-DSA on triangular grids [25]. In our implementation,  
 172 we have linked our code with the ML package [13] of the Trilinos library and with the AGMG library [14].  
 173 ML is a multigrid preconditioning package that uses a smoothed aggregation algebraic multigrid to build  
 174 a preconditioner for a Krylov method. AGMG is an aggregation-based algebraic multigrid library written

175 in Fortran 90. The first multigrid methods developed were geometric multigrid techniques, used as stand-  
 176 alone solvers. In many applications, they achieve the so-called “textbook multigrid efficiency”, i.e., “the  
 177 solution to the governing system of equations [is attained] in a computational work that is a small multiple  
 178 of the operation counts associated with discretizing the system” [26]. However, in many other applications,  
 179 multigrid methods, and particularly algebraic multigrid methods, may not achieve such efficiency, [27] and,  
 180 in such cases, they are often used as preconditioner for Krylov subspace methods.

181 Before describing a general multigrid method, we recall the process using two grids. Consider the following  
 182 system

$$\mathbf{A}_f u_f = b_f, \quad (21)$$

183 defined on a fine grid  $\Gamma_f$ . The two-grid algorithm is as follows:

- 184 1. Perform  $\nu_1$  pre-smoothing iterations using a smoother (Jacobi, Gauss-Seidel, or ILU) and an initial  
 185 guess  $u_0$ :  $u = S^{\nu_1}(u_0, b_f)$
- 186 2. Compute the residual on the fine grid  $\Gamma_f$  and restrict it to the coarse grid  $\Gamma_c$ :  $r_c = \mathbf{R}(b_f - \mathbf{A}_f u)$ ;
- 187 3. Solve the system on the coarse grid:  $v = \mathbf{A}_c^{-1} r_c$ ;
- 188 4. Interpolate the coarse grid correction to the fine grid and add the correction to  $u$ :  $u \leftarrow u + \mathbf{P}v$ ;
- 189 5. Perform  $\nu_2$  post-smoothing iterations:  $u = S^{\nu_2}(u, b_f)$ .

190 When using AMG, the matrix  $\mathbf{A}_c$  on the coarse grid is given by the Galerkin approximation:

$$\mathbf{A}_c = \mathbf{R} \mathbf{A}_f \mathbf{P}, \quad (22)$$

191 where  $\mathbf{P}$  is a prolongation matrix and  $\mathbf{R}$  is a restriction matrix. Generally, solving the system  $\mathbf{A}_c v = r_c$   
 192 on the coarse grid is still quite expensive, therefore this step is recursively replaced by  $n_\gamma$  sequences of the  
 193 two-grid method until the system can be efficiently inverted with a direct solver. When  $\gamma = 1$ , respectively  
 194  $\gamma = 2$ , the multigrid method is said to use a  $V$ -cycle, respectively a  $W$ -cycle. Figure 2 shows typical  $V$ -  
 195 and  $W$ -cycles. In Figure 2, a dot represents a smoothing operation and a square a direct inversion; the grid  
 196 transfer operators are symbolized by the line segments.

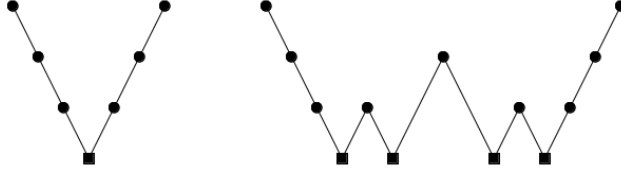


Figure 2:  $V$ - and  $W$ -cycles.

The main difference between geometric and algebraic multigrid techniques lies in the method used to coarsen the grid. Algebraic multigrid methods only use the properties of the matrix. Among the algebraic multigrid methods, there are three main categories: the classical Ruge-Stueben AMG, the plain aggregation AMG, and the smoothed aggregation AMG. ML uses smoothed aggregation AMG and AGMG uses plain aggregation AMG. Next, we briefly explain the coarsening step in the ML and AGMG implementations. The coarsening step is a crucial step because convergence rates will decrease if coarsening occurs too quickly. However, if coarsening is too slow, more memory may be required to solve the problem.

#### 4.2. The ML Package of Trilinos

When using a smoothed aggregation scheme, the smoothed interpolation operators  $\mathbf{P}_k$  are the transpose of the coarsening operators  $\mathbf{R}_k = \mathbf{P}_k^T$ . Therefore, when the  $\mathbf{P}_k$  matrices are built, the coarsening operator is also known. First, the graph of the matrix is constructed: if element  $(i, j)$  or  $(j, i)$  of the matrix is non-zero, an edge is built between the vertex  $i$  and the vertex  $j$  [13]. Second, the vertices are aggregated. When using ML on a single processor, two aggregation schemes can be used: the uncoupled scheme or the maximally independent sets (MIS) scheme. The uncoupled scheme tries to build aggregates of size  $3^d$  where  $d$  is the dimension of the problem; its algorithm proceeds as follows [28]:

**Step 1:** As long as there are points not adjacent to an aggregate:

1. Choose a point which is not adjacent to an aggregate. This point is a new root point.
2. Define a new aggregate as the root point and its neighbors .

**Step 2:** Add all the points left to the existing aggregates or form new aggregates with them.

The MIS scheme used in ML applies the MIS algorithm [29] to the graph of the matrix  $\mathbf{A}^2$ . These two coarsening schemes use a fixed ratio of coarsening between levels. Once the aggregation is done, a tentative

218 prolongation matrix,  $\tilde{\mathbf{P}}_k$  is constructed [28]. An example of  $\tilde{\mathbf{P}}_k$  is given by:

$$\tilde{\mathbf{P}}_k(i, j) = \begin{cases} 1 & \text{if the } i^{th} \text{ point is contained in the } j^{th} \text{ aggregate} \\ 0 & \text{otherwise} \end{cases} \quad (23)$$

219 This tentative prolongation operator could be used as is but smoothing it allows to have a more robust  
 220 scheme. Let  $\mathbf{S}_k$  be a smoother, for example damped Jacobi. Then, the prolongation matrix is given by:

$$\mathbf{P}_k = \mathbf{S}_k \tilde{\mathbf{P}}_k. \quad (24)$$

#### 221 4.3. The AGMG Package

222 Unlike ML, the prolongation operator in AGMG is not smoothed; this results in a cheaper setup and  
 223 a decrease in memory requirements [16]. However, such a scheme could be less robust. To counteract  
 224 this weakness, the aggregation scheme is more involved. Coarsening algorithms that control the size of  
 225 the aggregates tend to produce a few badly shaped aggregates. Since the convergence of AMG is bounded  
 226 by the worst aggregate, even a small number of badly shaped aggregates can have a huge impact on the  
 227 convergence. In AGMG, the aggregation algorithm has as input the upper bound of the two-grid condition  
 228 number. When the aggregates are constructed, their quality is checked. This increases the cost of the  
 229 coarsening and therefore, it is important that the coarsening is fast enough. Such an algorithm does not  
 230 control the size of the aggregates, and, therefore, it is difficult to control the speed of coarsening. However,  
 231 controlling the condition number rather than the coarsening speed can be an effective alternative approach.  
 232 By monitoring the condition number, bad aggregates will not be created and, instead, a few aggregates below  
 233 the target size may be generated. This does not affect the efficiency of the method in a noticeable way [16].  
 234 A simple way to create the aggregates would be to try to exhaust all of the possible combinations, compute  
 235 their quality, and then choose the optimal coarsening. In practice, this would be too costly, and, in AGMG,  
 236 the aggregation step is done by a few passes of a pairwise aggregation algorithm. Each pass aggregates the  
 237 variables two by two to allow a simple computation of the aggregate quality and to keep the cost per iteration  
 238 low.

## 5. Numerical Results

In this section, we present two series of results. First, Fourier analyses are carried out to analyze the performance of the MIP-DSA acceleration scheme for an homogeneous infinite medium meshed with rectangular cells and discretized with PWLD finite elements. The effects of the  $S_n$  order and the cell aspect ratio on the spectral radius of the iterative scheme are also studied. Second, the MIP-DSA technique is implemented in a 2D  $S_n$  code that uses arbitrary polygonal grids with a PWLD spatial discretization. Numerical examples employ several mesh types: arbitrary quadrilaterals; arbitrary polygons; a grid containing a regular layout of hexagons/triangles/rectangles; and a grid that mimics adaptive mesh refinement performed on rectangular cells. The MIP-DSA diffusion solves are performed using various linear solvers: Conjugate Gradient (CG); Conjugate Gradient Preconditioned with Symmetric Gauss-Seidel (PCG-SGS); Conjugate Gradient Preconditioned with ML using Uncoupled aggregation (PCG-ML/U); Conjugate Gradient Preconditioned with ML using MIS aggregation (PCG-ML/MIS); and Conjugate Gradient Preconditioned with AGMG (AGMG).

### 5.1. Fourier Analyses

Fourier analysis is often performed to assess some of the properties of DSA-accelerated transport solves [3, 4, 5]. In a Fourier analysis, the eigenvalues of the iteration matrix are analyzed, assuming a Fourier ansatz for the error modes. Specifically, the iteration matrices for the SI and SI+DSA schemes are given by

$$DL^{-1}M\Sigma \quad \text{and} \quad I - (I + A^{-1}\Sigma)(I - DL^{-1}M\Sigma), \quad (25)$$

respectively. The error modes are of the form  $\exp(i\mathbf{\Lambda} \cdot \mathbf{r})$  with the wave number  $\mathbf{\Lambda} = [\lambda_x, \lambda_y]^T$ . This expression for the error modes is inserted into the discretized equations and the spectral radius (largest eigenvalue in magnitude) of the iteration matrices are sought for  $0 \leq \lambda_x \leq 2\pi/X$  and  $0 \leq \lambda_y \leq 2\pi/Y$ , where  $X$  and  $Y$  are the dimensions of the rectangular domain.

#### 5.1.1. Spectral Radius as a Function of the $S_n$ Order

This Fourier analysis is carried on a square cell ( $X = Y$ ) using a Gauss-Legendre-Chebyshev (GLC) angular quadrature. The medium is homogeneous with a scattering ratio  $c = 0.9999$ ; periodic boundary conditions are used. The results are plotted on Figure 3, where the  $x$ -axis is the mesh size in mean free paths and the  $y$ -axis is the spectral radius. There are four curves corresponding to different  $S_n$  orders:  $S_2$ ,  $S_4$ ,  $S_8$ , and  $S_{16}$ . From Figure 3, we conclude that a PWLD discretization of the MIP-DSA equations is stable



for every cell size. The spectral radius is always less than 0.5, except for  $S_2$  where it is about 0.7. In the fine mesh limit, the continuum results are recovered [30]. Notably, the spectral radius of SI+DSA using an  $S_2$  quadrature in 2D undiscretized space is  $0.5c$ . Also, as the angular quadrature is refined, the standard limit value of  $0.2247c$  for the spectral result is reached.

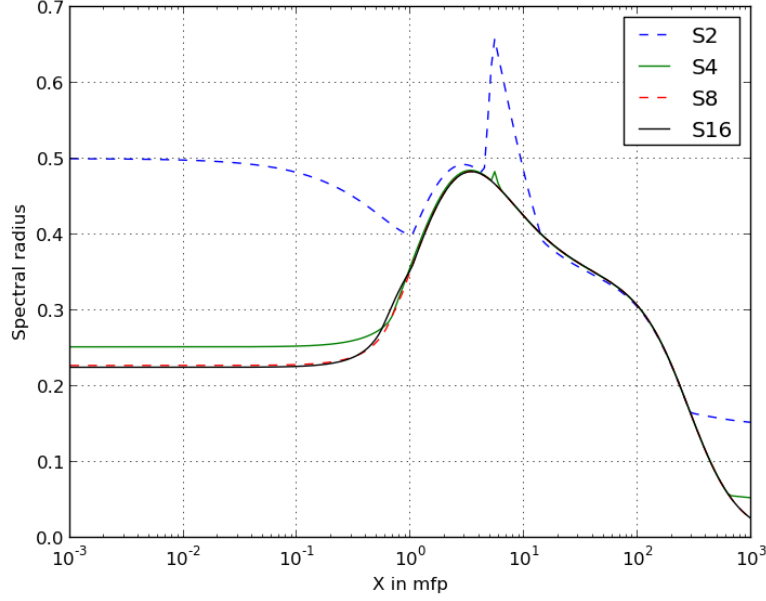


Figure 3: Spectral radius as a function of the mesh optical thickness for various  $S_n$  orders (Fourier analysis on a square cell).

### 5.1.2. Spectral Radius as a Function of the Cell Aspect Ratio

For this Fourier analysis, we use a  $S_{16}$  GLC quadrature. The medium is again homogeneous with  $c = 0.9999$ , and periodic boundary conditions apply. On Figure 4, the five curves correspond to the following cell aspect ratios:  $\frac{Y}{X} = \frac{1}{16}; \frac{1}{4}; 1; 4; 16$ ; and 100. We note that the MIP-DSA scheme is stable for every aspect ratio tested, including an aspect ratio value of 100, and that the maximum spectral radius shows little sensitivity to the aspect ratio.

### 5.2. Performance of MIP-DSA Implemented in a PWLD $S_n$ Transport Code

The MIP-DSA scheme has been implemented in a 2D  $S_n$  code that employs a PWLD discretization for arbitrary polygonal grids. Several test cases are presented.

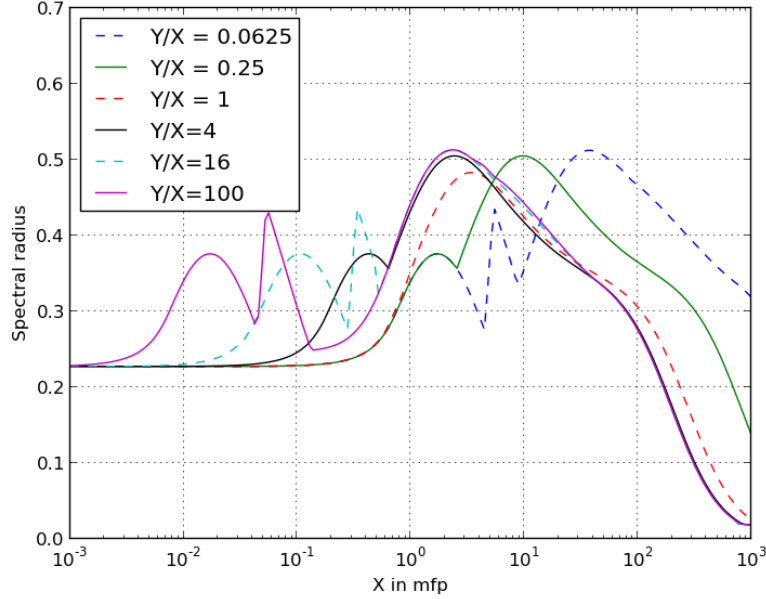


Figure 4: Spectral radius as a function of the mesh optical thickness for various cell aspect ratios (Fourier analysis carried out using an  $S_{16}$  angular quadrature).

### 5.2.1. Homogeneous Test Cases

We compare the different linear solvers employed for MIP-DSA using a homogeneous medium,  $100\text{cm} \times 100\text{cm}$ ,  $\Sigma_t = 1\text{cm}^{-1}$  and  $\Sigma_s = 0.999\text{cm}^{-1}$ , with vacuum boundary conditions and a unit source of intensity  $1\text{cm}^{-3}\text{s}^{-1}$ . We use an  $S_8$  angular quadrature. A relative tolerance of  $10^{-8}$  is used for the Source Iteration solver while the MIP-DSA solvers employs a relative tolerance of  $10^{-10}$ . The domain is discretized using two different meshes:

1. A quadrilateral grid composed of 49263 quadrilateral cells (197,052 degrees of freedom or spatial unknowns per angular direction). This grid was obtained from an unstructured triangular grid in which each triangle was split into 3 quadrangles.
2. A polygonal grid composed of 45,204 triangles, 823 quadrilaterals, 4,978 pentagons, 4,155 hexagons, 725 heptagons, and 24 octagons, for a total of 55,909 cells and 193,991 degrees of freedom. This grid was obtained from the same triangular grid as before. This time, vertices were removed and the triangles containing them were merged into polygons. Since an arbitrary number of triangles may contain a given point, deleting a vertex leads to various polygonal types. This example allows us to test MIP and the

different preconditioners on a mesh composed of different cell types.

The meshes and the numerical solutions are given in Figures 5 and 6. In Table 2, results obtained with the different linear solvers for MIP-DSA are compared for the quadrilateral grid. In Table 2, *SI iterations* is the number iterations of needed to solve the problem and *CG iterations* is the total number of CG iterations used to solve MIP. We note that the accelerated transport solves only require 24 SI iterations, regardless of the linear solver employed in MIP-DSA (as expected). Preconditioning CG significantly reduces the number of CG iterations. We observe that algebraic multigrid processes, PGC-ML and AGMG, require about the same number of iterations (two orders of magnitude less than unpreconditioned CG). The different linear solvers are compared for the polygonal grid in Table 3. We note that using different spatial cell types in the same grid does not affect the performance of MIP-DSA or that of its preconditioners.

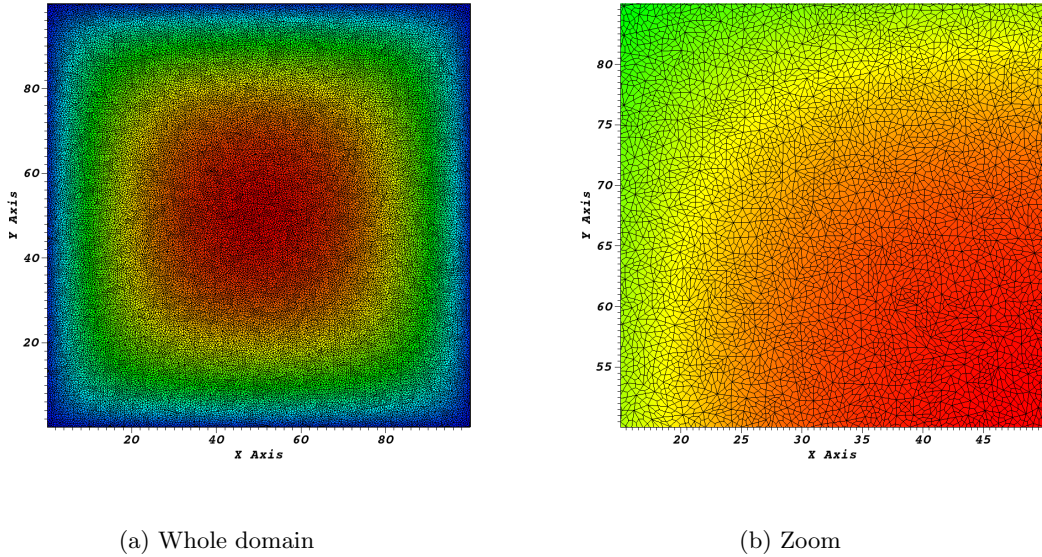
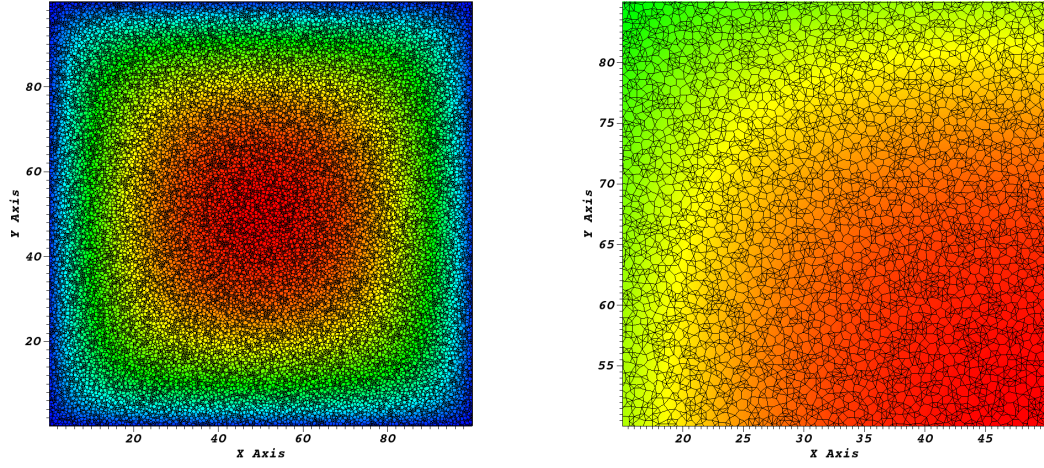


Figure 5: Grids and scalar flux solutions for the homogenous test problem using quadrilateral cells.

Table 2: Comparison of different preconditioners for the homogeneous test problem (quadrilateral cells).

	No-DSA	CG	PCG-SGS	PCG-ML/U	PCG-ML/MIS	AGMG
SI iterations	7311	24	24	24	24	24
CG iterations	NA	56649	17332	630	604	578



(a) Whole domain

(b) Zoom

Figure 6: Grids and scalar flux solutions for the homogenous test problem using arbitrary polygonal cells.

Table 3: Comparison of different preconditioners for the homogeneous test problem (polygonal cells).

	No-DSA	CG	PCG-SGS	PCG-ML/U	PCG-ML/MIS	AGMG
SI iterations	7311	23	23	23	23	23
CG iterations	NA	46262	16712	652	603	555

### 5.2.2. Test Case with Heterogeneous Media

In this example, a heterogeneous geometry with three materials is used. It is composed of 184 triangles, 3,720 quadrilaterals, and 2,791 regular hexagons of side  $0.05\text{cm}$  for a total of 6,695 cells and 32,178 degrees of freedom. The domain is  $5.28275\text{cm}$  by  $4.6\text{cm}$ . Reflective boundary conditions are used. A  $S_{16}$  angular quadrature is used. The SI solver has a relative tolerance of  $10^{-8}$ , and the relative tolerance for MIP-DSA is  $10^{-10}$ . Figure 7 shows the problem geometry and the material properties are in given Table 4. The different linear solvers for MIP-DSA are compared in Table 5. The remarks made in Section 5.2.1 for the homogeneous

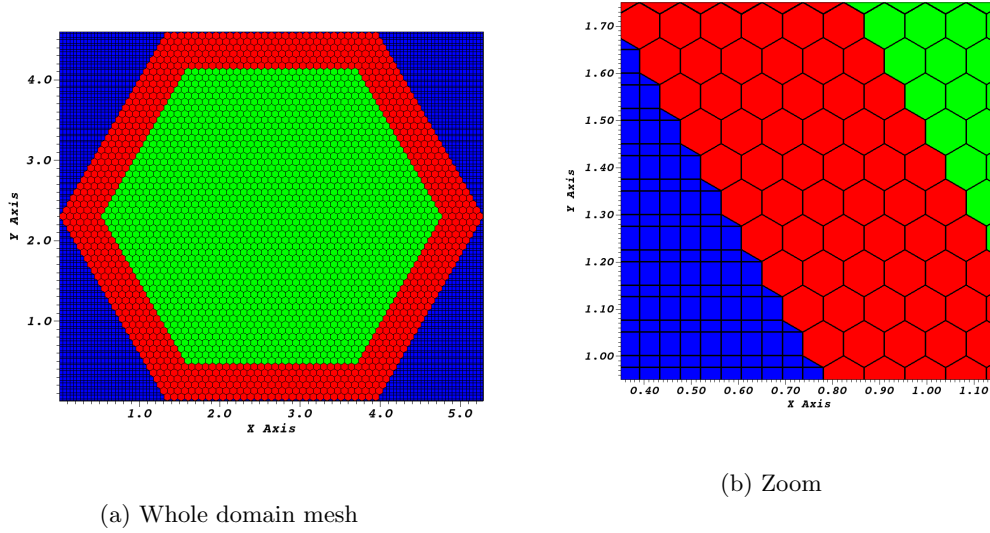


Figure 7: Material Zones for the Heterogeneous Test Problem

Table 4: Material Properties For the Different Regions of the Heterogeneous Test Problem

	Inner region	Intermediate region	Outer region
$\Sigma_t \text{ (cm}^{-1}\text{)}$	1.5	1.0	1.0
$\Sigma_s \text{ (cm}^{-1}\text{)}$	1.4999	0.999	0.3
source $\text{(cm}^{-3}\text{s}^{-1}\text{)}$	1.0	0.0	0.0

Table 5: Preconditioners Comparison (Heterogeneous Problem).

	No-DSA	CG	PCG-SGS	PCG-ML/U	PCG-ML/MIS	AGMG
SI iterations	278	17	17	17	17	17
CG iterations	NA	12214	6679	415	386	248

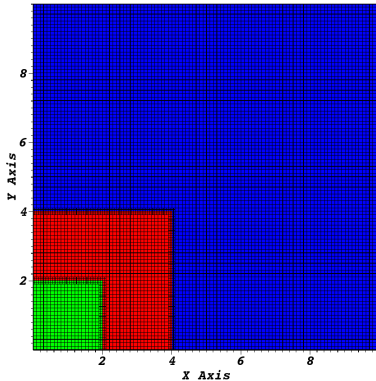
test problem remain mostly unchanged. MIP-DSA is effective for this heterogeneous test case, and AGMG is again the most efficient preconditioner. It is interesting to note that, contrary to the homogeneous tests where the number of CG iterations was about identical for all algebraic multigrid preconditioners, in this heterogeneous test case, AGMG requires significantly fewer iterations than both Trilinos implementations, PCG-ML/U and PCG-ML/MIS.

### 5.2.3. Test Case with Locally Refined Cells

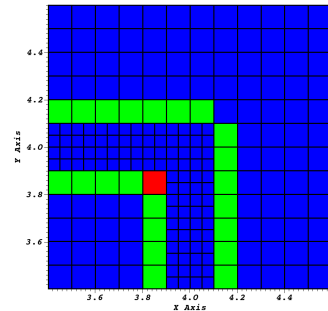
In this example, a 3-material domain of size  $10\text{cm} \times 10\text{cm}$  is used. Figure 8a shows the material zoning and the mesh used. The bottom and left sides of the domain have reflective boundary conditions, while the other two sides have vacuum boundary conditions. Material properties are given in Table 6.

The grid used mimics meshes obtained via adaptive mesh refinement: the rectangular cells at the interfaces between two materials are refined once more, leading to a grid composed of 10,482 quadrilaterals, 236 pentagons, and 2 hexagons for a total of 10,720 cells and 43120 degrees of freedom. The distribution of cells is given in Figure 8b.

A  $S_{16}$  angular quadrature is employed. The tolerance on SI is  $10^{-8}$ , and the tolerance on the CG solvers is  $10^{-10}$ . The different linear solvers for MIP-DSA are compared in Table 7. The conclusions drawn from this test case are similar to the ones made for our previous tests. This test case demonstrates that degenerate polygons (here, pentagons and hexagons) do not seem to affect the MIP-DSA acceleration.



(a) Material regions



(b) Polygons distribution (zoom): quadrilaterals (blue cells), pentagons (green cells), hexagons (red cells)

Figure 8: AMR-like test domain

Table 6: Material Properties (AMR-like Problem)

	Inner region	Intermediate region	Outer region
$\Sigma_t$ ( $cm^{-1}$ )	1.5	1.0	1.0
$\Sigma_s$ ( $cm^{-1}$ )	1.44	0.9	0.3
Source ( $cm^{-3}s^{-1}$ )	1.0	0.0	0.0

Table 7: Preconditioners Comparison for the AMR Mesh

	No-DSA	CG	PCG-SGS	PCG-ML/U	PCG-ML/MIS	AGMG
SI iterations	184	19	19	19	19	19
CG iterations	NA	11300	4734	361	361	264

#### 5.2.4. Test Cases High Aspect Ratio Grids

In these last two examples, a square domain of  $100cm \times 100cm$  with vacuum boundaries is employed. There are 10,000 cells (thus, 40,000 degrees of freedom). Again, the relative tolerance on SI is  $10^{-8}$  and the relative tolerance for CG is  $10^{-10}$ . An  $S_8$  GLC angular quadrature is used.  $\Sigma_t = 1cm^{-1}$ ,  $\Sigma_s = 0.999cm^{-1}$  and the source is  $1n/(cm^3s)$ . In the first run, the domain is discretized using 100 subdivisions of the  $x$  and  $y$  axes, i.e., 10,000 square cells with an aspect ratio of one. In the second run, the domain is discretized using 1,000 subdivision along  $x$  and 10 along  $y$  (the aspect ratio is then 100). As expected, solving the MIP-DSA equations requires more CG iterations when the aspect ratio increases. PCG-ML/U and PCG-ML/MIS are significantly more affected by the increase in the aspect ratio than the other methods. AGMG is the least affected by the change of aspect ratio and is again the best performing preconditioner.

Table 8: Preconditioners Comparison for a Rectangular Grid with an Aspect Ratio of 1

	No-DSA	CG	PCG-SGS	PCG-ML/U	PCG-ML/MIS	AGMG
SI iterations	7311	21	21	21	21	21
CG iterations	NA	8363	4853	376	375	221

Table 9: Preconditioners Comparison for a Rectangular Grid with an Aspect Ratio of 100

	No-DSA	CG	PCG-SGS	PCG-ML/U	PCG-ML/MIS	AGMG
SI iterations	7304	24	24	24	24	24
CG iterations	NA	84802	43466	14180	13896	821

## 6. Conclusions

We have extended the Modified Interior Penalty (MIP) form of Diffusion Synthetic Acceleration (DSA) to  $S_n$  radiation transport solves performed on arbitrary polygonal grids discretized with Piece-Wise Linear Discontinuous finite elements. The MIP-DSE equation employs the same discontinuous finite element trial spaces as the spatial discretization of the  $S_n$  transport equation. As such, only a few additional elementary matrices need to be implemented to an existing  $S_n$  code.

Fourier analyses show that the PWLD discretization of the MIP diffusion synthetic accelerator is stable and effective, including for cells with high-aspect ratios. Numerical experiments have been performed with grids containing various types of polygonal cells (random quadrilaterals, random polygons), including grids with different polygon types for a given mesh and tests with degenerate polygons that mimic grids obtained in adaptive mesh refinement strategies. In these tests, MIP-DSE always performed effectively, reducing significantly the number of Source Iterations needed for convergence. We noted that the effectiveness of MIP does not seem to be affected by the meshes employed.

The MIP-DSE matrix is Symmetric Positive Definite (SPD) and we solved the resulting linear system of equations using a standard Conjugate Gradient method with different preconditioners. Algebraic multigrid techniques (AGMG and ML) were found to be the most effective preconditioner.

We conclude that MIP-DSE is an effective alternative to carry out diffusion synthetic accelerations for  $S_n$  radiation transport problems on arbitrary polygonal grids and grids obtained through mesh adaptivity. Future extensions of this work will include the implementation of the MIP-DSE equations in a 3D parallel  $S_n$  code.

## References

- [1] Marvin L. Adams. Discontinuous Finite Element Transport Solutions in Thick Diffusive Problems. *Nuclear Science and Engineering*, 137:298–333, March 2001.
- [2] Youcef Saad and Martin H. Schultz. GMRES: A Generalized Minimal Residual Algorithm for Solving Nonsymmetric Linear Systems. *SIAM Journal on Scientific and Statistical Computing*, 7(3):856–869, 1986.
- [3] Marvin L. Adams and Edward W. Larsen. Fast iterative methods for discrete-ordinates particle transport calculations. *Progress in Nuclear Energy*, 40:3–159, 2002.



- [4] E. W. Larsen. Diffusion-Synthetic Acceleration Methods for Discrete-Ordinates Problems. *Transport Theory and Statistical Physics*, 13:107–126, 1984.
- [5] James S. Warsa, Todd A. Wareing, and Jim E. Morel. Fully Consistent Diffusion Synthetic Acceleration of Linear Discontinuous  $S_N$  Transport Discretizations on Unstructured Tetrahedral Meshes. *Nuclear Science and Engineering*, 141(3):236–251, July 2002.
- [6] E. W. Larsen. Unconditionally Stable Diffusion-Synthetic Acceleration Methods for Slab Geometry Discrete Ordinates Equations. Part I: Theory. *Nuclear Science and Engineering*, 82, 1982.
- [7] T. A. Wareing, E. W. Larsen, and M. L. Adams. Diffusion Accelerated Discontinuous Finite Element Schemes for the Sn Equations in Slab and X-Y Geometries. *Proc. Int. Topl. Mtg. Advances in Mathematics, Computations, and Reactor Physics, Pittsburgh, Pennsylvania*, April 28 - May 2 1991.
- [8] Y. Wang and J.C. Ragusa. Diffusion Synthetic Acceleration for High-Order Discontinuous Finite Element  $S_n$  Transport Schemes and Application to Locally Refined Unstructured Meshes. *Nuclear Science and Engineering*, 166:145–166, 2010.
- [9] James S. Warsa. A Continuous Finite Element-Based, Discontinuous Finite Element Method for Sn Transport. *Nuclear Science and Engineering*, 160:385–400, 2008.
- [10] Anthony P. Barbu and Marvin L. Adams. Semi-Consistent Diffusion Synthetic Acceleration for Discontinuous Discretizations of Transport Problems. *International Conference on Mathematics, Computational Methods and Reactor Physics, Saratoga Springs, New York*, May 3-7 2009.
- [11] William L. Briggs, Van Emden Henson, and Steve F. McCormick. *A Multigrid Tutorial*. Second edition, 2000.
- [12] Christian Wagner. *Introduction to Algebraic Multigrid*. 1998. Course Notes of an Algebraic Multigrid Course at the University of Heidelberg in the Winter Semester 1998/1999.
- [13] M.W. Gee, C.M. Siefert, J.J. Hu, R.S. Tuminaro, and M.G. Sala. ML 5.0 Smoothed Aggregation User’s Guide. Technical Report SAND2006-2649, Sandiad National Laboratories, 2006.
- [14] Yvan Notay. User’s Guide to AGMG. Technical report, Universite Libre de Bruxelles, 2011.

- [15] Y. Notay. An aggregation-based algebraic multigrid method. *Electronic Transactions on Numerical Analysis*, 37:123–146, 2010.
- [16] Artem Napov and Yvan Notay. An algebraic multigrid method with guaranteed convergence rate. *SIAM Journal of Scientific Computing*, 34(2):1079–1109, 2012.
- [17] Yvan Notay. Aggregation-based algebraic multigrid for convection-diffusion equations. *SIAM Journal of Scientific Computing*, 34(4):2288–2316, 2012.
- [18] Todd A. Wareing, John M. McGhee, Jim E. Morel, and Shawn D. Pautz. Discontinuous Finite Element  $S_N$  Methods on Three-Dimensional Unstructured Grids. *Nuclear Science and Engineering*, 138:256–268, January 2011.
- [19] Hiromi G. Stone and Marvin L. Adams. A piecewise linear finite element basis with application to particle transport. *Nuclear Mathematical and Computational Sciences : A Century in Review, A Century Anew, Gatlinburg, Tennessee*, April 2003.
- [20] J.E. Morel and T.A. Manteuffel. An Angular Multigrid Acceleration Technique for  $S_n$  Equations with Highly Forward-Peaked Scattering. *Nuclear Science and Engineering*, 107:330–342, 1991.
- [21] Bruno Turcksin, Jean C. Ragusa, and Jim E. Morel. Angular Multigrid Preconditioner for Krylov-based Solution Techniques applied to the  $S_n$  Equations with Highly Forward-Peaked Scattering. *Transp. Theor. & Stat. Phys.*, 41:1–22, 2012.
- [22] G. Kanschat. *Discontinuous Galerkin Methods for Viscous Incompressible Flow*. Advances in Numerical Mathematics. Teubner Research, Wiesbaden, Germany, 2007.
- [23] D. N. Arnold. An interior penalty finite element method with discontinuous elements. *SIAM J. Numer. Anal.*, 19:742–760, 1982.
- [24] Douglas N Arnold, Franco Brezzi, Bernardo Cockburn, and Donatella Marini. Discontinuous galerkin methods for elliptic problems. In *Discontinuous Galerkin Methods*, pages 89–101. Springer, 2000.
- [25] Yaqi Wang. Personal E-mail communication (2012-08-29).
- [26] James L. Thomas, Boris Diskin, and Achi Brandt. Textbook Multigrid Efficiency for Fluid Simulations. *Annual Review of Fluid Mechanics*, 35:317–340, January 2003.

- 415 [27] Yvan Notay and Panayot S. Vassilevski. Recursive Krylov-based multigrid cycles. *Numerical Linear*  
416 *Algebra with Applications*, 15:473–487, 2008.
- 417 [28] Ray S. Tuminaro and Charles Tong. Parallel Smoothed Aggregation Multigrid: Aggregation Strategies  
418 on Massively Parallel Machines. *SC Conference*, 0:5, 2000.
- 419 [29] Mark T. Jones and Paul E. Plassmann. A Parallel Graph Coloring Heuristic. *SIAM Journal on Scientific*  
420 *Computing*, 14:654, 1992.
- 421 [30] Marvin L. Adams and Todd A. Wareing. Diffusion-synthetic acceleration given anisotropic scattering,  
422 general quadratures, and multidimensions. *Conference: American Nuclear Society(ANS) annual meeting,*  
423 *San Diego*, 68, June 1993.

**LaTeX Source Files**

[Click here to download LaTeX Source Files: latex-source-files.zip](#)