

Vortexy fluid dynamics simulator

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Software documentation

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1 Introduction

Vortexy is a computational fluid dynamics (CFD) simulation package. It is written in C and uses the finite volume method with the SIMPLE algorithm to calculate flow of incompressible fluids, namely liquids.

The simulator is based on irregular tetrahedral meshes. These meshes can be computed from surfaces using the program Tetgen. The simulator takes a configuration file as input that contains paths to the simulation mesh and boundary conditions in addition to other settings. The state of the system is periodically written to an output file specified in the config. Included is also a renderer that uses OpenGL to visualize results.

2 Background

2.1 Navier-Stokes equations

The Navier-Stokes equations form the basis for all of fluid dynamics. The momentum equation is typically written as

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u} + g, \quad (1)$$

where $\mathbf{u} = (u, v, w)$ is velocity in [m/s], t time in [s], ρ density in [kg/m³], p pressure in [Pa], $\nu = \frac{\mu}{\rho}$ kinematic viscosity in [m²/s], g gravity in [m/s].

The continuity equation must be satisfied for incompressible fluids that have no sinks or sources

$$\nabla \cdot \mathbf{u} = 0, \quad (2)$$

where $\mathbf{u} = (u, v, w)$ is velocity in [m/s].

2.2 Turbulence

A simple way of predicting onset of turbulence is the Reynolds number:

$$\text{Re} = \frac{\mu u L}{\rho} = \frac{u L}{\nu}$$

Turbulence models in simulations include RANS (Reynolds Averaged), LES (Large Eddy) and DNS (Direct).

2.3 Finite volume method

The *finite volume method* (FVM) is based on a simulation mesh with volume elements. This enables evaluation of partial differential equations (PDEs) prevalent in physics. The divergence theorem allows us to convert volume integrals to surface integrals

$$\int_V \nabla \cdot \mathbf{F} \, dV = \oint_S \mathbf{F} \cdot d\mathbf{S},$$

so volume terms can be computed from fluxes at element faces.

2.4 Discretization

The momentum equation is written in a form conducive for discretization:

$$\frac{\partial \mathbf{u}}{\partial t} + \nabla \cdot (\mathbf{u} \otimes \mathbf{u}) = -\frac{\nabla p}{\rho} + \frac{\mu}{\rho} \nabla \cdot (\nabla \mathbf{u}) + \nabla \cdot (\nabla \mathbf{u})^T + \mathbf{f}_b$$

$$\mathbf{transient} + \mathbf{convective} = \mathbf{diffusion} + \mathbf{source}$$

Continuity equation:

$$\nabla \cdot \mathbf{u} = 0$$

This saddle-point problem can be represented in matrix form as:

$$A\mathbf{u} = \begin{pmatrix} F & B^T \\ B & 0 \end{pmatrix} \begin{pmatrix} \mathbf{u} \\ p \end{pmatrix} = \begin{pmatrix} \mathbf{f}_b \\ 0 \end{pmatrix}$$

In practice, the velocity and pressure fields are calculated using 4 matrices (vx, vy, vz, p).

Discretization for the one-dimensional momentum equation [1]:

$$a_v u_v + \sum_{f \in f_{nb}} a_f u_f = b_v, \quad (3)$$

which can be represented in matrix form as:

$$\begin{pmatrix} a_{00} & a_{01} & 0 & \dots & 0 \\ \vdots & \ddots & & & \vdots \\ 0 & a_f & a_v & a_f & 0 \\ \vdots & & & \ddots & \vdots \\ 0 & \dots & 0 & a_{n(n-1)} & a_{0n} \end{pmatrix} \begin{pmatrix} u_0 \\ \vdots \\ u_n \end{pmatrix} = \begin{pmatrix} b_0 \\ \vdots \\ b_n \end{pmatrix},$$

where $u_0 \dots u_n$ represents volume element velocities. In the pressure calculation these are referred to as u^* , i.e. momentum conserving. The coefficients a and b are calculated for each volume element and assembled into a matrix. Components X, Y and Z are calculated one after another.

Face fluxes:

$$\phi_f = \max(\dot{m}_f, 0) + \mu \frac{E_f}{d_{vf}} \quad (4)$$

$$\Phi_f = -\max(-\dot{m}_f, 0) - \mu \frac{E_f}{d_{vf}} \quad (5)$$

$$\vec{\psi}_f = -\mu(\nabla \mathbf{u}_f) \cdot \mathbf{T}_f + \dot{m}_f(\mathbf{u}_f^{\text{hr}} - \mathbf{u}_f^{\text{uw}}), \quad (6)$$

where $\mathbf{S}_f = A_f \hat{n}_f = \mathbf{E}_f + \mathbf{T}_f$ and $\dot{m}_f = \rho \mathbf{u}_f \cdot \mathbf{S}_f$. High-resolution model velocity $\mathbf{u}_f^{\text{hr}} = \mathbf{u}_f$ and upwind velocity \mathbf{u}_f^{uw} .

Volume fluxes:

$$\phi_v = \frac{\rho V}{\Delta t} \quad (7)$$

$$\vec{\psi}_v = \frac{\rho V}{\Delta t} \mathbf{u}_v - \mathbf{f}_b V \quad (8)$$

Coefficients:

$$a_v = \phi_v + \sum_{f \in f_{nb}} \phi_f = \frac{\rho V}{\Delta t} + \sum_{f \in f_{nb}} \left(\max(\dot{m}_f, 0) + \mu \frac{E_f}{d_{vf}} \right) \quad (9)$$

$$a_f = \Phi_f = -\max(-\dot{m}_f, 0) - \mu \frac{E_f}{d_{vf}} \quad (10)$$

$$\mathbf{b}_v = -\vec{\psi}_v - \sum_{f \in f_{nb}} \vec{\psi}_f + \sum_{f \in f_{nb}} (\mu (\nabla \mathbf{u}_f)^T \cdot \mathbf{S}_f) - V \nabla p_v \quad (11)$$

$$= -\frac{\rho V}{\Delta t} \mathbf{u}_v - \mathbf{f}_b V - \sum_{f \in f_{nb}} (-\mu (\nabla \mathbf{u}_f) \cdot \mathbf{T}_f + \dot{m}_f (\mathbf{u}_f - \mathbf{u}_f^{\text{uw}})) + \sum_{f \in f_{nb}} (\mu (\nabla \mathbf{u}_f)^T \cdot \mathbf{S}_f) - V \nabla p_v$$

Pressure equation:

$$u_v^* + \sum_{f \in f_{nb}} \frac{a_f}{f_v} \mathbf{u}_f^* = -\frac{V}{a_v} \nabla p_v + \frac{b_v + V \nabla p_v}{a_v} \quad (12)$$

Similarly, the pressure correction equation can be discretized [1]:

$$a_v p'_v + \sum_{f \in f_{nb}} a_f p'_f = b_v \quad (13)$$

With coefficients:

$$a_f = -\rho \frac{E_f}{d_{vf}} \quad (14)$$

$$a_v = \sum_{f \in f_{nb}} \rho \frac{E_f}{d_{vf}} \quad (15)$$

$$b_v = -\sum_{f \in f_{nb}} \dot{m}_f^* = -\sum_{f \in f_{nb}} \rho \mathbf{u}_f^* \cdot \mathbf{S}_f \quad (16)$$

All of these are $n \times n$ diagonally dominant matrices which can be solved using a numerical matrix solver, such as Gauss-Seidel. The off-diagonal coefficients represent neighbouring elements, so they will be mostly zero.

2.5 Gauss-Seidel method

2.6 SIMPLE algorithm

1. Set boundary conditions, set u and p
2. Compute gradients ∇u and ∇p
3. Compute mass fluxes j_m , flow rate $\dot{m} = j_m \cdot A$
4. Solve *momentum equation* using velocity guess u^0
5. Solve *pressure correction equation* to get p'

6. Correct pressure $p = p + p'$ and velocity
7. Increment time $t^{n+1} = t^n + \Delta t$
8. Repeat

2.7 Boundary conditions

No-slip wall

Slip wall

Inlet

Outlet

3 Implementation

3.1 Simulation mesh

Triangular/tetrahedral, closed, connected, Delaunay

3.2 Solver

Gauss-Seidel

3.3 Rendering

OpenGL

4 Software

4.1 Compilation

Deps: glibc, (OpenGL, GLEW, GLFW)

```
cmake .  
make
```

4.2 Configuration

Files: sim.cfg, data/fluid.cfg

4.3 Examples

4.3.1 Lid-driven cavity

4.3.2 Pipe flow

4.4 Problems

Checkerboard problem

References

- [1] Moukalled, F., Mangani, L. & Darwish M. (2016). The Finite Volume Method in Computational Fluid Dynamics: An Advanced Introduction with OpenFOAM and Matlab. Fluid Mechanics and Its Applications Volume 113. Springer International Publishing, Cham. <https://doi.org/10.1007/978-3-319-16874-6>
- [2] <https://quickersim.com/tutorial/tutorial-2-numerics-simple-scheme/>
- [3] <https://www.openfoam.com/documentation/guides/latest/doc/guide-applications-solvers-simple.html>
- [4] https://www.cfd-online.com/Wiki/SIMPLE_algorithm