

Gummel-Poon Model

MOD-4 ED-14

- Gummel-Poon Model is specially designed for very small BJT considering second-order effects.

It was described by Hermann Gummel and H.C. Poon in Bell Lab in 1970.

This model takes into account effects of base currents and at high level injection.

- This is a charge control model, rectifying drawbacks of Ebers-Moll model.

Considering the minority carrier movement in base for npn.

$$I_{EP} = qA \mu_p p(x_n) E - qA D_p \frac{dp(x_n)}{dx_n} \quad \text{--- (I)}$$

Assuming $n(x_n) = N_d(x_n)$

$$I_{EP} = qA \mu_p p \left(\frac{-kT}{q} \frac{1}{n} \frac{dn}{dx_n} \right) - qA D_p \frac{dp}{dx_n}$$

$$= - \frac{qA D_p}{n} \left(p \frac{dn}{dx_n} + n \frac{dp}{dx_n} \right) \quad \text{--- (II)}$$

$$= - \frac{qA D_p}{n} \frac{d(pn)}{dx_n} \quad \text{--- (III)}$$

$$\frac{-I_{EP} n}{qA D_p} = \frac{d(pn)}{dx_n} \quad \text{--- (IV)}$$

Integrating (IV) from $J_E (w_b=0)$ to $J_C (w_b=w_b)$, considering I_{EP} flowing from E to C.

$$-I_{EP} \int_0^{w_b} \frac{n dx_n}{q A D_p} = \int_0^{w_b} \frac{d(pn)}{dx_n} dx_n$$

$$= p(w_b) n(w_b) - p(0) n(0)$$

As we know $pn = n_i^2$ — (i) (for equilibrium values)
and for non equilibrium values.

$$pn = n_i^2 e^{\frac{F_n - F_p}{kT}} = n_i^2 e^{qV/kT} \quad \text{--- (VII)}$$

$$\left. \begin{aligned} p(w_b) n(w_b) &= n_i^2 e^{qV_{CB}/kT} \\ p(0) n(0) &= n_i^2 e^{qV_{EB}/kT} \end{aligned} \right\} \quad \text{--- (VIII)}$$

$$I_{EP} = \frac{-q A D_p n_i^2 \left(e^{\frac{qV_{CB}}{kT}} - e^{\frac{qV_{EB}}{kT}} \right)}{\int_0^{w_b} n dx_n} \quad \text{--- (IX)}$$

- The integral in denominator corresponds to the integrated majority carrier charge in the base and known as base Gummel number Q_B .

$$\text{Now } I_{EP} = \frac{q A D_p n_i^2 e^{\frac{qV_{EB}}{kT}}}{Q_B} \quad \text{--- (X)}$$

(-ve sign due to reverse bias)

$$\text{Similarly } I_{En} = \frac{q A D_p n_i^2 e^{\frac{qV_{EB}}{kT}}}{Q_E} \quad \text{--- (XI)}$$

(Base e^- current flowing back into emitter)

DE - Integrated majority carrier charge in emitter known as emitter Gummel number.

MOD-2 (20-15)

(V)

As for analysing secondary effect like early effect or high level injection effect

$$Q_B = \int_0^{W_B(V_{EB})} n(x_n) dx_n \quad \text{--- (XII)}$$

As eq (XII) is a biased dependent, so it justifying early effect.

Under high level injection the integrated majority carrier charge becomes greater than the integrated base dopant charge

$$\int_0^{W_B} n(x_n) dx_n > \int_0^{W_B} N_D(x_n) dx_n \quad \text{--- (XIII)}$$

(14)

- I_{EP} will increase less rapidly with emitter base volt. at high biases.

- As we know for low level injection

$$I_C \propto I_{EP} \propto e^{qV_{EB}/2kT} \quad \text{--- (XIV)}$$

$$I_B \propto I_{EN} \propto e^{qV_{EB}/kT}$$

Hence for high V_{EB}

$$\beta = \frac{I_C}{I_B} \propto \frac{e^{qV_{EB}/2kT}}{e^{qV_{EB}/kT}} \propto e^{-qV_{EB}/2kT} \propto I_C^{-1}$$

Hence common emitter gain decreases at high injection due to excess majority carrier in base.

Base current injected in emitter can be given by

$$I_B \propto I_{en} \propto e^{qV_{EB}/kT} \quad \text{--- (XV)}$$

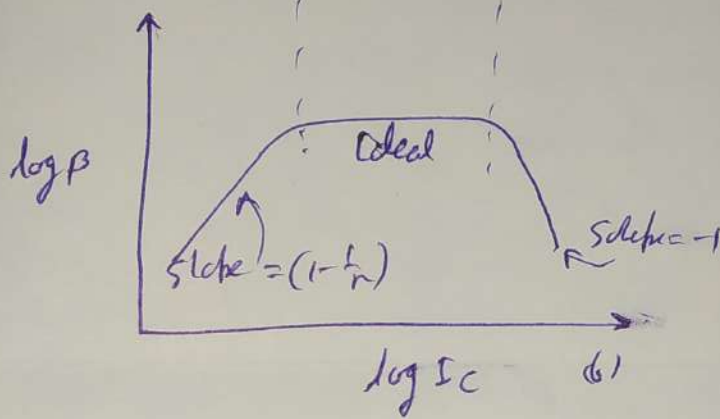
— Large emitter current injected into the base is not likely to be affected by the generation-recombination effects.

$$I_{Ep} \propto e^{qV_{EB}/kT} \quad \text{--- (XVI)}$$

— Hence for low V_{EB} or low I_C current gain is given

$$\beta = \frac{I_C}{I_B} \propto \frac{e^{qV_{EB}/kT}}{e^{qV_{EB}/kT}} \propto e^{qV_{EB}/kT(1-\frac{1}{n})} \propto I_C^{[1-\frac{1}{n}]} \quad \text{--- (XVII)}$$

— From the graph it can be concluded that at low injection levels β is degraded by poor emitter injection efficiency and at high current β decreases due to excess majority charge in base which degrades γ .



(a) \rightarrow Gummel plot of \log of collector and base currents as a function of emitter-base forward bias.

(b) \rightarrow DC current gain I_C/I_B .