

弹性力学 基础

$$\sigma_{eq} = \sqrt{\frac{2}{3} \sigma'_{ij} \sigma'_{ij}}$$

1. ()

$$[a, b, c] = \vec{a} \cdot (\vec{b} \times \vec{c})$$

$$\begin{vmatrix} \delta_{or} & \delta_{os} & \delta_{ot} \\ \delta_{pr} & \delta_{ps} & \delta_{pt} \\ \delta_{gr} & \delta_{gs} & \delta_{gt} \end{vmatrix} = e_{opq} e_{rst}$$

$$e_{ijk} e_{ist} = \delta_{js} \delta_{kt} - \delta_{jt} \delta_{ks}$$

$$e_{ijk} e_{ijkt} = 2 \delta_{kt}$$

$$e_{ijk} e_{ijkt} = 6$$

$$\beta_{ik} = \frac{\partial x_i}{\partial x_k}$$

$$T_{mn} = \beta_{mi} \beta_{nj} T_{ij}$$

$$T_{mn} = \beta_{im} \beta_{jn} T_{ij}$$

球形张量, 偏斜张量

$$(\lambda I - T) = 0$$

$$\Rightarrow \lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0$$

$$I_1 = \text{tr } T$$

I_2 : 所有二阶主子式之和
(顺序主子式)

$$I_3 = \det T$$

$$\frac{1}{2} (T_{ii} T_{jj} - T_{ij} T_{ji})$$

$$\int_V \nabla \cdot \vec{T} dV = \int_S \vec{n} \cdot \vec{T} ds$$

$$\int_S (\vec{n} \cdot \nabla \times \vec{T}) ds = \oint_L \vec{dl} \cdot \vec{T}$$

$$\text{名义应力: } \vec{\sigma} = \frac{\vec{F}}{A_0} \quad (\text{工程应力})$$

$$\text{柯西应力: } \vec{\sigma} = \frac{\vec{F}}{A}$$

正应力 ~ 切应力/剪应力

柯西公式/斜面应力公式:

$$\vec{\sigma}_{(n)} = \vec{n} \cdot \vec{\sigma}$$

应力分量转换公式/转轴公式

第二主应力: $\sigma_1 \geq \sigma_2 \geq \sigma_3$

(主根, 副主根对应作用平面)

任意正交方向

作正方向)

正应力之最大, 最小值: σ_1, σ_3

剪应力的六个极值: $\pm \frac{1}{2} (\sigma_1 - \sigma_2)$ 方向分别为对应主应力

T_{ii} $\pm \frac{1}{2} (\sigma_2 - \sigma_3)$ 方向之对角线

$\pm \frac{1}{2} (\sigma_1 - \sigma_3)$

八面体正应力 $\sigma_8 = \frac{1}{3} I_1$, 剪应力 $|T_8| = \frac{2}{3} \sqrt{T_{11}^2 + T_{22}^2 + T_{33}^2} = \frac{\sqrt{2}}{3} \sigma_{eq}$

$$\text{动量定理: } \rho \frac{d\vec{v}}{dt} = \vec{f} + \nabla \cdot \vec{\sigma} \xrightarrow{\text{平衡}} \vec{f} + \nabla \cdot \vec{\sigma} = 0$$

$$\text{柱坐标: } \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{\theta r}}{\partial \theta} + \frac{\partial \sigma_{zr}}{\partial z} + \frac{\sigma_{rr} - \sigma_{\theta\theta}}{r} + f_r = 0$$

$$\frac{\partial \sigma_{r\theta}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{\theta\theta}}{\partial \theta} + \frac{\partial \sigma_{z\theta}}{\partial z} + 2 \frac{\sigma_{r\theta}}{r} + f_\theta = 0$$

$$\frac{\partial \sigma_{rz}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{\theta z}}{\partial \theta} + \frac{\partial \sigma_{zz}}{\partial z} + \frac{\sigma_{rz}}{r} + f_z = 0$$

$$\epsilon_{ij} = \frac{1}{2} (q_{ij} - q_{ji}) \quad \text{工程剪应变 } \gamma_{ij} \quad \vec{F} = \frac{\partial x_i}{\partial x_j} \vec{e}_i \vec{e}_j$$

$$\text{格林应变张量 } E_{ij} = \frac{1}{2} \left(\frac{\partial x_m}{\partial x_i} \frac{\partial x_m}{\partial x_j} - \delta_{ij} \right)$$

$$\vec{E} = \frac{1}{2} (\vec{u} \cdot \nabla + \nabla \cdot \vec{u} + \vec{u} \cdot \nabla \cdot \vec{u}) \quad (\epsilon_{ij} \vec{e}_i \vec{e}_j)$$

$$\text{阿基曼西张量 } e_{ij} = \frac{1}{2} \left(\delta_{ij} - \frac{\partial u_m}{\partial x_i} \frac{\partial u_m}{\partial x_j} \right)$$

$$\vec{e} = \frac{1}{2} (\vec{u} \cdot \nabla + \nabla \cdot \vec{u} - \vec{u} \cdot \nabla \cdot \vec{u}) \quad (\epsilon_{ij} \vec{e}_i \vec{e}_j)$$

$$\text{对数应变: } \epsilon_{zz} = \ln \frac{L}{L_0}$$

$$\text{变形前后, 长度比 } \lambda = \frac{ds}{ds_0} = \sqrt{1 + 2 \vec{v} \cdot \vec{E} \cdot \vec{v}}$$

$$\text{单位向量 } \lambda \vec{v}' = (\vec{I} + \vec{u} \cdot \nabla) \cdot \vec{v} = \vec{F} \cdot \vec{v}$$

$$\text{夹角余弦: } \cos(\vec{v}', \vec{v}) = \frac{\vec{v} \cdot \vec{F} + 2 \vec{v} \cdot \vec{E} \cdot \vec{v}}{\lambda_v \lambda_t}$$

$$\text{线应变分量 } E_v = \frac{1}{2} (\lambda_v^2 - 1) \quad (\epsilon_{ii} = \epsilon_{ii})$$

$$\text{切应变分量 } E_{vt} = \frac{1}{2} (\lambda_v \vec{v}') \cdot (\lambda_t \vec{t}') \quad (\epsilon_{ij} = \frac{\pi}{4} - \frac{\psi_{ij}}{2})$$

小变形: $\lambda_v = \vec{v} \cdot \vec{E} \cdot \vec{v} =: \epsilon_v$, 工程正应变,

$$\gamma_{vt} = 2 \vec{v} \cdot \vec{E} \cdot \vec{t} = 2 \epsilon_{vt} \quad \text{工程剪应变}$$

$$\text{转动张量 } \vec{\Omega} = \frac{1}{2} (\nabla \vec{u} - \vec{u} \cdot \nabla) \quad \Omega_{ij} = e_{ijk} \omega_k$$

$$(-\vec{\Omega} \cdot d\vec{x} = \vec{\omega} \times d\vec{x})$$

$$\text{协调方程: } \epsilon_{ijkl} + \epsilon_{kl,ij} - \epsilon_{ik,jl} - \epsilon_{jl,ik} = 0$$

$$L_{mn} = e_{nijl} e_{mjkl} \epsilon_{ij,kl} = 0$$

$$\nabla \times \vec{\epsilon} \times \nabla = 0$$

$$\text{二维: } \epsilon_{\alpha\alpha, \beta\beta} - \epsilon_{\alpha\beta, \alpha\beta} = 0$$

单连通域上, 位移单值 \Leftrightarrow 满足协调方程

由应变求位移

直接积分: eg: $d(\frac{\partial u_1}{\partial x_2}) = \frac{\partial \epsilon_{11}}{\partial x_2} dx_1 + (\frac{\partial \gamma_{12}}{\partial x_2} - \frac{\partial \epsilon_{22}}{\partial x_1}) dx_2$
 $+ \frac{1}{2} (\frac{\partial \gamma_{12}}{\partial x_3} + \frac{\partial \gamma_{31}}{\partial x_2} - \frac{\partial \gamma_{23}}{\partial x_1}) dx_3$

有直接积分: eg: $\epsilon_{ij} = 0$

$\Rightarrow u_1 = f_1(x_2, x_3) \quad u_2 = f_2(x_1, x_3) \quad u_3 = f_3(x_1, x_2)$

$\Rightarrow \gamma_{12} = \dots = 0 \quad \gamma_{13} = \dots = 0 \quad \gamma_{23} = \dots = 0 \quad \sim \square$

观察应变, 再求导, 得: $\frac{\partial^2 f_1}{\partial x_2^2} = 0, \frac{\partial^2 f_1}{\partial x_3^2} = 0$

$\Rightarrow f_1 = g_1(x_3)x_2 + g_0(x_3) \Rightarrow g_1'' \cdot x_2 + g_0'' = 0$

$\Rightarrow g_1'' = g_0'' = 0 \Rightarrow f_1 = a_0 + a_1 x_2 + a_2 x_3 + a_3 x_2 x_3$

同理可得 u_2, u_3 .

设 $f_1 = a_0 + b_1 x_2 + a_2 x_3, \dots$

柱坐标系: $\epsilon_r = \frac{\partial u}{\partial r} \quad \epsilon_\theta = \frac{1}{r} \frac{\partial u}{\partial \theta} + \frac{u}{r} \quad \epsilon_z = \frac{\partial w}{\partial z}$

$\gamma_{r\theta} = \frac{\partial u}{\partial \theta} - \frac{u}{r} + \frac{\partial v}{\partial r} \quad \gamma_{\theta z} = \frac{\partial v}{\partial z} + \frac{1}{r} \frac{\partial w}{\partial \theta}$

$\gamma_{rz} = \frac{\partial w}{\partial r} + \frac{\partial u}{\partial z}$

协调方程: $\frac{\partial^2 \epsilon_r}{\partial r^2} - \frac{\partial \epsilon_r}{r \partial r} + \frac{1}{r} \frac{\partial^2}{\partial r^2} (r \epsilon_\theta) - \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} (r \gamma_{r\theta}) = 0$

$\epsilon_{ij} = \frac{1}{E} (\sigma_{ij} - \nu \sigma_{kk}) \quad \gamma_{ij} = \frac{1}{G} \tau_{ij} \quad G = \frac{E}{2(1+\nu)}$

$\epsilon_{ij} = \frac{1+\nu}{E} \sigma_{ij} - \frac{\nu}{E} \sigma_{kk} \delta_{ij}$

令 $\theta = \epsilon_{kk}, \quad \Phi = \sigma_{kk}, \quad k = \frac{E}{3(1-2\nu)}, \quad \lambda = \frac{E\nu}{(1+\nu)(1-2\nu)}$

$\theta = \frac{1}{3k} \Phi = \frac{1}{k} \sigma_0 \quad (\text{体积应变 vs 静水应力})$

令 $\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}, \quad \lambda, G, \text{第一、二、三}$

$\sigma_{ij} = 2G \epsilon_{ij} + \lambda \epsilon_{kk} \delta_{ij}$

梅蒂数

$\Rightarrow \sigma_0 = k\theta = 3k\epsilon_0 \quad \sigma'_{ij} = 2G \epsilon'_{ij}$

一般形式: $\sigma_{ij} = C_{ijkl} \epsilon_{kl} \quad W = \frac{1}{2} C_{ijkl} \epsilon_{ij} \epsilon_{kl} = \frac{1}{2} \sigma_{ij} \epsilon_{ij}$

应变能 $W = \int_0^{\tilde{\epsilon}} \tilde{\sigma} : d\tilde{\epsilon} \quad \frac{\partial W}{\partial \epsilon_{ij}} = \sigma_{ij}$

\Rightarrow 广义胡克公式 $\frac{\partial \sigma_{kk}}{\partial \epsilon_{ij}} = \frac{\partial \sigma_{kl}}{\partial \epsilon_{ij}}$

应变余能 $W_c = \int_0^{\tilde{\sigma}} \tilde{\epsilon} : d\tilde{\sigma} \quad \frac{\partial W_c}{\partial \sigma_{ij}} = \epsilon_{ij} \Rightarrow \frac{\partial \epsilon_{kl}}{\partial \sigma_{ij}} = \frac{\partial \epsilon_{ij}}{\partial \sigma_{kl}}$

$\hookrightarrow W + W_c = \sigma_{ij} \epsilon_{ij}$

正定性: $W = \frac{1}{2} k \epsilon_{ii}^2 + G \epsilon'_{ij} \epsilon'_{ij}$

$W_c = \frac{1}{2k} \sigma_{ii}^2 + \frac{1}{2 \cdot 2G} \sigma'_{ij} \sigma'_{ij}$

$W_c = \frac{1}{2k} \sigma_{ii}^2 + \frac{1}{2 \cdot 2G} \sigma'_{ij} \sigma'_{ij}$

Lame - Navier 方程: $G u_{ij,jj} + (\lambda + G) u_{j,j} i + f_i = 0 \quad (p_{ii})$

Beltrami - Michell 方程:

$\nabla^2 \sigma_{ij} + \frac{1}{1+\nu} \Phi_{,ij} = -(\bar{f}_{ij} + \bar{f}_{ji}) - \frac{\nu}{1-\nu} f_{k,k} \delta_{ij}$

圣维南定理