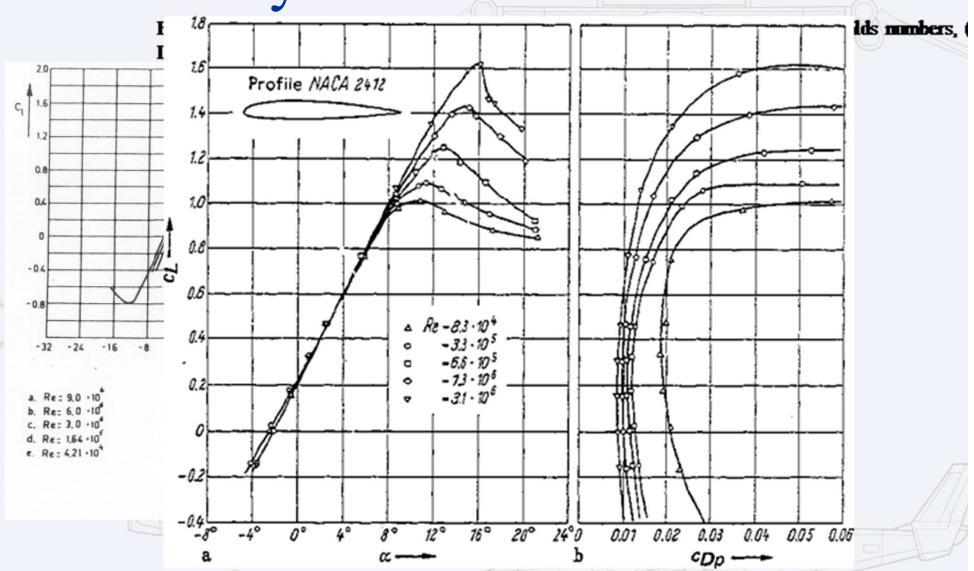


Steady flow around an airfoil



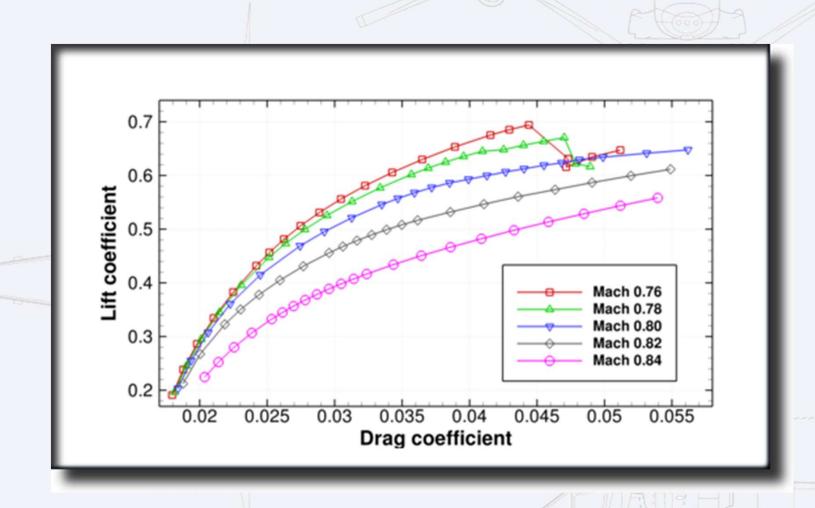


Reynolds influence





MACH influence



Slide 4



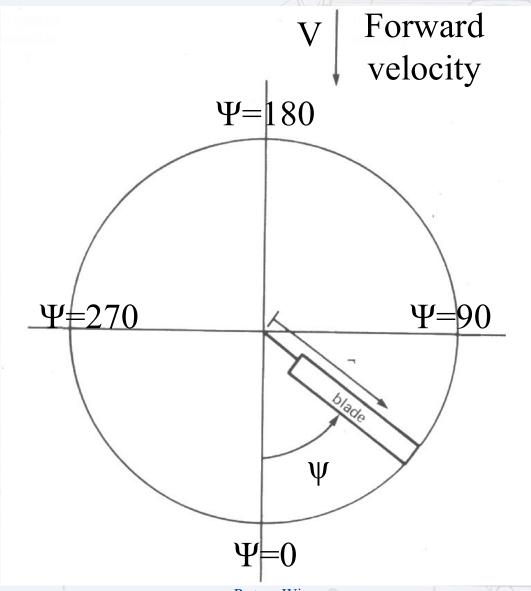
Rotary wing

- •The rotor provides 3 basic functions
 - Generation of lift
 - Generation of trust
 - Control the attitude and position



Azimuth Angle

Retreating Side

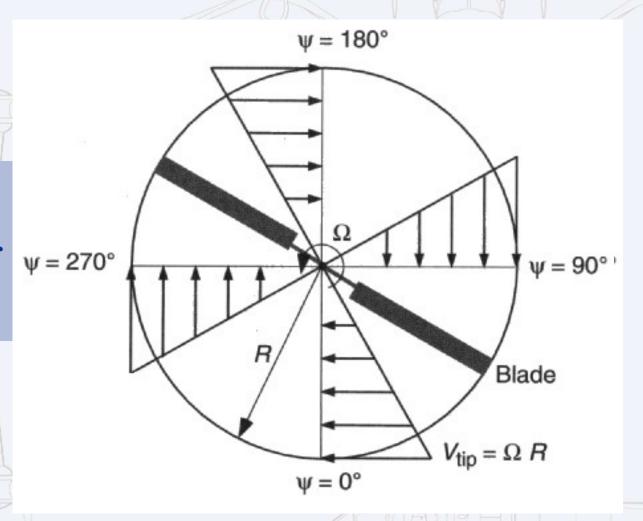


Advancing Side



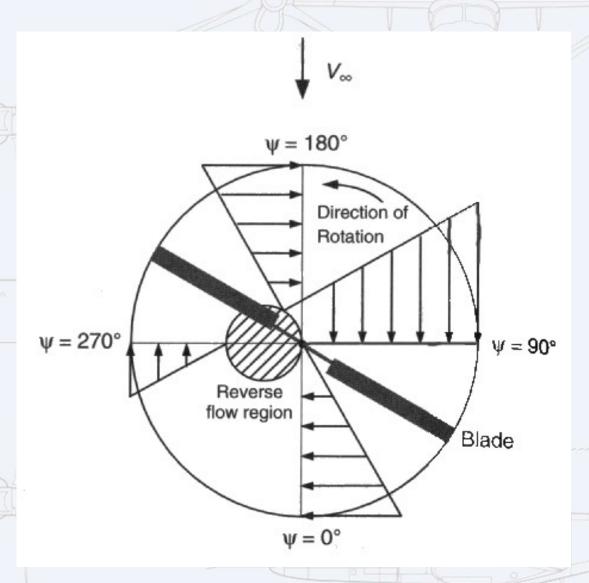
Velocity in hover

Velocity
Perpendicular
to the blade





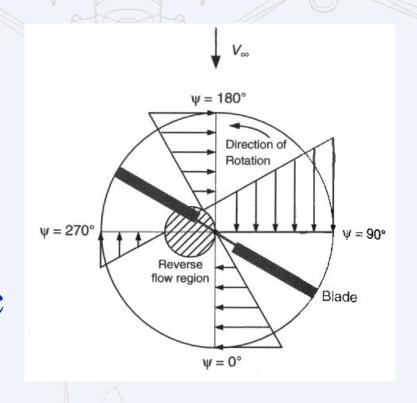
Velocity in Forward Flight





• Velocity distribution is not uniform

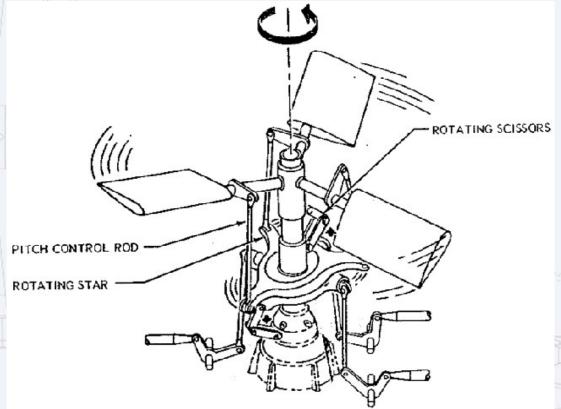
Generation of lift is asymmetric





To compensate de asymmetry

• It's necessary to change the blade AOA (lower on the advancing side, higher on the retreating side)

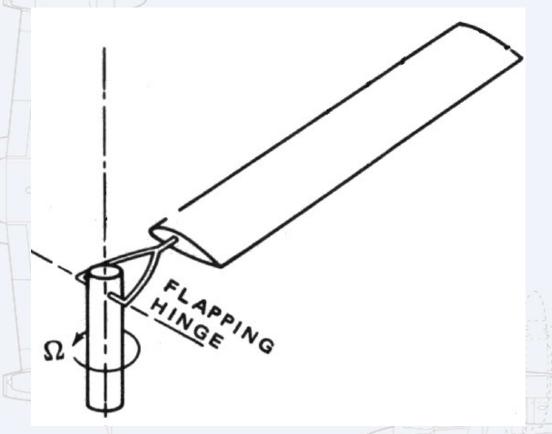






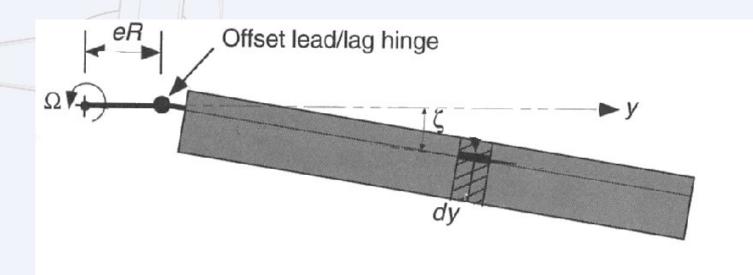


• To avoid high moments due to the blade on the rotor hub the blade is allowed to flap



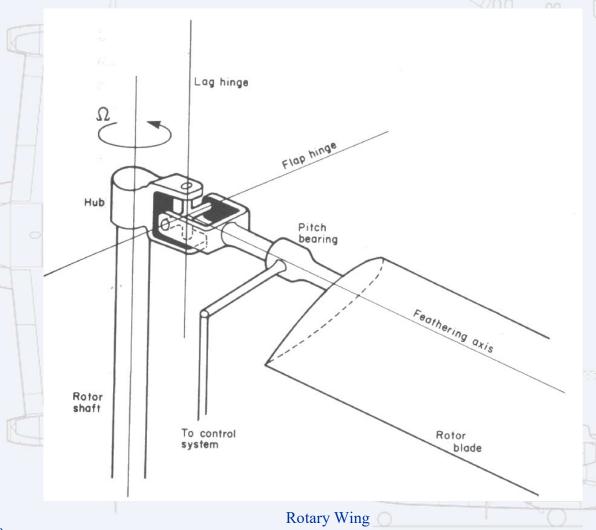


• To allow the blades movement at constant velocity the blade is allowed to lag.



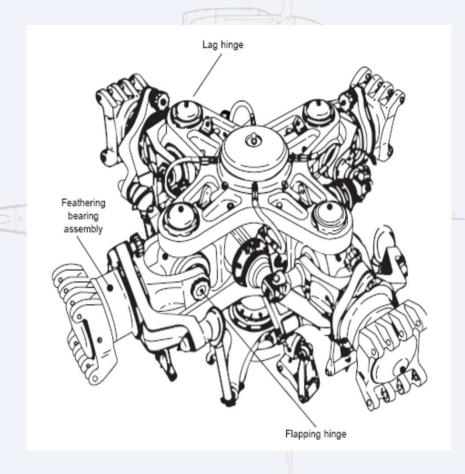


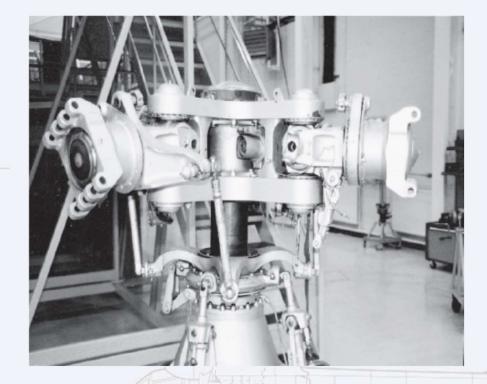
Rotor





The rotor Hub

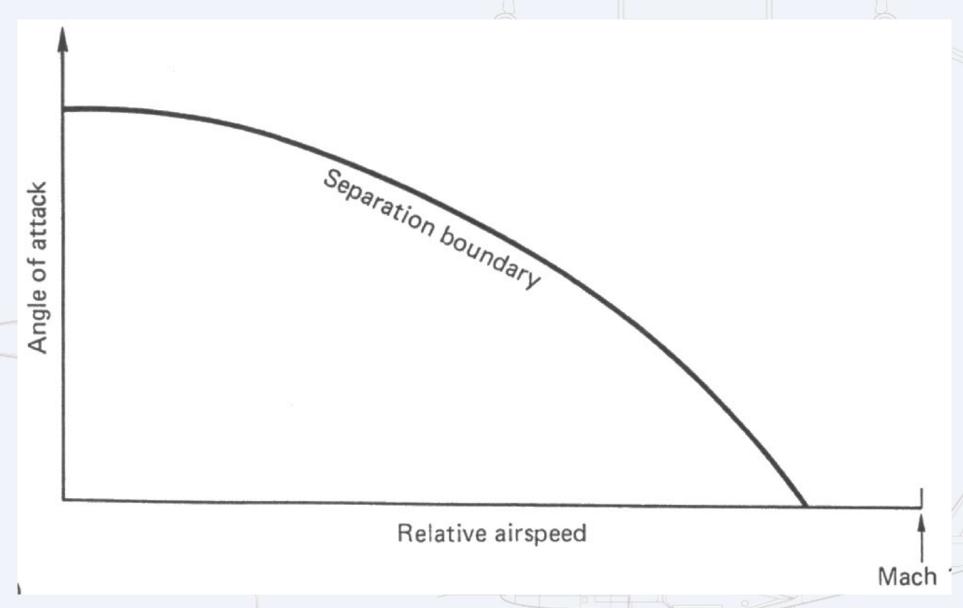








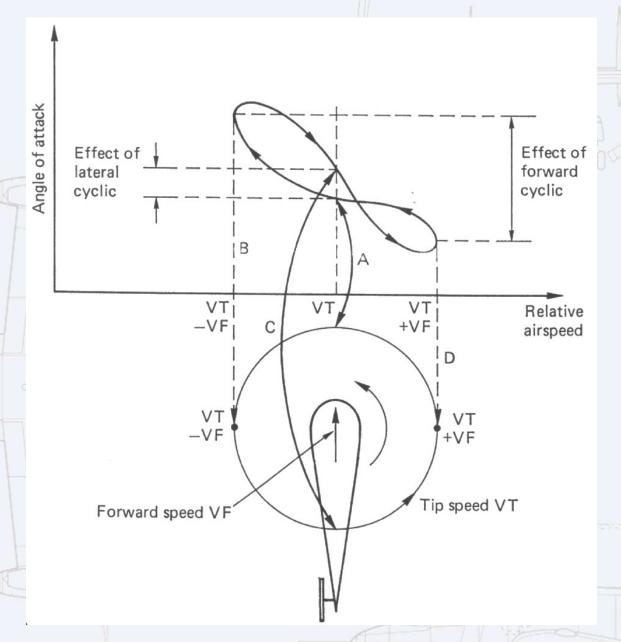




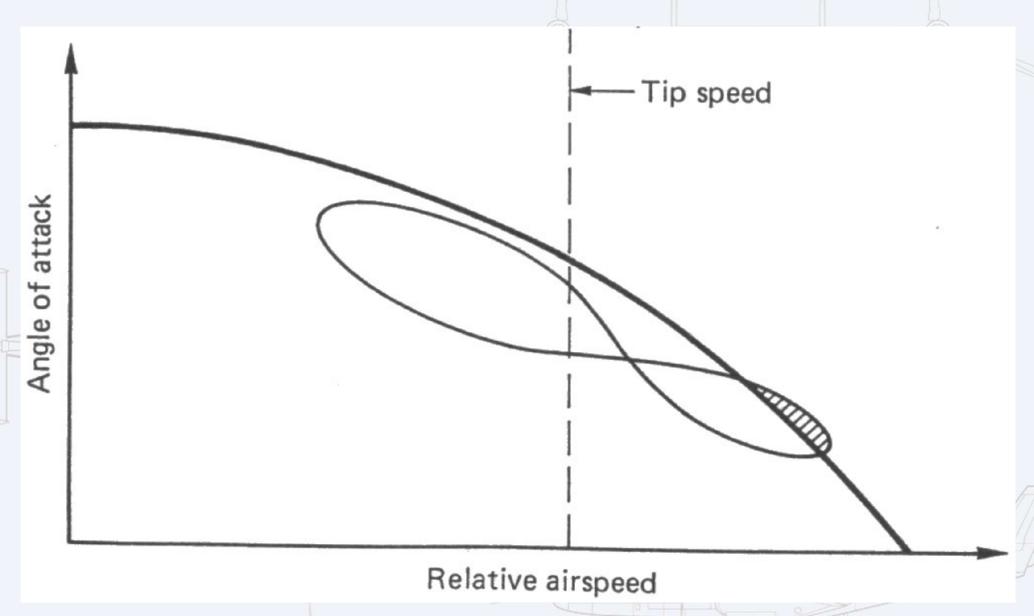
Rotary Wing

Slide 17



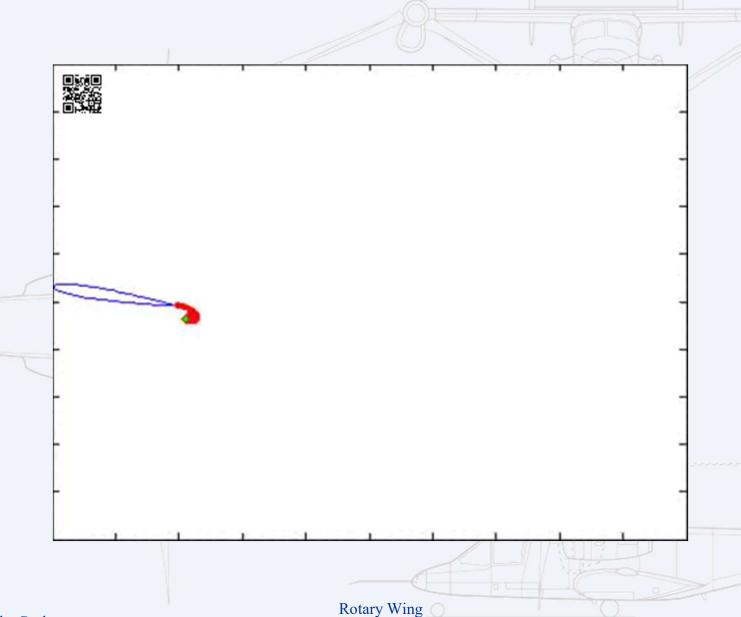








Airfoil Behaviour



NCAS / Filipe Szolnoky Cunha

Slide 20

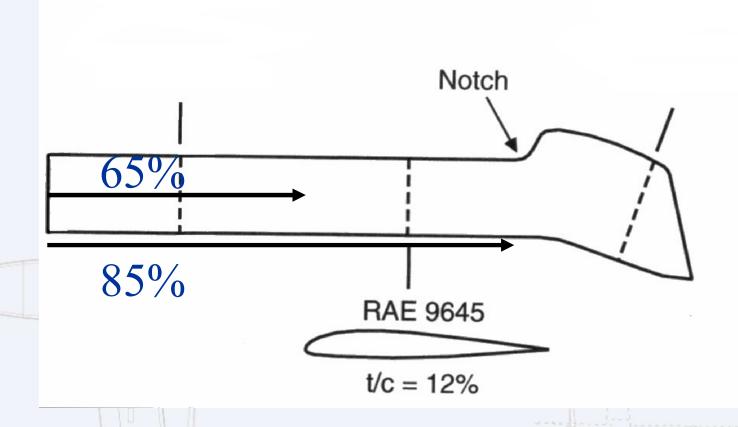


- During each revolution the blade encounters a wide variety of operating conditions.
 - Transonic flow on the advancing region
 - Blade must be thin enough to maximize the M_{dd} .
 - High AOA on the retreating side
 - Minimum thickness and incorporate some camber to give a relatively high C_{lmax} .
- No single airfoil will meet all requirements. The balance is between:
 - High M_{dd} .
 - High C_{lmax} at low M.



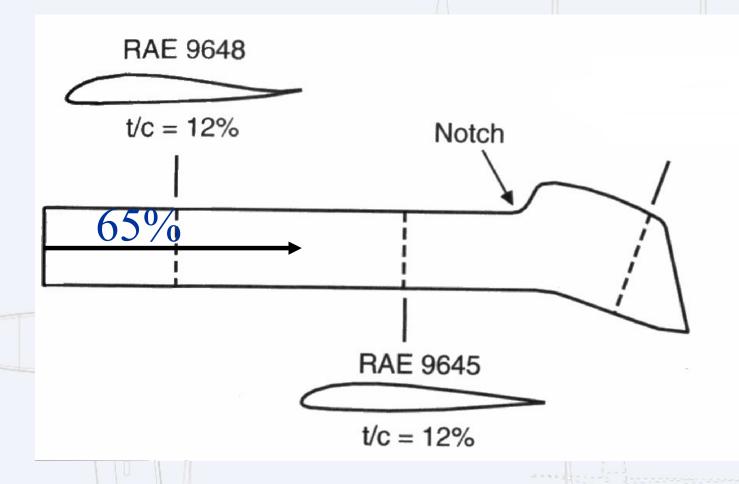
- British Experimental Rotor Program (BERP) rotor, was design to meet the conflicting requirements of the advancing and retreating blade conditions.
- The research paid off in 1986 when a Westland Lynx helicopter attained the absolute speed record for a conventional helicopter with a speed of 400.87 km/h





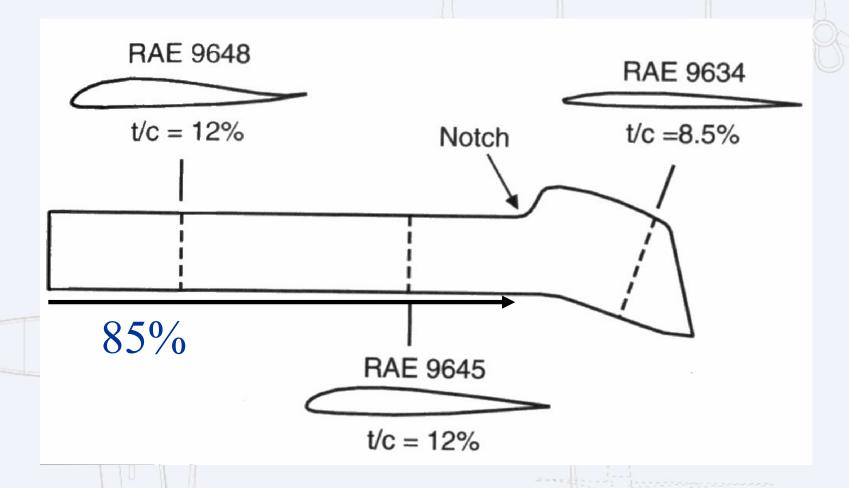
- High C_{lmax} (1.55).
- High pitching moments





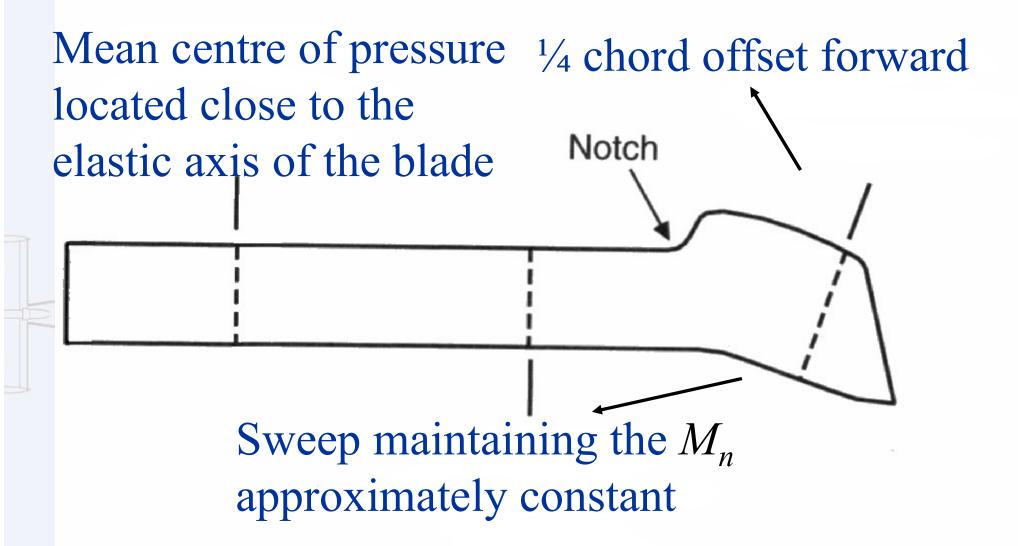
- Lower C_{lmax} .
- Offset of the high pitching moments (RAE 9645)



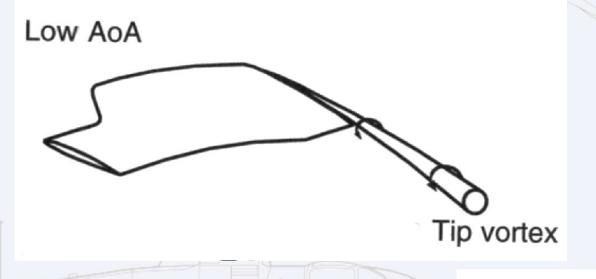


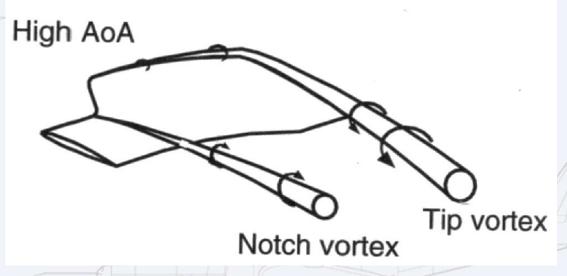
- Low t/c ratio (high M_{dd})
- Cambered to give weak shock wave and low pitching moments













Momentum Theory

• Let's simplify our first approach and develop a simple method capable of predicting the rotor thrust and power

Momentum Theory

• First developed by Rankine (1895) for marine propellers and developed further and generalized by several other authors

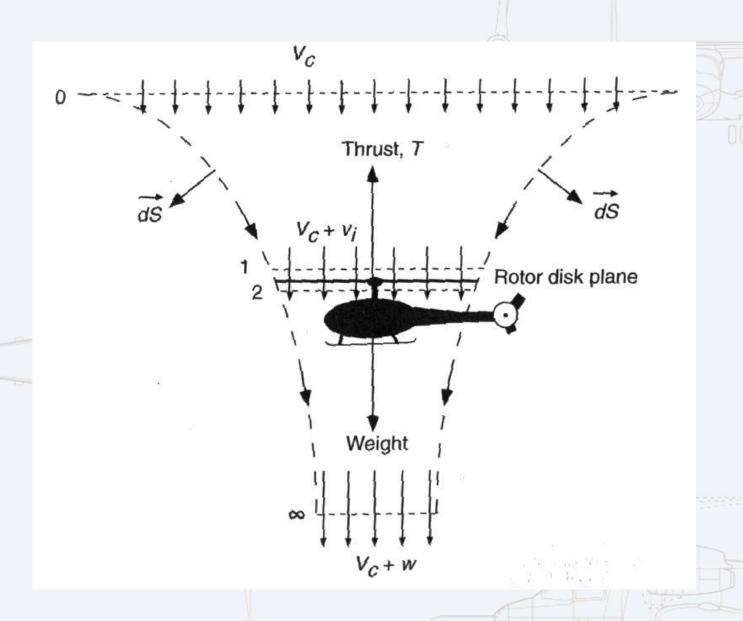


Assumptions

- Conditions in hover:
 - No forward speed
 - No vertical speed
 - The flow field is axisymetrical
 - There is a wake boundary with the flow outside this boundary being quiescent
 - The flow velocities inside this boundary can be quite high



Representation and notation





Ideal Power Hover

• Power consumed=Energy rate flow out-Energy rate flow in and the following expression can be obtained

$$P = T \sqrt{\frac{T}{2 A \rho}}$$

Or in terms of the induced velocity:

$$P = Tv_h = (2\rho A v_h^2) v_h = 2\rho A v_h^3$$



Non Ideal effects Hover

• The actual rotor power can then be expressed as:

$$C_{P_a} = C_{P_i} + C_{P_0} = \kappa \frac{C_T^{\frac{3}{2}}}{\sqrt{2}} + \frac{1}{8} \sigma C_{d0}$$

Where to profile power coefficient is

obtained:
$$P_0 = \frac{1}{8} \rho N_b \Omega^3 c C_{d0} R^4$$

$$C_{P_0} = \frac{1}{8} \left(\frac{N_b cR}{A} \right) C_{d0} = \frac{1}{8} \sigma C_{d0}$$



Power required for axial flight

• We can write the power ratio as:

$$\frac{P}{P_h} = \frac{V_c + v_i}{v_h} = \frac{V_c}{v_h} + \frac{v_i}{v_h}$$

• Using the previous equations obtained for climbing

$$\frac{P}{P_h} = \frac{V_c}{v_h} - \left(\frac{V_c}{2v_h}\right) + \sqrt{\left(\frac{V_c}{2v_h}\right)^2 + 1} = \left(\frac{V_c}{2v_h}\right) + \sqrt{\left(\frac{V_c}{2v_h}\right)^2 + 1}$$



Power required for axial flight

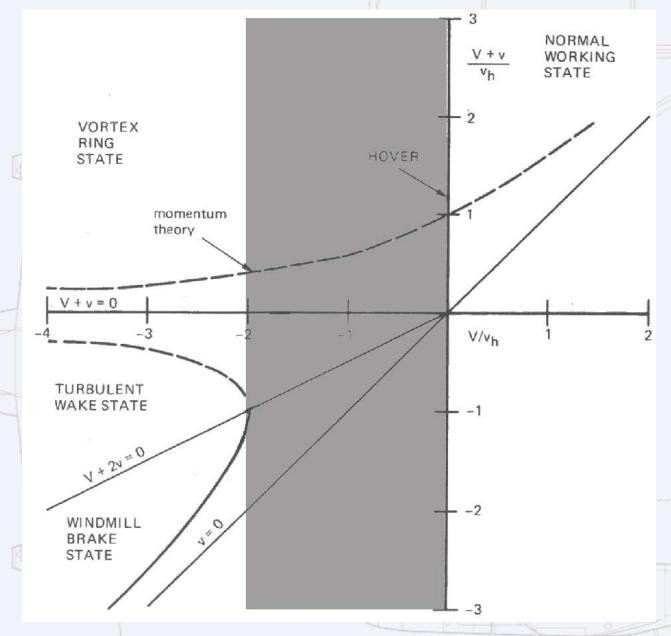
Using the previous equations obtained for descent

$$\frac{P}{P_h} = \frac{V_c}{2v_h} - \sqrt{\left(\frac{V_c}{2v_h}\right)^2 - 1}$$

Rotary Wing Slide 34

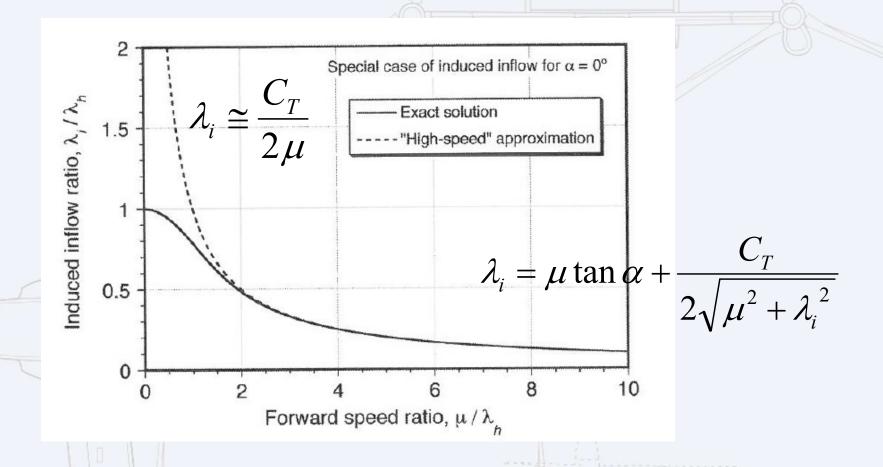


Power required for axial flight





Variation of Non-Dimensional Inflow with Advance Ratio



• Notice that inflow velocity rapidly decreases with advance ratio



Power Coefficient

$$C_{P_{\text{Uncorrected}}} = C_T \lambda_i + \frac{1}{2} \frac{f}{A} \mu^3 + \frac{\sigma C_{d0}}{8} \left[1 + 3\mu^2 \right]$$
Induced power Parasite Power

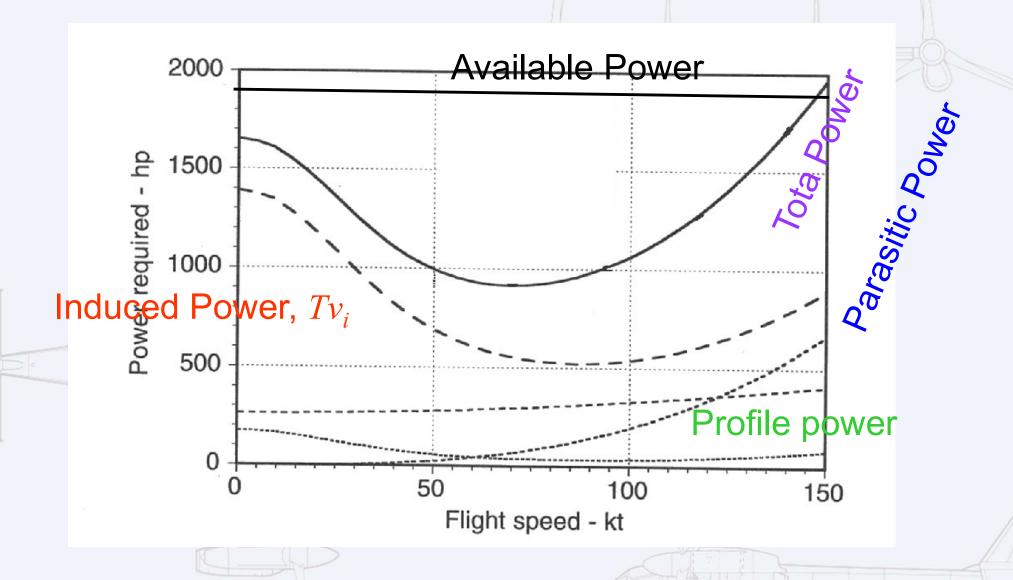
$$C_{P_{\text{Corrected}}} = \kappa C_T \lambda_i + \frac{1}{2} \frac{f}{A} \mu^3 + \frac{\sigma C_{d0}}{8} [1 + 4.6 \mu^2]$$

 C_D is the vehicle parasite drag coefficient and S the reference area. Because there is no agreement on a common reference area it is customary to supply the product $C_DS=f$ equivalent flat plate area

Rotary Wing



Power in Forward Flight





Steady aerodynamics

- What is the induced power in hover for a rotary wing aircraft with 2200kg, 11m diameter at sea level, using momentum theory?
- And with a forward speed of 15m/s?

Rotary Wing



Blade Element Theory

- The momentum theory does not take into account
 - Number of blades
 - Airfoil characteristics (lift, drag, angle of zero lift)
 - Blade planform (taper, sweep, root cut-out)
 - Blade twist distribution
 - Compressibility effects



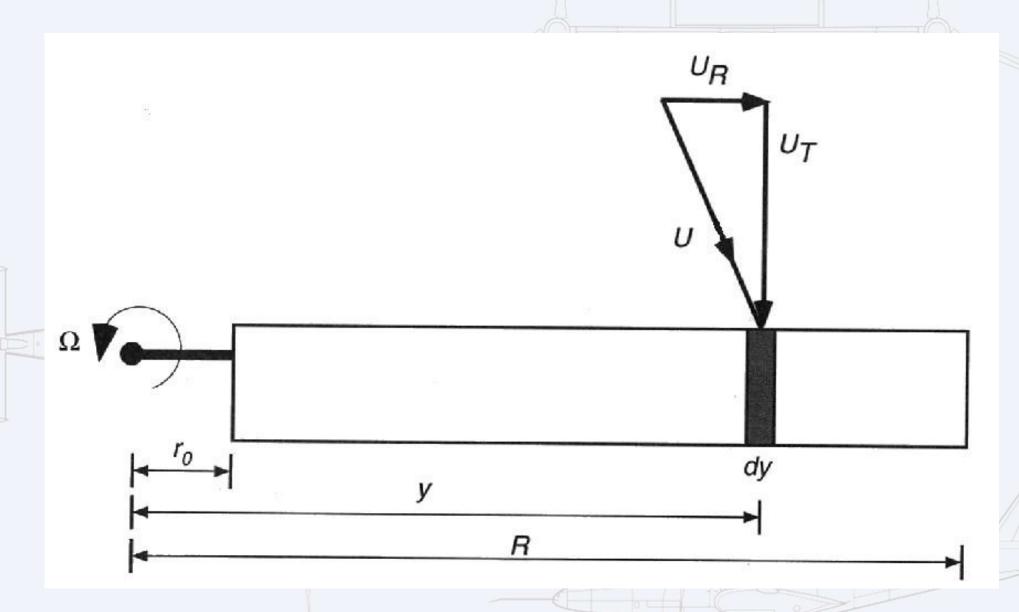
Blade Element Theory

- BET assumes that each blade section acts as a two-dimensional airfoil to produce aerodynamic forces
- The blade is then divided in non-interacting sections where all the computations are performed using 2-D aerodynamics
- An integration over the blade length gives the total thrust and total power

Slide 41

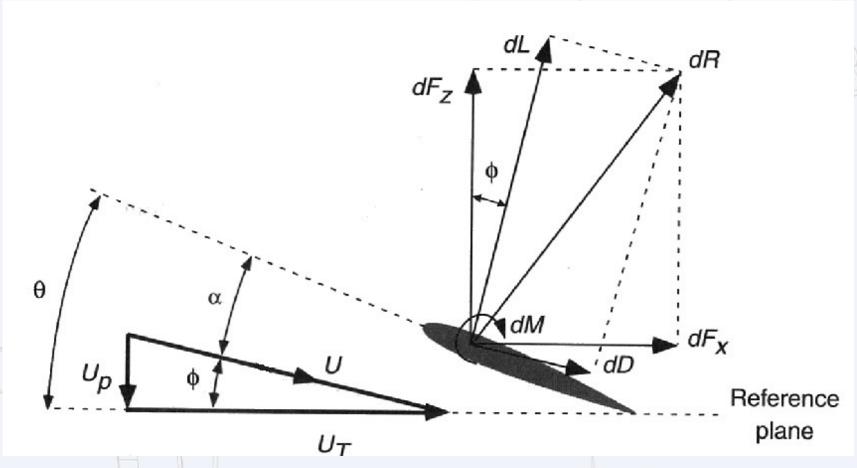


BET Model





BET Model



- The in plane Velocity $U_T = \Omega y$
- The out of plane Velocity $U_P = V_C + v_i$
- Therefore the total velocity is $U = \sqrt{U_T^2 + U_P^2}$



BET model

• The incremental lift per unit span:

$$dL = \frac{1}{2} \rho U^2 c C_l dy$$

• The incremental drag per unit span:

$$dD = \frac{1}{2}\rho U^2 cC_d dy$$

• Or in quantities parallel and perpendicular to the rotor disk plane:

$$\begin{cases} dF_z = dL \cos \phi - dD \sin \phi \\ dF_x = dL \sin \phi + dD \cos \phi \end{cases}$$



BET model

• We can then calculate the Thrust:

$$dT = N_b dF_z$$

• The Torque

$$dQ = N_b dF_x y$$

• The Power

$$dP = N_b dF_x \Omega y$$

• Remember N_b is the number of blades



BET model

• And we can relate all three with C_l and C_d

$$\begin{cases} dT = N_b (dL \cos \phi - dD \sin \phi) \\ dQ = N_b (dL \sin \phi + dD \cos \phi) y \\ dP = N_b (dL \sin \phi + dD \cos \phi) \Omega y \end{cases}$$



BET model assumptions

• The following assumptions are valid within the helicopter aerodynamics

$$U_{T} \gg U_{P} \Rightarrow U = \sqrt{U_{P}^{2} + U_{T}^{2}} \approx U_{T}$$

$$\phi = \tan^{-1}(U_{P}/U_{T}) \approx U_{P}/U_{T}$$

$$\phi \approx 0 \Rightarrow \begin{cases} \sin \phi = \phi \\ \cos \phi = 1 \end{cases}$$

$$dD \ll dL \implies dD \sin \phi \approx dD \phi \approx 0$$



Basic Equations

• The expression for Thrust, Torque and Power are:

$$\begin{cases} dT = N_b (dL \cos \phi - dD \sin \phi) = N_b (dL) \\ dQ = N_b (dL \sin \phi + dD \cos \phi) y = N_b (dL \phi + dD) y \\ dP = N_b (dL \sin \phi + dD \cos \phi) \Omega y = N_b (dL \phi + dD) \Omega y \end{cases}$$

• Let's now nondimensionalize using for length R and for speed $V_{tip} = \Omega R$



Nondimensional form

- r=y/R
- $U_T/\Omega R = \Omega y/\Omega R = y/R = r$
- And the thrust, torque and power coefficients already defined:

$$dC_T = \frac{dT}{\rho A(\Omega R)^2}, dC_Q = \frac{dQ}{\rho A(\Omega R)^2 R}, dC_P = \frac{dP}{\rho A(\Omega R)^3}$$

Now the inflow ratio is

$$\lambda = \frac{V_c + v_i}{\Omega R} = \frac{V_c + v_i}{\Omega y} \left(\frac{\Omega y}{\Omega R}\right) = \frac{U_P}{U_T} \left(\frac{y}{R}\right) = \phi r$$

Rotary Wing



Thrust coefficient (incremental)

• Substituting the previous equations in the Thrust coefficient equation:

$$dC_T = \frac{N_b dL}{\rho A(\Omega R)^2} = \frac{N_b \left(\frac{1}{2} \rho U_T^2 c C_l dy\right)}{\rho A(\Omega R)^2}$$

$$= \frac{1}{2} \left(\frac{N_b c}{\pi R} \right) C_l \left(\frac{y}{R} \right)^2 d \left(\frac{y}{R} \right) = \frac{1}{2} \sigma C_l r^2 dr$$



Power coefficient (incremental)

• Using the same analysis for the Power coefficient

$$dC_{P} = dC_{Q} = \frac{dQ}{\rho A(\Omega R)^{2} R} = \frac{N_{b}(\phi dL + dD)y}{\rho A(\Omega R)^{2} R}$$

$$= \frac{1}{2} \sigma(\phi C_{l} + C_{D}) \left(\frac{y}{R}\right)^{3} d\left(\frac{y}{R}\right)$$

$$= \frac{1}{2} \sigma(\phi C_{l} + C_{D}) r^{3} dr$$



Total Thrust and Power

• To find the total blade contribution for Thrust and power we have on integrate between the root and tip of the blade

$$C_{T} = \frac{1}{2} \int_{0}^{1} \sigma C_{l} r^{2} dr = \frac{1}{2} \sigma \int_{0}^{1} C_{l} r^{2} dr$$

- If the blade is rectangular c=const
- For the torque and power coefficient

$$C_{Q} = C_{P} = \frac{1}{2}\sigma \int_{0}^{1} (\phi C_{l} + C_{d})r^{3}dr = \frac{1}{2}\sigma \int_{0}^{1} (\lambda C_{l}r^{2} + C_{d}r^{3})dr$$



Total Thrust and Power

- To evaluate the previous expressions, we need:
- Inflow ratio $\lambda = \lambda(r)$
- Lift coefficient $C_l = C_l(\alpha, Re, M)$
- Drag coefficient $C_d = C_d(\alpha, Re, M)$
- AOA $\alpha = \alpha(V_C, \theta, v_i)$
- Induced Velocity $v_i = v_i(r)$

Numerical Solution needed



Approximations

- With certain assumptions and approximations it is possible to find closed form analytical solutions.
- The solutions are important because they serve to illustrate the fundamental form of the results in term of operational and geometric parameters of the rotor
- Let's the assume a rectangular blade c=const. From the definition $\sigma=const$. too.



Thrust approximation

• From the Steady linearized aerodynamics:

$$C_{l} = C_{l_{\alpha}}(\alpha - \alpha_{0}) = C_{l_{\alpha}}(\theta - \phi - \alpha_{0})$$

- We can consider C_l constant without serious loss of accuracy
- Let's also assume symmetric airfoils $\alpha_0 = 0$
- We can then write:

$$C_{T} = \frac{1}{2} \int_{0}^{1} \sigma C_{l} r^{2} dr = \frac{1}{2} \sigma C_{l\alpha} \int_{0}^{1} (\theta - \phi) r^{2} dr$$

$$C_{T} = \frac{1}{2} \sigma C_{l\alpha} \int_{0}^{1} (\theta r^{2} - \lambda r) dr$$

Rotary Wing



Untwisted Blades

- For a blade with zero twist $\theta = const. = \theta_0$.
- Let's also assume uniform inflow velocity, as assumed in the momentum theory $\lambda = const.$
- The Thrust coefficient can be written as:

$$C_T = \frac{1}{2} \sigma C_{l\alpha} \int_0^1 \left(\theta_0 r^2 - \lambda r \right) dr = \frac{1}{2} \sigma C_{l\alpha} \left[\theta_0 \frac{r^3}{3} - \lambda \frac{r^2}{2} \right]_0^1$$

$$C_T = \frac{1}{2} \sigma C_{l\alpha} \left[\frac{\theta_0}{3} - \frac{\lambda}{2} \right]$$



Power approximations

• We have seen that the incremental power coefficient (that is equal to the torque coefficient):

$$dC_{P} = \frac{1}{2}\sigma(\phi C_{l} + C_{d})r^{3}dr = \frac{1}{2}\sigma(\lambda C_{l}r^{2} + C_{d}r^{3})dr =$$

$$= \frac{1}{2}\sigma\lambda C_{l}r^{2}dr + \frac{1}{2}\sigma C_{d}r^{3}dr =$$

$$= dC_{P_{i}} + dC_{P_{0}}$$

Remembering that

$$dC_{P_i} = \lambda dC_T \Rightarrow dC_P = \lambda dC_T + dC_{P_0}$$



Power approximations

• Therefore the total power:

$$C_{P} = \int_{r=0}^{r=1} \lambda dC_{T} + \int_{0}^{1} \frac{1}{2} \sigma C_{d} r^{3} dr = \lambda C_{T} + \frac{1}{8} \sigma C_{d_{0}}$$

- Assuming uniform inflow and $C_d = C_{d0} = const.$
- Using once more the inflow expression obtained in hover: $C_T^{\frac{3}{2}}$

$$C_P = \frac{C_T^{\frac{3}{2}}}{\sqrt{2}} + \frac{1}{8}\sigma C_{d_0}$$

• Expression already obtained in the momentum theory



Example

- Given the inflow distribution over a rotor with rectangular untwisted blades is approximately triangular, and assuming no tip-loss
 - Compute the variation with radius of inflow angle and section angle of attack for such rotor.
 - Derive the relation between the blade element lift coefficient and C_T/σ
 - Derive an expression for the hovering power of such a rotor in terms of C_T/σ and C_{d0}