Dressed Graphene Model

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This thesis concerned with a specific model. Idea: Graphene with an added orbital on one of the lattice site with a low hopping, as to provide a flat band. I will call this model dressed Graphene from here on. This chapter reviews the lattice structure in section 1.1.

1.1 Lattice Structure

There exist a few different ways to define the lattice structure of Graphene which are all equivalent, but intermediate steps in calculating tight-binding models look different depending on the definition. This review on follows ref. [1].

Monolayer graphene forms a honeycomb lattice, which is a hexagonal Bravais lattice with a two atom basis, as can be seen in fig. 1.1a. The primitive lattice vectors of the hexagonal lattice are:

$$\mathbf{a}_1 = \frac{a}{2} \begin{pmatrix} 1 \\ \sqrt{3} \end{pmatrix}, \ \mathbf{a}_2 = \frac{a}{2} \begin{pmatrix} 1 \\ -\sqrt{3} \end{pmatrix}$$
 (1.1)

with lattice constant $a = \sqrt{3}a_0 \approx 2.46$ Å, using the nearest-neighbour distance a_0 . The vectors to the nearest-neighbor atoms B_i (i = 1, 2, 3,) from atom A are

Labels on vectors

$$\delta_{AB,1} = \begin{pmatrix} 0 \\ \frac{a}{\sqrt{3}} \end{pmatrix}, \ \delta_{AB,2} = \begin{pmatrix} \frac{a}{2} \\ -\frac{a}{2\sqrt{3}} \end{pmatrix}, \ \delta_{AB,3} = \begin{pmatrix} -\frac{a}{2} \\ -\frac{a}{2\sqrt{3}} \end{pmatrix}$$
(1.2)

and the vectors to the nearest-neighbor atoms A_i (i = 1, 2, 3,) from atom B are

$$\delta_{BA,1} = \begin{pmatrix} 0 \\ -\frac{a}{\sqrt{3}} \end{pmatrix}, \ \delta_{BA,2} = \begin{pmatrix} \frac{a}{2} \\ \frac{a}{2\sqrt{3}} \end{pmatrix}, \ \delta_{BA,3} = \begin{pmatrix} -\frac{a}{2} \\ \frac{a^2}{2\sqrt{3}} \end{pmatrix}. \tag{1.3}$$

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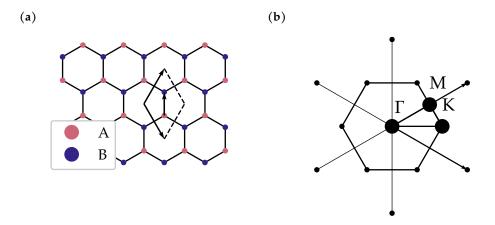


Figure 1.1: (a) Graphene lattice structure and (b) Brilluoin zone

The vectors between the Graphene A atom and the six neighbours on the same sub lattice are:

$$\delta_{AA,1} = \begin{pmatrix} 1 \\ \sqrt{3} \end{pmatrix}, \quad \delta_{AA,2} = a \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \delta_{AA,3} = a \begin{pmatrix} \frac{1}{2} \\ -\frac{\sqrt{3}}{2} \end{pmatrix}, \quad (1.4)$$

$$\delta_{AA,4} = a \begin{pmatrix} -\frac{1}{2} \\ -\frac{\sqrt{3}}{2} \end{pmatrix}, \, \delta_{AA,5} = a \begin{pmatrix} -1 \\ 0 \end{pmatrix}, \, \delta_{AA,6} = a \begin{pmatrix} -\frac{1}{2} \\ \frac{\sqrt{3}}{2} \end{pmatrix}$$
 (1.5)

The primitive reciprocal lattice vectors \mathbf{b}_1 , \mathbf{b}_2 fulfill

labels on vectors

$$\mathbf{a}_1 \cdot \mathbf{b}_1 = \mathbf{a}_2 \cdot \mathbf{b}_2 = 2\pi \tag{1.6}$$

$$\mathbf{a}_1 \cdot \mathbf{b}_2 = \mathbf{a}_2 \cdot \mathbf{b}_1 = 0 \,, \tag{1.7}$$

so we have:

$$\mathbf{b}_1 = \frac{2\pi}{a} \begin{pmatrix} 1\\ \frac{1}{\sqrt{3}} \end{pmatrix}, \ \mathbf{b}_2 = \frac{2\pi}{a} \begin{pmatrix} 1\\ -\frac{1}{\sqrt{3}} \end{pmatrix} \tag{1.8}$$

The first Brilluoin zone of the hexagonal lattice is shown in fig. 1.1b, with the points of high symmetry

$$\Gamma = \begin{pmatrix} 0 \\ 0 \end{pmatrix}, \ \mathbf{M} = \frac{\pi}{a} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{3}} \end{pmatrix}, \ \mathbf{K} = \frac{4\pi}{3a} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \ . \tag{1.9}$$

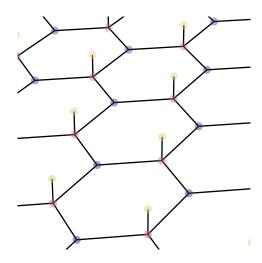


Figure 1.2: Dressed Graphene model

1.2 Dressed Graphene Model

The model I am concerned with in this thesis consists of a Hubbard Hamiltonian (as introduced in ??) on a Graphene lattice, with one additional atom at one of the two sites in a unit cell, which I will call X. This is shown in fig. 1.2. The kinetic term is

$$H_0 = -t_{\mathcal{X}} \sum_{\langle ij \rangle, \sigma} d^{\dagger}_{i,\sigma} d_{j,\sigma} - t_{\text{Gr}} \sum_{\langle ij \rangle, \sigma} c^{(\mathbf{A}),\dagger}_{i,\sigma} c^{(\mathbf{B})}_{j,\sigma} + V \sum_{i,\sigma\sigma'} d^{\dagger}_{i,\sigma} c^{(\mathbf{A})}_{i,\sigma'} + \text{h.c.}$$
 (1.10)

Work over image for dressed graphene lattice

with

- *d* operators on the X atom
- $c^{(\epsilon)}$ operators on the graphene sites $(\epsilon = A, B)$
- t_X nearest neighbour hopping for X
- ullet t_{Gr} nearest neighbour hopping between Graphene sites
- *V* hopping between X and Graphene A sites.

The (attractive) Hubbard interaction has the following form:

$$H_{\rm int} = -U_X \sum_i d^{\dagger}_{i,\uparrow} d^{\dagger}_{i,\downarrow} d_{i,\downarrow} d_{i,\uparrow} - U_{\rm Gr} \sum_{i,\epsilon=A,B} c^{(\epsilon)\dagger}_{i,\uparrow} c^{(\epsilon)\dagger}_{i,\downarrow} c^{\epsilon}_{i,\downarrow} c^{\epsilon}_{i,\downarrow}$$
(1.11)

The notation using different letters for the sites connects intuitively to the physical picture, but it is more economical and in line with the notation for mean field-theory established in ?? to write the Hamiltonian using a sublattice index

$$\alpha = 1, 2, 3 \tag{1.12}$$

with $1 \cong Gr_A$, $2 \cong Gr_B$, $3 \cong X$. Then we can write the non-interacting term as

$$H_0 = \sum_{\langle i,j\rangle,\alpha,\beta,\sigma} [\mathbf{t}]_{i\alpha,j\beta} c_{i\alpha}^{\dagger} c_{j\beta}$$
 (1.13)

with the matrix in the sublattice indices

$$\mathbf{t} = \begin{pmatrix} 0 & -t_{\mathrm{Gr}} & V\delta_{ij} \\ -t_{\mathrm{Gr}} & 0 & 0 \\ V\delta_{ij} & 0 & -t_{\mathrm{X}} \end{pmatrix}$$
 (1.14)

Also write the interaction part as

$$H_{\rm int} = -\sum_{i\alpha} U_{\alpha} c_{i\alpha\uparrow}^{\dagger} c_{i\alpha\downarrow}^{\dagger} c_{i\alpha\downarrow} c_{i\alpha\uparrow} . \qquad (1.15)$$

Using the Fourier transformation ??

Clean up the section from here

$$H_{0} = \sum_{\mathbf{k},\sigma,\sigma'} \begin{pmatrix} c_{k,\sigma}^{A,\dagger} & c_{k,\sigma}^{B,\dagger} & d_{k,\sigma}^{\dagger} \end{pmatrix} \begin{pmatrix} 0 & f_{Gr} & V \\ f_{Gr}^{*} & 0 & 0 \\ V & 0 & f_{X} \end{pmatrix} \begin{pmatrix} c_{k,\sigma}^{A} \\ c_{k,\sigma}^{B} \\ d_{k,\sigma} \end{pmatrix}$$
(1.16)

The band structure for the non-interacting dressed graphene model is easily obtained by diagonalising the matrix in eq. (1.16). This was done in fig. 1.3.

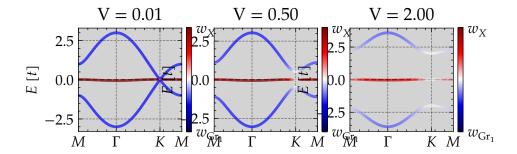


Figure 1.3: Bands of the non-interacting dressed Graphene model, with parameters $t_\chi = 0 \cdot t_{\rm Gr}$