## I EG-X Model with interactions

## I.1 BdG Hamiltonian and free energy

## I.1.1 BdG Hamiltonian

Define sublattice index

$$\alpha = 1, 2, 3 \tag{I.1}$$

with  $1 = Gr_1, 2 = Gr_2, 3 = X$ . Then we can write the non-interacting term as

$$H_0 = -\sum_{\langle i,j\rangle,\alpha,\beta,\sigma} [\mathbf{t}]_{i\alpha,j\beta} c_{i\alpha}^{\dagger} c_{j\beta}$$
 (I.2)

with the matrix

$$\mathbf{t} = \begin{pmatrix} 0 & t_{Gr} & 0 \\ t_{Gr} & 0 & -V\delta_{ij} \\ 0 & -V\delta_{ij} & t_{X} \end{pmatrix}$$
 (I.3)

Add chemical potential:

$$-\mu \sum_{i\alpha\sigma} n_{i\alpha\sigma} \tag{I.4}$$

Also write the interaction part with  $\alpha$  (with changed signs compared to Niklas, to keep in line with papers about the attractive Hubbard model):

$$H_{int} = -\sum_{i\alpha} U_{\alpha} c_{i\alpha\uparrow}^{\dagger} c_{i\alpha\downarrow}^{\dagger} c_{i\alpha\downarrow} c_{i\alpha\uparrow}$$
 (I.5)

Fourier transformation:

$$H_{int} = -\frac{1}{N^2} \sum_{\alpha, \mathbf{k}_{1,2,3,4}} U_{\alpha} e^{i(\mathbf{k}_1 + \mathbf{k}_4 - \mathbf{k}_1 - \mathbf{k}_3) r_{i\alpha}} c^{\dagger}_{\mathbf{k}_1 \alpha \uparrow} c^{\dagger}_{\mathbf{k}_3 \alpha \downarrow} c_{\mathbf{k}_2 \alpha \downarrow} c_{\mathbf{k}_4 \alpha \uparrow}$$
(I.6)

Impose zero-momentum pairing:  $\mathbf{k}_1 + \mathbf{k}_3 = 0$  and  $\mathbf{k}_2 + \mathbf{k}_4 = 0$ :

$$H_{int} = -\sum_{\alpha, \mathbf{k}, \mathbf{k'}} U_{\alpha} c_{\mathbf{k}\alpha\uparrow}^{\dagger} c_{-\mathbf{k}\alpha\downarrow}^{\dagger} c_{-\mathbf{k'}\alpha\downarrow} c_{\mathbf{k'}\alpha\uparrow}$$
(I.7)

Mean-field approximation:

$$H_{int} \approx \sum_{\alpha, \mathbf{k}} (\Delta_{\alpha} c_{\mathbf{k}\alpha\uparrow}^{\dagger} c_{-\mathbf{k}\alpha\downarrow}^{\dagger} + \Delta_{\alpha}^{*} c_{-\mathbf{k}\alpha\downarrow} c_{\mathbf{k}\alpha\uparrow})$$
 (I.8)

with

$$\Delta_{\alpha} = -U_{\alpha} \sum_{\mathbf{k}'} \langle c_{-\mathbf{k}'\alpha\downarrow} c_{\mathbf{k}'\alpha\uparrow} \rangle \tag{I.9}$$

$$\Delta_{\alpha}^{*} = -U_{\alpha} \sum_{\mathbf{k}'} \langle c_{\mathbf{k}'\alpha\uparrow}^{\dagger} c_{-\mathbf{k}'\alpha\downarrow}^{\dagger} \rangle \tag{I.10}$$

This gives the BCS mean field Hamiltonian:

$$H_{BCS} = \sum_{\mathbf{k}\alpha\beta\sigma} [H_{0,\sigma}(\mathbf{k})]_{\alpha\beta} c_{\mathbf{k}\alpha\sigma}^{\dagger} c_{\mathbf{k}\beta\sigma} - \mu \sum_{\mathbf{k}\alpha\sigma} n_{\mathbf{k}\alpha\sigma} + \sum_{\alpha,\mathbf{k}} (\Delta_{\alpha} c_{\mathbf{k}\alpha\uparrow}^{\dagger} c_{-\mathbf{k}\alpha\downarrow}^{\dagger} + \Delta_{\alpha}^{*} c_{-\mathbf{k}\alpha\downarrow} c_{\mathbf{k}\alpha\uparrow})$$
(I.11)

with Nambu spinor

$$\Psi_{\mathbf{k}} = \begin{pmatrix}
c_{1,\mathbf{k}\uparrow} \\
c_{2,\mathbf{k}\uparrow} \\
c_{3,\mathbf{k}\uparrow} \\
c_{1,-\mathbf{k}\downarrow}^{\dagger} \\
c_{2,-\mathbf{k}\downarrow}^{\dagger} \\
c_{3,-\mathbf{k}\downarrow}^{\dagger}
\end{pmatrix} (I.12)$$

we have:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \mathcal{H}(\mathbf{k}) \Psi_{\mathbf{k}}$$
 (I.13)

with

$$\mathcal{H}(\mathbf{k}) = \begin{pmatrix} H_{0,\uparrow}(\mathbf{k}) - \mu & \Delta \\ \Delta^{\dagger} & -H_{0,\downarrow}^{*}(-\mathbf{k}) + \mu \end{pmatrix}$$
(I.14)

with  $H_{0,\sigma}$  being the F.T. of the kinetic term and  $\Delta = diag(\Delta_1, \Delta_2, \Delta_3)$ .

## I.1.2 BdG Hamiltonian in band basis

Use transformation

$$c_{\mathbf{k}\alpha\sigma}^{\dagger} = \sum_{n} [\mathbf{G}]_{\alpha n}^{*} d_{n\mathbf{k}\sigma}^{\dagger} \tag{I.15}$$

where the columns are made up of the eigenvectors of  $\mathbf{H}_{0,\sigma}$  for a given  $\mathbf{k}$ :

$$\mathbf{G} = \begin{pmatrix} \mathbf{G}_1 & \mathbf{G}_2 & \mathbf{G}_3 \end{pmatrix} \tag{I.16}$$

with that:

$$\mathbf{G}_{\sigma}^{\dagger}(\mathbf{k})\mathbf{H}_{0,\sigma}(\mathbf{k})\mathbf{G}_{\sigma}(\mathbf{k}) = \begin{pmatrix} \epsilon_{1} & 0 & 0\\ 0 & \epsilon_{2} & 0\\ 0 & 0 & \epsilon_{3} \end{pmatrix}$$
(I.17)

So the kinetic part of the BdG Hamiltonian becomes:

$$\sum_{\mathbf{k}\alpha\beta\sigma} [H_{0,\sigma}(\mathbf{k})]_{\alpha\beta} \sum_{n} [\mathbf{G}(\mathbf{k})]_{\alpha n}^{*} d_{n\mathbf{k}\sigma}^{\dagger} \sum_{m} [\mathbf{G}(\mathbf{k})]_{\beta m} d_{m\mathbf{k}\sigma} - \mu \sum_{\mathbf{k}\alpha\sigma} n_{n\mathbf{k}\sigma}$$
(I.18)

$$= \sum_{mn\mathbf{k}\sigma} d_{n\mathbf{k}\sigma}^{\dagger} d_{m\mathbf{k}\sigma} \sum_{\alpha\beta} [\mathbf{G}(\mathbf{k})]_{\alpha n}^{*} [H_{0,\sigma}(\mathbf{k})]_{\alpha\beta} [\mathbf{G}(\mathbf{k})]_{\beta m} - \mu \sum_{\mathbf{k}\alpha\sigma} n_{n\mathbf{k}\sigma}$$
(I.19)

$$= \sum_{mn\mathbf{k}\sigma} d_{n\mathbf{k}\sigma}^{\dagger} d_{m\mathbf{k}\sigma} \epsilon_n \delta_{nm} - \mu \sum_{\mathbf{k}\alpha\sigma} n_{n\mathbf{k}\sigma}$$
(I.20)

$$= \sum_{n \mathbf{k}\sigma} \epsilon_n d_{n\mathbf{k}\sigma}^{\dagger} d_{n\mathbf{k}\sigma} - \mu \sum_{\mathbf{k}\sigma\sigma} n_{n\mathbf{k}\sigma}$$
 (I.21)

$$=: \sum_{n\mathbf{k}\sigma} \xi_{\mathbf{k}} d_{n\mathbf{k}\sigma}^{\dagger} d_{n\mathbf{k}\sigma} \tag{I.22}$$

with  $\xi_{\mathbf{k}} \coloneqq \epsilon_{\mathbf{k}} - \mu$ . The pairing terms become:

$$\sum_{\mathbf{k}\alpha} \Delta_{\alpha} c_{\mathbf{k}\alpha\uparrow}^{\dagger} c_{-\mathbf{k}\alpha\downarrow}^{\dagger} = \sum_{\mathbf{k}\alpha} \Delta_{\alpha} \sum_{n} [\mathbf{G}_{\uparrow}(\mathbf{k})]_{\alpha n}^{*} d_{n\mathbf{k}\uparrow}^{\dagger} \sum_{m} [\mathbf{G}_{\downarrow}(-\mathbf{k})]_{\beta m}^{*} d_{m-\mathbf{k}\downarrow}^{\dagger} \quad (I.23)$$

$$= (I.24)$$

So that:

$$\mathcal{H}(\mathbf{k}) = \begin{pmatrix} \epsilon_{\mathbf{k}} - \mu & G^{\dagger} \Delta G \\ G^{\dagger} \Delta^{\dagger} G & -\epsilon_{\mathbf{k}} + \mu \end{pmatrix}$$
(I.25)

with

$$\epsilon_{\mathbf{k}} = \begin{pmatrix} \epsilon_1(\mathbf{k}) & 0 & 0 \\ 0 & \epsilon_2(\mathbf{k}) & 0 \\ 0 & 0 & \epsilon_3(\mathbf{k}) \end{pmatrix}$$
(I.26)

Concrete example for transformation of gaps from orbital to band basis at  $K = \frac{4\pi}{3a} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ . There, the non-interacting part becomes simply:

$$\mathcal{H}_0 = \begin{pmatrix} 0 & 0 & V \\ 0 & 0 & 0 \\ V & 0 & 3t_X \end{pmatrix} \tag{I.27}$$

The eigenvalue problem can be solved e.g. via sympy:

$$G = \begin{pmatrix} \frac{-3t_X - \sqrt{4V^2 + 9t_X^2}}{\sqrt{4V^2 + \left(3t_X + \sqrt{4V^2 + 9t_X^2}\right)^2}} & 0 & \frac{-3t_X + \sqrt{4V^2 + 9t_X^2}}{\sqrt{4V^2 + \left(3t_X - \sqrt{4V^2 + 9t_X^2}\right)^2}} \\ 0 & 1 & 0 \\ \frac{2V}{\sqrt{4V^2 + \left(3t_X + \sqrt{4V^2 + 9t_X^2}\right)^2}} & 0 & \frac{2V}{\sqrt{4V^2 + \left(3t_X - \sqrt{4V^2 + 9t_X^2}\right)^2}} \end{pmatrix}$$
 (I.28)

So for  $V \to 0$ :

$$G = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \tag{I.29}$$

but for V > 0, there are off-diagonal elements, e.g. V = 0.1:

$$G = \begin{pmatrix} -0.7578 & 0 & 0.6526 \\ 0 & 1 & 0 \\ 0.6526 & 0 & 0.7578 \end{pmatrix}$$
 (I.30)

So the transformation of the gap from orbital to band space reads:

$$G^{\dagger}\Delta G = \begin{pmatrix} \frac{3\Delta_{1}t_{X} - 3\Delta_{3}t_{X} + (\Delta_{1} + \Delta_{3})\sqrt{4V^{2} + 9t_{X}^{2}}}{2\sqrt{4V^{2} + 9t_{X}^{2}}} & 0 & \frac{V(-\Delta_{1} + \Delta_{3})}{\sqrt{4V^{2} + 9t_{X}^{2}}} \\ 0 & \Delta_{2} & 0 \\ \frac{V(-\Delta_{1} + \Delta_{3})}{\sqrt{4V^{2} + 9t_{X}^{2}}} & 0 & \frac{-3\Delta_{1}t_{X} + 3\Delta_{3}t_{X} + (\Delta_{1} + \Delta_{3})\sqrt{4V^{2} + 9t_{X}^{2}}}{2\sqrt{4V^{2} + 9t_{X}^{2}}} \end{pmatrix}$$

$$(I.31)$$

So in particular there is no interband pairing for  $V \to 0$ :

$$G^{\dagger} \Delta G = \begin{pmatrix} \Delta_1 & 0 & 0 \\ 0 & \Delta_2 & 0 \\ 0 & 0 & \Delta_3 \end{pmatrix}$$
 (I.32)

But for V > 0, there is interband pairing (e.g. V = 0.1):

$$G^{\dagger}\Delta G = \begin{pmatrix} 0.5742\Delta_1 + 0.4258\Delta_3 & 0 & -0.4945\Delta_1 + 0.4945\Delta_3 \\ 0 & \Delta_2 & 0 \\ -0.4945\Delta_1 + 0.4945\Delta_3 & 0 & 0.4258\Delta_1 + 0.5742\Delta_3 \end{pmatrix}$$
 (I.33)