I Dressed Graphene Hamiltonian in Reciprocal Space

Clean up this section

In this chapter, the model Hamiltonian from **??**

$$H_0 = -t_X \sum_{\langle ij\rangle,\sigma} d^\dagger_{i,\sigma} d_{j,\sigma} - t_{Gr} \sum_{\langle ij\rangle,\sigma} c^{(A),\dagger}_{i,\sigma} c^{(B)}_{j,\sigma} + V \sum_{i,\sigma} d^\dagger_{i,\sigma} c^{(A)}_{i,\sigma} + \text{h.c.} \tag{I.1}$$

will be treated to obtain the band structure. The first step is to write out the sums over nearest neighbors $\langle i,j\rangle$ explicitly, writing $\delta_X, \delta_\epsilon$ ($\epsilon = A, B$) for the vectors to the nearest neighbors of the X atoms and Graphene A, B sites. Doing the calculation for example of the X atoms:

$$-t_X \sum_{\langle ij\rangle,\sigma} (d^\dagger_{i,\sigma} d_{j,\sigma} + d^\dagger_{j,\sigma} d_{i,\sigma}) = -\frac{t_X}{2} \sum_{i,\sigma} \sum_{\delta_X} d^\dagger_{i,\sigma} d_{i+\delta_X,\sigma} - \frac{t_X}{2} \sum_{j,\sigma} \sum_{\delta_X} d^\dagger_{j,\sigma} d_{j+\delta_X,\sigma} \tag{I.2}$$

$$= -t_X \sum_{i,\sigma} \sum_{\delta_X} d^\dagger_{i,\sigma} d_{i+\delta_X,\sigma} \tag{I.3}$$

The factor $1/2$ in eq. (I.2) is to account for double counting when going to the sum over all lattice sites i. By relabeling $j \to i$ in the second sum, the two sum are the same and eq. (I.3) is obtained. Using now the discrete Fourier transform

$$c_i = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{i\mathbf{k}\mathbf{r}_i} c_{\mathbf{k}}, \; c^\dagger_i = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{-i\mathbf{k}\mathbf{r}_i} c^\dagger_{\mathbf{k}} \tag{I.4}$$

with the completeness relation

$$\sum_i e^{i\mathbf{k}\mathbf{r}_i} e^{-i\mathbf{k}'\mathbf{r}_i} = N\delta_{\mathbf{k},\mathbf{k}'} \,, \tag{I.5}$$

eq. (I.3) reads:

$$-t_X \frac{1}{N} \sum_{i,\sigma} \sum_{\delta_X} d^\dagger_{i,\sigma} d_{i+\delta_X,\sigma} = -t_X \frac{1}{N} \sum_{i,\sigma} \sum_{\mathbf{k},\mathbf{k}',\delta_X} \left(e^{-\mathrm{i}\mathbf{k}\mathbf{r}_i} d^\dagger_{\mathbf{k},\sigma} \right) \left(e^{\mathrm{i}\mathbf{k}'\mathbf{r}_i} e^{\mathrm{i}\mathbf{k}'\delta_X} d_{\mathbf{k}',\sigma} \right) \tag{I.6}$$

$$= -t_X \frac{1}{N} \sum_{\mathbf{k},\mathbf{k}',\delta_X,\sigma} d^\dagger_{\mathbf{k},\sigma} d_{\mathbf{k}',\sigma} e^{\mathrm{i}\mathbf{k}'\delta_X} \sum_i e^{-\mathrm{i}\mathbf{k}\mathbf{r}_i} e^{\mathrm{i}\mathbf{k}'\mathbf{r}_i} \tag{I.7}$$

$$= -t_X \frac{1}{N} \sum_{\mathbf{k},\mathbf{k}',\sigma} d^\dagger_{\mathbf{k},\sigma} d_{\mathbf{k}',\sigma} \sum_{\delta_X} e^{\mathrm{i}\mathbf{k}'\delta_X} \left(N \delta_{\mathbf{k},\mathbf{k}'} \right) \tag{I.8}$$

$$= -t_X \sum_{\mathbf{k},\sigma} d^\dagger_{\mathbf{k},\sigma} d_{\mathbf{k},\sigma} \sum_{\delta_X} e^{\mathrm{i}\mathbf{k}\delta_X} \, . \tag{I.9}$$

This part is now diagonal in $\mathbf{k}$ space. The nearest neighbours vectors δ_X for the X atoms are the vectors $\delta_{AA,i}$ from **??**. With that, the sum over δ_X can be explicitly calculated:

Correct exp expressions

Example for a vector product

$$f_X(\mathbf{k}) = -t_X \sum_{\delta_X} e^{\mathrm{i}\mathbf{k}\delta_X} \tag{I.10}$$

$$= -t_X \left[\exp\left(\mathrm{i}a \left(\frac{k_x}{2} + \frac{\sqrt{3}k_y}{2} \right) \right) + e^{\mathrm{i}ak_x} + e^{\mathrm{i}a(\frac{k_x}{2} - \frac{\sqrt{3}k_y}{2})} \right. \tag{I.11}$$

$$\left. + \, e^{\mathrm{i}a(-\frac{k_x}{2} - \frac{\sqrt{3}k_y}{2})} + e^{-\mathrm{i}ak_x} + e^{\mathrm{i}a(-\frac{k_x}{2} + \frac{\sqrt{3}k_y}{2})} \right] \tag{I.12}$$

$$= -t_X \left(2\cos(ak_x) + 2e^{\mathrm{i}a\frac{\sqrt{3}k_y}{2}} \cos\left(\frac{a}{2}k_x\right) + 2e^{-\mathrm{i}a\frac{\sqrt{3}k_y}{2}} \cos\left(\frac{a}{2}k_x\right) \right) \tag{I.13}$$

$$= -2t_X \left(\cos(ak_x) + 2\cos\left(\frac{a}{2}k_x\right) \cos\left(\sqrt{3}\frac{a}{2}k_y\right) \right) \, . \tag{I.14}$$

The same can be done for the hopping between Graphene sites, for example :

$$-t_{\mathrm{Gr}} \sum_{\langle ij \rangle,\sigma\sigma'} c^{(A),\dagger}_{i,\sigma} c^{(B)}_{j,\sigma'} = -t_{\mathrm{Gr}} \sum_{i,\sigma\sigma'} \sum_{\delta_{AB}} c^{(A),\dagger}_{i,\sigma} c^{(B)}_{i+\delta_{AB},\sigma'} \tag{I.15}$$

$$= -t_{\mathrm{Gr}} \sum_{\mathbf{k},\sigma,\sigma'} c^{(A)\dagger}_{\mathbf{k},\sigma} c^{(B)}_{\mathbf{k},\sigma'} \sum_{\delta_{AB}} e^{\mathrm{i}\mathbf{k}\delta_{AB}} \tag{I.16}$$

Show that!

We note

$$\sum_{\delta_{AB}} e^{\mathrm{i}\mathbf{k}\delta_{AB}} = \left(\sum_{\delta_{BA}} e^{\mathrm{i}\mathbf{k}\delta_{BA}} \right)^* = \sum_{\delta_{BA}} e^{-\mathrm{i}\mathbf{k}\delta_{BA}} \tag{I.17}$$

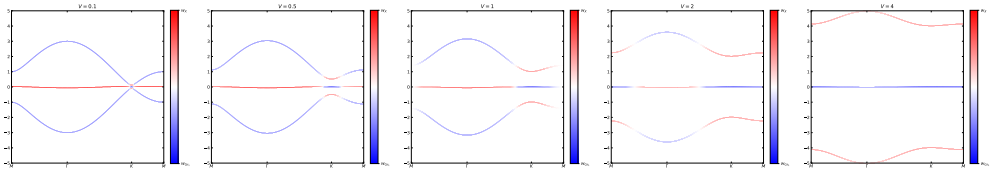

Figure I.1: *Bands of the non-interacting EG-X model. All the bands are spin-degenerate.*

and calculate

$$f_{Gr} = -t_{Gr} \sum_{\delta_{AB}} e^{\mathrm{i}\mathbf{k}\delta_{AB}} \tag{I.18}$$

$$= -t_{Gr} \left(e^{\mathrm{i}\frac{a}{\sqrt{3}}k_y} + e^{\mathrm{i}\frac{a}{2\sqrt{3}}(\sqrt{3}k_x - k_y)} + e^{\mathrm{i}\frac{a}{2\sqrt{3}}(-\sqrt{3}k_x - k_y)} \right) \tag{I.19}$$

$$= -t_{Gr} \left(e^{\mathrm{i}\frac{a}{\sqrt{3}}k_y} + e^{-\mathrm{i}\frac{a}{2\sqrt{3}}k_y} \left(e^{\mathrm{i}\frac{a}{2}k_x} + e^{-\mathrm{i}\frac{a}{2}k_x} \right) \right) \tag{I.20}$$

$$= -t_{Gr} \left(e^{\mathrm{i}\frac{a}{\sqrt{3}}k_y} + 2e^{-\mathrm{i}\frac{a}{2\sqrt{3}}k_y} \cos\left(\frac{a}{2}k_x\right) \right) \tag{I.21}$$

All together, we get:

$$H_0 = \sum_{\mathbf{k},\sigma,\sigma'} \begin{pmatrix} c^{A,\dagger}_{k,\sigma} & c^{B,\dagger}_{k,\sigma} & d^{\dagger}_{k,\sigma} \end{pmatrix} \begin{pmatrix} 0 & f_{Gr} & V \\ f^*_{Gr} & 0 & 0 \\ V & 0 & f_X \end{pmatrix} \begin{pmatrix} c^{A}_{k,\sigma} \\ c^{B}_{k,\sigma} \\ d_{k,\sigma} \end{pmatrix} \tag{I.22}$$

The band structure for the non-interacting EG-X model is easily obtained by diagonalising the matrix in eq. **??**. This was done in fig. I.1.

Values used for calculation:

- $a_0 = 1$
- $t_{\mathrm{Gr}} = 1$
- $t_{\mathrm{X}} = 0.01$

V is the control parameter. A range from $V = 0.1$ to $V = 2$ can be mapped onto materials in experiment.