Contents

Todo list

I Green's Function Formalism

Following [1]

Green's functions: method to encode influence of many-body effects on propagation of particles in a system.

Have different kinds of Green's functions, for example the retarded Green's function:

$$G^R(\mathbf{r}\sigma t, \mathbf{r}'\sigma' t') = -\mathrm{i}\Theta(t-t')\left\langle\{c_{\mathbf{r}\sigma}(t), c^\dagger_{\mathbf{r}\sigma}(t')\}\right\rangle \tag{I.1}$$

They give the amplitude of a particle inserted at point $\mathbf{r}'$ at time t' to propagate to position $\mathbf{r}$ at time t. In a translation invariant system: can use $\mathbf{k}$ as a natural basis set:

$$G^R(\mathbf{k}, \sigma t, \sigma' t') = -\mathrm{i}\Theta(t-t')\left\langle\{c_{\mathbf{k}\sigma}(t), c^\dagger_{\mathbf{k}\sigma}(t')\}\right\rangle \tag{I.2}$$

Can define the spectral function from this:

$$A(\mathbf{k}\sigma, \omega) = -2\Im G^R(\mathbf{k}\sigma, \omega) \tag{I.3}$$

Energy dependent retarded GF

Show GFs can be relatex to observables

Mathematical technique to calculate retarded GFs involves defining GFs on imaginary times τ:

$$t \to -\mathrm{i}\tau \tag{I.4}$$

where τ is real and has the dimension time. This enables the simultaneous expansion of exponential $e^{-\beta H}$ coming from the thermodynamic average and $e^{-\mathrm{i}Ht}$ coming from the time evolution of operators.

Definitions of Matsubara GFs

Define imaginary time/Matsubara GF $\mathcal{C}_{AB}(\tau, 0)$:

$$\mathcal{C}_{AB}(\tau, 0) = -\left\langle T_\tau(A(\tau)B(0))\right\rangle \tag{I.5}$$

with time-ordering operator in imaginary time:

ordering operator

Can prove from properties of Matsubara GF, that they are only defined for

$$-\beta < \tau < \beta \tag{I.6}$$

Due to this, the Fourier transform of the Matsubara GF is defined on discrete values:

$$\mathcal{C}_{AB}(\mathrm{i}\omega_n) = \int_0^\beta \mathrm{d}\tau \tag{I.7}$$

with fermionic/bosonic Matsubara frequencies

$$\omega_n = \begin{cases} \frac{2n\pi}{\beta} \text{ for bosons} \\ \frac{(2n+1)\pi}{\beta} \text{ for fermions} \end{cases} \tag{I.8}$$

How to resolve ambiguity at borders of integral

Connection to real-time GFs

single-particle Matsubara GF

equations of motion for Matsubara GF

A bit of perturbation theory -> Dyson equation and self-energy

More general introduction into NG GFs, how they look like, what they describe etc.

Sources for these?

How to get quasiparticle weight?

Extrapolation of the Matsubara GF to zero is proportional to the density of states at the chemical potential. Gapped: density is zero (Matsubara GF goes to 0), metal: density is finite (Matsubara GF goes to finite value) [1, p. 8.3.4].

I.1 Nambu-Gorkov GF

Order parameter can be chosen as the anomalous GF:

$$\Psi = F^{\mathrm{loc}}(\tau = 0^-) \tag{I.9}$$

or the superconducting gap

$$\Delta = Z\Sigma^{\mathrm{AN}} \tag{I.10}$$

that can be calculated from the anomalous self-energy Σ^{AN} and quasiparticle weight Z

II Superconductivity

II.0.1 Extracting T_C

From Niklas

Übrigens: Typische Varianten, um 'sauber(er)' Tc zu bestimmen, ist $|OP|^2$ gegen T aufzutragen, da das (als Phasenübergang 2. Ordnung) proportional zu T-Tc ist. Heißt, man kann Tc dann mittels linearem Fit finden - ist leider auch nicht immer der einfachste Weg, weil der Bereich, in dem diese lineare Näherung prinzipiell sehr klein um Tc herum sein kann. Aber pi-mal-Daumen Abschätzungen gehen damit ganz gut. Oder man macht es wie unten beschrieben mit einer daraus abgeleiteten Formel.

Find a source for that! Phase transitions

in the superconducting region in Fig. 2A. To estimate T_c, we assume that the order parameter as a function of temperature behaves like $P_{SC}(T) \propto \sqrt{T_c - T}$ for $T \lesssim T_c$. We pick the highest and the second highest temperatures (T_1 and T_2, respectively), which give stable superconducting solutions ($P_{SC} \gtrsim 0.01$). Then, T_c is determined by

$$T_c = T_1 + (T_1 - T_2)\frac{P_{SC}^2(T_1)}{P_{SC}^2(T_2) - P_{SC}^2(T_1)}. \qquad \text{(S4)}$$

Figure II.1: *Formula for extracting T_C*

III d-wave Superconductivity

Source: Coleman - *Introduction to Many-Body Physics* [2, ch. 15]

III.1 BCS theory with momentum dependent coupling

Starting point is a BCS-Hamiltonian with momentum-dependent coupling term $V_{\mathbf{k},\mathbf{k}'}$:

$$H = \sum_{\mathbf{k},\sigma} \epsilon_{\mathbf{k}} c^\dagger_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma} + \sum_{\mathbf{k},\mathbf{k}'} V_{\mathbf{k},\mathbf{k}'} c^\dagger_{\mathbf{k}\uparrow} c^\dagger_{-\mathbf{k}\downarrow} c_{-\mathbf{k}'\downarrow} c_{\mathbf{k}'\uparrow} \tag{III.1}$$

The original idea by Bardeen, Cooper and Schrieffer uses the coupling

$$V_{\mathbf{k},\mathbf{k}'} = \begin{cases} -\frac{g_0}{V} \ , & |\epsilon_{\mathbf{k}}| < \omega_D \\ 0 \end{cases} \tag{III.2}$$

Point out specific difference to BCS theory!

Then similar process as for BCS theory without the momentum-dependent term (Hubbard-Stratonovich decoupling, minimization of mean-field free energy). Gives self-consistent equation for the gap function:

$$\Delta_{\mathbf{k}} = -\sum_{\mathbf{k}'} V_{\mathbf{k},\mathbf{k}'} \frac{\Delta_{\mathbf{k}'}}{2E_{\mathbf{k}'}} \tanh\left(\frac{\beta E_{\mathbf{k}}}{2}\right) \tag{III.3}$$

or at $T = 0$:

$$\Delta_{\mathbf{k}} = -\sum_{\mathbf{k}'} V_{\mathbf{k},\mathbf{k}'} \frac{\Delta_{\mathbf{k}'}}{2E_{\mathbf{k}'}} \tag{III.4}$$

What is the E_k?

Important note: there is a minus sign in the front! If $V_{\mathbf{k},\mathbf{k}'} < 0$ (a uniformly attractive interaction), the equation is fulfilled by a uniformly positive gap function. In general $V_{\mathbf{k},\mathbf{k}'}$ contains repulsive (positive) terms (in particual stemming from the Coulomb interaction), so the gap function cannot be

uniformly positive, it acquires nodes in momentum space. Most satisfying solutions fulfill:

$$\operatorname{sign}(\Delta_{\mathbf{k}}) = -\operatorname{sign}(V_{\mathbf{k},\mathbf{k}'})\operatorname{sign}(\Delta_{\mathbf{k}'}) \tag{III.5}$$

So for an attractive interaction we have:

$$\operatorname{sign}(\Delta_{\mathbf{k}}) = -(-1)\operatorname{sign}(\Delta_{\mathbf{k}'}) \tag{III.6}$$

So areas in phase space linked by an attractive interaction have the same sign (and areas linked by repulsive interaction have opposite signs)! Solutions like this have the largest gaps and thus the largest mean-field transition temperature .

Why large gap?

Connection from gap to transition temperature?

Are there more?

Two cases :

- Electron-phonon superconductors: interaction is repulsive at high energies, $\Delta_{\mathbf{k}}$ is largely isotropic in momentum space, but changes sign at $\approx$ Debye frequency
- Anisotropic superconductors: $\Delta_{\mathbf{k}}$ is strongly momentum-dependent, acquires nodes in momentum space

The last mechanism is at work in heavy-fermion, high-temperature cuprate and iron-based superconductors.

III.2 Anisotropic pairing

III.2.1 Hubbard interaction

The goal in this section is to derive a BCS-like Hamiltonian with a term

$$V_{\mathbf{k},\mathbf{k}'}\Psi^{\dagger}_{\mathbf{k}}\Psi_{\mathbf{k}'} \tag{III.7}$$

We start from a Hubbard-like interaction term

$$V = \sum_{\mathbf{q}} V_{\mathbf{q}} : \rho_{-\mathbf{q}}\rho_{\mathbf{q}} := \frac{1}{2}\sum_{\mathbf{k}_1,\mathbf{k}_2,\mathbf{q},\sigma,\sigma'} V_{\mathbf{q}} c^{\dagger}_{\mathbf{k}_1+\mathbf{q}\sigma} c^{\dagger}_{\mathbf{k}_2-\mathbf{q}\sigma'} c_{\mathbf{k}_2\sigma'} c_{\mathbf{k}_1\sigma} \tag{III.8}$$

Proper implementation of normal-ordering

Hubbard-like would be $V_q = U$?

Cooper pairs have zero total momentum and the pairing potential is determined by the interaction on them, so we have

$$\mathbf{k}_1 + \mathbf{k}_2 = 0 \implies \mathbf{k}_1 = -\mathbf{k}_2 -: \mathbf{k}' \tag{III.9}$$

$$\mathbf{k}_1 + \mathbf{q} = -(\mathbf{k}_2 - \mathbf{q}) -: \mathbf{k} \implies \mathbf{k}' + \mathbf{q} = \mathbf{k} \implies \mathbf{q} = \mathbf{k} - \mathbf{k}' \tag{III.10}$$

and we can split up the interaction term

Show why the third line works!

$$V_{\mathrm{BCS}} = \frac{1}{2} \sum_{\mathbf{k},\mathbf{k}',\sigma,\sigma'} V_{\mathbf{k}-\mathbf{k}'} c^\dagger_{\mathbf{k}\sigma} c^\dagger_{-\mathbf{k}\sigma'} c_{-\mathbf{k}'\sigma'} c_{\mathbf{k}'\sigma} \tag{III.11}$$

$$= \frac{1}{2} \sum_{\mathbf{k},\mathbf{k}'} V_{\mathbf{k}-\mathbf{k}'} c^\dagger_{\mathbf{k}\uparrow} c^\dagger_{-\mathbf{k}\downarrow} c_{-\mathbf{k}'\downarrow} c_{\mathbf{k}'\uparrow} \qquad \left(= \frac{1}{2} V^{\uparrow\downarrow}_{\mathrm{BCS}}\right) \tag{III.12}$$

$$+ \frac{1}{2} \sum_{\mathbf{k},\mathbf{k}'} V_{\mathbf{k}-\mathbf{k}'} c^\dagger_{\mathbf{k}\downarrow} c^\dagger_{-\mathbf{k}\uparrow} c_{-\mathbf{k}'\uparrow} c_{\mathbf{k}'\downarrow} \qquad \left(= \frac{1}{2} V^{\downarrow\uparrow}_{\mathrm{BCS}} = \frac{1}{2} V^{\uparrow\downarrow}_{\mathrm{BCS}}\right) \tag{III.13}$$

$$+ \frac{1}{2} \sum_{\mathbf{k},\mathbf{k}'} V_{\mathbf{k}-\mathbf{k}'} c^\dagger_{\mathbf{k}\uparrow} c^\dagger_{-\mathbf{k}\uparrow} c_{-\mathbf{k}'\uparrow} c_{\mathbf{k}'\uparrow} \qquad \left(= V^{\uparrow\uparrow}_{\mathrm{BCS}}\right) \tag{III.14}$$

$$+ \frac{1}{2} \sum_{\mathbf{k},\mathbf{k}'} V_{\mathbf{k}-\mathbf{k}'} c^\dagger_{\mathbf{k}\downarrow} c^\dagger_{-\mathbf{k}\downarrow} c_{-\mathbf{k}'\downarrow} c_{\mathbf{k}'\downarrow} \qquad \left(= V^{\downarrow\downarrow}_{\mathrm{BCS}}\right) \tag{III.15}$$

$$= V^{\uparrow\downarrow}_{\mathrm{BCS}} + V^{\uparrow\uparrow}_{\mathrm{BCS}} + V^{\downarrow\downarrow}_{\mathrm{BCS}} \tag{III.16}$$

First we treat $V^{\uparrow\downarrow}_{\mathrm{BCS}}$. Pair of opposite spins are neither single nor triplet, because they are not appropiately symmetrised. If we have the pair wavefunction

Why do we define spatial parity? Only symmetrised wavefunctions physical?

$$F(\mathbf{k})_{\alpha\beta} = \langle \mathbf{k}\alpha, -\mathbf{k}\beta | \, |\mathbf{k}\rho\rangle \tag{III.17}$$

We define spatial parity of this wavefunction:

$$F(-\mathbf{k})_{\alpha\beta} = PF(\mathbf{k})_{\alpha\beta} \tag{III.18}$$

as well as the spin parity:

$$F(\mathbf{k})_{\beta\alpha} = XF(\mathbf{k})_{\alpha\beta} \, , \tag{III.19}$$

where we define singlets ($X = +1$) and triplets ($X = -1$). The join application of XP is an exchange of fermions, so it should have an eigenvalue -1. So we have

- even-parity pairs, $P = +1 \implies X = -1$, spin singlets, $(X, P) = (+, -)$
- odd-parity pairs, $P = -1 \implies X = +1$, spin triplets, $(X, P) = (-, +)$

How exactly?

We split up the interaction into the symmetric and asymmetric parts :

studied superconductors are mostly singlet, pure triplet not found? Thats why we split it up! Paper for that?

$$V_{\text{BCS}} = \sum_{\mathbf{k},\mathbf{k}'} \left(\frac{V_{\mathbf{k}-\mathbf{k}'} + V_{\mathbf{k}+\mathbf{k}'}}{2} + \frac{V_{\mathbf{k}-\mathbf{k}'} - V_{\mathbf{k}+\mathbf{k}'}}{2} \right) \Psi^\dagger_{\mathbf{k}} \Psi_{\mathbf{k}'} \tag{III.20}$$

$$:- \left(V^S_{\mathbf{k},\mathbf{k}'} + V^T_{\mathbf{k},\mathbf{k}'} \right) \Psi^\dagger_{\mathbf{k}} \Psi_{\mathbf{k}'} \,, \tag{III.21}$$

where we have defined the BCS pairing interaction in the singlet and triplet channel:

$$V^{S,T}_{\mathbf{k},\mathbf{k}'} = \frac{1}{2} \left(V_{\mathbf{k}-\mathbf{k}'} \pm V_{\mathbf{k}+\mathbf{k}'} \right) \tag{III.22}$$

explain last step here

The singlet channel is even in $\mathbf{k}, \mathbf{k}'$:

$$V^S_{-\mathbf{k},-\mathbf{k}'} = \frac{1}{2} \left(V_{-\mathbf{k}+\mathbf{k}'} \pm V_{-\mathbf{k}-\mathbf{k}'} \right) = \frac{1}{2} \left(V_{-(\mathbf{k}-\mathbf{k}')} \pm V_{-(\mathbf{k}+\mathbf{k}')} \right) = \frac{1}{2} \left(V_{\mathbf{k}-\mathbf{k}'} \pm V_{\mathbf{k}+\mathbf{k}'} \right) \,, \tag{III.23}$$

while the triplet channel is odd in $\mathbf{k}, \mathbf{k}'$. In the sum:

With everything we write the unequal spin pairing as:

$$V^{\uparrow\downarrow}_{\text{BCS}} = \frac{1}{4} \sum_{\mathbf{k}\mathbf{k}'} \left[V^S_{\mathbf{k},\mathbf{k}'} \Psi^{S\dagger}_{\mathbf{k}} \Psi^S_{\mathbf{k}'} + V^T_{\mathbf{k},\mathbf{k}'} \Psi^{T\dagger}_{\mathbf{k}} \Psi^T_{\mathbf{k}'} \right] \tag{III.24}$$

$$= \sum_{\mathbf{k}\mathbf{k}' \in \frac{1}{2}\text{BZ}} \left[V^S_{\mathbf{k},\mathbf{k}'} \Psi^{S\dagger}_{\mathbf{k}} \Psi^S_{\mathbf{k}'} + V^T_{\mathbf{k},\mathbf{k}'} \Psi^{T\dagger}_{\mathbf{k}} \Psi^T_{\mathbf{k}'} \right] \tag{III.25}$$

vector arrows over the psi (or bold)

The equal spin pairing also includes triplet pairing (these are wrapped up in the vectors) and all in all the BCS pairing potential is:

$$V_{\text{BCS}} = \sum_{\mathbf{k}\mathbf{k}' \in \frac{1}{2}\text{BZ}} \left[V^S_{\mathbf{k},\mathbf{k}'} \Psi^{S\dagger}_{\mathbf{k}} \Psi^S_{\mathbf{k}'} + V^T_{\mathbf{k},\mathbf{k}'} * \Psi^{T\dagger}_{\mathbf{k}} \cdot * \Psi^T_{\mathbf{k}'} \right] \tag{III.26}$$

How can we access that information in experiment?

Source for that?

In real materials we mostly see singlet pairing , in this case we can just write:

$$V_{\text{BCS}} = \sum_{\mathbf{k}\mathbf{k}' \in \frac{1}{2}\text{BZ}} V^S_{\mathbf{k},\mathbf{k}'} (c^\dagger_{\mathbf{k}\uparrow} c^\dagger_{-\mathbf{k}\downarrow})(c_{-\mathbf{k}'\downarrow} c_{\mathbf{k}'\uparrow}) \tag{III.27}$$

III.2.2 Magnetic interaction

Starting point here is a magnetic interaction:

$$V_{\text{mag}} = \frac{1}{2}\sum_{\mathbf{q}} J_{\mathbf{q}} \left[\mathbf{S}_{-\mathbf{q}} \cdot \mathbf{S}_{\mathbf{q}}\right] \tag{III.28}$$

$$= \frac{1}{2}\sum_{\mathbf{k}_1,\mathbf{k}_2,\mathbf{q}} J_{\mathbf{q}} c^{\dagger}_{\mathbf{k}_1+\mathbf{q}\alpha} c^{\dagger}_{\mathbf{k}_2-\mathbf{q}\gamma} \left(\frac{\boldsymbol{\sigma}}{2}\right)_{\alpha\beta} \left(\frac{\boldsymbol{\sigma}}{2}\right)_{\gamma\delta} c_{\mathbf{k}_2\delta} c_{\mathbf{k}_1\beta} \tag{III.29}$$

Important point: eigenvalues of $\mathbf{S}_1 \cdot \mathbf{S}_2$ are different for singlet and triplet states:

$$\mathbf{S}_1 \cdot \mathbf{S}_2 = \begin{cases} +\frac{1}{4} & \text{(triplet)} \\ -\frac{3}{4} & \text{(singlet)} \end{cases} \tag{III.30}$$

These eigenvalues enter as prefactors into the pairing potentials:

$$V^S_{\mathbf{k},\mathbf{k}'} = -\frac{3}{4}\left(\frac{J_{\mathbf{k}-\mathbf{k}'} + J_{\mathbf{k}+\mathbf{k}'}}{2}\right) \tag{III.31}$$

$$V^T_{\mathbf{k},\mathbf{k}'} = \frac{1}{4}\left(\frac{J_{\mathbf{k}-\mathbf{k}'} - J_{\mathbf{k}+\mathbf{k}'}}{2}\right) \tag{III.32}$$

So antiferromagnetic interactions ($J_{\mathbf{k}-\mathbf{k}'} > 0 \implies V^S_{\mathbf{k},\mathbf{k}'} < 0$) attract in the singlet channel, while ferromagnetic interactions ($J_{\mathbf{k}-\mathbf{k}'} < 0 \implies V^T_{\mathbf{k},\mathbf{k}'} < 0$) attracts in the triplet channel.

III.3 d-wave superconductivity in two dimensions - cuprates

Cuprate superconductors cannot be understood in Fermi liquid theory.

A bit more information on history, structure etc.

Three regimes :

How doped?

- Undoped: antiferromagnetic Mott insulators
- Doped: d-wave superconductors
- Over-doped: Fermi liquid behaviours reoccurs, BCS treatment is applicable

Why can we only treat BCS when we also have Fermi liquid?

Do we just treat this case in the following?

Approximate by 2D tight-binding lattice (with nearest-neighbour hopping strength t) with

$$\epsilon_{\mathbf{k}} = -2t(\cos(k_x a) + \cos(k_y a)) - \mu \tag{III.33}$$

interacting via onsite Coulomb repulsion and nearest-neighbour antiferromagnetic interaction:

$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c^\dagger_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma} + \sum_j U n_{j\uparrow} n_{j\downarrow} + J \sum_{\langle i,j\rangle} \mathbf{S}_i \cdot \mathbf{S}_j \tag{III.34}$$

In momentum space:

$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c^\dagger_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma} + \frac{1}{2} \sum_{\mathbf{q}} U \rho_{-\mathbf{q}} \rho_{\mathbf{q}} + J \sum_{\mathbf{q}} \mathbf{S}_{-\mathbf{q}} \cdot \mathbf{S}_{\mathbf{q}} \tag{III.35}$$

with $J_{\mathbf{q}} = 2J(\cos(q_x a) + \cos(q_y a))$. From the treatment of the Hubbard and magnetic interaction earlier we can get the singlet interaction

$V_q^{singlet}$ as well?

Put table here as well?

Calculate that fully

$$V_{\mathbf{k},\mathbf{k}'} = U - \frac{3J}{2}\left(c_x c_{x'} + c_y c_{y'}\right) \tag{III.36}$$

where we use the abbreviation $c_x = \cos(k_x a)$. So the mean-field BCS Hamiltonian is

$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c^\dagger_{\mathbf{k}\sigma} c_{\mathbf{k}\sigma} + \sum_{\mathbf{k}\mathbf{k}'} \left(U - \frac{3J}{2}\left(c_x c_{x'} + c_y c_{y'}\right)\right) \tag{III.37}$$

Looking at the gap equation

$$\Delta_{\mathbf{k}} = -\sum_{\mathbf{k}'} V_{\mathbf{k},\mathbf{k}'} \frac{\Delta_{\mathbf{k}'}}{2E_{\mathbf{k}'}} \tanh\left(\frac{\beta E_{\mathbf{k}}}{2}\right), \tag{III.38}$$

we see that the interaction preserves the symmetries of the pair ($\hat{=}$ symmetries of $\Delta_{\mathbf{k}}$). We divide the interaction into two parts:

Why is the symmetry preserved? And why are the symmetries of the pair conserved? Are these the same as of Δ_k?

$$V^S_{\mathbf{k},\mathbf{k}'} = U - \frac{3J}{4}(c_x + c_y)(c_{x'} + c_{y'}) \tag{III.39}$$

$$V^D_{\mathbf{k},\mathbf{k}'} = -\frac{3J}{2}(c_x - c_y)(c_{x'} - c_{y'}) \tag{III.40}$$

$$V^S_{\mathbf{k},\mathbf{k}'} + V^D_{\mathbf{k},\mathbf{k}'} = U - \frac{3J}{4}(c_x c_{x'} + c_x c_{y'} + c_{x'} c_y + c_y c_{y'}) \tag{III.41}$$

$$- \frac{3J}{4}(c_x c_{x'} - c_x c_{y'} - c_{x'} c_y + c_y c_{y'}) \tag{III.42}$$

$$= U - \frac{3J}{2}(c_x c_{x'} + c_y c_{y'}) = V_{\mathbf{k},\mathbf{k}'} \tag{III.43}$$

We call $\frac{3J}{4}(c_x + c_y)(c_{x'} + c_{y'})$ the extended s-wave term. The s-wave term is invariant under 90° rotations of $\mathbf{k}$ or $\mathbf{k}'$, whereas the d-wave term changes sign :

Calculate that

$$V^S_{\mathbf{k},\mathbf{k}'} = V^S_{\mathbf{k}R\mathbf{k}'} \tag{III.44}$$

$$V^D_{\mathbf{k},\mathbf{k}'} = -V^D_{\mathbf{k}R\mathbf{k}'} \tag{III.45}$$

with $R\mathbf{k} = (-k_y, k_x)$. Another point to note is that in the d-wave term, there is no onsite Coulomb interaction. So a condensate with d-wave symmetry ,

Can an s-wave condensate also appear? How is it decided what symmetry the condensate has?

$$\Delta^D_{\mathbf{k}} = \Delta_D(c_x - c_y) \tag{III.46}$$

$$\Delta^D_{R\mathbf{k}} = -\Delta^D_{\mathbf{k}} \tag{III.47}$$

couples to cooper pairs via d-wave interaction , because

What is the relationship between gap and interaction? aka where does this equation come from?

$$\sum_{\mathbf{k}'} V^S_{\mathbf{k},\mathbf{k}'} \Delta^D_{\mathbf{k}'}(\ldots) = 0 \tag{III.48}$$

(see gap equation, it preserves the symmetry of the pair). A condensate with extended s-wave symmetry

$$\Delta^S_{\mathbf{k}} = \Delta_1 + \Delta_2(c_x + c_y) \tag{III.49}$$

vanishes when integrated with the d-wave part of the interaction. This means the two types of pairing are symmetry decoupled and moreover, the symmetry of the d-wave pair decouples against the local Coulomb pseudopotential. The quasiparticle energy for the d-wave condensate is:

What quasiparticle?

$$E_{\mathbf{k}} = \sqrt{\epsilon^2_{\mathbf{k}} + \Delta^2_{\mathbf{k}}(c_y - c_x)^2} \tag{III.50}$$

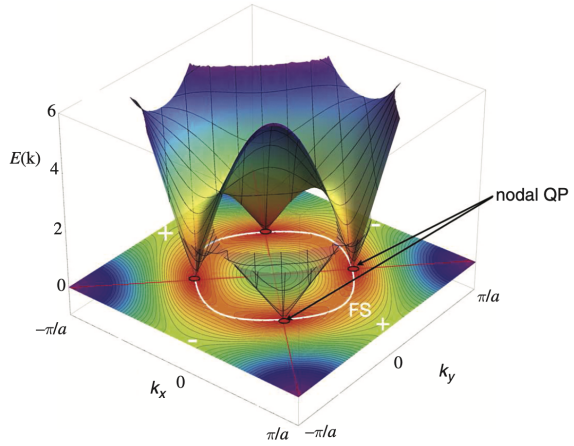

Figure III.1

What exactly is shown in the fig

It vanishes at intersections of nodes (where $\Delta_{\mathbf{k}} = 0$) and the Fermi surface (where $\epsilon_{\mathbf{k}} = 0$). At these points the dispersion can be linearized, they form Dirac cones of excitations with a relativistic dispersion . We can approximately solve the gap equation and get

What is the exact dispersion?

$$\Delta_D(c_y - c_x) = \Delta_D(k_x^2 - k_y^2) = \Delta_0 \cos(2\theta) \tag{III.51}$$

How exactly typical? $l = 2$?

The dependence $\Delta \propto \cos(2\theta)$ is typical for an $l = 2$ Cooper pair. The quasiparticle energy is then

Visualise that somehow?

$$E_{\mathbf{k}} = \sqrt{\epsilon_{\mathbf{k}}^2 + (\Delta_0 \cos(2\theta))^2} \tag{III.52}$$

The d-wave density of states does not have a clear gap, but instead a V-shaped structure. This linear DOS across the gap is due to the Dirac cones.

How does the DOS compare with real materials? Do we have the V-shaped structure?

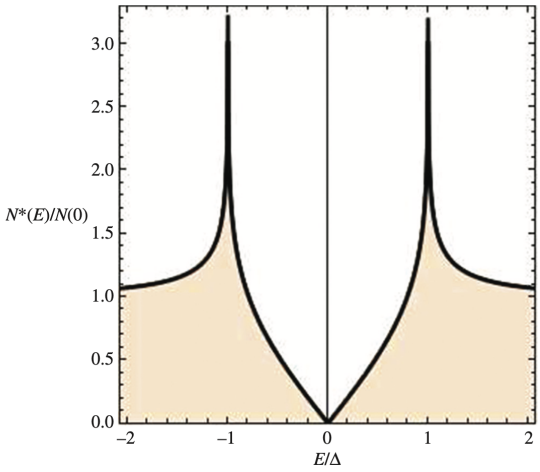

Figure III.2

IV Coherence length and penetration depth in strongly correlated superconductors

Put in source

Order parameter (OP) of a superconducting condensate with FMP has the form

$$\Psi_{\mathbf{q}}(\mathbf{r}) = |\Psi_{\mathbf{q}}| e^{i\mathbf{q}\mathbf{r}} \tag{IV.1}$$

where $\mathbf{q}$ is the center-of-mass momentum of Cooper pairs.

FMP is well known from Fulde-Ferrel-Larkin-Ovchinnikov (FFLO) theory, where the single-momentum phase used here corresponds to FF-type pairing.

What does that mean? More details on FFLO theory

IV.1 Ginzburg-Landau description

First: Motivate how the FMP constraint relates to λ_L and ξ_0.

GL low-order expansion of the free energy density f_{GL} in terms of the FMP-constrained OP reads

$$1 \tag{IV.2}$$

Fill in equation

The temperature dependent correlation length ξ appears as the natural length scale of the amplitude mode ($\propto \alpha$) and kinetic energy term

More details on GL theory in general

$$\xi(T) = \tag{IV.3}$$

with the zero temperature value ξ_0 being the coherence length.

Fill in equation

V Dynamical Mean-Field Theory

Source: Georges et al. - "Dynamical mean-field theory of strongly correlated fermion systems and the limit of infinite dimensions" [3]

With help from [4] and [5] to make it more concise.

Most general non-interacting electronic Hamiltonian in second quantization:

$$H_0 = \sum_{i,j,\sigma} \tag{V.1}$$

with lattice coordinates i, j and spin σ.

One particle Green's function (many-body object, coming from the Hubbard model):

$$G(\mathbf{k}, \mathrm{i}\omega_n) = \frac{1}{\mathrm{i}\omega_n + \mu - \epsilon_{\mathbf{k}} - \Sigma(\mathbf{k}, \mathrm{i}\omega_n)} \tag{V.2}$$

with the self energy $\Sigma(\mathrm{i}\omega_n)$ coming from the solution of the effect on-site problem:

The Dyson equation

$$G(\mathbf{k}, \mathrm{i}\omega_n) = (G_0(\mathbf{k}, \mathrm{i}\omega_n) - \Sigma(\mathbf{k}, \mathrm{i}\omega_n))^{-1} \tag{V.3}$$

relates the non-interacting Greens function $G_0(\mathbf{k}, \mathrm{i}\omega_n)$ and the fully-interacting Greens function $G(\mathbf{k}, \mathrm{i}\omega_n)$ (inversion of a matrix!).

VI Hubbard model

VII EG-X Model

VII.1 Lattice Structure of Graphene

Following [6].

Monolayer graphene forms a hexagonal lattice. This is formed by two triangular sub lattices. So in the unit cell of the hexagonal actually has two atoms.

Primitive lattice vectors of the hexagonal lattice:

$$\mathbf{a}_1 = \frac{a}{2}\begin{pmatrix} 1 \\ \sqrt{3} \end{pmatrix} \tag{VII.1}$$

$$\mathbf{a}_2 = \frac{a}{2}\begin{pmatrix} 1 \\ -\sqrt{3} \end{pmatrix} \tag{VII.2}$$

with lattice constant $a \approx 2.46\,\text{Å}$ (distance between unit cells). Have

$$a = \sqrt{3}a_0 \tag{VII.3}$$

with the nearest-neighbour distance a_0.

Vectors to the nearest-neighbor B_i $(i = 1, 2, 3,)$ atoms from atom A:

$$\boldsymbol{\delta}_{AB,1} = \begin{pmatrix} 0 \\ \frac{a}{\sqrt{3}} \end{pmatrix}, \boldsymbol{\delta}_{AB,2} = \begin{pmatrix} \frac{a}{2} \\ -\frac{a}{2\sqrt{3}} \end{pmatrix}, \boldsymbol{\delta}_{AB,3} = \begin{pmatrix} -\frac{a}{2} \\ -\frac{a}{2\sqrt{3}} \end{pmatrix} \tag{VII.4}$$

Vectors to the nearest-neighbor A_i $(i = 1, 2, 3,)$ atoms from atom B:

$$\boldsymbol{\delta}_{BA,1} = \begin{pmatrix} 0 \\ -\frac{a}{\sqrt{3}} \end{pmatrix}, \boldsymbol{\delta}_{BA,2} = \begin{pmatrix} \frac{a}{2} \\ \frac{a}{2\sqrt{3}} \end{pmatrix}, \boldsymbol{\delta}_{BA,3} = \begin{pmatrix} -\frac{a}{2} \\ \frac{a}{2\sqrt{3}} \end{pmatrix} \tag{VII.5}$$

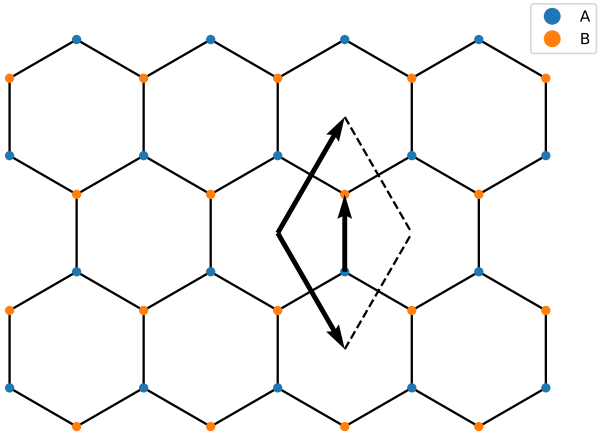

Figure VII.1: *Graphene lattice structure*

The vectors between the Graphene A atom and the six neighbours on the same sub lattice can be found by rotating $\mathbf{a}_1$ six times by $1/6 * 2\pi = \pi/3$:

$$\boldsymbol{\delta}_{AA,1} = \mathbf{a}_1 = \frac{a}{2}\begin{pmatrix}1\\ \sqrt{3}\end{pmatrix} = a\begin{pmatrix}\frac{1}{2}\\ \frac{\sqrt{3}}{2}\end{pmatrix} = a\begin{pmatrix}\sin\left(\frac{\pi}{6}\right)\\ \cos\left(\frac{\pi}{6}\right)\end{pmatrix} \tag{VII.6}$$

$$\boldsymbol{\delta}_{AA,2} = a\begin{pmatrix}\sin\left(\frac{3\pi}{6}\right)\\ \cos\left(\frac{3\pi}{6}\right)\end{pmatrix} = a\begin{pmatrix}1\\ 0\end{pmatrix} \tag{VII.7}$$

$$\boldsymbol{\delta}_{AA,3} = a\begin{pmatrix}\sin\left(\frac{5\pi}{6}\right)\\ \cos\left(\frac{5\pi}{6}\right)\end{pmatrix} = a\begin{pmatrix}\frac{1}{2}\\ -\frac{\sqrt{3}}{2}\end{pmatrix} \tag{VII.8}$$

$$\boldsymbol{\delta}_{AA,4} = a\begin{pmatrix}\sin\left(\frac{7\pi}{6}\right)\\ \cos\left(\frac{7\pi}{6}\right)\end{pmatrix} = a\begin{pmatrix}-\frac{1}{2}\\ -\frac{\sqrt{3}}{2}\end{pmatrix} \tag{VII.9}$$

$$\boldsymbol{\delta}_{AA,5} = a\begin{pmatrix}\sin\left(\frac{9\pi}{6}\right)\\ \cos\left(\frac{9\pi}{6}\right)\end{pmatrix} = a\begin{pmatrix}-1\\ 0\end{pmatrix} \tag{VII.10}$$

$$\boldsymbol{\delta}_{AA,6} = a\begin{pmatrix}\sin\left(\frac{11\pi}{6}\right)\\ \cos\left(\frac{11\pi}{6}\right)\end{pmatrix} = a\begin{pmatrix}-\frac{1}{2}\\ \frac{\sqrt{3}}{2}\end{pmatrix} \tag{VII.11}$$

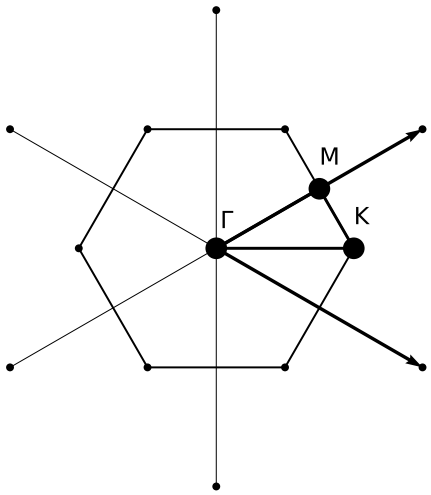

Figure VII.2: *Graphene Brillouin Zone*

The primitive reciprocal lattice vectors $\mathbf{b}_1$, $\mathbf{b}_2$ fulfill

$$\mathbf{a}_1 \cdot \mathbf{b}_1 = \mathbf{a}_2 \cdot \mathbf{b}_2 = 2\pi \tag{VII.12}$$
$$\mathbf{a}_1 \cdot \mathbf{b}_2 = \mathbf{a}_2 \cdot \mathbf{b}_1 = 0\ , \tag{VII.13}$$

so we have:

$$\mathbf{b}_1 = \frac{2\pi}{a} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{3}} \end{pmatrix} \tag{VII.14}$$
$$\mathbf{b}_2 = \frac{2\pi}{a} \begin{pmatrix} 1 \\ -\frac{1}{\sqrt{3}} \end{pmatrix} \tag{VII.15}$$

Points of high symmetry in the Brillouin zone are:

$$\Gamma = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \tag{VII.16}$$

$$\mathrm{M} = \frac{\pi}{a} \begin{pmatrix} 1 \\ \frac{1}{\sqrt{3}} \end{pmatrix} \tag{VII.17}$$

$$\mathrm{K} = \frac{4\pi}{3a} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \tag{VII.18}$$

VII.2 EG-X Model

Graphene lattice and a site X. Real-life motivation: layer of graphene on top

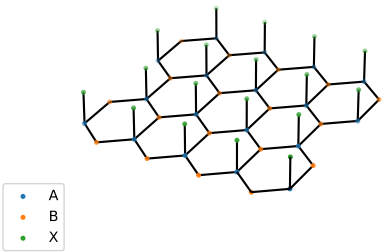

Figure VII.3: *EG-X model*

of a substrate of another material (which provides the additional X atoms). There is no spin-orbit coupling considered in the model (according to Niklas) when mapping to substrates Sn or Pb, it could be necessary (but does not the qualitative result?).

Spin-orbit coupling, drop second spin index?

Without interaction :

$$H_0 = -t_{\mathrm{X}} \sum_{\langle ij\rangle,\sigma\sigma'} d^\dagger_{i,\sigma} d_{j,\sigma'} - t_{\mathrm{Gr}} \sum_{\langle ij\rangle,\sigma\sigma'} \left(c^{(A),\dagger}_{i,\sigma} c^{(B)}_{j,\sigma'} + c^{(B),\dagger}_{j,\sigma'} c^{(A)}_{i,\sigma} \right) + V \sum_{i,\sigma\sigma'} \left(d^\dagger_{i,\sigma} c^{(A)}_{i,\sigma'} + c^{(A),\dagger}_{i,\sigma} d_{i,\sigma'} \right) \tag{VII.19}$$

with:

- d operators on the X atom
- $c^{(\epsilon)}$ operators on the graphene site ($\epsilon = A, B$)
- t_X NN hopping for X
- t_{Gr} NN hopping of Gr
- V hybridization between X and Graphene B sites

We can also introduce an onsite Hubbard interaction:

$$H_{\mathrm{int}} = U_{\mathrm{X}} \sum_i d^\dagger_{i,\uparrow} d^\dagger_{i,\downarrow} d_{i,\downarrow} d_{i,\uparrow} + U_{\mathrm{Gr}} \sum_{i,\epsilon=A,B} c^{(\epsilon)\dagger}_{i,\uparrow} c^{(\epsilon)\dagger}_{i,\downarrow} c^{\epsilon}_{i,\downarrow} c^{\epsilon}_{i,\uparrow} \tag{VII.20}$$

VII.2.1 Review: Hubbard model on the honeycomb lattice

Write review for Hubbard model on the honeycomb lattice

VII.2.2 Band structure of the non-interacting EG-X model

To treat eq. VII.19, we first write out the sums over nearest neighbours $\langle i, j\rangle$ explicitly, writing $\boldsymbol{\delta}_{\mathrm{X}}, \boldsymbol{\delta}_\epsilon$ ($\epsilon = A, B$) for the connections to the nearest neighbours of the X atoms and Graphene A, B sites. Doing the calculation for the example of the X atoms:

$$-t_{\mathrm{X}} \sum_{\langle ij\rangle,\sigma\sigma'} d^\dagger_{i,\sigma} d_{j,\sigma'} = -\frac{t_X}{2} \sum_{i,\sigma,\sigma'} \sum_{\delta_{\mathrm{X}}} d^\dagger_{i,\sigma} d_{i+\delta_{\mathrm{X}},\sigma'} \tag{VII.21}$$

$$\tag{VII.22}$$

(The factor $1/2$ is to account for double counting when going to the sum over all lattice sites i)

Now we can input the discrete Fourier transform (for both graphene and X operators) into eq. VII.21

$$c_i = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{\mathrm{i}\mathbf{k}\mathbf{r}_i} c_{\mathbf{k}} \tag{VII.23}$$

$$c_i^\dagger = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{-\mathrm{i}\mathbf{k}\mathbf{r}_i} c_{\mathbf{k}}^\dagger \tag{VII.24}$$

with the completeness relation:

$$\sum_i e^{\mathrm{i}\mathbf{k}\mathbf{r}_i} e^{-\mathrm{i}\mathbf{k}'\mathbf{r}_i} = N\delta_{\mathbf{k},\mathbf{k}'} \ . \tag{VII.25}$$

We get:

$$-\frac{t_X}{2}\frac{1}{N}\sum_{i,\sigma,\sigma'}\sum_{\boldsymbol{\delta}_\mathrm{X}} d_{i,\sigma}^\dagger d_{i+\boldsymbol{\delta}_\mathrm{X},\sigma'} = -\frac{t_X}{2}\frac{1}{N}\sum_{i,\sigma,\sigma'}\sum_{\boldsymbol{\delta}_\mathrm{X}}\sum_{\mathbf{k},\mathbf{k}'} e^{-\mathrm{i}\mathbf{k}\mathbf{r}_i} d_{\mathbf{k},\sigma}^\dagger e^{\mathrm{i}\mathbf{k}'\mathbf{r}_i} e^{\mathrm{i}\mathbf{k}'\boldsymbol{\delta}_\mathrm{X}} d_{\mathbf{k}',\sigma'} \tag{VII.26}$$

$$= -\frac{t_X}{2}\frac{1}{N}\sum_{\mathbf{k},\mathbf{k}',\sigma,\sigma'}\sum_{\boldsymbol{\delta}_\mathrm{X}} d_{\mathbf{k},\sigma}^\dagger e^{\mathrm{i}\mathbf{k}'\boldsymbol{\delta}_\mathrm{X}} d_{\mathbf{k}',\sigma'} \sum_i e^{-\mathrm{i}\mathbf{k}\mathbf{r}_i} e^{\mathrm{i}\mathbf{k}'\mathbf{r}_i} \tag{VII.27}$$

$$= -\frac{t_X}{2}\frac{1}{N}\sum_{\mathbf{k},\mathbf{k}',\sigma,\sigma'}\sum_{\boldsymbol{\delta}_\mathrm{X}} d_{\mathbf{k},\sigma}^\dagger e^{\mathrm{i}\mathbf{k}'\boldsymbol{\delta}_\mathrm{X}} d_{\mathbf{k}',\sigma'} N\delta_{\mathbf{k},\mathbf{k}'} \tag{VII.28}$$

$$= -\frac{t_X}{2}\sum_{\mathbf{k},\sigma,\sigma'} d_{\mathbf{k},\sigma}^\dagger d_{\mathbf{k},\sigma'} \sum_{\boldsymbol{\delta}_\mathrm{X}} e^{\mathrm{i}\mathbf{k}\boldsymbol{\delta}_\mathrm{X}} \tag{VII.29}$$

The nearest neighbours for X atoms are the vectors $\boldsymbol{\delta}_{AA,i}$ from section VII.1. With that, we can calculate:

$$f_{\mathrm{X}}(\mathbf{k}) = -\frac{t_X}{2}\sum_{\boldsymbol{\delta}_{\mathrm{X}}} e^{\mathrm{i}\mathbf{k}\boldsymbol{\delta}_{\mathrm{X}}} \tag{VII.30}$$

$$= -\frac{t_X}{2}\left(e^{\mathrm{i}a(\frac{k_x}{2}+\frac{\sqrt{3}k_y}{2})} + e^{\mathrm{i}ak_x} + e^{\mathrm{i}a(\frac{k_x}{2}-\frac{\sqrt{3}k_y}{2})}\right. \tag{VII.31}$$

$$\left. + e^{\mathrm{i}a(-\frac{k_x}{2}-\frac{\sqrt{3}k_y}{2})} + e^{-\mathrm{i}ak_x} + e^{\mathrm{i}a(-\frac{k_x}{2}+\frac{\sqrt{3}k_y}{2})}\right) \tag{VII.32}$$

$$= -\frac{t_X}{2}\left(2\cos(ak_x) + 2e^{\mathrm{i}a\frac{\sqrt{3}k_y}{2}}\cos(\frac{a}{2}k_x) + 2e^{-\mathrm{i}a\frac{\sqrt{3}k_y}{2}}\cos(\frac{a}{2}k_x)\right) \tag{VII.33}$$

$$= -t_X\left(\cos(ak_x) + 2\cos(\frac{a}{2}k_x)\cos(\sqrt{3}\frac{a}{2}k_y)\right) \tag{VII.34}$$

We can do the same for the hopping between Graphene sites, for example :

$$-t_{\mathrm{Gr}}\sum_{\langle ij\rangle,\sigma\sigma'} c_{i,\sigma}^{(A),\dagger}c_{j,\sigma'}^{(B)} = -\frac{t_{\mathrm{Gr}}}{2}\sum_{i,\sigma\sigma'}\sum_{\boldsymbol{\delta}_{AB}} c_{i,\sigma}^{(A),\dagger}c_{i+\boldsymbol{\delta}_{AB},\sigma'}^{(B)} \tag{VII.35}$$

$$= -\frac{t_{\mathrm{Gr}}}{2}\sum_{\mathbf{k},\sigma,\sigma'} c_{\mathbf{k},\sigma}^{(A)\dagger}c_{\mathbf{k},\sigma'}^{(B)}\sum_{\boldsymbol{\delta}_{AB}} e^{\mathrm{i}\mathbf{k}\boldsymbol{\delta}_{AB}} \tag{VII.36}$$

We note

$$\sum_{\boldsymbol{\delta}_{AB}} e^{\mathrm{i}\mathbf{k}\boldsymbol{\delta}_{AB}} = \left(\sum_{\boldsymbol{\delta}_{BA}} e^{\mathrm{i}\mathbf{k}\boldsymbol{\delta}_{BA}}\right)^* = \sum_{\boldsymbol{\delta}_{BA}} e^{-\mathrm{i}\mathbf{k}\boldsymbol{\delta}_{BA}} \tag{VII.37}$$

and calculate

$$f_{Gr} = -\frac{t_{Gr}}{2}\sum_{\boldsymbol{\delta}_{AB}} e^{\mathrm{i}\mathbf{k}\boldsymbol{\delta}_{AB}} \tag{VII.38}$$

$$= -\frac{t_{Gr}}{2}\left(e^{\mathrm{i}\frac{a}{\sqrt{3}}k_y} + e^{\mathrm{i}\frac{a}{2\sqrt{3}}(\sqrt{3}k_x-k_y)} + e^{\mathrm{i}\frac{a}{2\sqrt{3}}(-\sqrt{3}k_x-k_y)}\right) \tag{VII.39}$$

$$= -\frac{t_{Gr}}{2}\left(e^{\mathrm{i}\frac{a}{\sqrt{3}}k_y} + e^{-\mathrm{i}\frac{a}{2\sqrt{3}}k_y}\left(e^{\mathrm{i}\frac{a}{2}k_x} + e^{-\mathrm{i}\frac{a}{2}k_x}\right)\right) \tag{VII.40}$$

$$= -\frac{t_{Gr}}{2}\left(e^{\mathrm{i}\frac{a}{\sqrt{3}}k_y} + 2e^{-\mathrm{i}\frac{a}{2\sqrt{3}}k_y}\cos(\frac{a}{2}k_x)\right) \tag{VII.41}$$

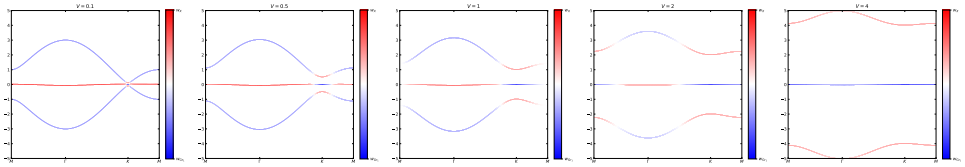

Figure VII.4: *Bands of the non-interacting EG-X model. All the bands are spin-degenerate.*

All together, we get:

$$H_0 = \sum_{\mathbf{k},\sigma,\sigma'} \tag{VII.42}$$

$$= \sum_{\mathbf{k},\sigma,\sigma'} \begin{pmatrix} c^{A,\dagger}_{k,\sigma} & c^{B,\dagger}_{k,\sigma} & d^{\dagger}_{k,\sigma} \end{pmatrix} \begin{pmatrix} 0 & f_{Gr} & V \\ f^{*}_{Gr} & 0 & 0 \\ V & 0 & f_X \end{pmatrix} \begin{pmatrix} c^{A}_{k,\sigma} \\ c^{B}_{k,\sigma} \\ d_{k,\sigma} \end{pmatrix} \tag{VII.43}$$

The band structure for the non-interacting EG-X model is easily obtained by diagonalising the matrix in eq. VII.43. This was done in fig. VII.4.

Values used for calculation:

- $a_0 = 1$
- $t_{\mathrm{Gr}} = 1$
- $t_{\mathrm{X}} = 0.01$

V is the control parameter. (According to Niklas), a range from $V = 0.1$ to $V = 2$ can be crudely onto real experimental materials.

VII.3 BCS Theory on the EG-X Model

VII.3.1 BdG Hamiltonian

Define sublattice index

$$\alpha = 1, 2, 3 \tag{VII.44}$$

with $1 \mathrel{\widehat{=}} \mathrm{Gr}_1, 2 \mathrel{\widehat{=}} \mathrm{Gr}_2, 3 \mathrel{\widehat{=}} \mathrm{X}$. Then we can write the non-interacting term as

$$H_0 = -\sum_{\langle i,j\rangle,\alpha,\beta,\sigma} [\mathbf{t}]_{i\alpha,j\beta}\, c^\dagger_{i\alpha} c_{j\beta} \tag{VII.45}$$

with the matrix

Does that make sense?

$$\mathbf{t} = \begin{pmatrix} 0 & t_{\mathrm{Gr}} & 0 \\ t_{\mathrm{Gr}} & 0 & -V\delta_{ij} \\ 0 & -V\delta_{ij} & t_{\mathrm{X}} \end{pmatrix} \tag{VII.46}$$

Add chemical potential (to control filling $n = \frac{N_\uparrow + N_\downarrow}{N_{\mathrm{sites}}}$):

$$-\mu \sum_{i\alpha\sigma} n_{i\alpha\sigma} \tag{VII.47}$$

Also write the interaction part with α (with changed signs -> to keep in line with papers about the attractive Hubbard model):

$$H_{int} = -\sum_{i\alpha} U_\alpha c^\dagger_{i\alpha\uparrow} c^\dagger_{i\alpha\downarrow} c_{i\alpha\downarrow} c_{i\alpha\uparrow} \tag{VII.48}$$

Fourier transformation:

$$H_{int} = -\frac{1}{N^2} \sum_{\alpha,\mathbf{k}_{1,2,3,4}} U_\alpha e^{\mathrm{i}(\mathbf{k}_1+\mathbf{k}_4-\mathbf{k}_1-\mathbf{k}_3)r_{i\alpha}} c^\dagger_{\mathbf{k}_1\alpha\uparrow} c^\dagger_{\mathbf{k}_3\alpha\downarrow} c_{\mathbf{k}_2\alpha\downarrow} c_{\mathbf{k}_4\alpha\uparrow} \tag{VII.49}$$

Imposing zero-momentum pairing: $\mathbf{k}_1 + \mathbf{k}_3 = 0$ and $\mathbf{k}_2 + \mathbf{k}_4 = 0$:

$$H_{int} = -\sum_{\alpha,\mathbf{k},\mathbf{k}'} U_\alpha c^\dagger_{\mathbf{k}\alpha\uparrow} c^\dagger_{-\mathbf{k}\alpha\downarrow} c_{-\mathbf{k}'\alpha\downarrow} c_{\mathbf{k}'\alpha\uparrow} \tag{VII.50}$$

Mean-field approximation:

$$H_{int} \approx \sum_{\alpha,\mathbf{k}} (\Delta_\alpha c^\dagger_{\mathbf{k}\alpha\uparrow} c^\dagger_{-\mathbf{k}\alpha\downarrow} + \Delta^*_\alpha c_{-\mathbf{k}\alpha\downarrow} c_{\mathbf{k}\alpha\uparrow}) \tag{VII.51}$$

with

$$\Delta_\alpha = -U_\alpha \sum_{\mathbf{k}'} \langle c_{-\mathbf{k}'\alpha\downarrow} c_{\mathbf{k}'\alpha\uparrow} \rangle \tag{VII.52}$$

$$\Delta^*_\alpha = -U_\alpha \sum_{\mathbf{k}'} \langle c^\dagger_{\mathbf{k}'\alpha\uparrow} c^\dagger_{-\mathbf{k}'\alpha\downarrow} \rangle \tag{VII.53}$$

This gives the BCS mean field Hamiltonian:

$$H_{BCS} = \sum_{\mathbf{k}\alpha\beta\sigma} [H_{0,\sigma}(\mathbf{k})]_{\alpha\beta} c^\dagger_{\mathbf{k}\alpha\sigma} c_{\mathbf{k}\beta\sigma} - \mu \sum_{\mathbf{k}\alpha\sigma} n_{\mathbf{k}\alpha\sigma} + \sum_{\alpha,\mathbf{k}} (\Delta_\alpha c^\dagger_{\mathbf{k}\alpha\uparrow} c^\dagger_{-\mathbf{k}\alpha\downarrow} + \Delta^*_\alpha c_{-\mathbf{k}\alpha\downarrow} c_{\mathbf{k}\alpha\uparrow}) \tag{VII.54}$$

with Nambu spinor

$$\Psi_{\mathbf{k}} = \begin{pmatrix} c_{1,\mathbf{k}\uparrow} \\ c_{2,\mathbf{k}\uparrow} \\ c_{3,\mathbf{k}\uparrow} \\ c^\dagger_{1,-\mathbf{k}\downarrow} \\ c^\dagger_{2,-\mathbf{k}\downarrow} \\ c^\dagger_{3,-\mathbf{k}\downarrow} \end{pmatrix} \tag{VII.55}$$

we have:

$$H_{MF} = \sum_{\mathbf{k}} \Psi^\dagger_{\mathbf{k}} \mathcal{H}(\mathbf{k}) \Psi_{\mathbf{k}} \tag{VII.56}$$

with

$$\mathcal{H}(\mathbf{k}) = \begin{pmatrix} H_{0,\uparrow}(\mathbf{k}) - \mu & \Delta \\ \Delta^\dagger & -H^*_{0,\downarrow}(-\mathbf{k}) + \mu \end{pmatrix} \tag{VII.57}$$

with $H_{0,\sigma}$ being the F.T. of the kinetic term and $\Delta = diag(\Delta_1, \Delta_2, \Delta_3)$.

VII.3.2 BdG Hamiltonian in band basis

Use transformation

$$c^\dagger_{\mathbf{k}\alpha\sigma} = \sum_n [\mathbf{G}]^*_{\alpha n} d^\dagger_{n\mathbf{k}\sigma} \tag{VII.58}$$

where the columns are made up of the eigenvectors of $\mathbf{H}_{0,\sigma}$ for a given $\mathbf{k}$:

$$\mathbf{G} = \begin{pmatrix} \mathbf{G}_1 & \mathbf{G}_2 & \mathbf{G}_3 \end{pmatrix} \tag{VII.59}$$

with that:

$$\mathbf{G}^\dagger_\sigma(\mathbf{k}) \mathbf{H}_{0,\sigma}(\mathbf{k}) \mathbf{G}_\sigma(\mathbf{k}) = \begin{pmatrix} \epsilon_1 & 0 & 0 \\ 0 & \epsilon_2 & 0 \\ 0 & 0 & \epsilon_3 \end{pmatrix} \tag{VII.60}$$

So the kinetic part of the BdG Hamiltonian becomes:

$$\sum_{\mathbf{k}\alpha\beta\sigma}[H_{0,\sigma}(\mathbf{k})]_{\alpha\beta}\sum_{n}[\mathbf{G}(\mathbf{k})]^*_{\alpha n}d^\dagger_{n\mathbf{k}\sigma}\sum_{m}[\mathbf{G}(\mathbf{k})]_{\beta m}d_{m\mathbf{k}\sigma}-\mu\sum_{\mathbf{k}\alpha\sigma}n_{n\mathbf{k}\sigma} \tag{VII.61}$$

$$=\sum_{mn\mathbf{k}\sigma}d^\dagger_{n\mathbf{k}\sigma}d_{m\mathbf{k}\sigma}\sum_{\alpha\beta}[\mathbf{G}(\mathbf{k})]^*_{\alpha n}[H_{0,\sigma}(\mathbf{k})]_{\alpha\beta}[\mathbf{G}(\mathbf{k})]_{\beta m}-\mu\sum_{\mathbf{k}\alpha\sigma}n_{n\mathbf{k}\sigma} \tag{VII.62}$$

$$=\sum_{mn\mathbf{k}\sigma}d^\dagger_{n\mathbf{k}\sigma}d_{m\mathbf{k}\sigma}\epsilon_n\delta_{nm}-\mu\sum_{\mathbf{k}\alpha\sigma}n_{n\mathbf{k}\sigma} \tag{VII.63}$$

$$=\sum_{n\mathbf{k}\sigma}\epsilon_n d^\dagger_{n\mathbf{k}\sigma}d_{n\mathbf{k}\sigma}-\mu\sum_{\mathbf{k}\alpha\sigma}n_{n\mathbf{k}\sigma} \tag{VII.64}$$

$$=:\sum_{n\mathbf{k}\sigma}\xi_{\mathbf{k}}d^\dagger_{n\mathbf{k}\sigma}d_{n\mathbf{k}\sigma} \tag{VII.65}$$

with $\xi_{\mathbf{k}} \coloneqq \epsilon_{\mathbf{k}}-\mu$. The pairing terms become (only looking at pairing of two electrons in the same band?):

$$\sum_{\mathbf{k}\alpha}\Delta_\alpha c^\dagger_{\mathbf{k}\alpha\uparrow}c^\dagger_{-\mathbf{k}\alpha\downarrow}=\sum_{\mathbf{k}\alpha}\Delta_\alpha\sum_{n}[\mathbf{G}_\uparrow(\mathbf{k})]^*_{\alpha n}d^\dagger_{n\mathbf{k}\uparrow}\sum_{m}[\mathbf{G}_\downarrow(-\mathbf{k})]^*_{\beta m}d^\dagger_{m-\mathbf{k}\downarrow} \tag{VII.66}$$

$$=-\sum_{n\mathbf{k}}\Delta_n d^\dagger_{n\mathbf{k}\uparrow}d^\dagger_{n-\mathbf{k}\downarrow} \tag{VII.67}$$

with gap $\Delta_n(\mathbf{k})=-\sum_\alpha[\mathbf{G}_\uparrow(\mathbf{k})]^*_{\alpha n}\Delta_\alpha[\mathbf{G}_\downarrow(-\mathbf{k})]^*_{\alpha n}$ for band n.

$$\sum_{\mathbf{k}\alpha}\Delta^*_\alpha c_{-\mathbf{k}\alpha\downarrow}c_{\mathbf{k}\alpha\uparrow}=-\sum_{n\mathbf{k}}\Delta^*_n d_{n-\mathbf{k}\downarrow}d_{n\mathbf{k}\uparrow} \tag{VII.68}$$

VII.3.3 Self-consistent calculation of the superconducting gaps

Compare [1, ch. 10]. Notable here: Multiple bands, and the gaps in each band depend in a complicated manner on the parameters U_α and the orbital Green's functions.

Define normal Green's function:

$$\mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau)=-\langle T_\tau d_{n\mathbf{k}\uparrow}(\tau)d^\dagger_{n\mathbf{k}\uparrow}(0)\rangle \tag{VII.69}$$

Anomalous Green's function:

$$\mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau)=-\langle T_\tau d_{n-\mathbf{k}\downarrow}(\tau)d^\dagger_{n\mathbf{k}\uparrow}(0)\rangle \tag{VII.70}$$

Equations of motion (Heisenberg equation), follow [1, ch. 17]:

$$\partial_\tau \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) = -\delta(\tau) + \langle T_\tau \left[d_{n\mathbf{k}\uparrow}, H\right](\tau) d^\dagger_{n\mathbf{k}\uparrow}(0)\rangle \tag{VII.71}$$

$$\partial_\tau \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) = \langle T_\tau \left[d_{n-\mathbf{k}\downarrow}, H\right](\tau) d^\dagger_{n\mathbf{k}\uparrow}(0)\rangle \tag{VII.72}$$

To calculate the commutators, use the relation (for operators A, B, C):

$$[A, BC] = ABC - BCA = (\{A,B\} - BA)C - B(\{C,A\} - AC) \tag{VII.73}$$

$$\left[d^\dagger_{n-\mathbf{k}\downarrow}, H_0\right] = \sum_{n'\mathbf{k}'\sigma'} \xi_{n'\mathbf{k}'} \left[d^\dagger_{n-\mathbf{k}\downarrow}, d^\dagger_{n'\mathbf{k}'\sigma'} d_{n'\mathbf{k}'\sigma'}\right] \tag{VII.74}$$

$$= \sum_{n'\mathbf{k}'\sigma'} \xi_{n'\mathbf{k}'} \left(\{d^\dagger_{n-\mathbf{k}\downarrow}, d^\dagger_{n'\mathbf{k}'\sigma'}\} - d^\dagger_{n'\mathbf{k}'\sigma'} d^\dagger_{n-\mathbf{k}\downarrow}\right) d_{n'\mathbf{k}'\sigma'} \tag{VII.75}$$

$$- d^\dagger_{n'\mathbf{k}'\sigma'} \left(\{d_{n'\mathbf{k}'\sigma'}, d^\dagger_{n-\mathbf{k}\downarrow}\} - d^\dagger_{n-\mathbf{k}\downarrow} d_{n'\mathbf{k}'\sigma'}\right) \tag{VII.76}$$

$$= \sum_{n'\mathbf{k}'\sigma'} \xi_{n'\mathbf{k}'} \left(-d^\dagger_{n'\mathbf{k}'\sigma'} d^\dagger_{n-\mathbf{k}\downarrow} d_{n'\mathbf{k}'\sigma'} - d^\dagger_{n'\mathbf{k}'\sigma'} \delta_{n'\mathbf{k}'\sigma',n-\mathbf{k}\uparrow} + d^\dagger_{n'\mathbf{k}'\sigma'} d^\dagger_{n-\mathbf{k}\downarrow} d_{n'\mathbf{k}'\sigma'}\right) \tag{VII.77}$$

$$= -\xi_{n\mathbf{k}} d^\dagger_{n\mathbf{k}\uparrow} \tag{VII.78}$$

$$\left[d_{n-\mathbf{k}\downarrow}, -\sum_{m\mathbf{k}'} \Delta^*_m d_{m-\mathbf{k}'\downarrow} d_{m\mathbf{k}'\uparrow}\right] \tag{VII.79}$$

$$= -\sum_{m\mathbf{k}'} \Delta^*_m \left(\{d_{n-\mathbf{k}\downarrow}, d_{m-\mathbf{k}'\downarrow}\} - d_{m-\mathbf{k}'\downarrow} d_{n-\mathbf{k}\downarrow}\right) d_{m\mathbf{k}'\uparrow} \tag{VII.80}$$

$$- d_{m-\mathbf{k}'\downarrow} \left(\{d_{m\mathbf{k}'\uparrow}, d_{n-\mathbf{k}\downarrow}\} - d_{n-\mathbf{k}\downarrow} d_{m\mathbf{k}'\uparrow}\right) \tag{VII.81}$$

$$= -\sum_{m\mathbf{k}'} \Delta^*_m \left(\delta_{n-\mathbf{k}\downarrow,m-\mathbf{k}'\downarrow} - d_{m-\mathbf{k}'\downarrow} d_{n-\mathbf{k}\downarrow}\right) d_{m\mathbf{k}'\uparrow} + d_{m-\mathbf{k}'\downarrow} d_{n-\mathbf{k}\downarrow} d_{m\mathbf{k}'\uparrow} \tag{VII.82}$$

$$= -\Delta^*_n d_{n\mathbf{k}\uparrow} \tag{VII.83}$$

$$\partial_\tau \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) = -\xi_{n\mathbf{k}} \left\langle T_\tau (d^\dagger_{n-\mathbf{k}\downarrow}(\tau) d^\dagger_{n\mathbf{k}\uparrow}(0))\right\rangle - \Delta_n^* \left\langle T_\tau (d_{n\mathbf{k}\uparrow}(\tau) d^\dagger_{n\mathbf{k}\uparrow}(0))\right\rangle \tag{VII.84}$$

$$= \xi_{n\mathbf{k}} \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) + \Delta_n^* \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) \tag{VII.85}$$

Similarly:

$$[d_{n-\mathbf{k}\uparrow}, H_0] = \sum_{n'\mathbf{k}'\sigma'} \xi_{n'\mathbf{k}'} \left[d^\dagger_{n-\mathbf{k}\downarrow}, d^\dagger_{n'\mathbf{k}'\sigma'} d_{n'\mathbf{k}'\sigma'} \right] \tag{VII.86}$$

$$= \xi_n d^\dagger_{n\mathbf{k}\uparrow} \tag{VII.87}$$

$$\left[d_{n-\mathbf{k}\uparrow}, -\sum_{m\mathbf{k}'} \Delta_m d^\dagger_{m-\mathbf{k}'\uparrow} d^\dagger_{m-\mathbf{k}'\downarrow} \right] \tag{VII.88}$$

$$= -\Delta_n d^\dagger_{n-\mathbf{k}\downarrow} \tag{VII.89}$$

$$\partial_\tau \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) = -\delta(\tau) + \xi_{n\mathbf{k}} \left\langle T_\tau d_{n\mathbf{k}\uparrow}(\tau) d^\dagger_{n\mathbf{k}\uparrow}\right\rangle - \Delta_n \left\langle T_\tau d_{n-\mathbf{k}\downarrow}(\tau) d^\dagger_{n\mathbf{k}\uparrow}(0)\right\rangle \tag{VII.90}$$

$$= -\delta(\tau) - \xi_{n\mathbf{k}} \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) + \Delta_n \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) \tag{VII.91}$$

$$\tag{VII.92}$$

All in all:

$$\partial_\tau \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) = -\delta(\tau) - \xi_{n\mathbf{k}} \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) + \Delta_n \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) \tag{VII.93}$$

$$\partial_\tau \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) = \xi_{n\mathbf{k}} \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\tau) + \Delta_n^* \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\tau) \tag{VII.94}$$

Fourier transform:

$$(-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}) \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\mathrm{i}\omega_n) = -1 + \Delta_n \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\mathrm{i}\omega_n) \tag{VII.95}$$

$$(-\mathrm{i}\omega_n - \xi_{n\mathbf{k}}) \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\mathrm{i}\omega_n) = \Delta_n^* \mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k},\mathrm{i}\omega_n) \tag{VII.96}$$

This algebraic expression can be easily solved:

$$(-\mathrm{i}\omega_n - \xi_{n\mathbf{k}}) \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\mathrm{i}\omega_n) = \frac{\Delta_n^*}{-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}} (-1 + \Delta_n \mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k},\mathrm{i}\omega_n)) \tag{VII.97}$$

$$(-\mathrm{i}\omega_n - \xi_{n\mathbf{k}} - \frac{|\Delta_n|^2}{-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}})\mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k}, \mathrm{i}\omega_n) = \frac{-\Delta_n^*}{-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}} \tag{VII.98}$$

$$(\frac{(-\mathrm{i}\omega_n - \xi_{n\mathbf{k}})(-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}) - |\Delta_n|^2}{-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}})\mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k}, \mathrm{i}\omega_n) = \frac{-\Delta_n^*}{-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}} \tag{VII.99}$$

$$\mathcal{F}_{n\downarrow n\uparrow}(\mathbf{k}, \mathrm{i}\omega_n) = \frac{-\Delta_n^*}{(-\mathrm{i}\omega_n - \xi_{n\mathbf{k}})(-\mathrm{i}\omega_n + \xi_{n\mathbf{k}}) - |\Delta_n|^2} \tag{VII.100}$$

$$= \frac{-\Delta_n^*}{(\mathrm{i}\omega_n)^2 - \xi_{n\mathbf{k}}^2 - |\Delta_n|^2} \tag{VII.101}$$

$$= \frac{-\Delta_n^*}{(\mathrm{i}\omega_n)^2 - E_{n\mathbf{k}}} \tag{VII.102}$$

$$(-\mathrm{i}\omega_n + \xi_{n\mathbf{k}})\mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k}, \mathrm{i}\omega_n) = -1 + \frac{-|\Delta_n|^2}{(\mathrm{i}\omega_n)^2 - \xi_{n\mathbf{k}}^2 - |\Delta_n|^2} \tag{VII.103}$$

$$= \frac{-(\mathrm{i}\omega_n)^2 + \xi_{n\mathbf{k}}^2 + |\Delta_n|^2 - |\Delta_n|^2}{(\mathrm{i}\omega_n)^2 - \xi_{n\mathbf{k}}^2 - |\Delta_n|^2} \tag{VII.104}$$

$$= \frac{-(\mathrm{i}\omega_n)^2 + \xi_{n\mathbf{k}}^2}{(\mathrm{i}\omega_n)^2 - \xi_{n\mathbf{k}}^2 - |\Delta_n|^2} \tag{VII.105}$$

$$= \frac{(\mathrm{i}\omega_n + \xi_{n\mathbf{k}})(-\mathrm{i}\omega_n + \xi_{n\mathbf{k}})}{(\mathrm{i}\omega_n)^2 - \xi_{n\mathbf{k}}^2 - |\Delta_n|^2} \tag{VII.106}$$

$$\mathcal{G}_{n\uparrow n\uparrow}(\mathbf{k}, \mathrm{i}\omega_n) = \frac{\mathrm{i}\omega + \xi_{n\mathbf{k}}}{(\mathrm{i}\omega_n)^2 - \xi_{n\mathbf{k}}^2 - |\Delta_n|^2} \tag{VII.107}$$

$$= \frac{\mathrm{i}\omega + \xi_{n\mathbf{k}}}{(\mathrm{i}\omega_n)^2 - E_{n\mathbf{k}}} \tag{VII.108}$$

with the energies $E_{n\mathbf{k}} = \pm\sqrt{\xi_{n\mathbf{k}}^2 + |\Delta_n|^2}$.

To calculate the band gap in band n:

$$\Delta_n(\mathbf{k}) = -\sum_{\alpha}[G_{k\uparrow}]^*_{\alpha n}\Delta_\alpha[G_{-k\downarrow}]^*_{\alpha n} \tag{VII.109}$$

$$= -\sum_{\alpha\mathbf{k}'} U_\alpha[G_{k\uparrow}]^*_{\alpha n}\left\langle c_{-k'\alpha\downarrow}c_{k'\alpha\uparrow}\right\rangle[G_{-k\downarrow}]^*_{\alpha n} \tag{VII.110}$$

$$= -\sum_{\alpha\mathbf{k}'} U_\alpha[G_{k\uparrow}]^*_{\alpha n}[G_{-k\downarrow}]^*_{\alpha n}\sum_m[G_{-k'\downarrow}]_{\alpha m}[G_{k'\uparrow}]_{\alpha m}\left\langle d_{-k'm\downarrow}d_{k'm\uparrow}\right\rangle \tag{VII.111}$$

Can now use $\mathcal{F}$ and fourier-transform:

$$\left\langle d_{-k'm\downarrow}d_{k'm\uparrow}\right\rangle = \mathcal{F}^*_{m\downarrow m\uparrow}(\mathbf{k}', \tau = 0^+) \tag{VII.112}$$

$$= \frac{1}{\beta}\sum_{\mathrm{i}\omega_n} e^{-\mathrm{i}\omega_n 0^+}\mathcal{F}^*_{m\downarrow m\uparrow}(\mathbf{k}', \mathrm{i}\omega_n) \tag{VII.113}$$

The summation over the Matsubara frequencies can be solved via the Residue theorem (the poles z_0 of $\mathcal{F}$ are the energies $\pm E_{m\mathbf{k}}$):

$$\frac{1}{\beta}\sum_{\mathrm{i}\omega_n} e^{-\mathrm{i}\omega_n 0^+}\mathcal{F}^*_{m\downarrow m\uparrow}(\mathbf{k}', \mathrm{i}\omega_n) \tag{VII.114}$$

$$= \sum_{z_0 \mathrm{poles of} \mathcal{F}} e^{-z_0 0^+} n_F(z_0) Res_{z_0}\mathcal{F}^*_{m\downarrow m\uparrow}(\mathbf{k}', z_0) \tag{VII.115}$$

$$= e^{-E_{m\mathbf{k}}0^+}n_F(E_{m\mathbf{k}})Res_{E_{m\mathbf{k}}}\frac{-\Delta_m}{(\mathrm{i}\omega_n)^2 - E_{m\mathbf{k}}} + e^{E_{m\mathbf{k}}0^+}n_F(-E_{m\mathbf{k}})Res_{-E_{m\mathbf{k}}}\frac{-\Delta_m}{(\mathrm{i}\omega_n)^2 - E_{m\mathbf{k}}} \tag{VII.116}$$

with residue:

$$Res_{E_{m\mathbf{k}}}\frac{1}{(\mathrm{i}\omega_n)^2 - z_0^2} = \frac{1}{\partial_z|_{z_0=E_{m\mathbf{k}}}((\mathrm{i}\omega)^2 - z_0^2)} = \frac{1}{2E_{m\mathbf{k}}} \tag{VII.117}$$

So we have

$$\left\langle d_{-k'm\downarrow}d_{k'm\uparrow}\right\rangle = -\Delta_m\left(\frac{n_F(E_{m\mathbf{k}})}{2E_{m\mathbf{k}}} - \frac{n_F(-E_{m\mathbf{k}})}{2E_{m\mathbf{k}}}\right) \tag{VII.118}$$

The n_F term can be written as:

$$n_F(E_{m\mathbf{k}'}) - n_F(-E_{m\mathbf{k}'}) = \frac{1}{e^{\beta E_{m\mathbf{k}'}}+1} - \frac{1}{e^{-\beta E_{m\mathbf{k}'}}+1} \tag{VII.119}$$

$$= \frac{e^{-\frac{1}{2}\beta E_{m\mathbf{k}'}}}{e^{-\frac{1}{2}\beta E_{m\mathbf{k}'}}}\frac{1}{e^{\beta E_{m\mathbf{k}'}}+1} - \frac{e^{\frac{1}{2}\beta E_{m\mathbf{k}'}}}{e^{\frac{1}{2}\beta E_{m\mathbf{k}'}}}\frac{1}{e^{-\beta E_{m\mathbf{k}'}}+1} \tag{VII.120}$$

$$= \frac{e^{-\frac{1}{2}\beta E_{m\mathbf{k}'}} - e^{\frac{1}{2}\beta E_{m\mathbf{k}'}}}{e^{\frac{1}{2}\beta E_{m\mathbf{k}'}} + e^{-\frac{1}{2}\beta E_{m\mathbf{k}'}}} \tag{VII.121}$$

$$= -\tanh\left(\frac{\beta E_{m\mathbf{k}'}}{2}\right) \tag{VII.122}$$

This results in the self-concistency equation for the gap:

$$\Delta_n(\mathbf{k}) = \sum_{\alpha m \mathbf{k}'} U_\alpha [G_{k\uparrow}]^*_{\alpha n}[G_{-k\downarrow}]^*_{\alpha n}[G_{-k'\downarrow}]_{\alpha m}[G_{k'\uparrow}]_{\alpha m}\Delta_m(\mathbf{k}')\frac{\tanh\left(\frac{\beta E_{m\mathbf{k}'}}{2}\right)}{E_{m\mathbf{k}'}} \tag{VII.123}$$

Using time-reversal symmetry $[G_{-\mathbf{k}\downarrow}]^*_{\alpha m} = [G_{\mathbf{k}\uparrow}]_{\alpha m}$ this expression gets a bit simpler:

$$\Delta_n(\mathbf{k}) = \sum_{\alpha m \mathbf{k}'} U_\alpha |[G_{k\uparrow}]_{\alpha n}|^2 |[G_{k'\uparrow}]_{\alpha m}|^2 \Delta_m(\mathbf{k}')\frac{\tanh\left(\frac{\beta E_{m\mathbf{k}'}}{2}\right)}{E_{m\mathbf{k}'}} \tag{VII.124}$$

VII.3.4 Computational Implementation

Use scipys `fixed_point` solver to solve the gap equation self-consistently.

Flatten $\Delta_n(\mathbf{k})$ the following way, to put it into the solver ($\mathbf{k}$ discretized in some way):

$$x = \begin{pmatrix} \Re(\Delta_1(\mathbf{k}_1)) \\ \Re(\Delta_1(\mathbf{k}_2)) \\ \vdots \\ \Re(\Delta_2(\mathbf{k}_1)) \\ \vdots \\ \Re(\Delta_3(\mathbf{k}_1)) \\ \vdots \\ \Im(\Delta_1(\mathbf{k}_1)) \\ \vdots \\ \Im(\Delta_2(\mathbf{k}_1)) \\ \vdots \\ \Im(\Delta_3(\mathbf{k}_1)) \\ \vdots \end{pmatrix} \tag{VII.125}$$

so that accessing a certain element takes the form:

$$\Re\Delta_n(\mathbf{k}) = x\left[\text{index}(\mathbf{k}) + \frac{\text{len}(x)\cdot n}{6}\right] \tag{VII.126}$$

$$\Im\Delta_n(\mathbf{k}) = x\left[\text{index}(\mathbf{k}) + \frac{\text{len}(x)\cdot n}{6} + \frac{1}{2}\text{len}(x)\right] \tag{VII.127}$$

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List of Symbols

Symbol	Meaning	Definition
τ	general imaginary time variable	1
$\mathcal{C}_{AB}(\tau, 0)$	Correlation function in imaginary time	1