I EG-X Hamiltonian in Reciprocal Space

In the following chapter, the model Hamiltonian

$$H_0 = -t_{\mathcal{X}} \sum_{\langle ij \rangle, \sigma} d_{i,\sigma}^{\dagger} d_{j,\sigma} - t_{\mathcal{G}r} \sum_{\langle ij \rangle, \sigma} c_{i,\sigma}^{(A),\dagger} c_{j,\sigma}^{(B)} + V \sum_{i,\sigma} d_{i,\sigma}^{\dagger} c_{i,\sigma}^{(A)} + \text{h.c.}$$
 (I.1)

will be treated to obtain the band structure. The first step is to write out the sums over nearest neighbors $\langle i,j \rangle$ explicitly, writing $\boldsymbol{\delta}_{\mathrm{X}}, \boldsymbol{\delta}_{\epsilon}$ ($\epsilon=A,B$) for the vectors to the nearest neighbors of the X atoms and Graphene A,B sites. Doing the calculation for example of the X atoms:

$$-t_{\mathbf{X}} \sum_{\langle ij \rangle, \sigma} (d_{i,\sigma}^{\dagger} d_{j,\sigma} + d_{j,\sigma}^{\dagger} d_{i,\sigma}) = -\frac{t_{\mathbf{X}}}{2} \sum_{i,\sigma} \sum_{\boldsymbol{\delta}_{\mathbf{X}}} d_{i,\sigma}^{\dagger} d_{i+\boldsymbol{\delta}_{\mathbf{X}},\sigma} - \frac{t_{\mathbf{X}}}{2} \sum_{j,\sigma} \sum_{\boldsymbol{\delta}_{\mathbf{X}}} d_{j,\sigma}^{\dagger} d_{j+\boldsymbol{\delta}_{\mathbf{X}},\sigma}$$
(I.2)

$$= -t_X \sum_{i,\sigma} \sum_{\delta_{\mathcal{X}}} d_{i,\sigma}^{\dagger} d_{i+\delta_{\mathcal{X}},\sigma}$$
 (I.3)

The factor 1/2 in eq. (I.2) is to account for double counting when going to the sum over all lattice sites i. By relabeling $j \to i$ in the second sum, the two sum are the same and eq. (I.3) is obtained. Using now the discrete Fourier transform

$$c_i = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{i\mathbf{k}\mathbf{r}_i} c_{\mathbf{k}}, \ c_i^{\dagger} = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{-i\mathbf{k}\mathbf{r}_i} c_{\mathbf{k}}^{\dagger}$$
 (I.4)

with the completeness relation

$$\sum_{i} e^{i\mathbf{k}\mathbf{r}_{i}} e^{-i\mathbf{k}'\mathbf{r}_{i}} = N\delta_{\mathbf{k},\mathbf{k}'} , \qquad (I.5)$$

eq. (I.3) reads:

$$-t_{X}\frac{1}{N}\sum_{i,\sigma}\sum_{\boldsymbol{\delta}_{X}}d_{i,\sigma}^{\dagger}d_{i+\boldsymbol{\delta}_{X},\sigma} = -t_{X}\frac{1}{N}\sum_{i,\sigma}\sum_{\mathbf{k},\mathbf{k}',\boldsymbol{\delta}_{X}}\left(e^{-i\mathbf{k}\mathbf{r}_{i}}d_{\mathbf{k},\sigma}^{\dagger}\right)\left(e^{i\mathbf{k}'\mathbf{r}_{i}}e^{i\mathbf{k}'\boldsymbol{\delta}_{X}}d_{\mathbf{k}',\sigma}\right)$$

$$= -t_{X}\frac{1}{N}\sum_{\mathbf{k},\mathbf{k}',\boldsymbol{\delta}_{X},\sigma}d_{\mathbf{k},\sigma}^{\dagger}d_{\mathbf{k}',\sigma}e^{i\mathbf{k}'\boldsymbol{\delta}_{X}}\sum_{i}e^{-i\mathbf{k}\mathbf{r}_{i}}e^{i\mathbf{k}'\mathbf{r}_{i}}$$

$$= -t_{X}\frac{1}{N}\sum_{\mathbf{k},\mathbf{k}',\sigma}d_{\mathbf{k},\sigma}^{\dagger}d_{\mathbf{k}',\sigma}\sum_{\boldsymbol{\delta}_{X}}e^{i\mathbf{k}'\boldsymbol{\delta}_{X}}\left(N\boldsymbol{\delta}_{\mathbf{k},\mathbf{k}'}\right) \qquad (I.8)$$

$$= -t_{X}\sum_{\mathbf{k},\sigma}d_{\mathbf{k},\sigma}^{\dagger}d_{\mathbf{k},\sigma}\sum_{\boldsymbol{\delta}_{X}}e^{i\mathbf{k}\boldsymbol{\delta}_{X}}. \qquad (I.9)$$

This part is now diagonal in **k** space. The nearest neighbours vectors $\boldsymbol{\delta}_{\mathrm{X}}$ for the X atoms are the vectors $\boldsymbol{\delta}_{AA,i}$ from ??. With that, the sum over $\boldsymbol{\delta}_{\mathrm{X}}$ can be explicitly calculated:

Correct exp expressions

Example for a vector product

$$f_{\mathcal{X}}(\mathbf{k}) = -t_X \sum_{\delta_{\mathcal{X}}} e^{i\mathbf{k}\delta_{\mathcal{X}}}$$
(I.10)

$$= -t_X \left[\exp \left(ia \left(\frac{k_x}{2} + \frac{\sqrt{3}k_y}{2} \right) \right) + e^{iak_x} + e^{ia(\frac{k_x}{2} - \frac{\sqrt{3}k_y}{2})} \right)$$
(I.11)

$$+ e^{\mathrm{i}a(-\frac{k_x}{2} - \frac{\sqrt{3}k_y}{2})} + e^{-\mathrm{i}ak_x} + e^{\mathrm{i}a(-\frac{k_x}{2} + \frac{\sqrt{3}k_y}{2})}$$
 (I.12)

$$= -t_X \left(2\cos(ak_x) + 2e^{ia\frac{\sqrt{3}k_y}{2}}\cos(\frac{a}{2}k_x) + 2e^{-ia\frac{\sqrt{3}k_y}{2}}\cos(\frac{a}{2}k_x) \right)$$
(I.13)

$$= -2t_X \left(\cos(ak_x) + 2\cos(\frac{a}{2}k_x)\cos(\sqrt{3}\frac{a}{2}k_y)\right). \tag{I.14}$$

The same can be done for the hopping between Graphene sites, for example :

$$-t_{Gr} \sum_{\langle ij\rangle,\sigma\sigma'} c_{i,\sigma}^{(A),\dagger} c_{j,\sigma'}^{(B)} = -t_{Gr} \sum_{i,\sigma\sigma'} \sum_{\delta_{AB}} c_{i,\sigma}^{(A),\dagger} c_{i+\delta_{AB},\sigma'}^{(B)}$$
(I.15)

$$= -t_{Gr} \sum_{\mathbf{k}, \sigma, \sigma'} c_{\mathbf{k}, \sigma}^{(A)\dagger} c_{\mathbf{k}, \sigma'}^{(B)} \sum_{\delta_{AB}} e^{i\mathbf{k}\delta_{AB}}$$
 (I.16)

We note

$$\sum_{\delta_{AB}} e^{i\mathbf{k}\delta_{AB}} = \left(\sum_{\delta_{BA}} e^{i\mathbf{k}\delta_{BA}}\right)^* = \sum_{\delta_{BA}} e^{-i\mathbf{k}\delta_{BA}}$$
(I.17)

and calculate

$$f_{Gr} = -t_{Gr} \sum_{\delta_{AB}} e^{i\mathbf{k}\delta_{AB}} \tag{I.18}$$

$$= -t_{Gr} \left(e^{i\frac{a}{\sqrt{3}}k_y} + e^{i\frac{a}{2\sqrt{3}}(\sqrt{3}k_x - k_y)} + e^{i\frac{a}{2\sqrt{3}}(-\sqrt{3}k_x - k_y)} \right)$$
(I.19)

$$= -t_{Gr} \left(e^{i\frac{a}{\sqrt{3}}k_y} + e^{-i\frac{a}{2\sqrt{3}}k_y} \left(e^{i\frac{a}{2}k_x} + e^{-i\frac{a}{2}k_x} \right) \right)$$
 (I.20)

$$= -t_{Gr} \left(e^{i\frac{a}{\sqrt{3}}k_y} + 2e^{-i\frac{a}{2\sqrt{3}}k_y} \cos\left(\frac{a}{2}k_x\right) \right)$$
 (I.21)

All together, we get:

$$H_0 = \sum_{\mathbf{k}, \sigma, \sigma'} \begin{pmatrix} c_{k, \sigma}^{A, \dagger} & c_{k, \sigma}^{B, \dagger} & d_{k, \sigma}^{\dagger} \end{pmatrix} \begin{pmatrix} 0 & f_{Gr} & V \\ f_{Gr}^* & 0 & 0 \\ V & 0 & f_X \end{pmatrix} \begin{pmatrix} c_{k, \sigma}^A \\ c_{k, \sigma}^B \\ d_{k, \sigma} \end{pmatrix}$$
(I.22)