Order Statistics of the SIR and Interference Cancellation in a Narrow-Beam LEO Uplink

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Abstract—We investigate the factorial moment measure of the signal-to-interference ratios (SIR) at the typical low Earth orbit base station (LEO BS) with a narrow Gaussian antenna serving an urban area, with a Gaussian mixture shadowing model. This SIR process is characterized by a Poisson-Dirichlet distribution $PD(0,\cdot)$, which allows us to derive the density of the factorial moment measure. We analyze the coverage probability at the typical LEO BS receiving the three strongest signals and implement successive interference cancellation (SIC). Our results demonstrate that SIC can notably reduce the variance of the SIR while maintaining robust performance.

I. INTRODUCTION

HILE THE ORDER STATISTICS of the signal-tointerference (SIR) and interference cancellation have been studied for terrestrial networks [1] by using stochastic geometry, they are yet to be explored in low Earth orbit (LEO) networks. We study the SIR of user equipments (UEs) at the typical LEO base station (BS) by utilizing the narrow-beam LEO uplink system model from [2] in an urban environment. We utilize the Gaussian Mixture shadowing model, similar to [3] and [4], using the parameters presented in [5]. The density of the factorial moment measure of the signal-tototal-interference (STIR) process follows a Poisson-Dirichlet distribution $PD(0,\cdot)$ (contrary to $PD(\cdot,0)$ in [1]), of which the factorial moment measure is well-known. We derive the joint pdf of the STIR and SIR and study the performance metrics of the three strongest UE signals with and without successive interference cancellation (SIC) schemes. In [2], it was observed that the system parameters optimizing the average throughput, corresponding to mean log(2) user equipments (UEs) inside a LEO BS -3 dB footprint, leads to a high variation in the SIR over the LEO BSs. We demonstrate that the link is more stable with interference cancellation. We show that with the SIC, the number of UEs inside a LEO BS -3 dB footprint can be doubled while maintaining the average performance of the strongest UE but profoundly reducing the variance in the SIR. Furthermore, the coverage probability of UEs with less strong signals drastically improves.

II. SYSTEM MODEL

A narrow-band LEO uplink is considered. The UEs follow a homogeneous PPP $\Phi \subset \mathbb{R}^2$ of density λ . The LEO BSs form

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TABLE I: Principal symbols the values and units in the numerical results. We denote (approximate) proportionality " \propto " or equality "=" to a variable. For a dimensional number, we denote the units "SI" and "[non-SI]."

Symbol	Explanation	Values and units
h	Receiving LEO BS altitudes.	1000 km
α	Power path loss exponent.	4
λ	The density of Φ and Θ , <i>i.e.</i> , the mean	$\{0.83, 13.3\}$
	number of UEs inside an unit area.	10^{-4} /km ²
ϵ	The typical LEO BS elevation angle	$(\pi/2, \pi/6)$ rad
	w.r.t o.	$= (90, 30)[^{\circ}]$
p_{LoS}	LoS probability; $p_{\text{LoS}} \propto \sin(\epsilon)$.	(0.992, 0.493)
$\mu_{ ext{LoS}}$	Mean of the LoS component of	0 [dB]
	the Gaussian mixture shadow fading	
	model.	
$\mu_{ m NLoS}$	Mean of the NLoS component.	-26 [dB]
σ_{LoS}^2	Variance of the LoS component.	$4^{2} [dB]$
$\sigma_{\rm NLoS}^2$	Variance of the NLoS component.	6^2 [dB]
φ_{RX}	Half-width of the -3 dB antenna gain.	0.028 = 1.6[°]
κ	Average number of UEs inside a -3 dB	(1.63, 6.51)
	footprint; $\kappa = \pi \lambda \left(\varphi_{\rm RX} h / \sin^2(\epsilon) \right)^2$.	
v	Fraction of effective UEs; $v \propto \sin(\epsilon)$.	(0.85, 0.426)
κv	Average number of effective UEs inside	$\{2,4\} \times \log(2)$
	a - 3 dB footprint.	
θ	SIR threshold of a uccessful transmis-	(0.2, 10) =
	sion.	(-7, 10) [dB]
au	SIR threshold of a successful interfer-	0.2 =
	ence cancellation.	-7 [dB]

a homogeneous point process (p.p.) possibly with a different density than the UEs. Because of the translation invariancy of the PPP, all locations are statistically equivalent, and we define the origin *o* to represent *the typical LEO BS* footprint focus point.

The path loss, representing the antenna gain at the typical LEO BS, is given over the planar distance $r \in [0, \infty)$ as a Gaussian function

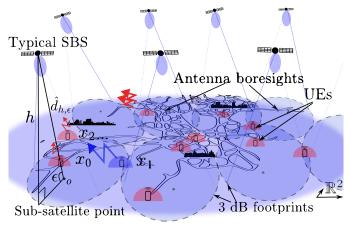
$$G(r) = 2^{-(D_{h,\epsilon}r)^2/\varphi_{RX}^2}.$$
 (1)

The angle $\varphi_{\rm RX}$ denotes the -3 dB antenna gain width. Furthermore, the scaling constant $D_{h,\epsilon} \triangleq \sin^2(\epsilon)/h$ is a first-order coefficient of the Taylor expansion of the angle φ_r w.r.t. the boresight of the typical antenna pattern. (Thorough details in [2].)

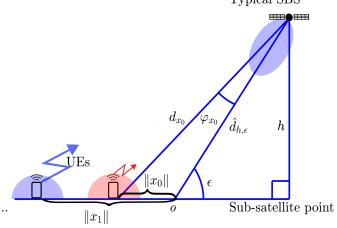
A. Shadowing

1) Gaussian mixture shadowing model

Consider a two-tier {LoS, NLoS} (line-of-sight and non-line-of-sight) Gaussian mixture shadow fading model with the parameters $\mu_{\text{LoS}} = 0$ dB, $\sigma_{\text{LoS}} = 4$ dB, $\mu_{\text{NLoS}} = -26$ dB, and $\sigma_{\text{NLoS}} = 6$ dB, which correspond to an urban environment [5]. Assuming i.i.d. power shadow fading for all UEs, the



(a) Interpretation of the *planar* system model with the satellites in adjacent orbits serving an urban area and a realization of the UEs. The altitudes are not to scale. $Typical\ SBS$



(b) The typical LEO BS as seen from the side. The transmitters are projected into line $(0,\infty)$ according to their norm.

Fig. 1: The simplified narrow-beam LEO uplink system model. The satellite antenna boresight is oriented towards o, the focus point of the elliptical footprint. The omnidirectionally transmitting UEs $\{x_i\}$ are located according to the homogeneous PPP on the plane. The transmitter with the strongest signal is the first-served UE.

typical shadowed transmit power $H_{\mathcal{MLN}}$ follows a log-normal mixture distribution;

$$H_{\mathcal{MLN}} \sim p_{\text{LoS}} \mathcal{LN}(\rho \mu_{\text{LoS}}, (\rho \sigma_{\text{LoS}})^2) + p_{\text{NLoS}} \mathcal{LN}(\rho \mu_{\text{NLoS}}, (\rho \sigma_{\text{NLoS}})^2), \tag{2}$$

where $p_{\text{LoS}} = 1 - p_{\text{NLoS}}$ is the LoS probability as in Figure 2b. Considering a natural base for the log-normal distribution, the constant $\rho \triangleq \log(10)/10$ normalizes the parameters $\mu_{\text{LoS}}, \sigma_{\text{LoS}}, \mu_{\text{NLoS}},$ and $\sigma_{\text{NLoS}},$ ensuring that the conditioned r.v.'s $10\log_{10}(H_{\mathcal{MLN}}|\text{LoS})$ and $10\log_{10}(H_{\mathcal{MLN}}|\text{NLoS})$ evaluate to r.v.'s following the normal distributions $\mathcal{N}(\mu_{\text{LoS}}, \sigma_{\text{LoS}}^2)$ and $\mathcal{N}(\mu_{\text{NLoS}}, \sigma_{\text{NLoS}}^2)$, respectively.

2) Defective exponential shadowing distribution

As a compromise between analytical tractability and realism, we introduce a *defective* exponential power fading distribution for the UEs, described by the distribution function

$$F_{H_{\text{Exp}}}(t) = ve^{-t}, t > 0.$$
 (3)

Essentially, this is a mixture distribution. Namely, $0 \le 1-v < 1$ denotes the probability that the shadowed signal is entirely

attenuated and takes the value of zero, otherwise, the power follows the exponential distribution.

We introduce a scaling term, Υ , to ensure that the means of the log-normal mixture distribution and the defective exponential distribution match: $\mathbb{E}(\Upsilon H_{\mathcal{MLN}}) = \mathbb{E}(H_{\text{Exp}}) = v$. By equating the first two moments of H_{Exp} (l.h.s.) and $\Upsilon H_{\mathcal{MLN}}$ (r.h.s.)

$$\begin{cases}
v = \Upsilon \left(p_{\text{LoS}} e^{\mu_{\text{LoS}} + \sigma_{\text{LoS}}^2/2} + p_{\text{NLoS}} e^{\mu_{\text{NLoS}} + \sigma_{\text{NLoS}}^2/2} \right) \\
2v = \Upsilon^2 \left(p_{\text{LoS}} e^{2(\mu_{\text{LoS}} + \sigma_{\text{LoS}}^2)} + p_{\text{NLoS}} e^{2(\mu_{\text{NLoS}} + \sigma_{\text{NLoS}}^2)} \right),
\end{cases} (4)$$

we can solve for the parameter v:

$$v = \frac{2\left(p_{\text{LoS}}e^{\mu_{\text{LoS}} + \sigma_{\text{LoS}}^2/2} + p_{\text{NLoS}}e^{\mu_{\text{NLoS}} + \sigma_{\text{NLoS}}^2/2}\right)^2}{p_{\text{LoS}}e^{2(\mu_{\text{LoS}} + \sigma_{\text{LoS}}^2)} + p_{\text{NLoS}}e^{2(\mu_{\text{NLoS}} + \sigma_{\text{NLoS}}^2)}.$$
 (5)

The parameter $v=v(\epsilon)$ varies with the elevation angle, influencing the shadow fading characteristics. The parameter Υ holds no significance in an interference-limited scenario, as the equal scaling of all UE powers neutralizes its effect.

Remark. The variable $0 < v \le 1$ is not generally solvable for all log-normal distribution parameters (by matching the first two moments). Broadly said, the variance of the shadowing has to be large enough. However, the variable $0 < v \le 1$ is solvable for almost every shadowing scenario in [5], particularly for the urban scenario.

B. The spherical Earth model in the Monte Carlo simulations

The planar model approximates the spherical system model, where the UEs are located on the spherical Earth surface of radius $R_{\oplus}=6378$ km according to a homogeneous PPP Θ of density λ represented in spherical coordinates. This p.p. can be constructed from $\Phi\subset\mathbb{R}^2$ by a preserving mapping. Namely, for $x=(x_1,x_2)\in[-\pi,\pi]\times[-1,1]\cap\Phi/R_{\oplus}$,

$$(x_1, x_2) \mapsto (R_{\oplus}, x_1, \sin^{-1}(x_2)) = (R_{\oplus}, \theta_{x_1}, \theta_{x_2}) \in \Theta.$$
 (6)

The total interference at the typical LEO BS from the UEs in the PPP $\Theta \cap E$, with E denoting the area above the horizon of the typical LEO BS, is defined as

$$\mathring{I} \triangleq \sum_{x \in \Theta \cap E} \frac{H_{\mathcal{MLN}_x} G(\varphi_x)}{(d_x/d_0)^{\alpha}},\tag{7}$$

where d_0 is a normalizing constant. The simulated values are based on the spherical model, the angle φ_x and distance d_x calculated precisely for each $x \in \Theta \cap E$, and accurate i.i.d. Gaussian mixture shadowing $H_{\mathcal{MLN}_x}$.

III. ANALYSIS

For a more detailed analysis and planar model comparison to the spherical model, please refer to [2].

Given i.i.d. shadowing variables $\{H_x\}_{x\in\Phi}$, we define the process of the received signal powers at the typical location, referred to as the gain process (GP), by

$$\mathcal{G} \triangleq \{H_x G(\|x\|) : x \in \Phi\}, \tag{8}$$

where ||x|| is the Euclidean distance from o.

The GP is a *projection process* mapping the points from \mathbb{R}^2 into $(0, \infty)$ and, as such, forms a nonhomogeneous PPP [6, Section 4.2.5].

Since the variables $\{H_x\}_{x\in\Phi}$ are i.i.d., we can denote the typical shadowing variable simply as H without the subscript.

Proposition 1 (Density of the GP). Let $F_H(\cdot)$ be the (possibly degenerate) complementary cumulative distribution function (ccdf) of a fading variable H. The density function of $\mathcal G$ is given by

$$\lambda_{\mathcal{G}}(t) = \tilde{\kappa} F_H(t)/t, \quad t \in (0, \infty),$$
(9)

where $\tilde{\kappa} = \kappa/\log(2)$ and

$$\kappa \triangleq \pi \lambda \left(\frac{\varphi_{RX} h}{\sin^2(\epsilon)} \right)^2 \tag{10}$$

is approximately the average number of UEs inside $a-3 \ dB$ footprint.

Proof. Let $f_H(\cdot)$ be the pdf of H. Denote $G^{-1}(\cdot)$ as the generalized inverse of G, defined as $G^{-1}(y) = \inf\{x : G(x) < y\}$. According to [6, Eq. 4.55],

$$\begin{split} &\int_{t}^{\infty} \lambda_{\mathcal{G}}(y) \, dy = \pi \lambda \mathbb{E} \left[\left(G^{-1}(t/H) \right)^{2} \right] \\ &= \pi \lambda \int_{t}^{\infty} \left(-\frac{\varphi_{\text{RX}} \sqrt{-\log(t/h)}}{D_{h,\epsilon} \sqrt{\log(2)}} \right)^{2} f_{H}(h) \, dh \\ &= -\tilde{\kappa} \int_{t}^{\infty} \log(t/h) f_{H}(h) \, dh \\ &\stackrel{(a)}{=} -\tilde{\kappa} \left[\log(t/h) F_{H}(h) \big|_{t}^{\infty} + \int_{t}^{\infty} \frac{F_{H}(h)}{h} \, dh \right]. \end{split}$$

In (a), we use integration by parts. The result follows by differentiating with respect to t and applying the negative sign. Note that a necessary condition for this procedure is that $\int_t^\infty \log(t/h) f_H(h) \, dh$ converges for all t > 0.

Please refer to [2, Lemma 1] for throughout explanation of the interpretation of κ .

The total interference, or total received power, is defined as the sum of the GP at the footprint location o of the typical LEO BS:

$$I \triangleq \sum_{x \in \Phi} H_x G(\|x\|) = \sum_{x \in \mathcal{G}} x. \tag{11}$$

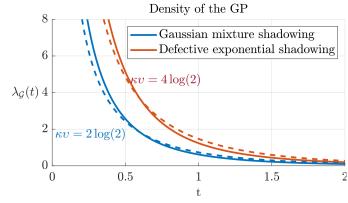
The mean and the variance of I are respectively given by

$$\mathbb{E}(I) = \int_0^\infty t \lambda_{\mathcal{G}}(t) dt = \tilde{\kappa} \int_0^\infty F_H(t) dt = \tilde{\kappa} \mathbb{E}(H), \quad (12)$$

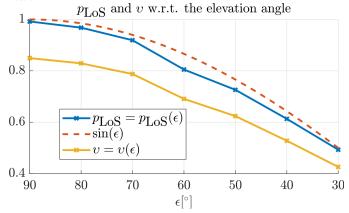
and

$$\operatorname{Var}(I) = \int_{0}^{\infty} t^{2} \lambda_{\mathcal{G}}(t) dt = \tilde{\kappa} \int_{0}^{\infty} t F_{H}(t) dt$$
$$= \tilde{\kappa} \frac{\operatorname{Var}(H) + \mathbb{E}(H)^{2}}{2} = \tilde{\kappa} \mathbb{E}[H^{2}]/2. \tag{13}$$

Note that matching the first two moments of the fading distributions (2) and (3) is equivalent to matching the mean and the variance of the total interference.



(a) The density of the GP for $\kappa \upsilon \in \{2\log(2), 4\log(2)\}$ using $\epsilon = \pi/2$. Interestingly, the elevation angle did not have visible effect on the density in the Gaussian mixture model.



(b) The dependence of the shadowing parameters $p_{\rm LoS}$ and υ on the elevation angle. The parameters are approximately proportional to the sine function.

Fig. 2: The density of the GP and shadowing parameters in the Gaussian mixture and defective exponential shadowing models.

A. Laplace transform of the total received power

With defective exponential shadowing H_{exp} , for Re(s) > 1,

$$\mathcal{L}_{I}(s) \triangleq \mathbb{E}\left(e^{-sI}\right) = \exp\left\{-\int_{0}^{\infty} (1 - e^{-sr})\lambda_{\mathcal{G}}(r)dr\right\}$$

$$= \exp\left\{-\tilde{\kappa}\int_{0}^{\infty} (1 - e^{-sr})F_{H_{\exp}}(r)/rdr\right\}$$

$$= \exp\left\{-\tilde{\kappa}v\int_{0}^{\infty} (1 - e^{-sr})e^{-r}/rdr\right\} = (1 + s)^{-\tilde{\kappa}v},$$
(14)

which is the Laplace transform of the gamma distribution with the shape parameter $\tilde{\kappa}v$.

B. Order statistics of the STIR and SIR processes

At the typical LEO BS, we denote the signal-to-interference ratio (SIR) process of the UEs as follows:

$$\Psi = \{ \mathsf{Z} : \mathsf{Z} \in \Psi \} \triangleq \left\{ \frac{u}{I - u} : u \in \mathcal{G} \right\}
= \left\{ \frac{H_x G(D_{h,\epsilon} || x ||)}{I - H_x G(D_{h,\epsilon} || x ||)} : x \in \Phi \right\},$$
(15)

where I is defined in (11). Similarly, the signal-to-total-interference ratio (STIR) process is defined as

$$\Psi' = \left\{ \mathsf{Z}' : \mathsf{Z}' \in \Psi' \right\} \triangleq \left\{ \frac{u}{I} : u \in \mathcal{G} \right\}. \tag{16}$$

We can always recover the process from another:

$$\Psi = \left\{ \frac{\mathsf{Z}'}{1 - \mathsf{Z}'} : \mathsf{Z}' \in \Psi' \right\}, \quad \Psi' = \left\{ \frac{\mathsf{Z}}{1 + \mathsf{Z}} : \mathsf{Z} \in \Psi \right\}. \tag{17}$$

Let θ denote the SIR threshold for successful transmission. The event $\Psi \ni \mathsf{Z} > \theta$ is equivalent to $\Psi' \ni \mathsf{Z}' > \theta'$ with $\theta' \triangleq \theta/(1+\theta)$ and $\theta \triangleq \theta'/(1-\theta')$.

We denote $Z'_{(1)} > Z'_{(2)} > Z'_{(3)} \dots$ as the order statistics of the STINR process Ψ' , such that $Z'_{(1)}$ is the largest value in Ψ' . Through the monotonicity of the relations (17), the order statistics of the STIR process are equivalent to the order statistics of the SIR process.

Proposition 2. The density of the nth factorial moment measure of the STIR process at the typical LEO BS with a narrow Gaussian antenna beam and Gaussian mixture shadow fading is approximately given by

$$\mu'^{(n)}(t'_1, \dots, t'_n) = (\tilde{\kappa}v)^n \prod_{j=1}^n t'_j^{-1} \left(1 - \sum_{j=1}^n t'_j\right)^{\kappa v - 1}, (18)$$

whenever $t'_1 > \cdots > t'_n$ and $\sum_{i=1}^n t'_i \leq 1$, and 0 otherwise.

Proof. The total interference can be characterized by the gamma process at time $\tilde{\kappa}v$ [7, Eq. 8] (recall (14)). Hence, the STIR process Ψ' can be characterized by a Poisson-Dirichlet distribution PD $(0, \tilde{\kappa}v)$ that has the given density [8, Eq. 2.3].

The partial densities can be derived from the density of the nth factorial moment measure as [1, Eq. 62]

$$\mu'_{n}^{(n+i)}(z'_{1}, \dots, z'_{n}) = \int_{z'_{n}}^{1} \dots \int_{z'_{n}}^{1} \mu'^{(n+i)}(z'_{1}, \dots, z'_{n}, \zeta'_{1}, \dots, \zeta'_{i}) d\zeta'_{1} \dots d\zeta'_{i},$$
 (19)

the support of the density being in the region $\sum_{j=1}^{n} z'_j + iz'_n \le 1$

The joint pdf of the n strongest values of the STIR process $(\mathsf{Z}'_{(1)},\ldots,\mathsf{Z}'_{(n)})$ is given as a series expansion involving the partial densities [1, Eq. 64]

$$f'_{(n)}(z'_1, \dots, z'_n) = \sum_{i=0}^{i_{\text{max}}} \frac{(-1)^i}{i!} {\mu'}_n^{(n+i)}(z'_1, \dots, z'_n), \quad (20)$$

for $z_1'>z_2'>\cdots>z_n'$ and $f_{(n)}'(z_1',\ldots,z_n')=0$ otherwise. The upper bound for the index $i_{\max}<1/z_n'-n$ corresponds to the non-zero terms of the series expansion.

The n-coverage probability that the first n strongest signals reach the threshold θ is given by

$$\mathcal{P}^{(n)}(\theta) \triangleq \int_{\theta'}^{1} \cdots \int_{\theta'}^{1} f'_{(n)}(z'_1, \dots, z'_n) dz'_1 \dots dz'_n, \quad (21)$$

with $\theta' = \theta/(1+\theta)$ and $i_{max} < 1/\theta' - n$.

The density of the *n*th factorial moment measure of the SIR process can be extracted from $\mu'^{(n)}$ [6, Corollary 6.1.3]:

$$\mu^{(n)}(z_1, \dots, z_n) = \prod_{j=1}^n \frac{1}{(1+z_j)^2} \mu'^{(n)}\left(\frac{z_1}{1+z_1}, \dots, \frac{z_n}{1+z_n}\right)$$
(22)

C. SIR under interference cancellation

Let $(u_{(1)},\ldots,u_{(k)})\subset\mathcal{G}$ represent an ordered set of points in the GP, where $u_{(1)}$ denotes the strongest signal at the typical LEO BS. The signals with indices in the set $[k] \triangleq (1,\ldots,k), k\geq n$, are canceled from the total interference. We denote the SIR with interference cancellation as

$$\operatorname{SIR}_{n,[k]} \triangleq \frac{u_{(n)}}{I - \sum_{j \in [k]} u_{(j)}}.$$
 (23)

Let us first study SIR_{1,[1]}. Combining (18), (20), and (22), we can derive a closed-form for the SIR pdf of the strongest signal in the *simple coverage region* $z \geq 1$: $f_{(1)}(z) = \tilde{\kappa} v (z+1)^{-\tilde{\kappa} v}/z^{-1}$. The second moment of the SIR is bounded by

$$\mathbb{E}(SIR_{1,[1]}^2) \ge \int_1^\infty f_{(1)}(z)z^2 dz = \frac{2^{1-\tilde{\kappa}v}(\tilde{\kappa}v)^2}{(\tilde{\kappa}v - 1)(\tilde{\kappa}v - 2)}, \quad (24)$$

which is divergent for $\tilde{\kappa} \upsilon \leq 2$, *i.e.*, for less than $2\log(2)$ effective UEs inside a -3 dB footprint on average, the first and second moments—hence, also the variance—are infinite (or undefined). Despite the strong average SIR, the infinite variance for $\tilde{\kappa} \upsilon \leq 2$ is not desirable if we want a consistent user experience in the link quality. We demonstrate that successive interference cancellation (SIC) can improve user fairness.

Under interference cancellation, we have the following identity in terms of the STIR process [1, Eq. 69]:

$$\mathbb{P}(\operatorname{SIR}_{n,[k]} > \theta) = \mathbb{P}\left(\mathsf{Z}'_{(n)} + \theta' \sum_{j \in [k] \setminus \{n\}} \mathsf{Z}'_{(j)} > \theta'\right). \tag{25}$$

Following the Poisson-Dirichlet order statistics of $Z'_{(1)} > Z'_{(2)} > \ldots$, if $Z'_{(1)}$ has a finite variance, each $\{Z'_{(j)}\}_{j \in [k]}$ also has a finite variance, hence $SIR_{n,[k]}$ has a finite variance.

Finally, we consider the SIR under the (perfect) successive signal cancellation (SIC-SIR). A necessary condition for the successful reception of the nth strongest UE at the typical LEO BS is that the preceding n signals are successively decoded and removed from the interference. Formally, $\{Z'_m + \tau' \sum_{j=1}^{m-1} Z'_j > \tau'\}$ for all $m \in \{1 \dots n\}$, where the signal detection threshold is denoted as $\tau = \tau'/(1-\tau') \leq \theta$. When the first n signals are successfully removed from the interference, the nth UE is considered covered if $\mathrm{SIR}_{n,[n]} > \theta$. If not, the SIC continues until the $\mathrm{SIR}_{n,[k]} > \theta$ or the maximum number of interference cancellation stages K is reached.

¹The SIR has a heavy-tailed distribution, cf. [2, Eq. (32)].

Proposition 3. Consider the SIC with at most $K \ge n$ interference cancellation stages. The coverage probability of the UE with nth strongest signal is given by

$$\mathcal{P}_{\text{SIC}}^{(n,K)}(\theta,\tau) \triangleq \sum_{k=-n}^{K} \Delta_{\text{SIC}}^{(n,k)}(\theta,\tau), \tag{26}$$

where

$$\Delta_{\text{SIC}}^{(n,k)}(\theta,\tau) \triangleq \\
\times \int_{0}^{1} \cdots \int_{0}^{1} f_{(k)}(z'_{1},\ldots,z'_{k}) \prod_{m=1}^{k} \mathbb{1} \left(z'_{m} + \tau' \sum_{j=1}^{m-1} z'_{j} > \tau' \right) \\
\times \left(\mathbb{1}(k > n) \mathbb{1} \left(z'_{n} + \theta' \sum_{j \in [k-1] \setminus \{n\}} z'_{j} < \theta' \right) + \mathbb{1}(k = n) \right) \\
\times \mathbb{1} \left(z'_{n} + \theta' \sum_{j \in [k] \setminus \{n\}} z'_{j} > \theta' \right) dz'_{1} \ldots dz'_{k} \tag{27}$$

with the upper summation limit bounded by $i_{max} < 1/\tau' - 1 = 1/\tau$.

Proof. The expression follows using the joint pdf of the order statistics (20). Furthermore, the first condition (with the form $\prod \mathbbm{1}(\cdot)$) allows the relaxation of i_{\max} . Namely, a necessary condition is $z_k'+\tau'\sum_{j=1}^{k-1}z_j'>\tau'$. By simple algebra, $\sum_{j=1}^{k-1}z_j>1-z_k/\tau'$. Recall the condition on the non-zero terms of $\mu'^{(k+i)}\colon \sum_{j=1}^k z_j'+iz_k'=\sum_{j=1}^{k-1}z_j'+z_k'+iz_k'\leq 1$. The condition certainly does *not* hold if $1-z_k/\tau'+z_k'+iz_k'>1$. We arrive at the inequality $z_k'(-1/\tau'+1+i)>0$. Divide both sides by $z_k'>0$, and the general upper bound of i follows. \square

IV. NUMERICAL RESULTS AND CONCLUCIONS

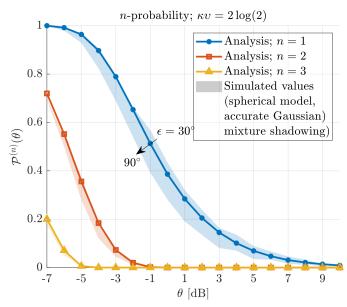


Fig. 3: The n-probabilities for $\kappa v = 2\log(2)$ (the average number of effective UEs inside a -3 dB footprint).

Figures 3 and 4 depict the *n*-probabilities (21) and SIC-SIR for $\tilde{\kappa}v = 2$ and $\tilde{\kappa}v = 4$ (27) with K = 3, respectively.

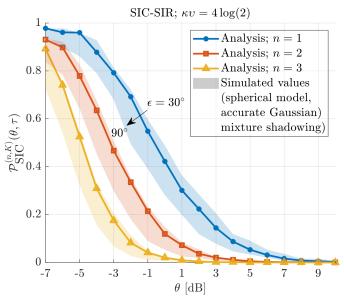


Fig. 4: The SIC-SIR for the transmitters $n \in \{1, 2, 3\}$ by doubling the average number of effective UEs inside the -3 dB footprints compared to Figure 3; $\kappa v = 4 \log(2)$. The SIC parameter K = 3.

The signal detection threshold $\tau=-7$ dB. We use the values presented in Table I in the simulations. However, the crucial system parameter is $\tilde{\kappa}v$. Hence, for example, instead of scaling λ , we could adjust the width of the antenna gain for each elevation angle according to (5) and (10) to match the corresponding $\tilde{\kappa}v\equiv\kappa v/\log(2)$.

The figures illustrate that SIC-SIR can achieve comparable coverage probabilities for the strongest UE within the region $\theta \in (-7,10)$ dB while doubling the average number of effective UEs inside a -3 dB footprint, denoted as κv . Additionally, the performance of the 2nd and 3rd UEs is significantly enhanced. Consequently, a single LEO BS could potentially serve multiple UEs effectively.

Further, similar to (24), we can calculate an upper bound for the variance of the SIR of the strongest signal before interference cancellation for $\tilde{\kappa} \upsilon = 4$: $\text{var}(\text{SIR}_{1,[1]}) = \mathbb{E}(\text{SIR}_{1,[1]}^2) - \mathbb{E}(\text{SIR}_{1,[1]})^2 \leq 1.2$. This represents a significant improvement compared to the infinite variance for $\tilde{\kappa} \upsilon = 2$.

We conclude that interference cancellation, particularly successive interference cancellation (SIC), is a viable solution for mitigating the considerable variability in link quality experienced by users in a narrow-beam low Earth orbit (LEO) uplink. These findings are also relevant to the downlink, given that the LEO footprint locations follow a Poisson distribution on the Earth's surface.

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