

Wide Post-Mass Transfer White-Dwarf Binaries

Runqiu Ye

June, 2024

1 Introduction

Common envelope (CE) is the outcome of unstable mass transfer. During CE, both stars orbit inside an envelope and spiral inward. If the energy liberated in this process is enough to eject the CE, then the result of CE process is a close binary. If, on the contrary, the liberated energy is not enough to eject the envelope, then the result of the process is a merger. In observation results, a number of WD+MS binaries are wider than expected, and CE is expected to be the main channel to form these white dwarf binaries [Yamaguchi et al., 2024a, Yamaguchi et al., 2024b]. In the MESA model utilized by [Yamaguchi et al., 2024a, Yamaguchi et al., 2024b], wide post-CE WD+MS binaries can be formed in a certain range of initial separation with only gravitational and internal energy included in the calculation.

In this paper, we are going to explore the formation possibility of these wide post-mass transfer WD+MS binaries with COSMIC model. We perform COSMIC simulation on both individual binary star systems and sampled population to explore the formation of WD+MS binaries under different conditions. In section 2, we are going to investigate formation of individual wide WD+MS binaries through CE for a variety of variables, including initial mass, initial separation, CE efficiency, and the energy budget of the CE. In section 3, we are going to explore the formation of individual wide WD+MS binaries through stable mass transfer. In section 4, we will research on how different COSMIC model affect simulation results of a sample population of binary star systems, and how the formation of post WD+MS binaries depends on various parameters.

Throughout the paper, the stellar type and evolution type follows the BSE conventions. `kstar` values specify the evolutionary state of the star, and `evol_type` values specify the evolutionary changes of the binary systems:

kstar	evolutionary state	evol_type	evolutionary change
0	MS < 0.7 M _⊙	1	initial state
1	MS > 0.7 M _⊙	2	kstar change
2	Hertzsprung Gap	3	begin Roche lobe overflow
3	First Giant Branch	4	end Roche lobe overflow
4	Core Helium Burning	5	contact
5	Early AGB	6	coalescence
6	TP-AGB	7	begin common envelope
7	Naked Helium Star MS	8	end common envelope
8	Naked Helium Star Hertzsprung Gap	9	no remnant leftover
9	Naked Helium Star Giant Branch	10	max evolution time
10	Helium White Dwarf	11	binary disruption
11	Carbon/Oxygen White Dwarf	12	begin symbiotic phase
12	Oxygen/Neon White Dwarf	13	end symbiotic phase
13	Neutron Star	14	blue straggler
14	Black Hole	15	supernova of primary
15	Massless Remnant	16	supernova of secondary

2 Formation Through Common Envelope

In [Yamaguchi et al., 2024a, Yamaguchi et al., 2024b], the author discussed the formation of WD+MS binaries through CE for both high-mass and low-mass systems, and it was found that wide WD+MS systems can be formed for certain energy budgets and certain initial separation in the MESA model. We hope to explore similar models in COSMIC.

In [Yamaguchi et al., 2024a, Yamaguchi et al., 2024b], to explore the dependency of the final separation a_f in the post-CE binary on the initial separation a_i in the pre-MT binary, the author first ran a stellar evolution model of the primary star using MESA. Its evolution up to AGB phase is followed. With the model, the binding energy of the CE is calculated as

$$E_{\text{bind}} = E_{\text{grav}} + E_{\text{int}} = \int_{M_{\text{core}}}^{M_{\text{tot}}} -\frac{Gm}{r(m)} + U(m)dm, \quad (1)$$

or

$$E_{\text{bind}} = E_{\text{grav}} + \alpha_{\text{th}} E_{\text{th}} + \alpha_{\text{rec}} E_{\text{rec}}. \quad (2)$$

After that, the change of separation through CE evolution is calculated using

$$E_{\text{bind}} = \alpha_{\text{CE}} \left(-\frac{GM_{\text{WD}}M_{\star}}{2a_f} + \frac{GM_iM_{\star}}{2a_i} \right). \quad (3)$$

In COSMIC, the CE is modeled using a structural parameter λ , where

$$E_{\text{bind}} = \frac{GM_iM_{\text{env}}}{\lambda R_i}. \quad (4)$$

This λ is represented in COSMIC as the `lambdaf` flag in `BSEDict`. To compare MESA model with COSMIC, we intend to compare how the energy budget utilized in past papers matches with the CE model in COSMIC. We approach this by calculating the effective λ value for each of the common envelope energy budget used in [Yamaguchi et al., 2024a, Yamaguchi et al., 2024b].

2.1 High Mass Systems ($7 M_{\odot} + 1 M_{\odot}$)

2.1.1 Methods and results

According to [Yamaguchi et al., 2024a], the lower limit of WD mass $M_{\text{MD,min}}$ in the observational sample is calculated to be $1.244 M_{\odot} \sim 1.418 M_{\odot}$. The corresponding mass of MS progenitor is expected to be in the range of $6 \sim 9 M_{\odot}$, and the companion mass has median around $1 M_{\odot}$. MESA results for $7 M_{\odot} + 1 M_{\odot}$ systems is documented in [Yamaguchi et al., 2024a]. Hence, we also consider a $7 M_{\odot} + 1 M_{\odot}$ evolution model in COSMIC and compare the results with those in [Yamaguchi et al., 2024a].

We ran COSMIC binary evolution simulation of a $7 M_{\odot} + 1 M_{\odot}$ system and investigate how the final separation a_f depends on the initial separation a_i and the CE parameter λ . The initial separation is a linear grid in range $2 \sim 8 \text{ au}$ with 400 points. The CE flag `lambdaf` is a linear grid in range $0 \sim -100$ with 2000 steps, which corresponds to $0 \sim 100$ for CE parameter λ . Also, we ran these parameters with four different CE efficiencies — $\alpha_{\text{CE}} = 1$, $\alpha_{\text{CE}} = 0.9$, $\alpha_{\text{CE}} = 0.6$, and $\alpha_{\text{CE}} = 0.3$, the same as recorded in [Yamaguchi et al., 2024a]. In total, for each of the 4 different CE efficiency α_{CE} , we ran a grid of 800000 binary star systems with different initial separation a_i and CE parameter λ .

After simulation is finished, we select the time-step when the system finished CE for the first time (`evol_type = 8`). At this moment, if the system is a WD+MS system, we record their separation as the final separation a_f . Otherwise, we set $a_f = 0$ and consider this initial condition unable to produce a WD+MS system. We create four heat maps for the four different CE efficiencies. In Figure 1, we plot the final separation a_f of the WD+MS systems against initial separation a_i and the CE flag $-\lambda$.

2.1.2 Discussions

Among the 3200000 binary star systems we simulated, a typical change of `evol_type` through the evolution is **3-7-8-4-3-7-8-4** (refer to Section 1 for correspondence between `evol_type` and evolutionary changes, `kstar` changes and other evolutionary changes before and after the mass transfer are omitted), where WD+MS is reached during the second mass transfer process. Another possible change of `evol_type` through the evolution is **3-7-8-7-8-4**. It is worth noting that most of these systems experience two mass transfer phases in COSMIC before reaching WD+MS. That is, during the evolution process, two CE processes take place, between which there is a period of time. Statistically, 3146556 of the 3200000 binaries (98%) we simulated reach WD+MS in their evolutionary process. In these 3146556 binaries, 2660287 of them went through two CE processes before reaching WD+MS (85%). To explain this qualitatively, we notice that when the primary loses too much mass during the first CE, the CE process ends, as the radius of the primary becomes smaller than the Roche lobe. After the primary fills its Roche lobe and has a large enough envelope again, the CE process resumes and causes two CE processes in total. On the contrary, for the rest of the population that reached WD+MS with only one CE, the mass loss of the primary is not enough for the first CE to stop, and the evolution reaches WD+MS directly.

There are several interesting points of the results. First of all, comparing the four panels, we found that as α_{CE} decreases, the final separation decreases in general. This can be explained by Equation 3. Qualitatively, note that more energy from the orbit is needed to eject the envelope for low α_{CE} , so the resulting a_f is expected to be smaller.

From Figure 1, we also notice that final separation a_f jumps obviously at about 2.9 au, and 7.1 au. After investigating the `bpp` array of these simulations, we find out the jump at about 2.9 au and 7.1 au is likely due to the difference of `kstar_1` at the start of RLOF. For $a_i < 2.9 \text{ au}$, RLOF starts when `kstar_1 = 3`. For $a_i = 2.9 \text{ au}$, RLOF starts when `kstar_1 = 4`. For $2.9 \text{ au} < a_i < 7.1 \text{ au}$, RLOF starts when `kstar_1 = 5`. For $a_i > 7.1 \text{ au}$, RLOF starts when `kstar_1 = 6`. For different star type,

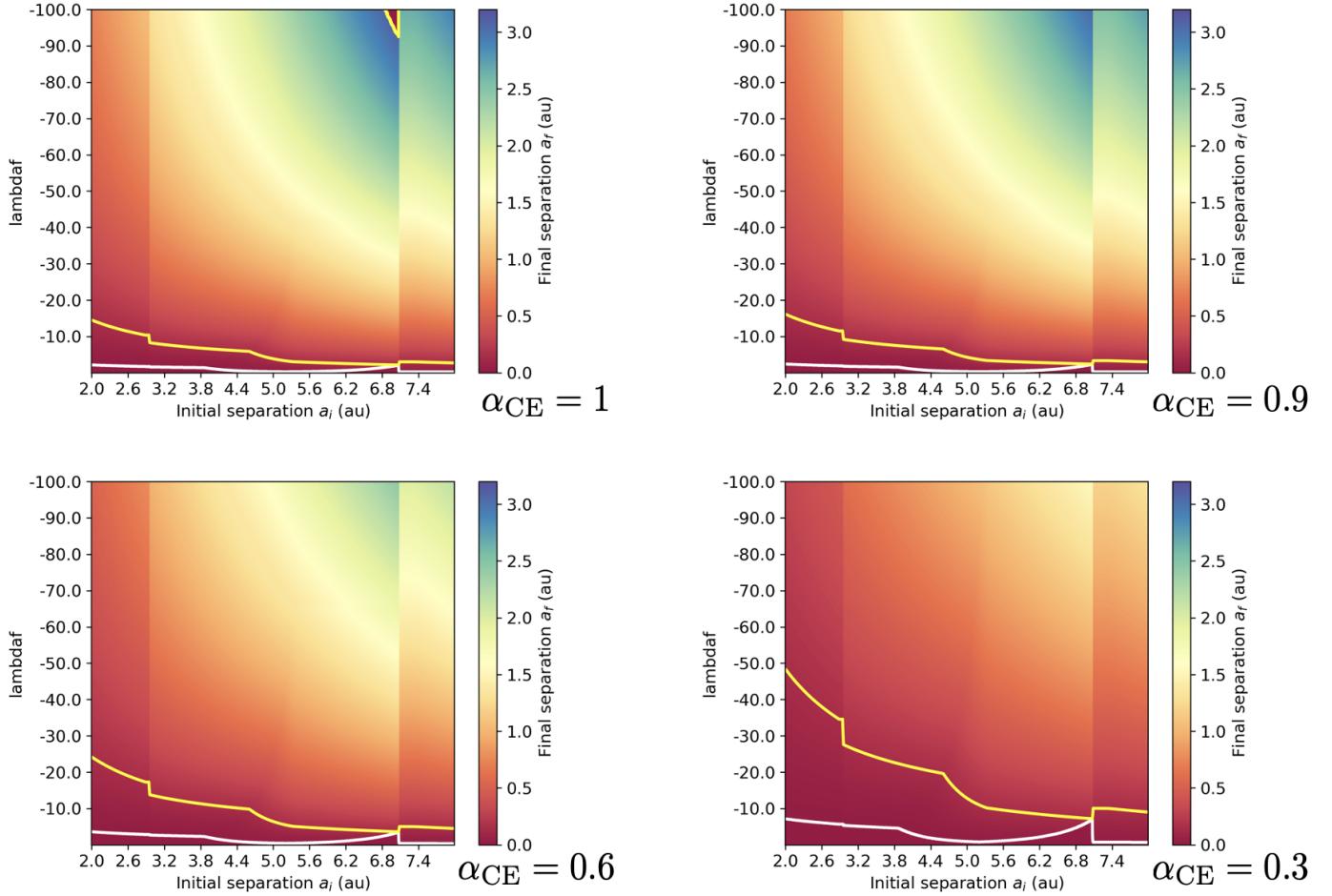


Figure 1: Heat map of final separation a_f against initial separation a_i and CE flag $-\lambda$. Different panels represent four different CE efficiencies $\alpha_{CE} = 1, 0.9, 0.6$, and 0.3 . The white and yellow contour corresponds to $a_f = 0.01$ au and $a_f = 0.15$ au respectively. WD+MS binaries is not possible for separation smaller than 0.01 au, and 0.15 au is the minimum separation of wide WD+MS binaries in observational results. Outside the white contour no desired system forms.

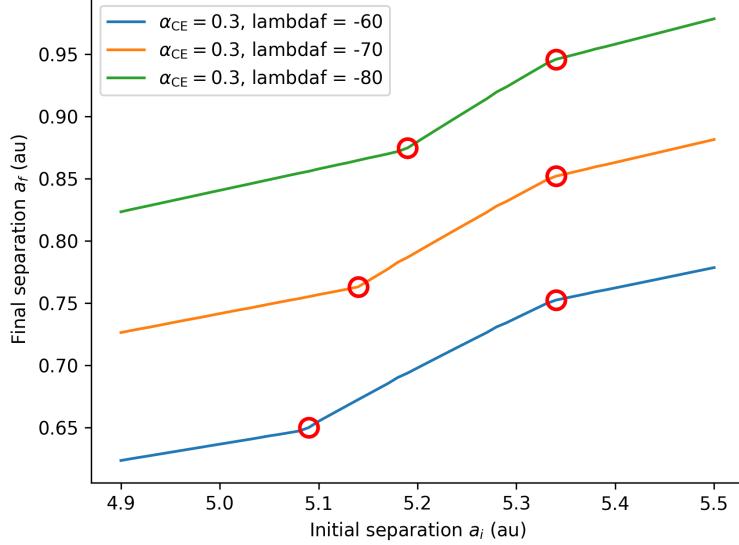


Figure 2: Dependence of final separation a_f on initial separation a_i for $\alpha_{CE} = 0.3$. Simulation results of $\lambda_f = 60, 70$, and 80 are shown. The red circles emphasize the jump of final separation as initial separation increases.

the structure and mass of the envelope is different, leading to different binding energy E_{bind} and thus different final separation a_f .

Inspecting more closely, we found that the final separation a_f jumps at about $5.2 \sim 5.4$ au. This is especially obvious with $\alpha_{CE} = 0.3$. To illustrate, we zoom in the simulation results of $\alpha_{CE} = 0.3$, and present the change of final separation a_f as initial separation a_i increases in Figure 2. From here, we can clearly see two jumps — one between 5.1 au and 5.2 au, the other at 5.34 au. To explain this, we again go to the `bpp` array of these systems. The first jump is because the evolution process changes from [3-7-8-4-3-7-8-4] to [3-7-8-7-8-4]. This means that for relatively larger initial separation a_i , the primary fills its Roche lobe all the time, leading to more mass loss during a shorter timescale and thus larger final separation. For the second jump, it is because `kstar_1` changes from 8 to 9 at the start of the second common envelope. This changes the binding energy of the common envelope and thus changes the final separation.

Finally, notice that there exists a small triangular gap in the figure for $\alpha_{CE} = 1$. This is because `kstar_1` becomes a neutron star and no rows in the `bpp` gets selected.

In conclusion, with a large λ , it is possible to create wide post-mass transfer WD+MS systems with final separation $a_f > 0.15$ au, which is the minimal separation of our observed objects. In next section, we will compare these results with default settings in COSMIC and with results in [Yamaguchi et al., 2024a], in order to investigate the possibility of forming wide WD+MS systems through CE.

2.1.3 Comparison

After we have obtained the simulation results for different initial conditions, now we can calculate the effective λ that matches with results in [Yamaguchi et al., 2024a]. For each fixed initial separation a_i , we loop through the COSMIC results with the same initial separation, and find the λ value that results in a final separation closest to MESA model results. We record this λ as the effective λ that matches

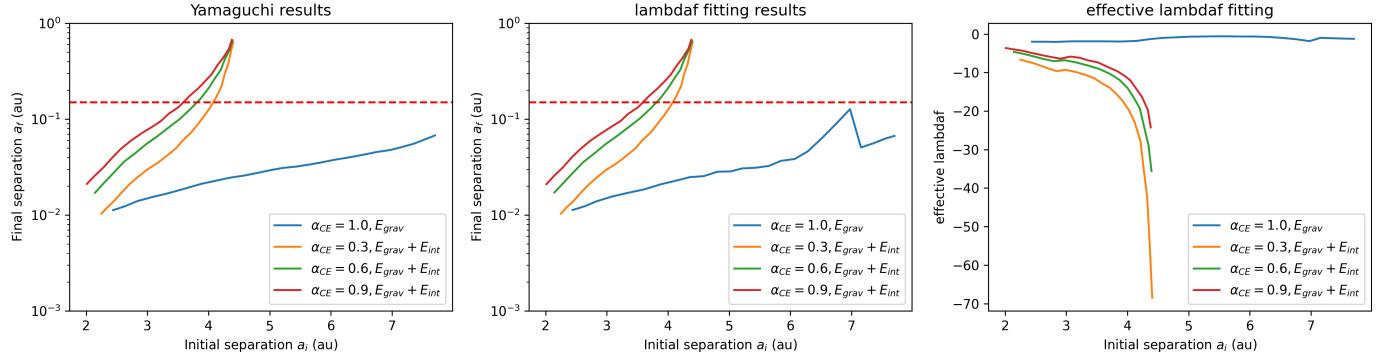


Figure 3: Effective λ fitting results of a $7 M_\odot + 1 M_\odot$ system for common envelope structure in MESA. Different lines represent different α_{CE} and different common envelope energy budget, as shown in the legend. Left panel: recreation of MESA results in [Yamaguchi et al., 2024a] of how final separation a_f depends on initial separation a_i . Central panel: COSMIC results of how final separation a_f depends on initial separation a_i , produced by the effective λ calculated. Right panel: effective λ in COSMIC that corresponds to the MESA results for each initial separation a_i . The red dashed line in the left and middle panel represents 0.15 au, which is the minimal separation for observed wide post-CE WD+MS binaries.

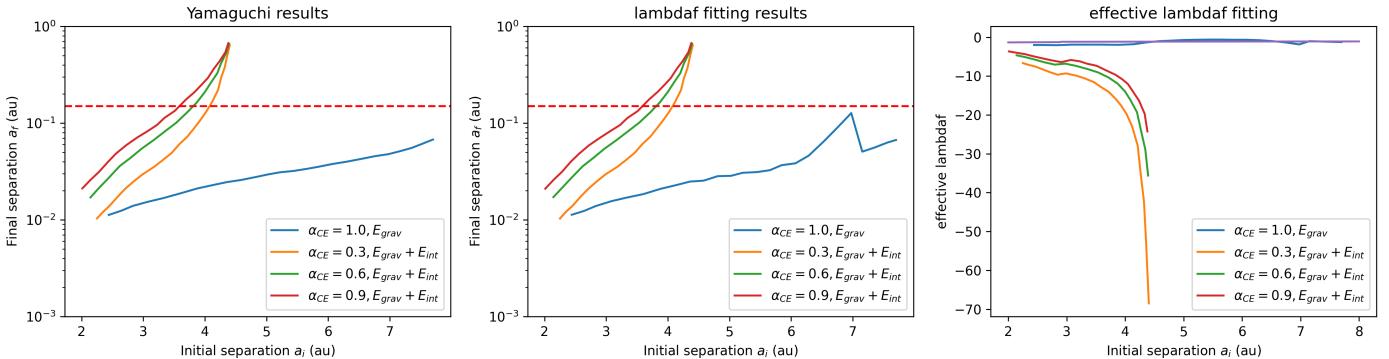


Figure 4: Same as Figure 3, but with default λ in COSMIC model included in the right panel with the purple line.

the binding energy formalism at initial separation a_i .

We present the fitting results for energy budget $E_{bind} = E_{grav} + E_{int}$ in Figure 3. Notice that in the central panel, there is a small bump of the blue line. This is because at $6 \sim 7$ au, we cannot reach such low final separation in COSMIC while forming a WD+MS binary at the same time. From the left panel, we can see that it is possible to create WD+MS binaries with $a_f > 0.15$ au in MESA as long as internal energy E_{int} is included, with initial separation $a_i > 3.5$ au. However, in the right panel we show that this correspond to very large λ in COSMIC ($\lambda > \sim 20$).

To further compare with the default λ value in COSMIC, we include in the right panel an extra line which represents the default λ value in COSMIC. This default λ in the COSMIC model is calculated following Appendix A of [Claeys et al., 2014]. For all of our case, we have $M_{env} > 1$. The results is shown in Figure 4. Notice that the default λ value is at order ~ 1 , far smaller than the effective λ needed to recreate results in [Yamaguchi et al., 2024a]. Hence, we conclude that in COSMIC, it is not practical to form wide high mass WD+MS system through common envelope processes, since the required λ is too large.

In Figure 5, we present the same results for another energy budget of the CE process. Rather than

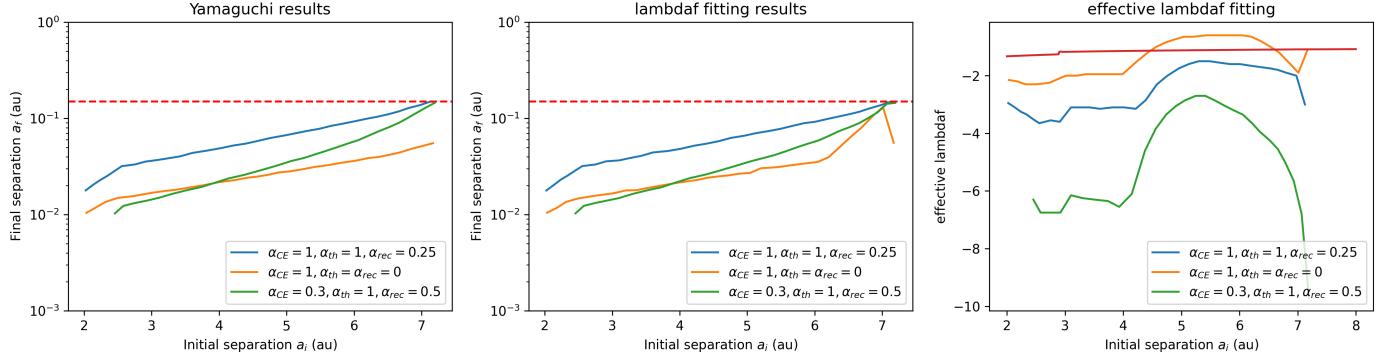


Figure 5: Same as Figure 4, but with a different common envelope energy budget. For results shown in these figures, $E_{\text{bind}} = E_{\text{grav}} + \alpha_{\text{th}}E_{\text{th}} + \alpha_{\text{rec}}E_{\text{rec}}$, where the coefficients α_{th} and α_{rec} for different line is shown in the legend. The red line in the right panel indicates default λ in COSMIC.

including both thermal and recombination energy into the binding energy, we use two coefficients α_{th} and α_{rec} to consider their contribution to E_{bind} separately. That is,

$$E_{\text{bind}} = E_{\text{grav}} + \alpha_{\text{th}}E_{\text{th}} + \alpha_{\text{rec}}E_{\text{rec}}.$$

We notice that the effective λ now gets with in the order of ~ 10 . However, none of these initial separation and λ results in wide WD+MS system with final separation $a_f > 0.15$ au. Therefore, we conclude that with the new energy budget, it is not practical to form wide WD+MS binaries either.

2.2 Low Mass Systems ($1.5 M_{\odot} + 0.85 M_{\odot}$)

2.2.1 Methods and Results

According to [Yamaguchi et al., 2024b], the WD mass in the Gaia sample ranges from $0.5 M_{\odot}$ to $0.8 M_{\odot}$, which corresponds to progenitor mass of $1 M_{\odot}$ to $3 M_{\odot}$. The companions mass has a median of $0.85 M_{\odot}$. MESA results for $1.5 M_{\odot} + 0.85 M_{\odot}$ systems is documented in [Yamaguchi et al., 2024a]. Hence, we consider a $1.5 M_{\odot} + 0.85 M_{\odot}$ evolution model in COSMIC and compare the results with those in [Yamaguchi et al., 2024a]. For the common envelope, we include all thermal energy E_{th} into the binding energy E_{bind} , while the recombination energy E_{rec} is included in part with a parameter α_{rec} . That is,

$$E_{\text{bind}} = E_{\text{grav}} + E_{\text{th}} + \alpha_{\text{rec}}E_{\text{rec}}.$$

We take the same approach as high mass systems to investigate the effective value of λ in COSMIC that matches with the results in MESA. The initial separation is a linear grid in range $0.5 \sim 6$ au with 400 points. The CE flag `lambdadaf` is a linear grid in range $0 \sim -100$ with 2000 steps, which corresponds to $0 \sim 100$ for CE parameter λ . Also, we ran these parameters with four different CE efficiencies — $\alpha_{\text{CE}} = 1$, $\alpha_{\text{CE}} = 0.9$, $\alpha_{\text{CE}} = 0.6$, and $\alpha_{\text{CE}} = 0.3$, the same as recorded in [Yamaguchi et al., 2024b]. In total, for each of the 4 different CE efficiency α_{CE} , we ran a grid of 800000 binary star systems with different initial separation a_i and CE parameter λ .

After simulation is finished, we again select the time-step when the system finished CE for the first time, and record the final separation if the system is a desired WD+MS system. The result is presented in Figure 6.

To compare the results with [Yamaguchi et al., 2024b], we use the same approach as that for high mass systems. In Figure 7, we compare COSMIC with MESA model that includes binding energy of the entire envelope. In Figure 8, we compare COSMIC with MESA model that includes only binding

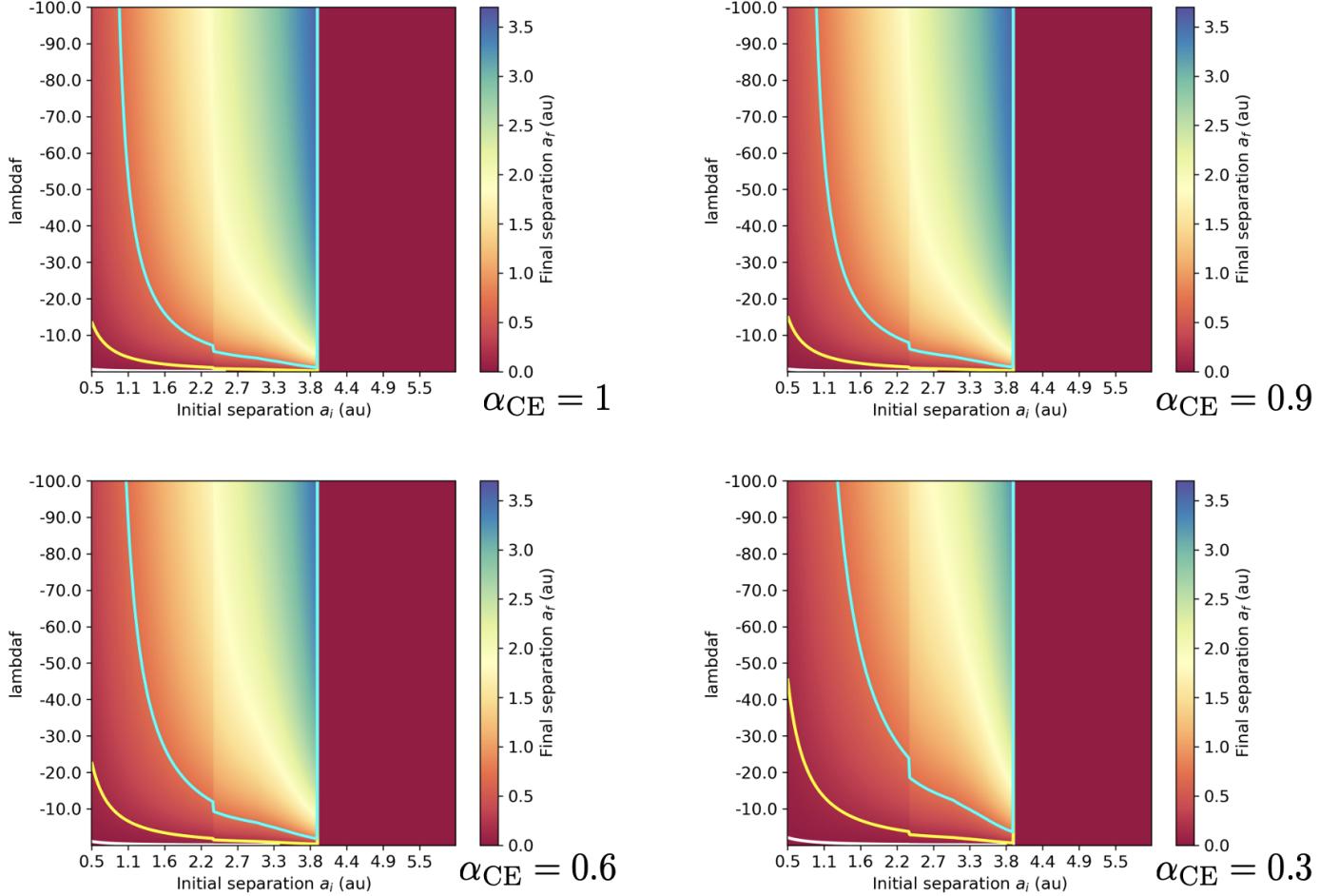


Figure 6: Heat map of final separation a_f against initial separation a_i and CE flag $-\lambda$. Different panels represent four different CE efficiencies $\alpha_{\text{CE}} = 1, 0.9, 0.6$, and 0.3 . The white and yellow contour corresponds to $a_f = 0.01$ au and $a_f = 0.15$ au respectively. WD+MS binaries is not possible for separation smaller than 0.01 au, and 0.15 au is the minimum separation of wide WD+MS binaries in observational results. Outside the white contour no desired system forms.

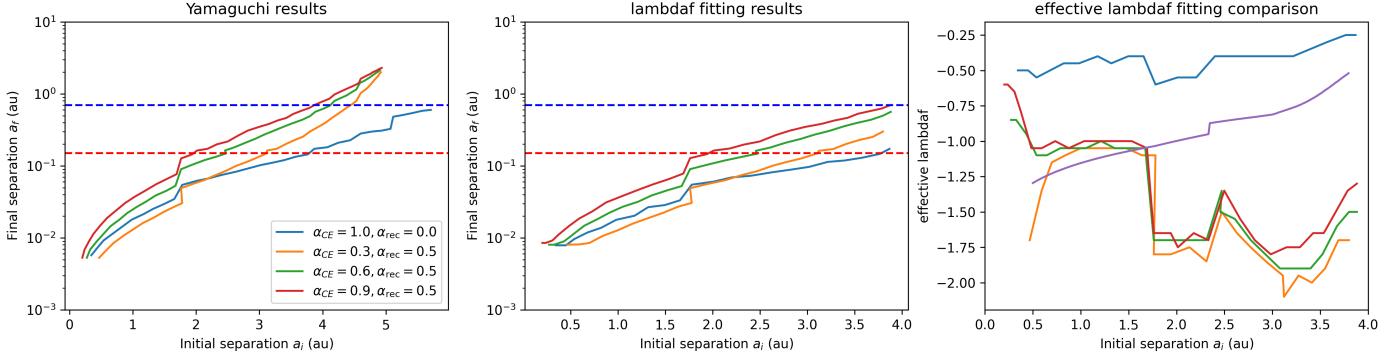


Figure 7: Effective λ fitting results of a $1.5 M_\odot + 0.85 M_\odot$ system for common envelope structure in MESA. Binding energy of the *entire envelope* is considered. Left Panel: MESA results in [Yamaguchi et al., 2024b] of how final separation a_f depends on initial separation a_i . Central Panel: COSMIC results of how final separation a_f depends on initial separation a_i , produced by the effective λ . Right Panel: effective λ in COSMIC that corresponds to the MESA results for each initial separation a_i . The purple line represents default λ calculated using [Claeys et al., 2014]. The red dashed line and the blue dashed line in the left and central panel represents 0.15 au and 0.7 au respectively. 0.15 au is the minimal separation for observed wide post-CE WD+MS binaries, while 0.7 au is the typical separation for wide post-CE WD+MS binaries.

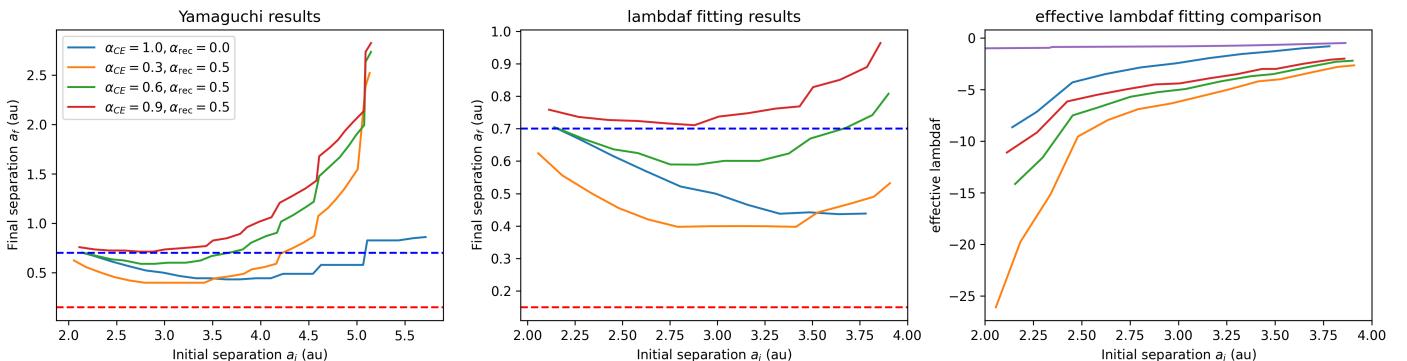


Figure 8: Same as Figure 8 but only binding energy of the *outer envelope* is considered.

energy of the outer envelope. Notice that we do not have fitting results for $a_i > 4$ au. This is because no WD+MS binaries form in COSMIC for such initial separation range.

From the left panel of Figure 7, we can see that wide WD+MS systems can form in MESA model when $a_i > 2.5$ au. This is reproduced in COSMIC with $\lambda \sim 1$. Hence, wide WD+MS can form for $a_i > 2.5$ au with parameter λ about the same as default. Recall that for high mass system, wide WD+MS systems cannot form without $\lambda \sim 10^2$. This is a significant difference between high mass systems and low mass systems.

3 Formation Through Stable Mass Transfer

After exploring the formation of WD+MS binaries through common envelope, we investigate the possibility of forming WD+MS binaries through stable mass transfer. In [Yamaguchi et al., 2024b], the author plot the $M_{WD} - P_{\text{orb}}$ relation of observed wide WD+MS binaries, and argues that it is not possible to form wide WD+MS binaries through stable mass. The plot is recreated in Figure 9, and from the figure, we can see that the observed wide post-mass transfer WD+MS binaries do not match the

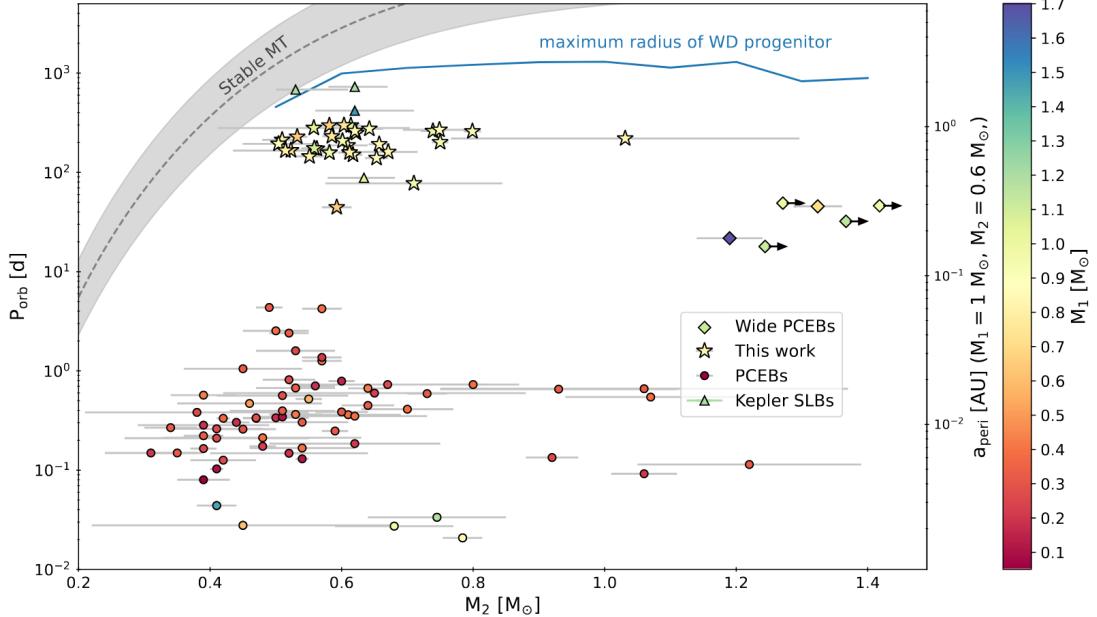


Figure 9: $M_{\text{WD}} - P_{\text{orb}}$ relation of observed wide post-mass transfer binaries. White dwarf mass M_{WD} correspond to The colors of the points represent the masses of the luminous companion, M_* . On the right axis, we show the minimum orbital separation a_{peri} , corresponding to the period assuming companion mass $M_* = 1 M_\odot$ and white dwarf mass $M_{\text{WD}} = 0.6 M_\odot$. The dashed line and the grey region is the track along which binaries undergoing stable MT are expected to evolve according to [Rappaport et al., 1995]

theoretical $M_{\text{WD}} - P_{\text{orb}}$ relation for stable mass transfer calculated in [Rappaport et al., 1995]. Also, the author argues that stable mass transfer requires donors on SGB at the onset of mass transfer. However, this will not lead to WDs with masses observed [Yamaguchi et al., 2024b]. We will verify this idea in this section using $7 M_\odot + 1 M_\odot$ binaries.

3.1 Methods and Results

First of all, to ensure stable mass transfer in COSMIC, we modify every entry of `qcrit_array` in `BSEDict` from the default 0.0 to 20.0. This changes the critical mass ratio for onset of unstable mass transfer to 20.0. Based on our initial mass ratio 7 : 1, this will make the actual mass ratio of our system below this critical throughout the evolutionary process.

Secondly, a parameter on which stable mass transfer process depends is the accretion efficiency β , which in COSMIC is represented by `acc_lim` in `BSEDict`. Accretion efficiency is the fraction β of all transferred materials that is accreted by the accretor. That is, $\dot{M}_a = -\beta \dot{M}_d$ where M_d is the mass of donor and M_a is the mass of accretor.

In our simulation, we investigate how accretion efficiency β and initial separation a_i affects the final separation a_f . We run a grid of three different accretion efficiency — $\beta = 0.0, 0.5$ and 1.0 , and a linear grid of a_i from 2 au to 8 au with 400 steps. In total, we simulate 1200 binaries with different accretion limit and initial separation. For each system, after the simulation is finished, we select the first timestep, if any, that shows a WD+MS binary. We record the separation at this moment as the final separation a_f , and we compare them with the observed minimum separation 0.15 au for post-mass transfer WD binaries.

The results are presented in Figure 10. Notice the gap for a_i ranges in $3 \sim 4.5$ au. This means no WD+MS forms for this initial separation. In fact, for these initial separations and accretion efficiencies,

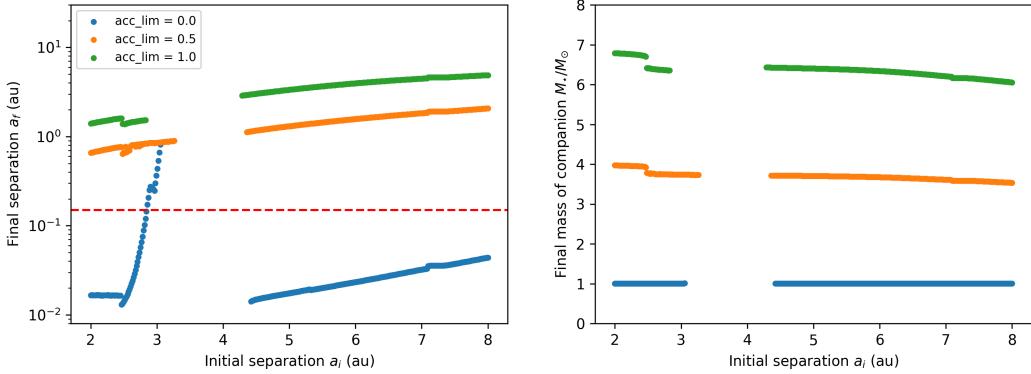


Figure 10: COSMIC results of WD+MS binaries after stable mass transfer for $7M_{\odot} + 1M_{\odot}$ systems, with accretion efficiency $\beta = 0$, $\beta = 0.5$, and $\beta = 1$ respectively. Left panel: dependence of final separation a_f of the WD+MS binary on initial separation a_i . Right panel: the corresponding companion mass in solar units in the WD+MS binary. The gap in both figure represents that no WD+MS system forms for that initial separation.

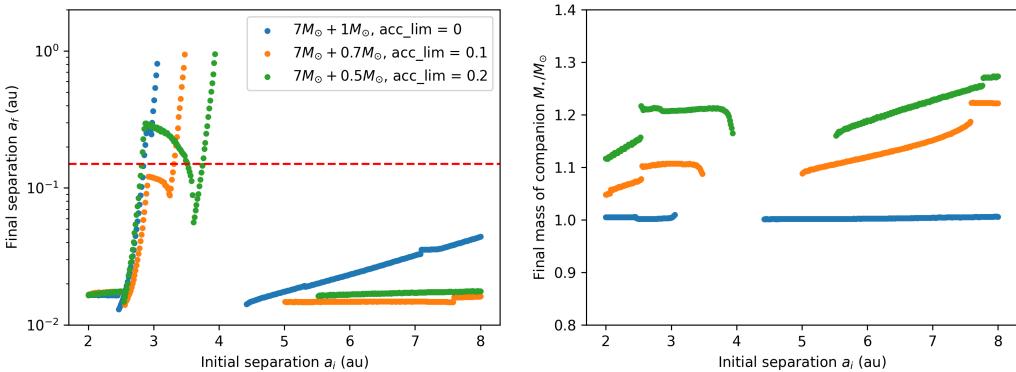


Figure 11: COSMIC results of WD+MS binaries after stable mass transfer for two systems — $7M_{\odot} + 0.7M_{\odot}$ with accretion efficiency $\beta = 0.1$, and $7M_{\odot} + 0.5M_{\odot}$ with accretion efficiency $\beta = 0.2$. The left and right panel presents the same content as in Figure 10.

two star merges.

Notice that the majority of these systems have final separation $a_f > 0.15$ au. However, this is **not the correct approach**, since the companion mass in our observed wide WD+MS binaries has median mass $\sim 1 M_{\odot}$. However, in the simulation results shown in Figure 10, the final companion mass is about $4 M_{\odot}$ and $6.5 M_{\odot}$ for $\beta = 0.5$ and $\beta = 1.0$ respectively. Hence, we need to adjust the initial companion mass $M_{\star,i}$ and the accretion efficiency β , in order to make the final companion mass $M_{\star,f}$ close to $1 M_{\odot}$. Therefore, we lower the initial companion mass $M_{\star,i}$ and adjust the accretion efficiency β and re-run the simulation. The results for $7M_{\odot} + 0.7M_{\odot}$ with $\beta = 0.1$, and $7M_{\odot} + 0.5M_{\odot}$ with $\beta = 0.2$ is shown in Figure 11.

After adjusting the initial mass of companion and the accretion efficiency, now the final mass of the companion is about $1 M_{\odot}$ and we get systems similar to our observed samples. Also, in Figure 11, we note that when initial separation a_i lies in the range of $2.5 \sim 4.5$ au, the final separation increases sharply. Also, it is only in this region that a_f goes larger than the observed minimum 0.15 au for wide WD binaries. Hence, in the next subsection, we will try to explain this sharp increases qualitatively.

3.2 Discussions

In the previous section, we mentioned how $7 M_{\odot} + 1 M_{\odot}$ systems with $\beta = 0$, $7 M_{\odot} + 0.7 M_{\odot}$ systems with $\beta = 0.1$, and $7 M_{\odot} + 0.5 M_{\odot}$ systems with $\beta = 0.2$ produces WD+MS binaries with final separation $a_f > 0.15$ au and companion mass $\sim 1 M_{\odot}$ in certain range of a_i .

From figure 11, we can see that the dependence of a_f on a_i can be roughly divided into four parts. The **pre-spike** part before the sharp increase with $a_i < 2.5$ au, the **spike** part where the a_f increases sharply with $2.5 \text{ au} < a_i < 3.5 \text{ au}$, the gap where the two stars merge with $3.5 \text{ au} < a_i < 4.5 \text{ au}$, and the part after that with $a_i > 4.5$ au.

Typical **pre-spike** and **spike** systems both follows the evolution path **3-5-7-8-4-3-4** before reaching WD+MS. Note that the binaries come into contact (`evol_type = 5`), which follows by a common envelope phase. After inspecting the evolutionary process more closely, we found that the second mass transfer (the stable one) tends to widens the orbit significantly.

To qualitatively explain the widening of the orbits, notice that for stable and non-conservative mass transfer with $\dot{M}_a = -\beta \dot{M}_d$, we have

$$\frac{\dot{J}}{J} \propto (1 - \beta) q_d \frac{\dot{M}_d}{M_T},$$

where $q_d = M_d/M_a$ is the mass ratio between the donor and the accretor. Also,

$$\frac{\dot{a}}{a} = 2 \frac{\dot{J}}{J} - 2 \frac{\dot{M}_d}{M_d} - 2 \frac{\dot{M}_a}{M_a} + \frac{\dot{M}_T}{M_T}.$$

Hence, for stable and non-conservative mass transfer, the widening of orbit can be evaluated with

$$\frac{\dot{a}}{a} \propto - \left(\frac{2}{q_d} - 2q_d - \beta \right) \frac{\dot{M}_d}{M_T}. \quad (5)$$

Also note that the mass transfer begins when the primary has `evol_type = 8`, a naked Helium star on the Hertzsprung gap, whose wind mass loss is significant. For a fast wind mass loss, we have

$$\frac{\dot{a}}{a} = - \frac{\dot{M}_d}{M_T}. \quad (6)$$

This means that strong wind stripping mass from the primary will cause the orbit to expand even further. To illustrate this, we change the `windflag` in `BSEDict` to 0, which correspond to weak wind. Now the $a_i = 2.80$ au system has final separation $a_f = 0.05$ au, and the $a_i = 2.90$ au system has final separation $a_f = 0.14$ au, both smaller than the original 0.09 au and 0.27 au that correspond to stronger wind.

For our systems, the mass ratio $q_d > 1$, and the value of q_d increases as the initial separation increases. Hence, based on Equation 5, the stable mass transfer process actually shrinks the orbit, while the wind mass loss compensates it by widening the orbit. For **spike** binaries, although q_d is slightly higher, causing the orbit to shrink more according to Equation 5, the mass loss rate is significantly higher than **pre-spike** binaries, leading to a more significant increase in orbital separation. Also, **spike** binaries have larger orbital separation a before mass transfer, which also leads to more significant increase in separation during the mass transfer.

For **merge** binaries, we notice that the separation of binaries at the start of CE is smaller than **spike** binaries. This lack of orbital energy to eject the CE explains the merger event. As initial separation continues to increase, the binaries possess more energy to eject the CE after coming into contact, explaining the reappearance of WD+MS binaries instead of merger.

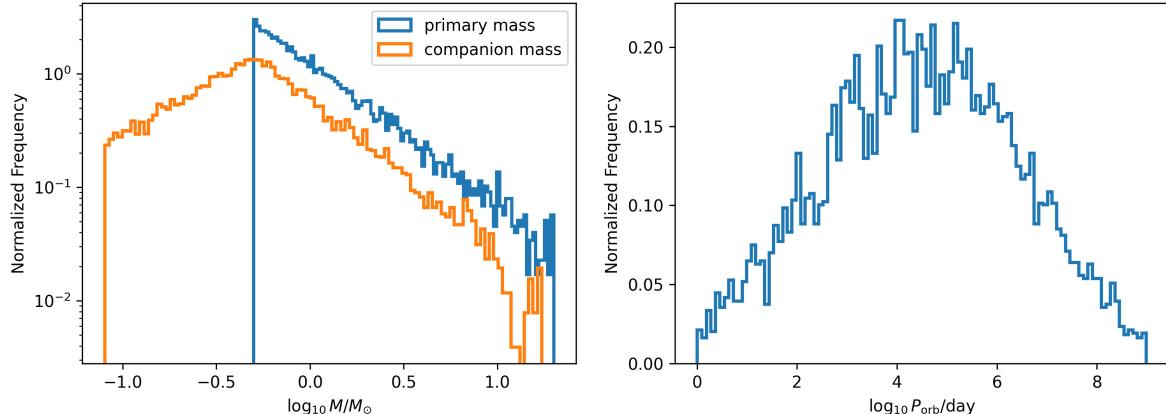


Figure 12: The distribution of primary mass, companion mass, and orbital period in the sampled population using independent sampler. Left panel: distribution of primary mass and companion mass in log scale. Right panel: distribution of orbital period P_{orb} in log scale.

This concludes our qualitative analysis on the evolution results of $7 M_\odot + 1 M_\odot$ through stable mass transfer. Again, as shown in Figure 10, for stable mass transfer, the observed minimum separation of wide post-mass transfer WD+MS systems 0.15 au can only be achieved in the **spike** region. That is, in our results, all wide post-mass transfer systems are due to the strong wind mass loss from the primary that widens the orbit significantly.

4 Population Synthesis Results

After investigating how the final separation depends on various parameters in individual binary star systems, we now shift our focus onto the distribution of final orbital period in a population of binaries (equivalent to final separation since $P_{\text{orb}}^2 \sim a^3$).

4.1 Methods

First of all, we sample the binary population using the independent sampler in COSMIC. We set the primary mass model to `korupa01`, eccentricity model to `uniform`, and orbital period model, together with the binary fraction model to `moe19`. We sample a total number of about 10000 binaries with solar metallicity. The distribution of primary mass and the orbital period in the sampled population is shown in the Figure 12.

To investigate how the final white dwarf mass and orbital period depends on `lambdaf`, we simulate the same sampled population with five different λ values, $\lambda = 5, 15, 25, 35, 45$. After that, we select the first time-step that satisfies `evol_type = 8`, `kstar_1` is WD, and `kstar_2` is MS. We record the white dwarf mass M_{WD} and the final orbital period $P_{\text{orb},f}$ from this moment. We also select the moment when the first Roche lobe overflow starts, and record the initial orbital separation $P_{\text{orb},i}$ and `kstar_1` type at this moment. These variables will later help us study how the orbit widens, with respect to different parameter λ and initial separation a_i .

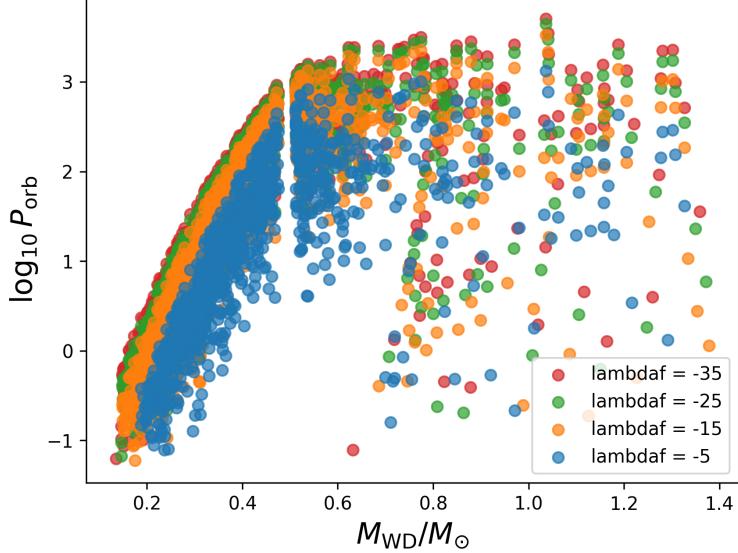


Figure 13: Scatter plot of white dwarf mass M_{WD} and final orbital period after common envelope in log scale $\log_{10} P_{\text{orb}}$. Different colors indicates results for different λ values. For each different λ value the same initial population is used.

4.2 Results

To show how the change in λ affects the final orbital period, we plot final orbital period P_{orb} against white dwarf mass M_{WD} in Figure 13. From Figure 13, we note that the final orbital period is larger in general for larger λ . This can be clearly seen from the figure for M_{WD} ranges in $0.2 \sim 0.7$ au. This is also expected based on Equation 4, since for larger λ , less energy is required to unbind the common envelope, leading to more energy after mass transfer and a larger orbital period.

To also show how the final separation P_{orb} depends on `kstar_1` type at the start of first RLOF, we plot P_{orb} against M_{WD} for different λ separately, and use different color to distinguish between different `kstar_1` type at the start of mass transfer. The plot is shown in Figure 14. In general, if the star is at later stages of evolution (larger `kstar`) when mass transfer starts, then the resulting white dwarf mass M_{WD} and P_{orb} are also both larger. Note also that few systems start the mass transfer at `kstar_1` = 2 (Hertzsprung gap) and `kstar_1` = 4 (Core Helium Burning). This is because stars crosses Hertzsprung in a short timescale, leaving little time for interaction. Also, stellar radius is small when going through Core Helium Burning, reducing the possibility of interactions that starts when `kstar_1` = 4.

To visualize the difference in final orbital period P_{orb} , and how they corresponds to different `kstar_1` type at the start of RLOF, we create plot the difference in orbital period $\Delta P = P_{\text{orb}}^{\lambda=35} - P_{\text{orb}}^{\lambda=5}$ against the white dwarf mass M_{WD} in Figure 15. From Figure 15, we can see that for large `kstar_1` type at the start of RLOF, $\Delta P = P_{\text{orb}}^{\lambda=35} - P_{\text{orb}}^{\lambda=5}$ is also larger. To explain this, we plot the how final white dwarf mass M_{WD} , difference in final orbital period ΔP_{orb} , and orbital period before mass transfer $P_{\text{orb},i}$ related to each other in Figure 16. We see that the larger initial orbital period $P_{\text{orb},i}$, the more significant effect does λ have on the final orbital period. This is not unexpected since during mass transfer, the separation changes as $\dot{a}/a = \text{const}$. Hence $P_{\text{orb}}/P_{\text{orb},i} = \text{const}$. Therefore, the larger the initial orbital period, the larger it changes, and thus the larger effect does λ have.

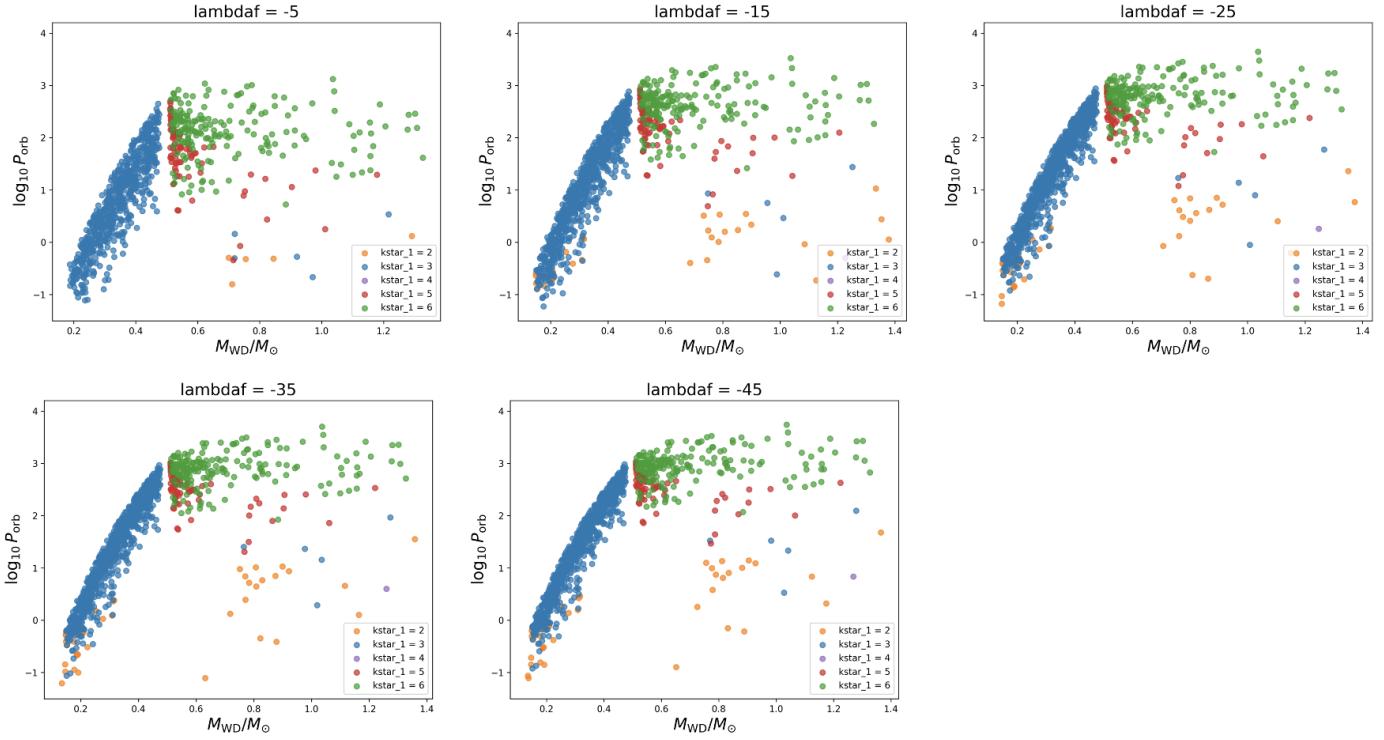


Figure 14: Each panel: for this λ value in the title, scatter plot of white dwarf mass M_{WD} and final orbital period after common envelope in log scale $\log_{10} P_{\text{orb}}$. Different colors represent different kstar type of the primary at the start of mass transfer.

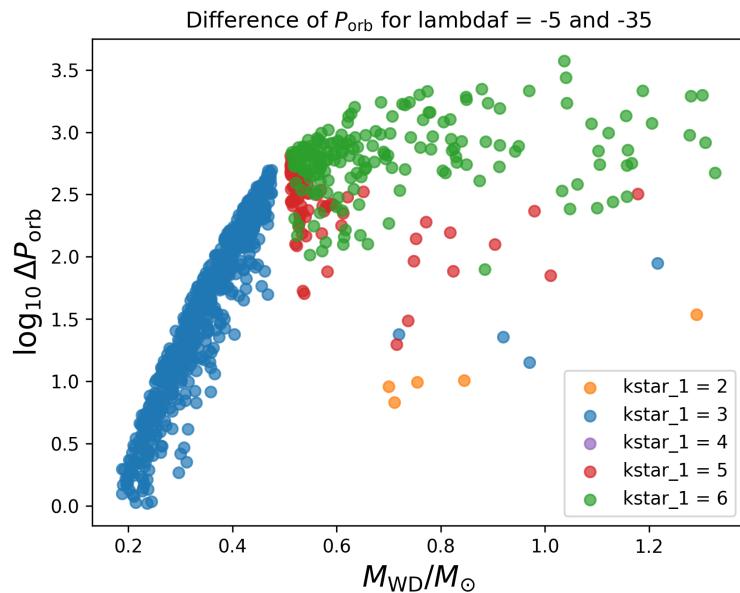


Figure 15: Scatter plot of difference in orbital period $\Delta P = P_{\text{orb}}^{\lambda=35} - P_{\text{orb}}^{\lambda=5}$ against the white dwarf mass M_{WD} . Different colors indicates different kstar type of the primary at the start of mass transfer.

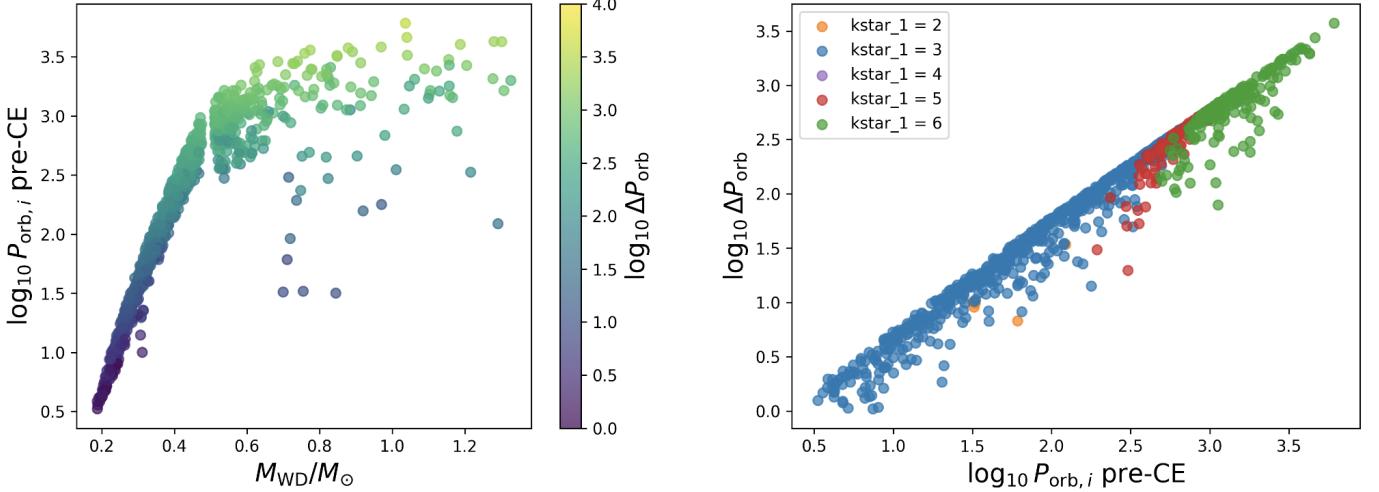


Figure 16: Left panel: scatter plot of initial orbital period at the start of RLOF $P_{\text{orb},i}$ against the final white dwarf mass M_{WD} . The color of the dots indicates how the final orbital period $P_{\text{orb},f}$ changes with respect to $\lambda = 35$ and $\lambda = 5$. That is, the color indicates the magnitude of $\Delta P = P_{\text{orb}}^{\lambda=35} - P_{\text{orb}}^{\lambda=5}$. Right panel: scatter plot of $\log_{10} \Delta P_{\text{orb}}$ against the initial orbital period in log scale $\log_{10} P_{\text{orb},i}$ at the start of RLOF. Different colors indicates different kstar type of the primary at the start of RLOF.

5 METISSE

To resolve inconsistencies of results across different population synthesis simulation, METISSE, acronym for METhod of Interpolation for Single Star Evolution is developed. As documented in documented in [Agrawal et al., 2020, Agrawal et al., 2023], it is a code package designed for rapidly computing the evolution of stellar populations by interpolating within a set of pre-computed stellar evolutionary tracks. To make simulation of binary star systems using METISSE more convenient and improve the simulation accuracy of binary evolution, we modify METISSE to interpolate and save the envelope binding energy of stars. In this way, when we need to calculate envelope binding energy in COSMIC, we no longer need to first calculate common envelope λ using the fit formula in [Claeys et al., 2014], and then use Equation 4 for binding energy. Instead, we can use the interpolated binding energy obtained from MESA tracks in METISSE directly, achieving improved accuracy and efficiency of simulation results.

Notice that binding energy is stored in CGS units in the input tracks of METISSE, so the value of energy is very large. Hence, it is the best for us to interpolate the log of binding energy per unit envelope mass, instead of binding energy directly. Also, for different tracks and different phase of evolution, the binding energy may either be positive or negative. Therefore, to keep the sign of binding energy accessible after interpolation, we preprocess the binding energy field in input tracks into

$$u = \text{sign } E_{\text{bind}} \cdot \frac{\log_{10} |E_{\text{bind}}|}{M_{\text{env}}}, \quad (7)$$

if the binding energy column is provided in the input tracks. Then, we interpolate u , just as other parameters are interpolated in METISSE. Finally, for non-degenerate stars, we recover the binding energy

$$E_{\text{bind}} = \text{sign } \hat{u} \cdot 10^{\hat{u} M_{\text{env}}}, \quad (8)$$

where \hat{u} is the interpolated value of u , and save E_{bind} in the `binding_energy` field in `star_parameters`. In `comenv_lambda.f90`, a file called by COSMIC when common envelope binding enrgy is needed, we

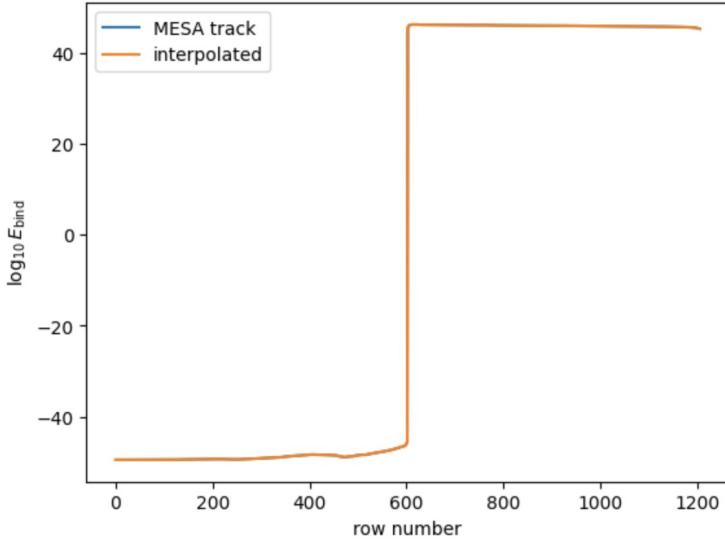


Figure 17: Comparison of interpolated binding energy using METISSE with binding energy stored in input track for initial mass $5.26 M_{\odot}$.

check if the binding energy column is interpolated. If so, we return the value that corresponds to the interpolated E_{bind} directly.

For a small caveat, path to input tracks must be provided in file `metissee.input`. Both hydrogen tracks and helium tracks is required, and the path must be a full path returned by the `pwd` command in terminal. That is, paths should not start with `~/foo`, but something similar to `/Users/john-doe/foo` instead. Also, there should not be spaces anywhere in the path, for example `Foo/foo__bar/Bar` for a directory `foo__bar`. This will cause the simulation to behave incorrectly.

After modifying the code as described above, we change the initial condition of the simulation in `main.input` and run the program to test it. We set the initial mass to be exactly the same initial mass as one input track. Then we compare the interpolated binding energy using our method with the binding energy stored in the input track. One example of this is shown in Figure 17. From the figure, we can see that the interpolated binding energy and the values in input MESA tracks has little deviation, validating our approach.

References

- [Agrawal et al., 2023] Agrawal, P., Hurley, J., Stevenson, S., Rodriguez, C. L., Szécsi, D., and Kemp, A. (2023). Modelling stellar evolution in mass-transferring binaries and gravitational-wave progenitors with <scp>metisse</scp>. *Monthly Notices of the Royal Astronomical Society*, 525(1):933–951.
- [Agrawal et al., 2020] Agrawal, P., Hurley, J., Stevenson, S., Szécsi, D., and Flynn, C. (2020). The fates of massive stars: exploring uncertainties in stellar evolution with metisse. *Monthly Notices of the Royal Astronomical Society*, 497(4):4549–4564.
- [Claeys et al., 2014] Claeys, J., Pols, O., Izzard, R., Vink, J., and Verbunt, F. (2014). Theoretical uncertainties of the type ia supernova rate. *Astronomy & Astrophysics*, 563:A83.
- [Rappaport et al., 1995] Rappaport, S., Podsiadlowski, P., Joss, P., Di Stefano, R., and Han, Z. (1995). The relation between white dwarf mass and orbital period in wide binary radio pulsars. *Monthly Notices of the Royal Astronomical Society*, 273(3):731–741.
- [Yamaguchi et al., 2024a] Yamaguchi, N., El-Badry, K., Fuller, J., Latham, D. W., Cargile, P. A., Mazeh, T., Shahaf, S., Bieryla, A., Buchhave, L. A., and Hobson, M. (2024a). Wide post-common envelope binaries containing ultramassive white dwarfs: evidence for efficient envelope ejection in massive asymptotic giant branch stars. *Monthly Notices of the Royal Astronomical Society*, 527(4):11719–11739.
- [Yamaguchi et al., 2024b] Yamaguchi, N., El-Badry, K., Rees, N., Shahaf, S., Mazeh, T., and Andrae, R. (2024b). Wide post-common envelope binaries from gaia: orbit validation and formation models. *arXiv preprint arXiv:2405.06020*.