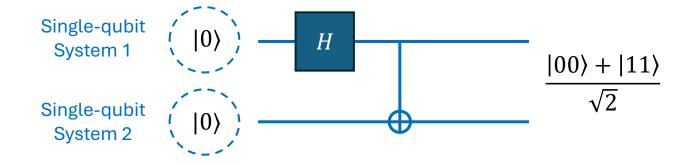
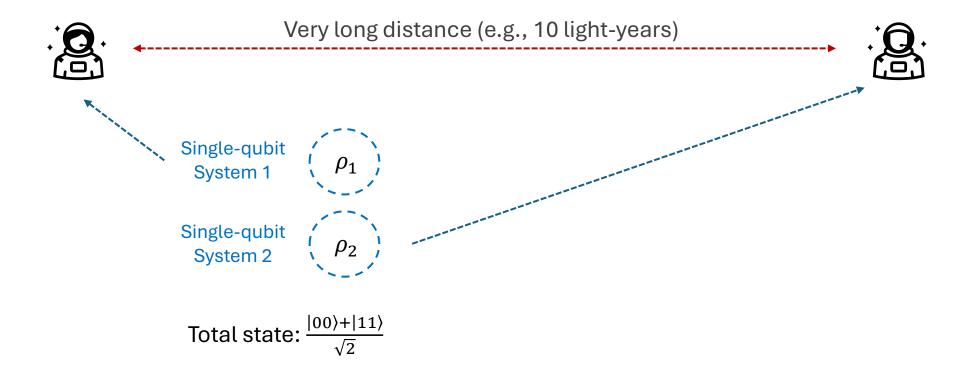
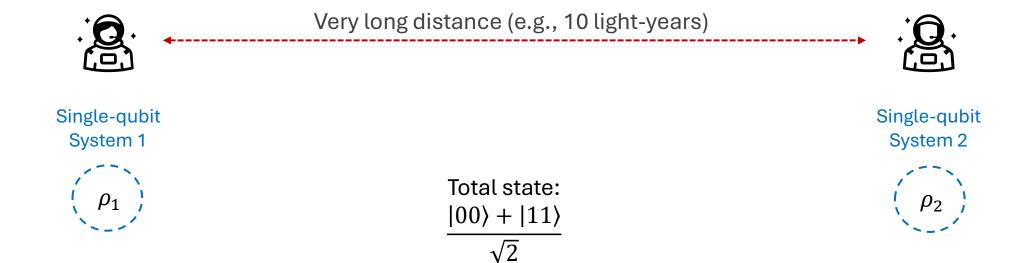
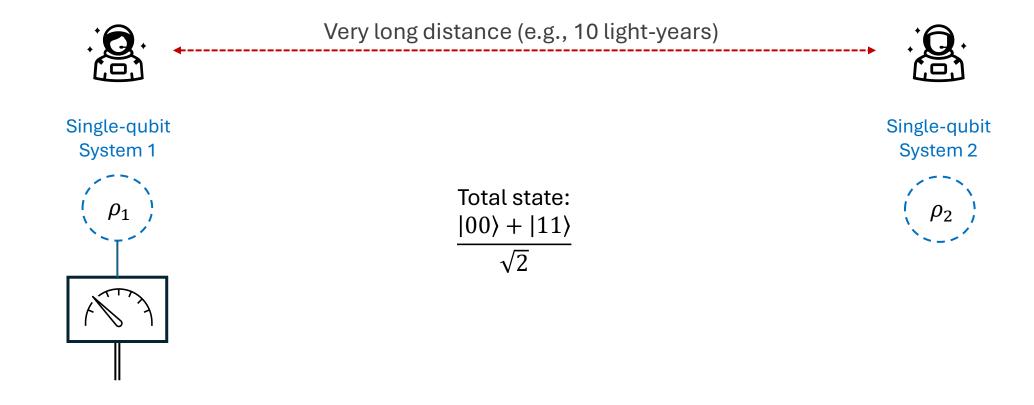
# **Quantum Computing**

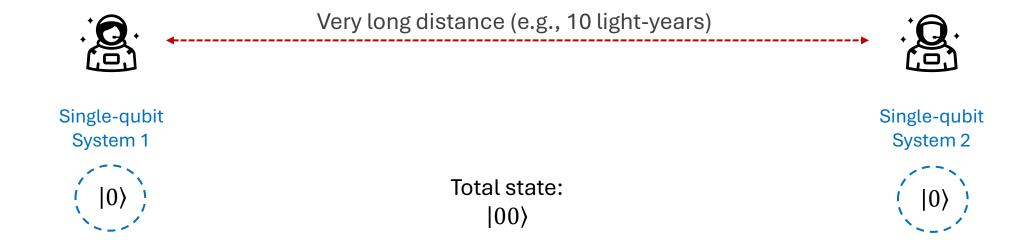
- Lectures 9 and 10 (June 4, 2025)
- Today:
  - Superdense coding
  - Quantum teleportation











• Application: Superdense coding

#### **Pauli Matrices**

Pauli matrices:

$$\mathbf{X} = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$
 (Pauli- $\mathbf{X}$ )

$$m{Y} = \sigma_2 \coloneqq egin{bmatrix} 0 & -i \ i & 0 \end{bmatrix}$$
 (Pauli- $m{Y}$ )

$$\mathbf{X} = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$
  $\mathbf{Y} = \sigma_2 \coloneqq \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$   $\mathbf{Z} = \sigma_3 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$  (Pauli- $\mathbf{X}$ ) (Pauli- $\mathbf{Z}$ )

- Some facts:
  - Pauli-X is the qNOT gate (in the computational basis)
  - $\sigma_i^2 = I \text{ for } j = 1,2,3$

#### **Pauli Matrices**

• (Extended) Pauli matrices:

$$I = \sigma_0 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$
  $X = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$   $Y = \sigma_2 \coloneqq \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$   $Z = \sigma_3 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$  (Pauli-X) (Pauli-Z)

- Some facts:
  - Pauli-X is the qNOT gate (in the computational basis)
  - $\sigma_i^2 = I \text{ for } j = 0,1,2,3$

Consider the four matrices

$$I = \sigma_0 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$
  $X = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$   $Z = \sigma_3 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$   $Z \cdot X = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$ 

$$\mathbf{Z} = \sigma_3 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

$$\mathbf{Z} \cdot \mathbf{X} = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

Let's define:

$$U_{b_1b_2} \coloneqq Z^{b_2}X^{b_1} egin{cases} U_{00} &= I \ U_{01} &= X \ U_{10} &= Z \ U_{11} &= Z \cdot X \end{cases}$$

Small Exercise:

$$\frac{|00\rangle + |11\rangle}{\sqrt{2}}$$

Consider the four matrices

$$I = \sigma_0 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$
  $X = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$   $Z = \sigma_3 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$   $Z \cdot X = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$ 

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Small Exercise:

$$\frac{|00\rangle + |11\rangle}{\sqrt{2}} \qquad \qquad |B_{b_1b_2}|$$

Consider the four matrices

$$I = \sigma_0 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$X = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

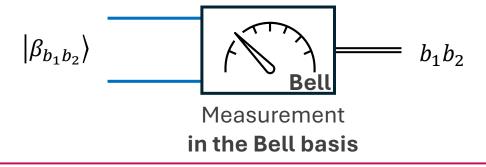
$$I = \sigma_0 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$
  $X = \sigma_1 \coloneqq \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$   $Z = \sigma_3 \coloneqq \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$   $Z \cdot X = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$ 

$$\mathbf{Z} \cdot \mathbf{X} = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

Let's define:

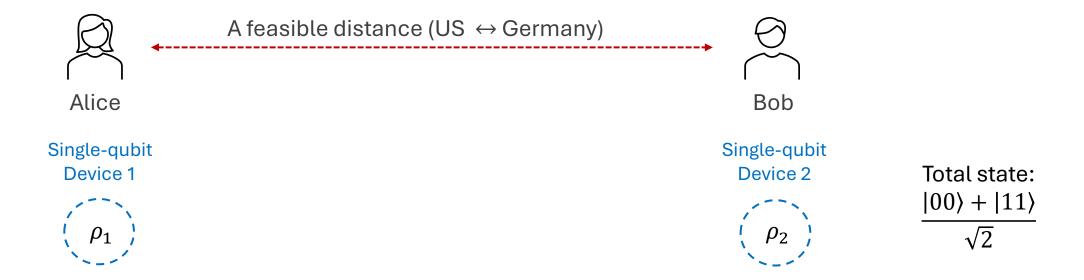
$$egin{aligned} egin{aligned} m{U_{01}} &= m{I} \ m{U_{01}} &= m{X} \ m{U_{10}} &= m{X} \ m{U_{10}} &= m{Z} \ m{U_{11}} &= m{Z} \cdot m{X} \end{aligned}$$

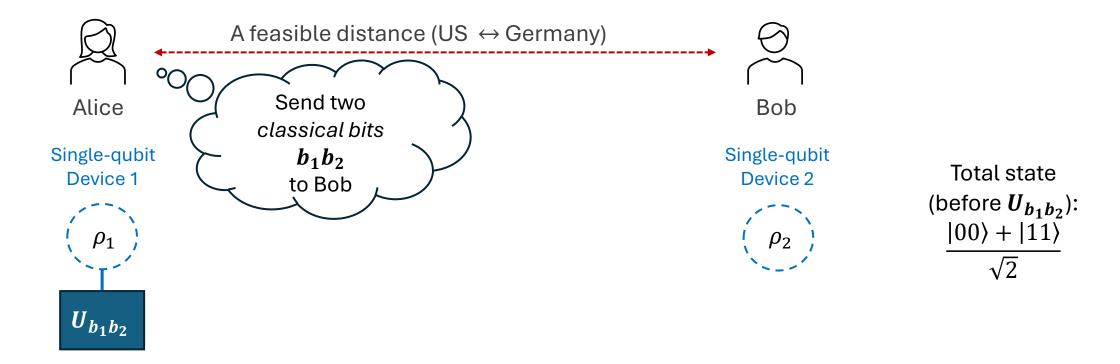
Small Exercise:

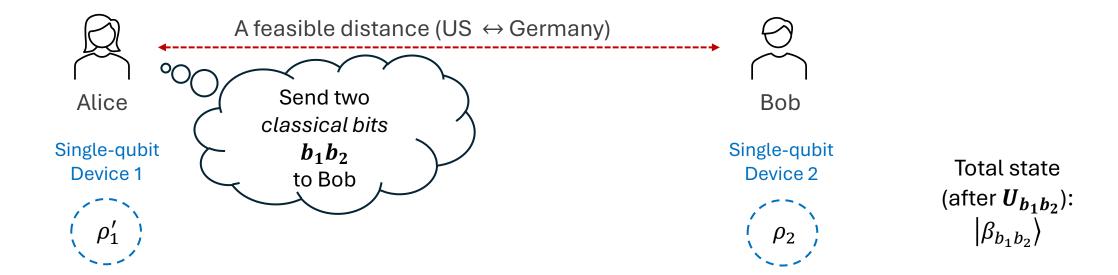


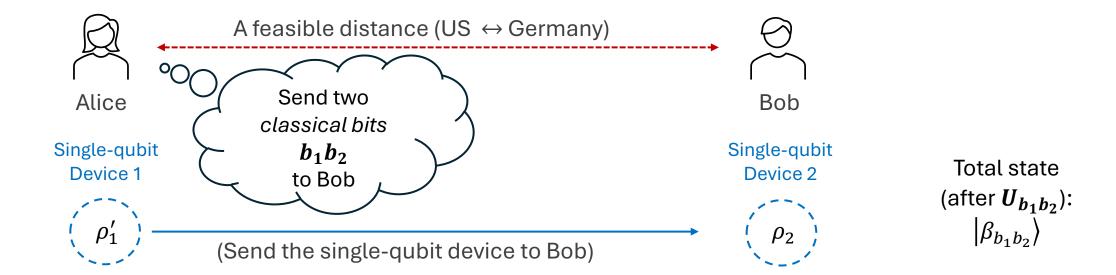
#### The Bell basis:

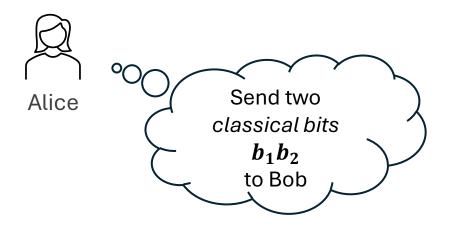
$$\begin{cases} M_{00} \colon |\beta_{00}\rangle\langle\beta_{00}|, \\ M_{01} \colon |\beta_{01}\rangle\langle\beta_{01}|, \\ M_{10} \colon |\beta_{10}\rangle\langle\beta_{10}|, \\ M_{11} \colon |\beta_{11}\rangle\langle\beta_{11}|, \end{cases}$$





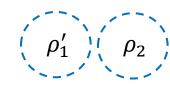




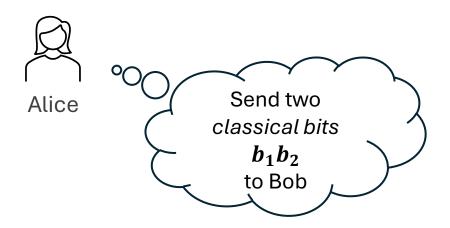


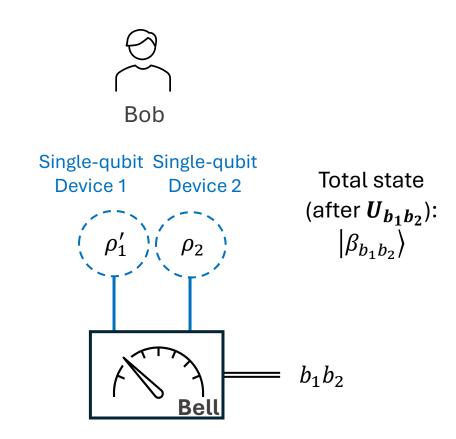


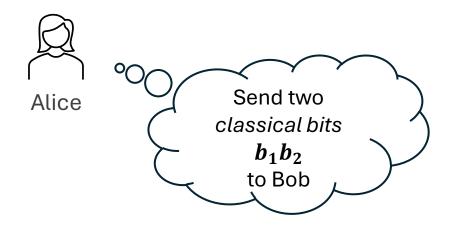
Single-qubit Single-qubit
Device 1 Device 2



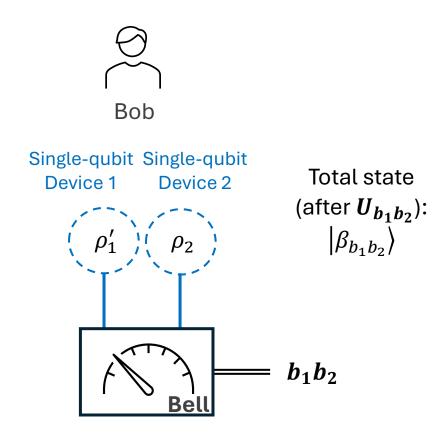
Total state (after  $m{U_{b_1b_2}}$ ):  $ig|m{eta_{b_1b_2}}ig>$ 



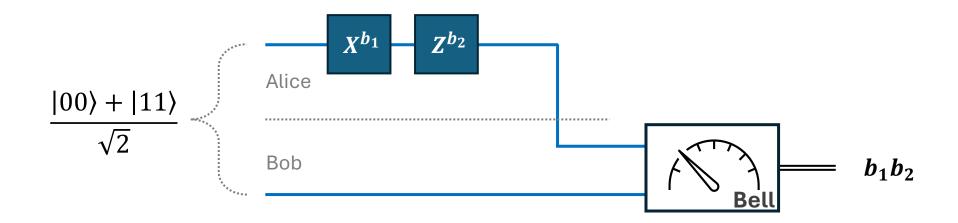




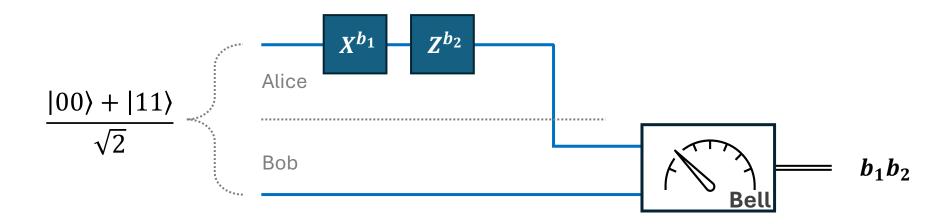
- One physical qubit can transmit two classical bits of information
  - (Require prior entanglement)
- Superdense coding: One qubit "encodes" two classical bits...



• A more compact description of the experiment:



• A more compact description of the experiment:

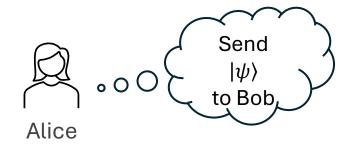


- Quick questions:
- (1) We wrote  $U_{b_1b_2}\coloneqq Z^{b_2}X^{b_1}$  before, but why does X come first here?
- (2) How can we transmit 2n classical bits using only n qubits?

- Superdense coding: Transmit classical bits via qubits
- Quantum teleportation: Transmit qubits via classical bits

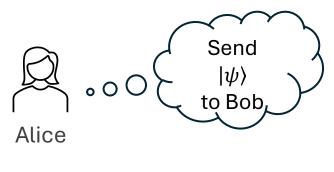


- A trivial approach (transportation): Physically deliver the device carrying  $|\psi\rangle$  to Bob
- But here we consider **Teleportation**: Transmit via (quantum) channel



$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$







Single-qubit Single-qubit Device 1 Device 2

Device 1

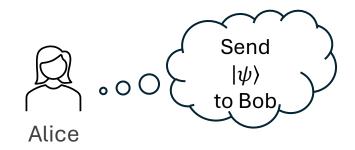


pre-sharing  $|\beta_{00}\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}}$ 

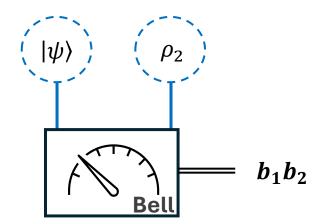
Single-qubit Device 3



Total state  $|\psi\rangle\otimes|\beta_{00}\rangle$ 



Single-qubit Single-qubit Device 1 Device 2

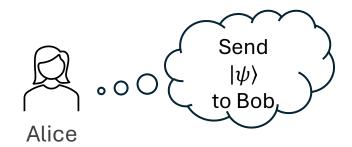




Single-qubit Device 3



Total state (before measurement)  $|\psi\rangle\otimes|eta_{00}
angle$ 



Single-qubit Single-qubit
Device 1 Device 2







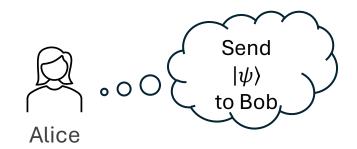
Single-qubit Device 3



Total state (after measurement)



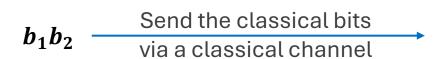
 $b_1b_2$  Send the classical bits via a classical channel



Single-qubit Single-qubit
Device 1 Device 2









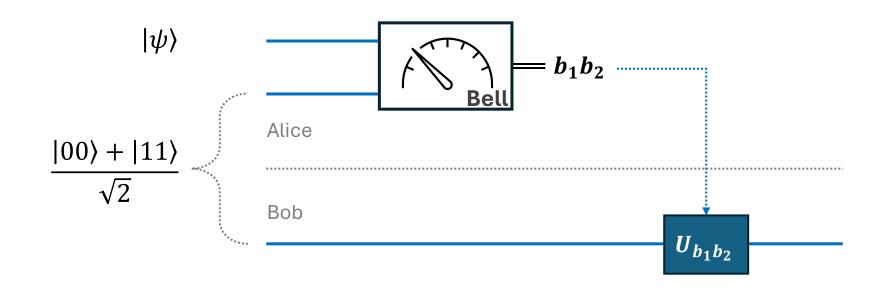
Single-qubit Device 3



Total state (after measurement)

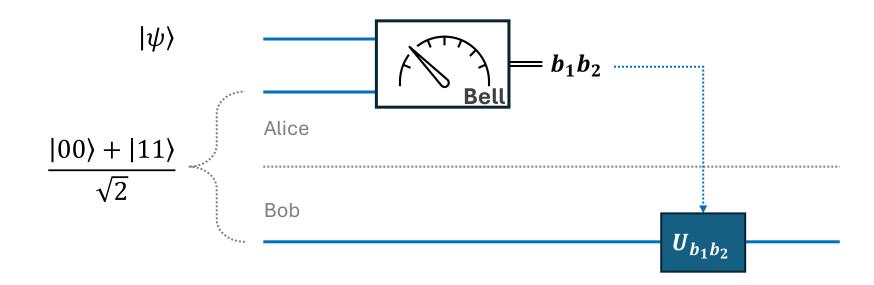


A more compact description (for analyzing states):



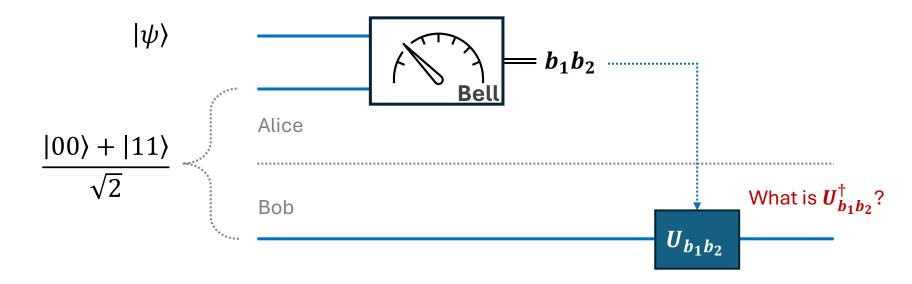
Total states:  $|\phi_0
angle$ 

A more compact description (for analyzing states):

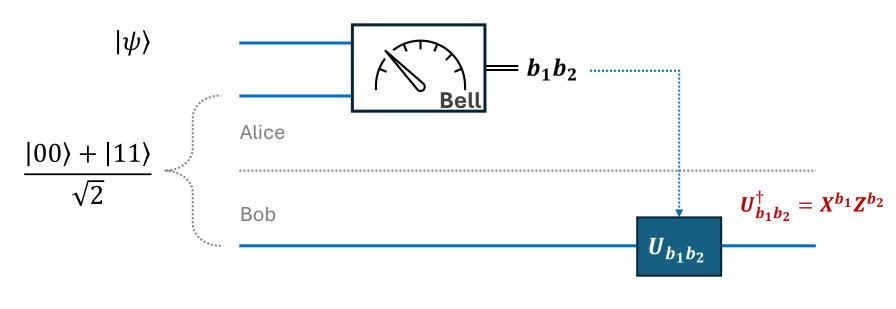


Total states:  $|\phi_0\rangle$ 

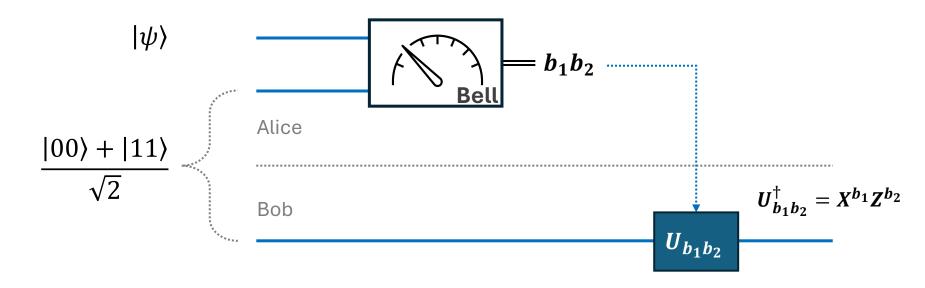
Can we re-write the state of the first two systems using the Bell basis?



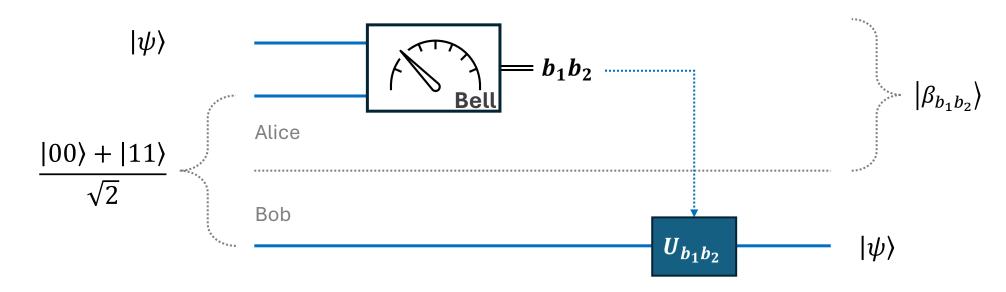
Total states: 
$$|\phi_0\rangle \\ = \sum_{b_1,b_2 \in \{0,1\}} \! \big|\beta_{b_1b_2}\big\rangle U_{b_1b_2}^\dagger |\psi\rangle$$



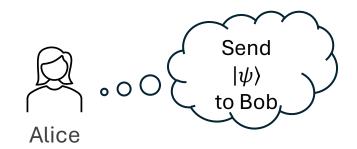
Total states: 
$$\begin{aligned} |\phi_0\rangle & |\phi_1\rangle \\ &= \sum_{b_1,b_2 \in \{0,1\}} |\beta_{b_1b_2}\rangle U_{b_1b_2}^\dagger |\psi\rangle \end{aligned}$$



Total states: 
$$\begin{aligned} |\phi_0\rangle & |\phi_1\rangle \\ &= \sum_{b_1,b_2 \in \{0,1\}} |\beta_{b_1b_2}\rangle U_{b_1b_2}^\dagger |\psi\rangle & = |\beta_{b_1b_2}\rangle U_{b_1b_2}^\dagger |\psi\rangle \end{aligned}$$



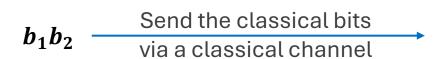
Total states: 
$$\begin{aligned} |\phi_0\rangle & |\phi_1\rangle & |\phi_2\rangle \\ &= \sum_{b_1,b_2 \in \{0,1\}} \left|\beta_{b_1b_2}\right\rangle U_{b_1b_2}^{\dagger} |\psi\rangle & = \left|\beta_{b_1b_2}\right\rangle U_{b_1b_2}^{\dagger} |\psi\rangle & = \left|\beta_{b_1b_2}\right\rangle \otimes |\psi\rangle \end{aligned}$$



Single-qubit Single-qubit
Device 1 Device 2







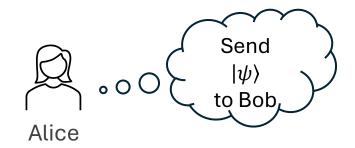


Single-qubit Device 3

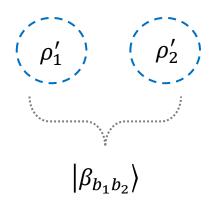


Total state (after measurement):

$$\left|oldsymbol{eta}_{b_1b_2}
ight>\otimes U_{b_1b_2}^{\dagger}|oldsymbol{\psi}
angle$$



Single-qubit Single-qubit
Device 1 Device 2





Single-qubit Device 3



Total state (Final):

$$\ket{oldsymbol{eta}_{b_1b_2}}\otimes\ket{\psi}$$

#### Reference

- **[NC00]:** Sections 1.3.6 and 1.3.7
- **[KLM07]:** Chapter 5

## **Next topic**

- Quantum Circuits
  - Controlled operations and Measurement