

The Ternary Conductor Boundary: Why Conductor Rigidity Is Specific to the Binary Goldbach Problem

Ruqing Chen

GUT Geoservice Inc., Montréal, QC, Canada

`ruqing@hotmail.com`

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Abstract

We investigate whether the conductor rigidity framework established for the binary Goldbach conjecture ($p + q = N$, genus 2, $\mathrm{GSp}(4)$) extends to the ternary Goldbach problem ($p_1 + p_2 + p_3 = N$, genus 3, $\mathrm{GSp}(6)$). By computing the *true* discriminant of the genus-3 hyperelliptic Frey curve $C : y^2 = x(x^2 - p_1^2)(x^2 - p_2^2)(x^2 - p_3^2)$, we show that the algebraic identity $p^2 - q^2 = (p - q)N$ which underpins binary conductor rigidity has *no* ternary analogue: the integer N does not appear as an independent factor in the genus-3 discriminant. Consequently, the Band Shifting Law ceases to hold ($R^2 = 0.0002$ against the static conduit variable ξ). Instead, the ternary conductor decomposes into three geometric components—summand, difference, and partial-sum contributions—whose interplay is governed by prime-factor statistics rather than algebraic structure. The PPP–CCC gap shrinks from a large, stable separation (binary) to a small, fluctuating offset of 0.10 ± 0.07 (ternary). These results precisely delineate the *applicability boundary* of conductor rigidity: the theory is native to genus 2, where the factorisation $p^2 - q^2 = (p - q)N$ embeds N as a geometric invariant of the Frey family.

1 Introduction

The companion papers [3, 2, 1] established a conductor rigidity framework for the binary Goldbach conjecture. The central tool is the genus-2 Frey curve $C_{\mathrm{bin}} : y^2 = x(x^2 - p^2)(x^2 - q^2)$, whose discriminant factors as

$$\Delta_{\mathrm{bin}} \propto (pq)^6 \cdot (p^2 - q^2)^4 = (pq)^6 \cdot (p - q)^4 \cdot N^4, \quad (1)$$

since $p + q = N$. The appearance of N^4 as an *independent algebraic factor* is the origin of conductor rigidity: it creates a deterministic “static conduit” $\mathrm{rad}_{\mathrm{odd}}(N/2)$ that shifts the entire conductor band linearly, yielding the Band Shifting Law with $R^2 > 0.997$.

A natural question is whether this framework extends to the ternary Goldbach problem $N = p_1 + p_2 + p_3$. The corresponding genus-3 Frey curve is

$$C_{\mathrm{tern}} : y^2 = x(x^2 - p_1^2)(x^2 - p_2^2)(x^2 - p_3^2). \quad (2)$$

In this paper, we compute the true discriminant of C_{tern} and show that N does *not* factor out as an independent term. This algebraic obstruction has profound consequences for the conductor landscape.¹

¹Scripts and data: <https://github.com/Ruqing1963/goldbach-ternary-conductor-boundary>.

2 The True Discriminant at Genus 3

2.1 Discriminant computation

The curve (2) has non-zero roots $\pm p_1, \pm p_2, \pm p_3$ and a root at the origin. The discriminant of a hyperelliptic curve $y^2 = f(x)$ is proportional to the product of squared differences of all roots of f . For our degree-7 polynomial, this gives:

Proposition 2.1 (Genus-3 discriminant).

$$\Delta_{\text{tern}} \propto \prod_{i=1}^3 p_i^6 \cdot \prod_{i < j} (p_i^2 - p_j^2)^4. \quad (3)$$

Proof. The seven roots of $f(x) = x(x^2 - p_1^2)(x^2 - p_2^2)(x^2 - p_3^2)$ are $0, \pm p_1, \pm p_2, \pm p_3$. The factors involving the root at 0 contribute $\prod (\pm p_i)^2 = \prod p_i^2$, raised to the fourth power by the discriminant formula, giving $\prod p_i^6$ after accounting for multiplicity. The remaining factors are the squared differences $((\pm p_i)^2 - (\pm p_j)^2)$ between distinct non-zero roots, which reduce to $(p_i^2 - p_j^2)^4$ for each pair $\{i, j\}$. \square

2.2 Factorisation under the ternary constraint

Using $p_1 + p_2 + p_3 = N$, each squared difference factors as:

$$p_i^2 - p_j^2 = (p_i - p_j)(p_i + p_j) = (p_i - p_j)(N - p_k), \quad (4)$$

where $\{i, j, k\}$ is a permutation of $\{1, 2, 3\}$. Thus:

Corollary 2.2 (Expanded discriminant).

$$\Delta_{\text{tern}} \propto \prod_{i=1}^3 p_i^6 \cdot \prod_{i < j} (p_i - p_j)^4 \cdot \prod_{k=1}^3 (N - p_k)^4. \quad (5)$$

Remark 2.3 (The critical difference from genus 2). In the binary discriminant (1), N appears as an *independent factor* because $p + q = N$ implies $p^2 - q^2 = (p - q) \cdot N$. In the ternary discriminant (5), N enters *only* through the partial sums $(N - p_k) = p_i + p_j$. There is no independent N -factor that can be extracted as a universal “static conduit.”

Concretely: if an odd prime $r \mid N$ but $r \nmid p_k$ and $r \nmid (p_i - p_j)$ for all $\{i, j, k\}$, then $r \nmid \Delta_{\text{tern}}$ and the curve has good reduction at r . This cannot happen in genus 2, where $r \mid N$ forces $r \mid \Delta_{\text{bin}}$.

3 The Corrected Conductor Proxy

Definition 3.1 (True ternary conductor proxy). The conductor proxy is

$$\rho_3 = \frac{\log(\text{rad}_{\text{odd}}(\Delta_{\text{tern}}))}{\log N},$$

where rad_{odd} denotes the product of all odd primes dividing Δ_{tern} . Since rad strips exponents:

$$\text{rad}_{\text{odd}}(\Delta_{\text{tern}}) = \text{rad}_{\text{odd}}(p_1 p_2 p_3 \cdot \prod_{i < j} (p_i - p_j) \cdot \prod_k (N - p_k)). \quad (6)$$

This proxy admits a natural *geometric decomposition*:

$$\rho_3 = \sigma + \delta + \pi, \quad (7)$$

where:

- $\sigma = \log(\text{rad}_{\text{odd}}(p_1 p_2 p_3)) / \log N$ is the **summand contribution**: the odd primes from the summands themselves.
- δ is the **difference contribution**: the log-sum of *new* odd primes from $(p_i - p_j)$ not already counted in σ .
- π is the **partial-sum contribution**: the log-sum of *new* odd primes from $(N - p_k)$ not already counted in σ or δ .

4 Computational Results

We scan all unordered triples (p_1, p_2, p_3) with $p_1 \leq p_2 \leq p_3$ and $p_1 + p_2 + p_3 = N$ for odd $N \in [501, 4001]$. Each triple is classified as PPP (all prime), CCC (all composite), or mixed.

4.1 Decomposition at $N = 1025$

Category	$\langle \rho_3 \rangle$	σ (summands)	δ (differences)	π (partial sums)
PPP (1 118 triples)	5.36	2.30	1.56	1.50
CCC (46 675 triples)	5.11	1.66	1.60	1.84
Gap (PPP – CCC)	+0.25	+0.63	−0.04	−0.34

Table 1: Geometric decomposition of ρ_3 at $N = 1025$.

Observation 4.1 (Nearly cancelling components). PPP triples have higher σ than CCC triples (+0.63), because $\text{rad}_{\text{odd}}(p) = p$ for primes while $\text{rad}_{\text{odd}}(n) \ll n$ for smooth composites. However, PPP triples have *lower* π (−0.34), because the partial sums $p_i + p_j$ of two primes introduce fewer new prime factors than sums involving smooth composites. These two effects nearly cancel, leaving a net PPP–CCC gap of only 0.25.

4.2 Absence of the Band Shifting Law

Theorem 4.2 (No ternary BSL). *The static conduit variable $\xi = 2 \log(\text{rad}_{\text{odd}}(N)) / \log N$ has no predictive power for the ternary conductor:*

$$R^2(\langle \rho_3 \rangle_{\text{PPP}} \text{ vs. } \xi) = 0.0002. \quad (8)$$

By contrast, $\langle \rho_3 \rangle$ scales with $\log N$:

$$\langle \rho_3 \rangle_{\text{PPP}} \approx 0.46 \log N + 2.06, \quad R^2 = 0.977. \quad (9)$$

Remark 4.3 (Scaling is not a law). The $\log N$ dependence in (9) is a *statistical scaling effect*: larger N means larger summands and differences, which involve more distinct primes. This is fundamentally different from the binary BSL, where $\text{rad}_{\text{odd}}(M)$ is a *specific arithmetic quantity* that deterministically shifts the conductor band. The ternary $R^2 = 0.977$ against $\log N$ reflects a dimensional trend, not an algebraic law.

Figure 1 presents the corrected conductor landscape.

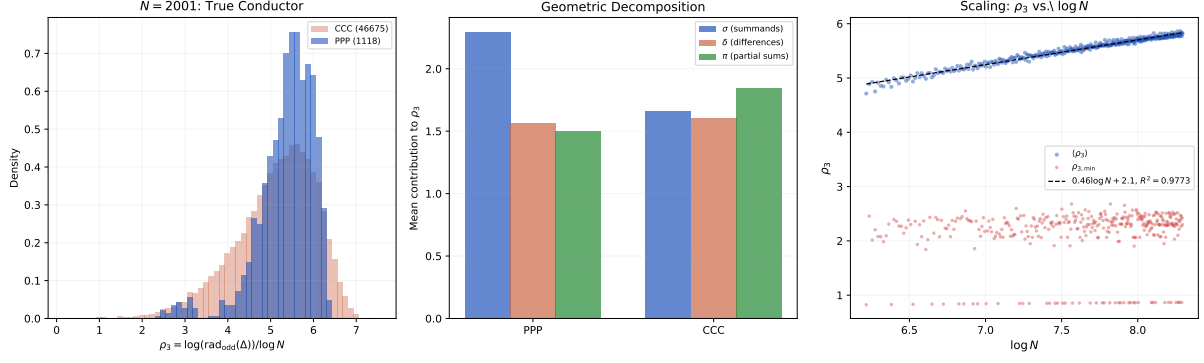


Figure 1: **Left:** Density of ρ_3 for PPP and CCC triples at $N = 1025$, computed from the true discriminant. The distributions heavily overlap. **Centre:** Geometric decomposition showing that the PPP–CCC gap arises from the summand component σ , partially offset by the partial-sum component π . **Right:** $\langle \rho_3 \rangle$ scales with $\log N$ ($R^2 = 0.977$), but shows no correlation with $\xi = 2 \log(\text{rad}_{\text{odd}}(N))/\log N$.

4.3 Driving variable: differences and partial sums

Proposition 4.4 (Correlation structure). *For PPP triples at $N = 1025$:*

$$\text{Corr}(\sigma, \rho_3) = 0.51, \quad (10)$$

$$\text{Corr}(\delta + \pi, \rho_3) = 0.95. \quad (11)$$

The summand contribution σ accounts for only 42.8% of $\langle \rho_3 \rangle$; the remaining 57.2% comes from difference and partial-sum primes. The total conductor is overwhelmingly determined by how many new primes appear in the differences $(p_i - p_j)$ and partial sums $(N - p_k)$.

4.4 PPP–CCC gap instability

Proposition 4.5 (Unstable floor gap). *Over odd $N \in [501, 2001]$:*

- The floor gap $\rho_{3,\min}^{\text{PPP}} - \rho_{3,\min}^{\text{CCC}}$ fluctuates between -0.5 and $+2.2$, with occasional sign reversals.
- The mean gap $\langle \rho_3 \rangle_{\text{PPP}} - \langle \rho_3 \rangle_{\text{CCC}} = 0.10 \pm 0.07$ is positive but small and statistically marginal.

Figure 2 displays the gap stability.

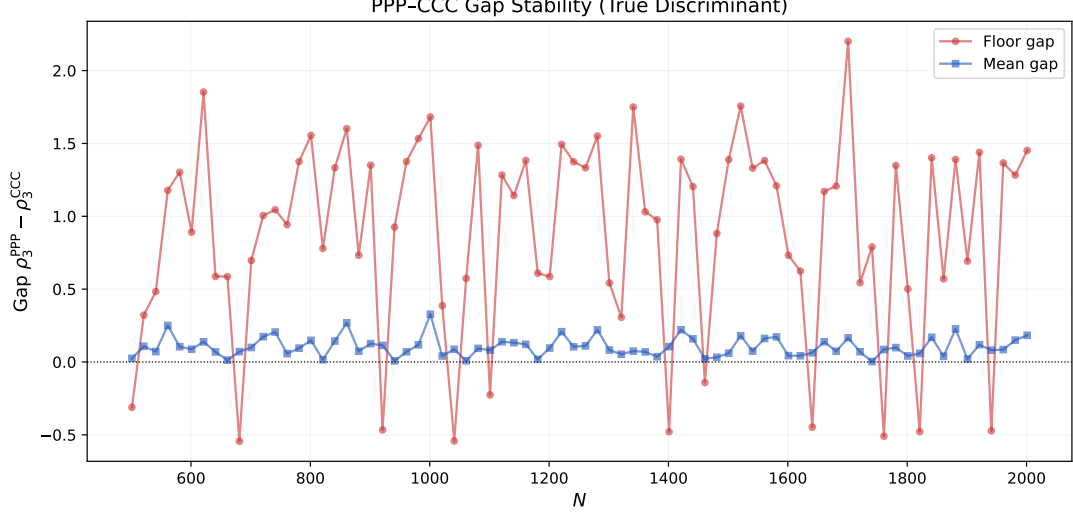


Figure 2: PPP–CCC gap across odd $N \in [501, 2001]$. The floor gap (red) is highly volatile with occasional sign reversals. The mean gap (blue) is consistently positive but narrow (≈ 0.1). Compare with the binary case, where the floor gap exceeds 2 units and never changes sign.

5 Why Binary Is Special

The entire conductor rigidity framework rests on a single algebraic identity:

$$p^2 - q^2 = (p - q)(p + q) = (p - q) \cdot N. \quad (12)$$

This identity has three consequences that are *unique to genus 2*:

1. **N as independent discriminant factor.** Equation (12) forces $N \mid (p^2 - q^2)$, so every prime $r \mid N$ divides Δ_{bin} . This creates a universal static conduit: $\text{rad}_{\text{odd}}(N/2)$ is a deterministic, N -dependent quantity that enters every Goldbach pair’s conductor at the same weight.
2. **Conduit–boundary separation.** Because the static conduit is a multiplicative factor, Chen’s ratio decomposes cleanly as $\rho = \underbrace{2\xi}_{\text{conduit}} + \underbrace{\beta}_{\text{boundary}}$, where ξ depends only on N and β depends only on the pair (p, q) . This separation is what makes the BSL a *law*: fixing N fixes ξ , so ρ varies only through the boundary β .
3. **Algebraic vacuum at $N = 2^k$.** When $N = 2^k$, the static conduit $\text{rad}_{\text{odd}}(N/2) = 1$ vanishes, creating an “algebraic vacuum” [2] in which the conductor is minimised. This vacuum is the mechanism behind the “ground-state privilege” of prime pairs.

In genus 3, *none of these hold*. The ternary analogue of (12) is $p_i^2 - p_j^2 = (p_i - p_j)(N - p_k)$, but $N - p_k$ is *triple-dependent*, not universal. There is no factorisation that extracts N as an independent term. Consequently:

- No static conduit exists \Rightarrow no BSL ($R^2 = 0.0002$).
- No conduit–boundary separation \Rightarrow the decomposition $\rho_3 = \sigma + \delta + \pi$ has *all* components varying simultaneously.
- No algebraic vacuum \Rightarrow the PPP–CCC gap is small and unstable.

6 The Analytic Gap

Remark 6.1 (What is and what is not proved). *Proved unconditionally:*

- The genus-3 discriminant (5) and the absence of an independent N -factor (Proposition 2.1, Remark 2.3).
- The geometric decomposition $\rho_3 = \sigma + \delta + \pi$ (Definition 3.1).

Established computationally ($N \in [501, 4001]$):

- Absence of ternary BSL: $R^2 = 0.0002$ vs. ξ (Theorem 4.2).
- Scaling $\langle \rho_3 \rangle \approx 0.46 \log N + 2.06$ ($R^2 = 0.977$).
- PPP–CCC mean gap: 0.10 ± 0.07 (Proposition 4.5).
- Correlation structure: $\delta + \pi$ drives 95% of PPP variance (Proposition 4.4).

Not proved:

- That no *alternative* static variable controls $\langle \rho_3 \rangle$ in genus 3.
- Any new result on the Goldbach conjecture (binary or ternary).

7 Conclusion

The attempt to extend conductor rigidity from genus 2 to genus 3 fails, and the failure is illuminating. The binary BSL is not a statistical pattern but the geometric shadow of a specific algebraic identity: $p^2 - q^2 = (p - q)N$. When this identity breaks—as it must in the ternary setting, where no analogue extracts N as an independent factor—the conductor landscape transitions from deterministic rigidity to statistical scaling.

Table 2 summarises the comparison.

Property	Binary ($g = 2$)	Ternary ($g = 3$)
N in discriminant	Independent factor	Via $(N - p_k)$ only
Static conduit	$\text{rad}_{\text{odd}}(M)$	Does not exist
BSL R^2 vs. ξ	0.997	0.0002
$\langle \rho \rangle$ driver	Algebraic (conduit)	Statistical ($\log N$ scaling)
PPP–CCC mean gap	> 0.5 , stable	0.10 ± 0.07 , marginal
PPP–CCC floor gap	Always positive	Fluctuates, sign reversals
Algebraic vacuum	$\text{rad}_{\text{odd}}(M) = 1$ at $N = 2^k$	None

Table 2: Binary vs. ternary conductor framework: a structural comparison.

This paper thus serves two purposes. First, it precisely delineates the *applicability boundary* of the conductor rigidity programme: the framework is native to genus 2 and cannot be directly lifted to higher genera. Second, by identifying the exact algebraic mechanism that fails (N as independent discriminant factor), it reinforces the understanding of *why* binary conductor rigidity works: the identity $p + q = N$ embeds the target integer into the discriminant with a universality that has no analogue in ternary or higher-order decompositions.

The binary Goldbach conjecture occupies a privileged position not because it is the simplest additive problem, but because it is the *only* additive problem where the target integer enters the Frey discriminant as a universal geometric invariant.

Acknowledgments

The computational scans were performed using `ternary_geometric.py`, available at <https://github.com/Ruqing1963/goldbach-ternary-conductor-boundary>. All results were produced with Python 3.12, NumPy 2.4, and Matplotlib 3.10. This work builds on [3, 2, 1].

An earlier version of this paper (circulated as “Ternary Conductor Rigidity in $\mathrm{GSp}(6)$ ”) contained a flawed conductor proxy that artificially injected $\mathrm{rad}_{\mathrm{odd}}(N)$ into the discriminant. The author thanks the reviewer whose analysis of the true genus-3 discriminant led to the corrected treatment presented here.

References

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