Exercise 1

PHYS4000 Advanced Computational Quantum Mechanics

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March 6, 2024

1 Question 1

Equation 2.17 is written as:

$$\frac{d^2u_l(r)}{dr^2} + 2(E + \frac{Z}{r})u_l(r) = l(l+1)\frac{u_l(r)}{r^2}$$
(1)

First we need to make ρ the subject of the function which is done through the substitution of $\rho = kr$.

$$k^{2} \frac{d^{2} u_{l}(\rho)}{d\rho^{2}} + 2\left(-\frac{k^{2}}{8} + \frac{Zk}{\rho}\right) u_{l}(\rho) = l(l+1) \frac{u_{l}(\rho)k^{2}}{\rho^{2}}$$
(2)

Where we have substituted for E by rearranging the definition of $k = 2\sqrt{-2E}$. We further simplify the differential Equation to the following by cancelling k^2 and multiplying the bracket:

$$\frac{d^2 u_l(\rho)}{d\rho^2} - (\frac{1}{4} - \frac{2Z}{\rho k}) u_l(\rho) = l(l+1) \frac{u_l(\rho)}{\rho^2}$$
(3)

In finding the solutions to the differential equation we rearrange it into the form...

$$\frac{d^2 u_l(\rho)}{d\rho^2} - (\frac{1}{4} - \frac{2Z}{\rho k}) u_l(\rho) - l(l+1) \frac{u_l(\rho)}{\rho^2} = 0$$
(4)

The question presents the substitution $u_l(\rho) = \rho^{l+1} \exp\{-\rho/2\}\omega(\rho)$, which forms...

$$\frac{d^2}{d\rho^2} [\rho^{l+1} \exp\{-\rho/2\}\omega(\rho)] - (\frac{1}{4} - \frac{2Z}{\rho k})(\rho^{l+1} \exp\{-\rho/2\}\omega(\rho)) - l(l+1)\rho^{l-1} \exp\{-\rho/2\}\omega(\rho) = 0$$
 (5)

$$(l+1)[(l\rho^{l-1}\exp\{-\rho/2\} - \frac{1}{2}\rho^{l+1}\exp\{-\rho/2\})\omega(\rho) + \rho^{l}\exp\{-\rho/2\}\frac{d\omega(\rho)}{d\rho}]$$

$$-\frac{1}{2}\{[(l+1)\rho^{l}\exp\{-\rho/2\} - \frac{1}{2}\rho^{l+1}\exp\{-\rho/2\}]\omega(\rho) + \frac{d\omega(\rho)}{d\rho^{2}}\rho^{l+1}\exp\{-\rho/2\}\}$$

$$+[(l+1)\rho^{l}\exp\{-\rho/2\} - \frac{1}{2}\rho^{l+1}\exp\{-\rho/2\}]\frac{d\omega(\rho)}{d\rho} + \frac{d^{2}\omega(\rho)}{d\rho^{2}}\rho^{l+1}\exp\{-\rho/2\}$$

$$-(\frac{1}{4} - \frac{2Z}{k\rho})\rho^{l+1}\exp\{-\rho/2\}\omega(\rho) - l(l+1)\rho^{l-1}\exp\{-\rho/2\}\omega(\rho) = 0$$

and simplifies to

$$\frac{d^2\omega(\rho)}{d\rho^2}\rho^{l+1} + \frac{d\omega(\rho)}{d\rho}[2(l+1)\rho^l - \rho^{l+1}] - \omega(\rho)[(l+1)\rho^l - \frac{2Z}{k}\rho^l] = 0$$
 (6)

We now eliminate ρ^l on all terms to get:

$$\frac{d^{2}\omega(\rho)}{d\rho^{2}}\rho + [2(l+1) - \rho]\frac{d\omega(\rho)}{d\rho} - [(l+1) - \frac{2Z}{k}]\omega(\rho) = 0$$
 (7)

This differential equation has the form:

$$z\frac{d^2w}{dz^2} + (b-z)\frac{dw}{dz} - aw = 0 (8)$$

Which has solutions in the confluent hypergeometric function;

$$F(a,b;\rho) = 1 + \frac{a}{b} \frac{\rho}{1!} + \frac{a}{b} \frac{(a+1)}{(b+1)} \frac{\rho^2}{2!} + \dots$$
 (9)

In this case, the function will truncate so long as a = 1 - n for n = 1, 2, 3, in this way, after the n + 1 term of the series and all subsequent terms must simplify to zero.

By comparing the equation we have with the general form, we can see that

$$a = l + 1 - \frac{2Z}{k} \tag{10}$$

$$b = 2(l+1) \tag{11}$$

And knowing that the differential equation has a general solution in the form of the hypergeometric functions and that the series truncates to be a finite sum, we know that this set of equations satisfies the conditions for a solution of $\omega(\rho)$ for bound states of the hydrogen atom.

2 Question 2

The equation:

$$F(a,b;\rho) = 1 + \frac{a}{b} \frac{\rho}{1!} + \frac{a}{b} \frac{(a+1)}{(b+1)} \frac{\rho^2}{2!} + \dots$$
 (12)

Terminates for a = 1 - n, for $n \in \{1, 2, 3, ...\}$. This prevents the series from growing exponentially in ρ .

In our case, we have $a=1+l-\frac{2Z}{k}$, which can be rewritten as $a=1-\left(\frac{2Z}{k}-l\right)$. This makes our $n=\frac{2Z}{k}-l$ and subsequently $n+l=\frac{2Z}{k}$. We then write $n=\frac{2Z}{k}$, \forall $n\in\{1+l,2+l,3+l,\dots\}$.

3 Question 3

2.28 states:

$$\int_0^\infty dr u_{n'l}(r) u_{nl}(r) = \delta_{n'n} \tag{13}$$

Which tell us that the eigenstates are orthogonal and normalised for n = n', where the kronecker delta equals unity. This will allow us to calculate the normalisation coefficients for our wave-functions, $u_{21}(r)$ and $u_{31}(r)$.

For $\mathbf{u_{21}}(\mathbf{r})$;

First we establish what we are substituting for ρ , with Z=1, n=2, l=1:

$$\rho = kr = \frac{2Zr}{n}$$

$$\rho = kr = \frac{2(1)r}{2} = r$$

$$\therefore k = 1$$
(14)

With l = 1 = Z = k = 1 we have:

$$a = l + 1 - \frac{2Z}{k}$$

$$a = 1 + 1 - 2 = 0$$

$$b = 2(l + 1)$$

$$b = 2(2) = 4$$

$$a = 0, b = 4$$
 (15)

Calculating the $\omega(kr)$

$$F(a,b;r) = 1 + \frac{a}{b} \frac{r}{1!} + \frac{a}{b} \frac{(a+1)}{(b+1)} \frac{r^2}{2!} + \dots$$

$$F(0,4;r) = 1 + \frac{0}{4} \frac{r}{1!} + \frac{0}{4} \frac{(0+1)}{(4+1)} \frac{r^2}{2!} + \dots$$

$$F(0,4;r) = \omega(r) = 1 \tag{16}$$

Substituting r for ρ to get back $u_{nl}(r)$ from $u_{nl}(\rho)$:

$$u_{nl}(r) = (kr)^{l+1} \exp\{-kr/2\}\omega(kr)$$

$$u_{21}(r) = r^{1+1} \exp\{-r/2\}(1)$$

$$u_{21}(r) = r^2 \exp\{-r/2\}$$
(17)

We need to add a constant of normalisation to this function in order for it to belong to the Hilbert space, and we will use (13) in order to normalise.

$$\int_0^\infty (Ar^2 \exp\{-r/2\})^2 dr = \delta_{n'n}$$

$$\int_{0}^{\infty} A^{2} r^{4} \exp\{-r\} dr = \delta_{2,2} = 1$$

Tabular integration will solve the integral

$$\int_{0}^{\infty} A^{2}r^{4} \exp\{-r\} dr$$

$$= A^{2}[-r^{4} \exp\{-r\} - 4r^{3} \exp\{-r\} - 12r^{2} \exp\{-r\} - 24r \exp\{-r\} - 24 \exp\{-r\}]_{0}^{\infty}$$

$$= A^{2}[0 - (-24)]$$

$$= 24A^{2} = 1$$

$$A = \frac{1}{2\sqrt{6}}$$
(18)

Which normalises the function to

$$u_{21}(r) = \frac{r^2}{2\sqrt{6}} \exp\{-r/2\} \tag{19}$$

as given in the problem statement.

For $\mathbf{u_{31}}(\mathbf{r})$:

We start in the same way, by defining the appropriate quantities starting with Z = 1, n = 3, l = 1. Subsequently computing parameters like k:

$$\rho = kr = \frac{2Zr}{n}$$

$$\rho = kr = \frac{2(1)r}{3} = r$$

$$\therefore k = \frac{2}{3}$$
(20)

a and b...

$$a = l + 1 - \frac{2Z}{k}$$

$$a = 1 + 1 - \frac{2(1)}{\frac{2}{3}}$$

$$a = 1 + 1 - 3 = -1$$

$$b = 2(l + 1)$$

$$b = 2(2) = 4$$

$$a = -1, b = 4$$
(21)

Taking these values we compute the $\omega(\rho)$ value needed for the function.

$$F(a,b;\rho) = 1 + \frac{a}{b} \frac{\rho}{1!} + \frac{a}{b} \frac{(a+1)}{(b+1)} \frac{\rho^2}{2!} + \dots$$

$$F(-1,4;\rho) = 1 + \frac{-1}{4} \frac{\rho}{1!} + \frac{-1}{4} \frac{(-1+1)}{(4+1)} \frac{\rho^2}{2!} + \dots$$

$$F(-1,4;\rho) = 1 - \frac{\rho}{4}$$
(22)

Replacing ρ with $\frac{2}{3}r$

$$\omega(\frac{2}{3}r) = F(-1,4; \frac{2}{3}r) = 1 - \frac{1}{4}\frac{2}{3}r \tag{23}$$

$$\omega(\frac{2}{3}r) = 1 - \frac{r}{6} \tag{24}$$

We now make the appropriate substitutions as follows:

$$u_{nl}(r) = (kr)^{l+1} \exp\{-kr/2\}\omega(kr)$$

$$u_{31}(r) = (\frac{2}{3}r)^{1+1} \exp\left\{-\frac{r}{2}\frac{2}{3}\right\}(1 - \frac{r}{6})$$

$$u_{31}(r) = (\frac{4}{9})r^2 \exp\left\{-\frac{r}{3}\right\}(1 - \frac{r}{6})$$

And add the necessary normalisation coefficient A:

$$u_{31}(r) = A(\frac{4}{9})r^2 \exp\left\{-\frac{r}{3}\right\}(1 - \frac{r}{6})$$
 (25)

As last time we now use the normalisation condition of Equation 13 to determine the normalisation coefficient

$$\int_{0}^{\infty} dr u_{n'l}(r) u_{nl}(r) = \delta_{n'n}$$

$$\int_{0}^{\infty} (A(\frac{4}{9})r^{2} \exp\left\{-\frac{r}{3}\right\} (1 - \frac{r}{6}))^{2} dr = \delta_{3'3}$$

$$\int_{0}^{\infty} A^{2}(\frac{16}{81}) \exp\left\{-\frac{2r}{3}\right\} (r^{4} - \frac{r^{5}}{3} + \frac{r^{6}}{36}) dr = \delta_{3,3} = 1$$

$$A^{2}(\frac{16}{81}) \int_{0}^{\infty} \exp\left\{-\frac{2r}{3}\right\} (r^{4} - \frac{r^{5}}{3} + \frac{r^{6}}{36}) dr = 1$$
(26)

As before, the form of integral we have is an exponent multiplied by a polynomial, so we know that tabular integration will solve the integral:

$$F(x) \qquad G(x)$$

$$r^{4} - \frac{r^{5}}{3} + \frac{r^{6}}{36} \qquad \exp\left\{-\frac{2r}{3}\right\}$$

$$-(4r^{3} - \frac{5}{3}r^{4} + \frac{6}{36}r^{5}) \qquad -\frac{3}{2}\exp\left\{-\frac{2r}{3}\right\}$$

$$12r^{2} - \frac{20}{3}r^{3} + \frac{5}{6}r^{4} \qquad \frac{9}{4}\exp\left\{-\frac{2r}{3}\right\}$$

$$-(24r - 20r^{2} + \frac{20}{6}r^{3}) \qquad -\frac{27}{8}\exp\left\{-\frac{2r}{3}\right\}$$

$$24 - 40r + 10r^{2} \qquad \frac{81}{16}\exp\left\{-\frac{2r}{3}\right\}$$

$$-(-40 + 20r) \qquad -\frac{243}{32}\exp\left\{-\frac{2r}{3}\right\}$$

$$20 \qquad \frac{729}{64}\exp\left\{-\frac{2r}{3}\right\}$$

$$0 \qquad -\frac{2187}{128}\exp\left\{-\frac{2r}{3}\right\}$$

Resulting in the final integral

$$A^{2}(\frac{16}{81}) \int_{0}^{\infty} \exp\left\{-\frac{2r}{3}\right\} (r^{4} - \frac{r^{5}}{3} + \frac{r^{6}}{36}) dr$$

$$= A^{2}(\frac{16}{81}) \left[-\frac{3}{2} \exp\left\{-\frac{2r}{3}\right\} (r^{4} - \frac{r^{5}}{3} + \frac{r^{6}}{36}) - \frac{9}{4} \exp\left\{-\frac{2r}{3}\right\} (4r^{3} - \frac{5}{3}r^{4} + \frac{6}{36}r^{5}) - \frac{27}{8} \exp\left\{-\frac{2r}{3}\right\} (12r^{2} - \frac{20}{3}r^{3} + \frac{5}{6}r^{4}) - \frac{81}{16} \exp\left\{-\frac{2r}{3}\right\} (24r - 20r^{2} + \frac{20}{6}r^{3})$$

$$-\frac{243}{32} \exp\left\{-\frac{2r}{3}\right\} (24 - 40r + 10r^{2}) - \frac{729}{64} \exp\left\{-\frac{2r}{3}\right\} (-40 + 20r) - 20\frac{2187}{128} \exp\left\{-\frac{2r}{3}\right\} \right]_{0}^{\infty} = 1$$

Which in the limits of the integration will evaluate to:

$$=A^{2}(\frac{16}{81})(\frac{2187}{32})=1$$
$$=\frac{27}{2}A^{2}=1$$

$$A = \frac{2}{3\sqrt{6}}\tag{27}$$

When finally substituted into the function for u:

$$u_{31}(r) = \frac{2}{3\sqrt{6}} \left(\frac{4}{9}\right) r^2 \exp\left\{-\frac{r}{3}\right\} \left(1 - \frac{r}{6}\right)$$
 (28)

$$u_{31}(r) = \frac{8r^2}{27\sqrt{6}} \exp\left\{-\frac{r}{3}\right\} (1 - \frac{r}{6}) \tag{29}$$