

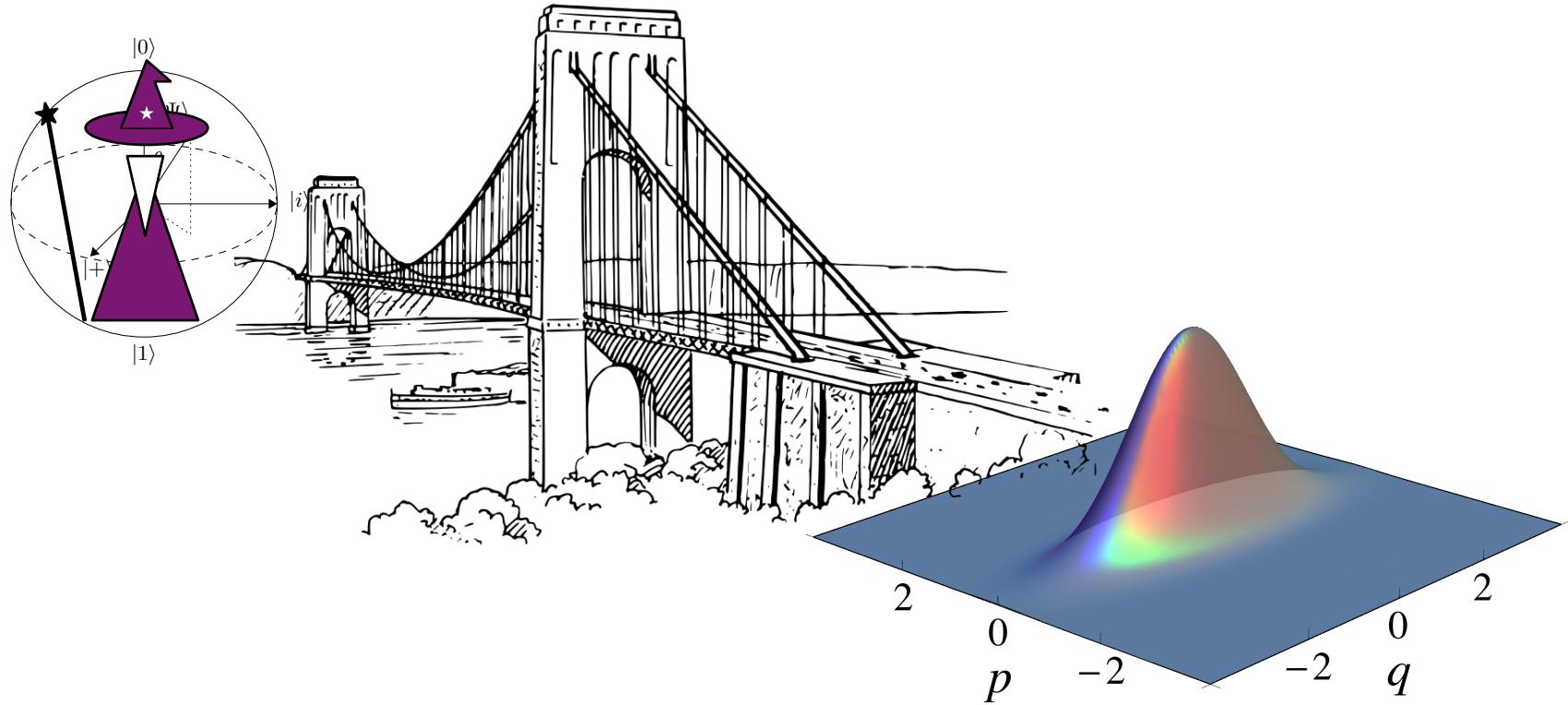


QUANTIFYING QUBIT MAGIC RESOURCE WITH GOTTESMAN-KITAEV-PRESKILL ENCODING

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QUANTUM COMPUTING WITH QUBITS

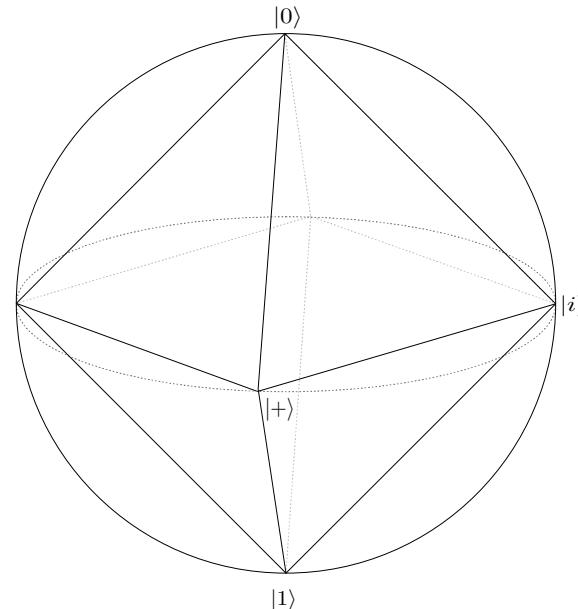
- We consider an array of two-level systems (qubits) $|\psi\rangle = \sum_{\mathbf{i} \in \mathcal{F}_2^n} c_{\mathbf{i}} |\mathbf{i}\rangle$
- Operations are implemented using quantum gates
- A Universal gate set is $\underbrace{\{ CX, H, S \}}_{\text{Clifford}}, T \}$

$$C_2 \equiv \{ \hat{U} | \hat{U} \hat{P} \hat{U}^\dagger \subseteq C_1, \forall \hat{P} \in C_1 \}$$



STABILIZER AND MAGIC STATES

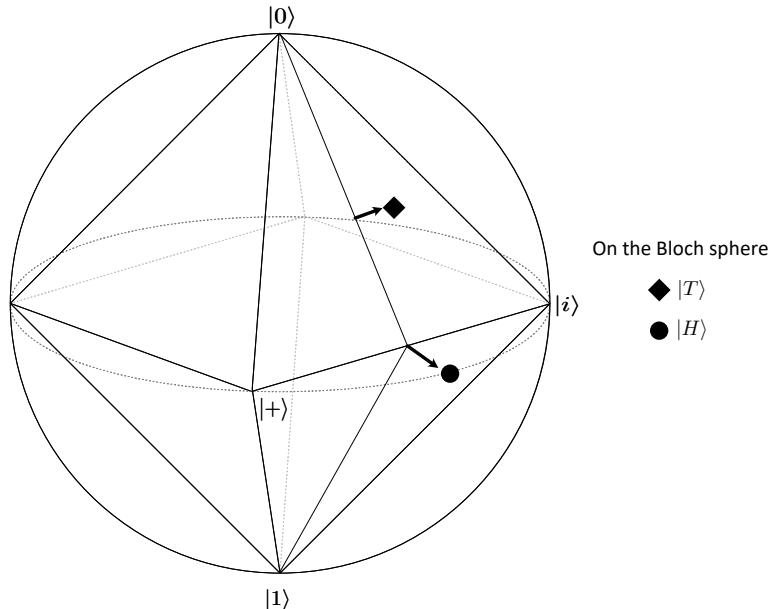
- Stabilizer states are closed under Clifford
- Magic states enable non-Clifford operations through teleportation





STABILIZER AND MAGIC STATES

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- Magic states enable non-Clifford operations through teleportation



$$|H\rangle = (|0\rangle + e^{i\pi/4} |1\rangle)/\sqrt{2}$$

$$|T\rangle = \cos(\beta) |0\rangle + \sin(\beta) e^{i\frac{\pi}{4}} |1\rangle$$

$$\cos(2\beta) = \frac{1}{\sqrt{3}}$$





MAGIC AND STABILISER STATES

- **Gottesman-Knill theorem:**

A quantum computer based only on:

1. Qubits initialized in a Pauli eigenstate
2. Clifford group operations
3. Pauli measurements

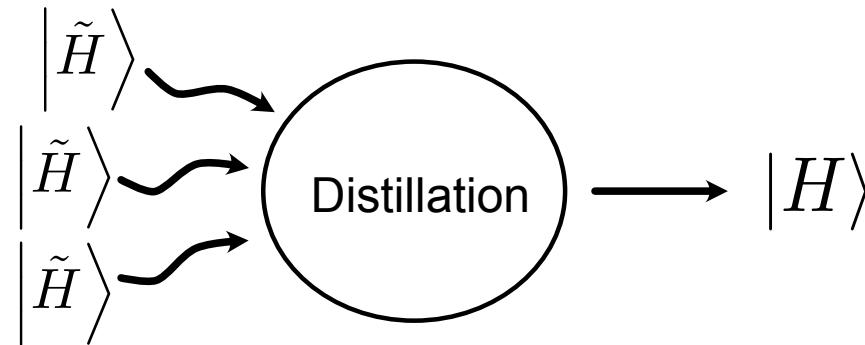
Can be simulated efficiently with a classical computer



D. Gottesman, arXiv:quant-ph/9807006 (1998)

RESOURCES IN FAULT-TOLERANCE

- 2D topological codes: Only Clifford operations are natively fault-tolerant
- Magic-state distillation is a way for non-Clifford gates



- Requires most of quibts/gates used

Resource: Magic/ non-stabilizerness



S. Bravyi, et al, PRA 71 022316 (2005)



MAGIC MEASURES

- Relative entropy of magic, V. Veitch *et al*, *New Journal of Physics* 16, 013009 (2014)
- Robustness of magic, Howard *et al*, *Phys. Rev. Lett.* 118, 090501 (2017)
- Stabilizer rank, Bravyi *et al*, *Quantum* 3, 181 (2019)
- Stabilizer nullity and the dyadic monotone, Beverland *et al*, *Quantum Science and Technology* 5, 035009 (2020)
- Dyadic negativity, mixed state extend, generalised robustness, Seddon *et al*, *PRX Quantum* 2, 010345 (2021)

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- Stabilizer Rény Entropy, Leone *et al*, *PRL* 128, 050402 (2022)
 - Bell magic, Haug *et al*, *arXiv:2204.10061* (2022)

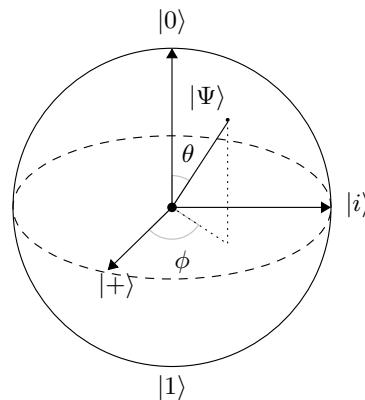




DISCRETE VS CONTINUOUS

Discrete variables:

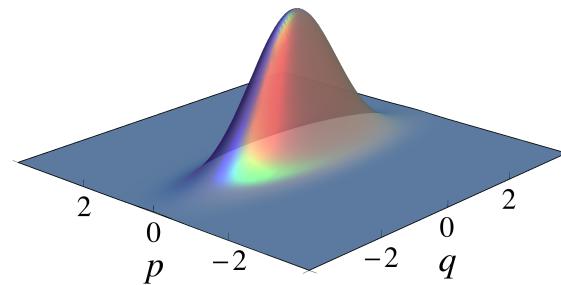
- Discrete basis
- Information encoded in Qubits
- Finite dimensional Hilbert space



Bloch Sphere

Continuous variables:

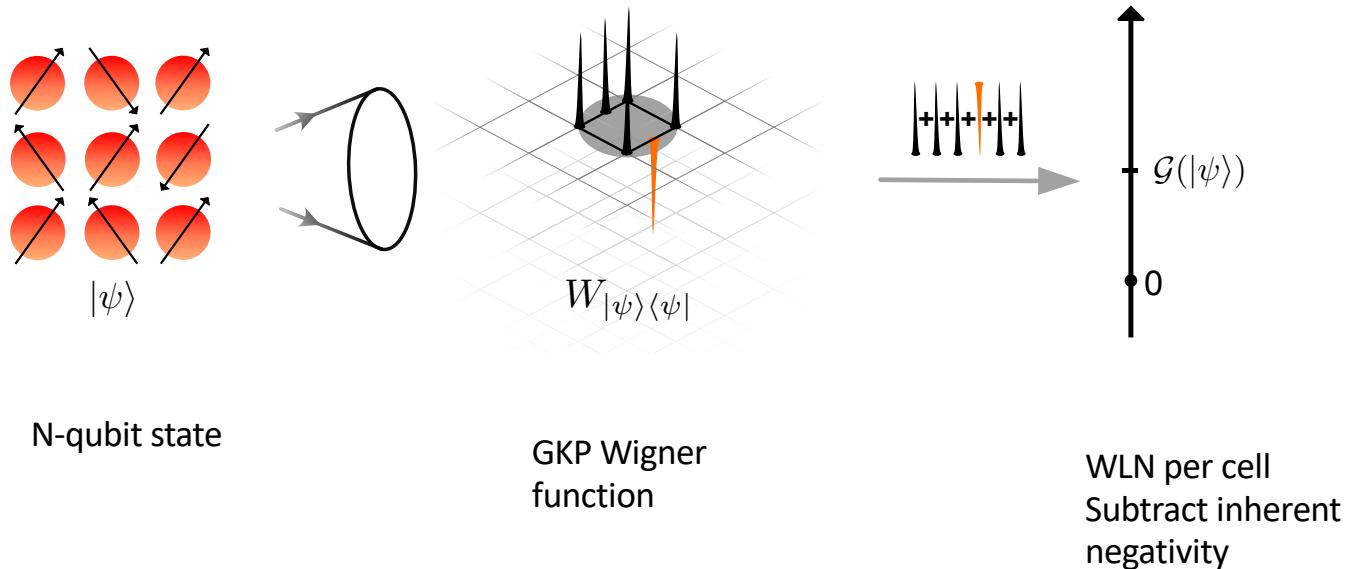
- Information encoded in q. modes
- Relevant observables have continuous spectra
- Infinite dimensional Hilbert space



Phase space representation



NEW MEASURE

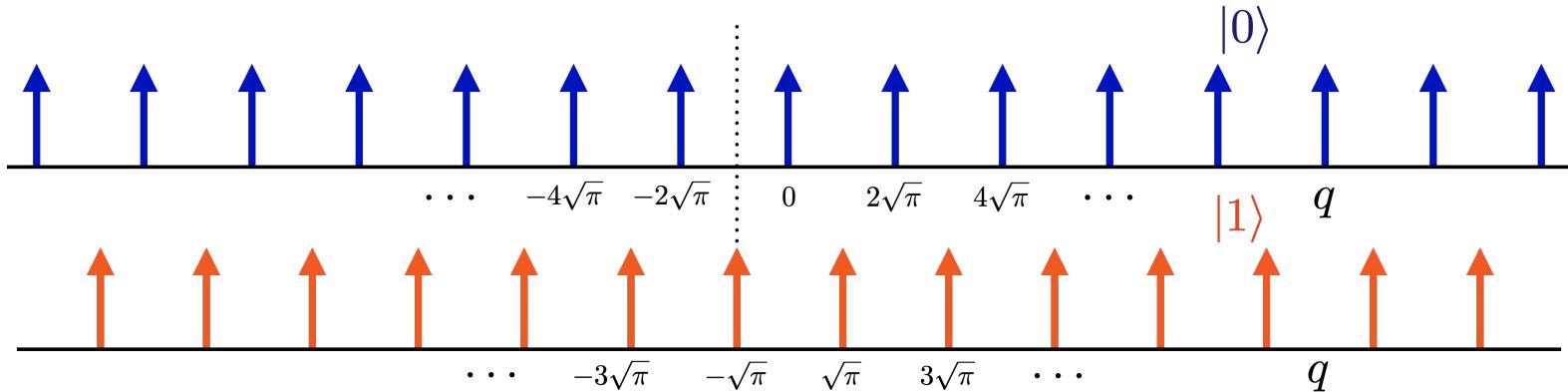




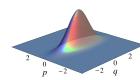
GOTTESMAN-KITAEV-PRESKILL CODE

- Error correction code for continuous variable systems to encode qubits

$$|u_i\rangle = \sum_{s_i=-\infty}^{\infty} |x_i = \sqrt{\pi}(u_i + 2s_i)\rangle_{\hat{q}}$$



D. Gottesman, et al, Phys. Rev. A 64, 012310 (2001)



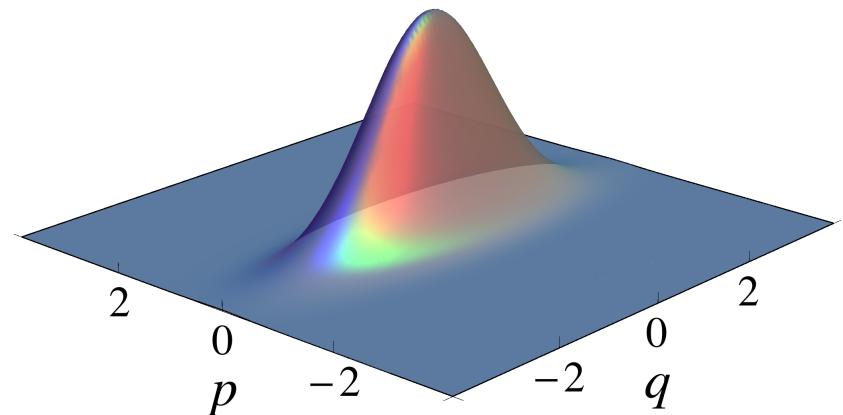


WIGNER FUNCTION

- Wigner Function represents a state in a phase-space

$$W_{\hat{\rho}}(\mathbf{q}, \mathbf{p}) \equiv \frac{1}{(2\pi)^n} \int d^n x e^{i\mathbf{p}\mathbf{x}} \left\langle \mathbf{q} + \frac{\mathbf{x}}{2} \right| \hat{\rho} \left| \mathbf{q} - \frac{\mathbf{x}}{2} \right\rangle_{\hat{q}}$$

- Is Quasi-probability distribution
- Resource theory related to Wigner negativity





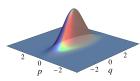
WIGNER LOGARITHMIC NEGATIVITY

- Resource theory of Wigner logarithmic negativity (WLN)
 - Gaussian states are the free states
 - Gaussian operations are the free operations

$$\mathcal{W}(\hat{\sigma}) = \log_2 \left(\int d^n q d^n p |W_{\hat{\sigma}}(q, p)| \right)$$

- Feature of GKP: All Clifford operations are Gaussian
- Transfer properties using GKP to define our measure

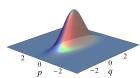
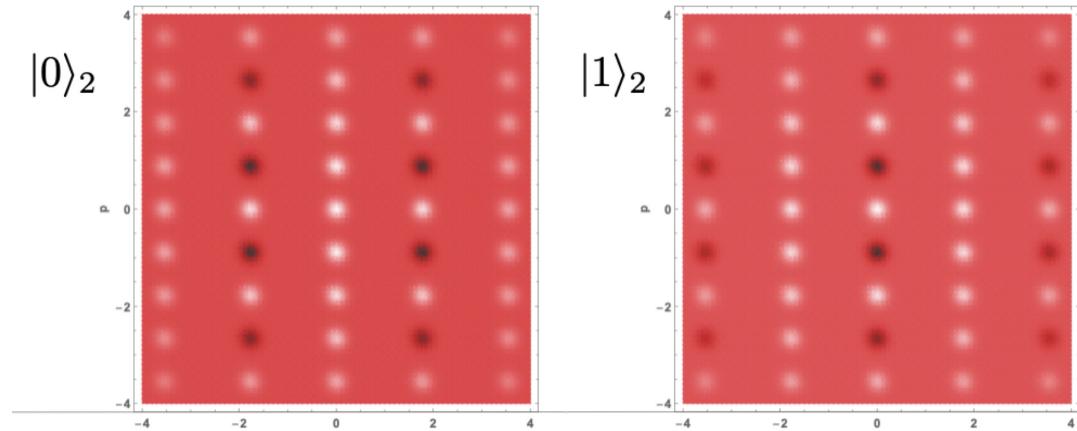
F. Albarelli, et al, Phys. Rev. A 98.5, 052350 (2018)
R. Takagi, et al, Phys. Rev. A 97, 062337 (2018)





WIGNER FUNCTION: GKP STATE

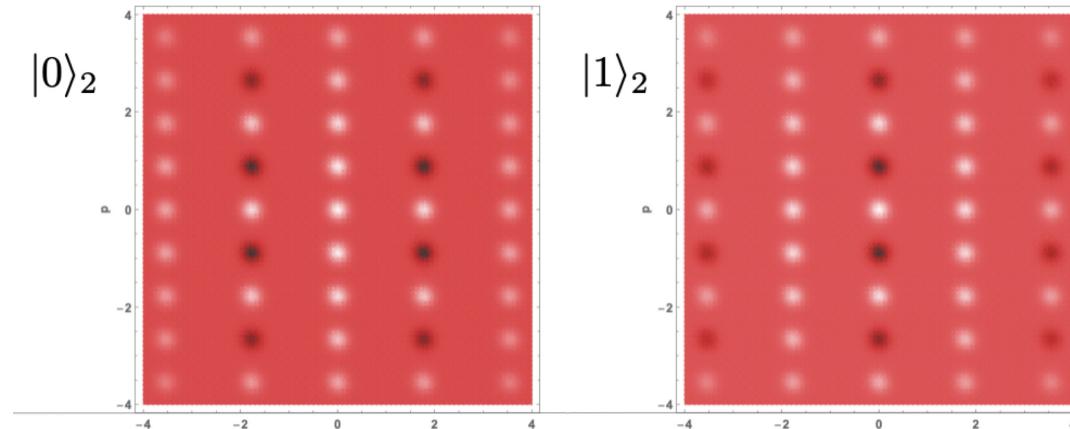
- **Stabilizer states:** $\frac{1}{4}$ of peaks are negative



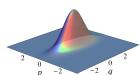


WIGNER LOG. NEGATIVITY: GKP STATE

- WLN is infinite negativity even for stabilizer states



- **Lattice:** We restrict to one unit cell
- WLN is non-zero for stabilizer states
 - But constant!

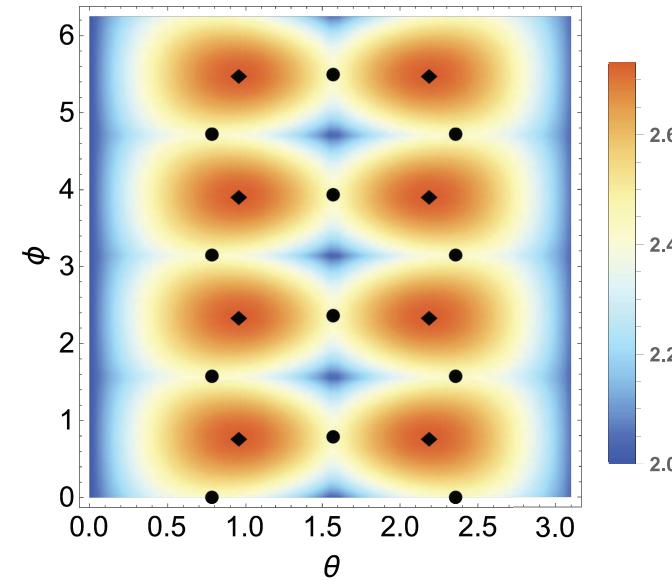
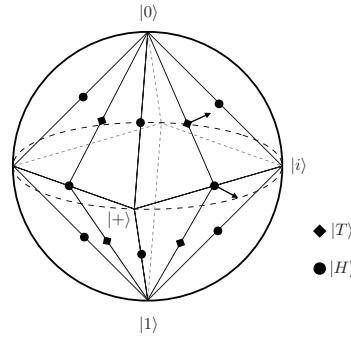




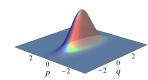
WIGNER LOG. NEGATIVITY: GKP STATE

- WLN has constant value for stabiliser states
- WLN has maximal values for T/H-type states

$$|\Psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle$$

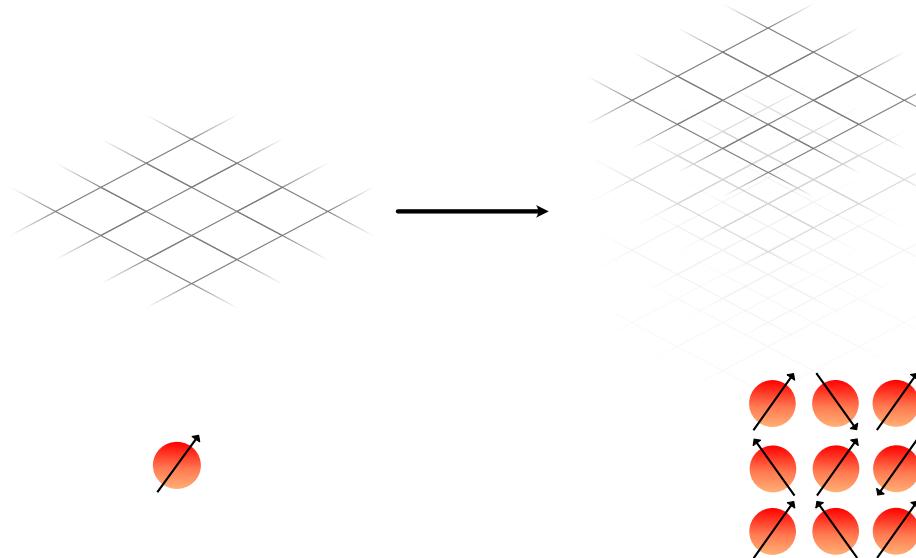


L. García-Álvarez, et . al, *Int. Symposium on Mathematics, Quantum Theory, and Cryptography*, 2021



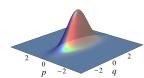


EXTENSION TO N-QUBITS



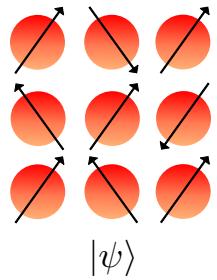
$$|\Psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle$$

$$|\psi\rangle = \sum_{\mathbf{i} \in \mathcal{F}_2^n} c_{\mathbf{i}} |\mathbf{i}\rangle$$



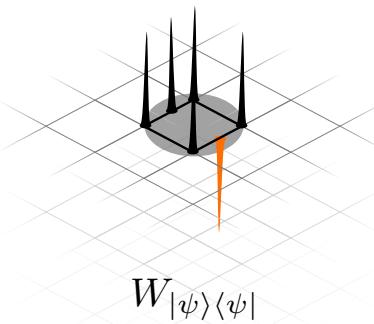
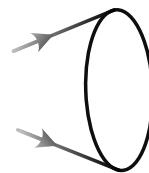


NEW MEASURE

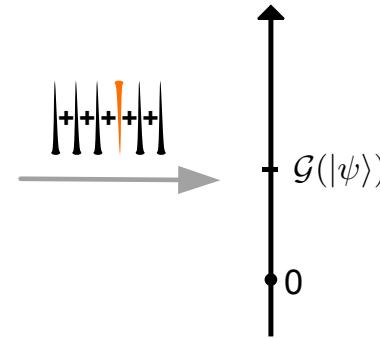


N-qubit state

$$|\psi\rangle = \sum_{i \in \mathbb{F}_2^n} c_i |i\rangle$$



GKP Wigner function



WLN per cell

$$\mathcal{W}_C(\hat{\rho}) = \log_2 \left(\int_0^{2\sqrt{\pi-\epsilon}} d^n q d^n p |W_{\hat{\rho}}(q, p)| \right)$$

Subtract inherent negativity



THE MEASURE

$$\mathcal{G}(|\psi\rangle) = \log_2 \left(\sum_{\mathbf{i}, \mathbf{j} \in \mathbb{F}_2^n} \left| \sum_{\mathbf{k} \in \mathbb{F}_2^n} \frac{(-1)^{\mathbf{i} \cdot \mathbf{k}}}{2^n} c_{\mathbf{k}} c_{\mathbf{k} + \mathbf{j}} \right| \right)$$

- No obvious connection to continuous variables left
- We recover the **st-norm!**
 - And prove that it is a magic measure for pure states
 - And the stabilizer Rényi entropy for $\alpha = \frac{1}{2}$
- Easier to compute than other measures (up to 12 qubits on a laptop)





THE MEASURE

$$\mathcal{G}(|\psi\rangle) = \log_2 \left[\frac{1}{2^n} \sum_{\hat{P} \in \mathcal{P}_n^+} \left| \text{Tr} \left[\hat{P} |\psi\rangle \langle \psi| \right] \right| \right]$$

- No obvious connection to continuous variables left
- We recover the **st-norm!**
 - And prove that it is a magic measure for pure states
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PROPERTIES

- Faithfulness $\mathcal{G}(|\psi_S\rangle) = 0$
- Invariance under Clifford unitaries $\mathcal{G}(\hat{U}_C |\psi\rangle) = \mathcal{G}(|\psi\rangle)$
- Invariance under composition $\mathcal{G}(|\psi\rangle \otimes |\phi_S\rangle) = \mathcal{G}(|\psi\rangle)$
- Additivity $\mathcal{G}(|\psi\rangle_A \otimes |\phi\rangle_B) = \mathcal{G}(|\psi\rangle) + \mathcal{G}(|\phi\rangle)$





BOUNDS ON OPERATOR COST

Upper Bound

Given by existing algorithm

Lower Bound

No better algorithm is possible

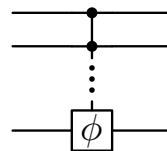
- Allows for optimization
- Use teleportation circuits and magic measures to find lower bounds
- Lower bound gate synthesis
 - e.g. how many T-gate needed at least to implement
- Form of our measures allows for analytical expressions





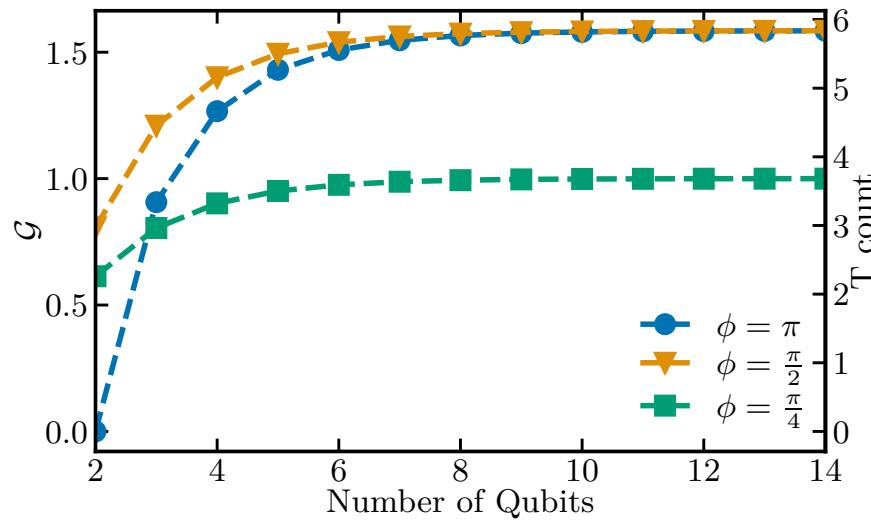
MULTIPLY CONTROLLED PHASE GATE

- Analytical solution for



$$M_\phi = \text{diag}(1, \dots, 1, e^{i\phi})$$

The multiply controlled Phase-Gate



$$\phi = \pi : C^n Z$$

$$\phi = \frac{\pi}{2} : C^n S$$

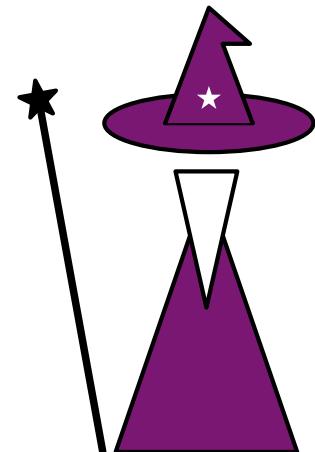
$$\phi = \frac{\pi}{4} : C^n T$$





CONCLUSION

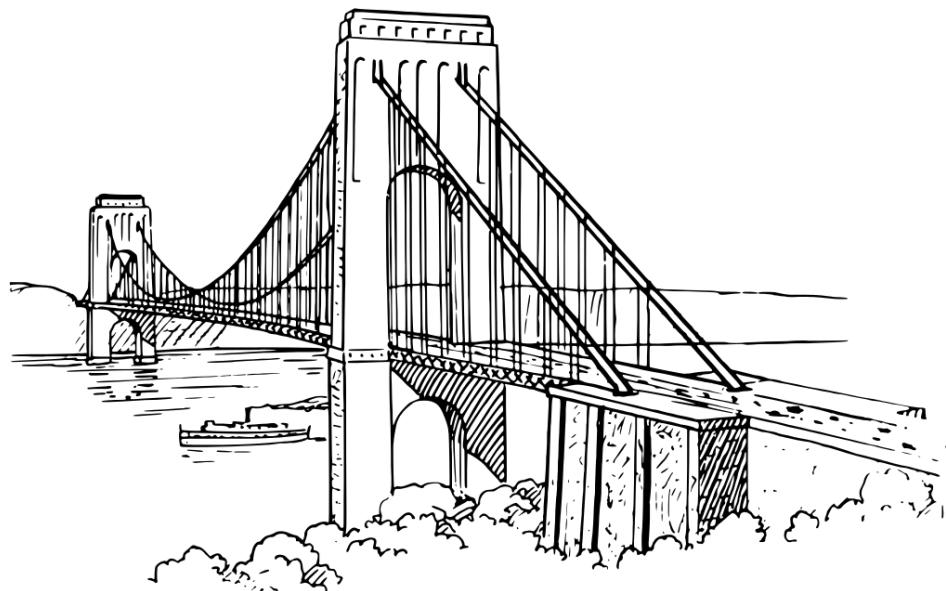
- Defined a magic measure
- Based on CV techniques and resources, we recover the **st-norm** and **upgrade** its status
- Its structure allows easier numerical computation and analytical values





OUTLOOK

- Generalise to qudits
- Transfer more results from CV to DV via error correction codes





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BOUNDS

- Coincide with the ones obtain by the robustness

| U_{target} | Robustness of Magic | GKP Magic | T -count |
|---------------------|---------------------|-----------|------------|
| T_1 | 1.41421 | 0.272 | 1 |
| $T_{1,2}$ | 1.74755 | 0.543 | 2 |
| CS_{12} | 2.2 | 0.807 | 3(2.967) |
| $T_{1,2,3}$ | 2.21895 | 0.815 | 3(2.996) |
| $CS_{12,13}$ | 2.55556 | 0.907 | 4(0.907) |
| $T_1 CS_{23}$ | 2.80061 | 1.079 | 4(3.966) |



MIXED QUBIT STATES

- The Wigner log. Negativity can be smaller than the one for pure stabilizer states
- Need attribute for that

$$\tilde{\mathcal{G}}(\hat{\rho}) = \max \left[0, \log_2 \left(\frac{1}{\sqrt{\pi}^n} \sum_{\mathbf{i}, \mathbf{j} \in \mathcal{F}_2^n} \left| \sum_{\mathbf{k} \in \mathcal{F}_2^n} (-1)^{\mathbf{i} \cdot \mathbf{k}} \rho_{\mathbf{k}, \mathbf{k} + \mathbf{j}} \right| \right) - \log_2 (\mathcal{N}_0(n)) \right]$$

- But for $\hat{\rho} = \hat{\rho}_I \otimes \hat{\rho}_O$

$$\tilde{\mathcal{G}}(\hat{\rho}_I \otimes \hat{\rho}_O) = \max \left[0, \log_2 \left(\int_{\mathcal{C}_1} d^n \mathbf{r}_1 |W_{\hat{\rho}_I}| \int_{\mathcal{C}_2} d^n \mathbf{r}_2 |W_{\hat{\rho}_O}| \right) - \log_2 (\mathcal{N}_0(n)) \right] = 0$$



CHOI-JAMIOLKOWSKI ISOMORPHISM

- Bijection between channel and state

$$\hat{\varphi} = (\hat{\Phi} \otimes 1) |\psi\rangle\langle\psi| = \frac{1}{2} \sum_{j,k=0}^1 \Phi(|j\rangle\langle k|) \otimes |j\rangle\langle k|$$

$$|\varphi_U\rangle = (\hat{U} \otimes 1) \frac{1}{\sqrt{2^n}} \sum_{\mathbf{j} \in \mathcal{F}_2^n} |\mathbf{j}, \mathbf{j}\rangle$$

$$|\psi\rangle = \frac{1}{\sqrt{2}} \sum_{i=0}^1 |i, i\rangle$$