

Computational Methods and Modelling

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Lecture 4

Ordinary Differential Equations (ODEs), Euler and Runge-Kutta methods



In the present and following lectures, we will study how to deal with Differential Equations using numerical methods

- ▶ How to solve Ordinary Differential Equations using numerical methods
 - ▶ How to approximate derivatives of different orders that appear in Ordinary and Partial Differential Equation
 - ▶ Learn about the stability of numerical schemes
-
- ▶ In this lecture, we will discuss
 - ▶ Derivatives and Taylor series
 - ▶ Euler Method
 - ▶ Runge-Kutta Methods

Differential equations are omnipresent in mathematical models of natural phenomena and engineering applications.

Examples are countless and with complexity that might range from

- Simple, linear Ordinary Differential Equations (ODE):
Pendulum equation for in the small-angle approximation:

$$\frac{d^2\theta}{dt^2} + \frac{g}{l}\theta = 0$$

- To non-linear Ordinary Differential Equations:
Logistic differential equation (applications in machine learning, population dynamics, virus spread)

$$\frac{df}{dt} = rf - \frac{rf^2}{k}$$

where f is the population (or number of infected)

- And non-linear systems of Partial Differential Equations (PDE):
Navier-Stokes equations of fluid-dynamics (weather forecasting, energy production devices like gas and wind turbines)

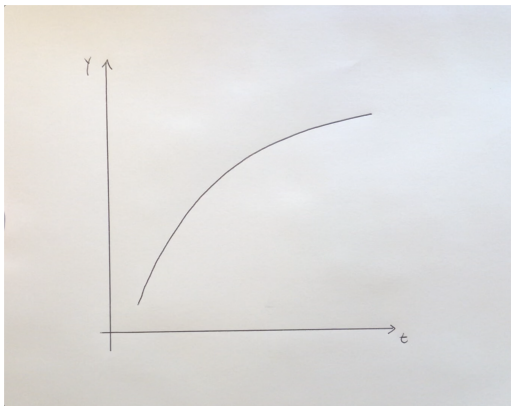
$$\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} = -\frac{1}{\rho} \frac{\partial p}{\partial x_i} + \nu \frac{\partial \tau_{ij}}{\partial x_j}$$

where u_i is the velocity vector, p the pressure, ν the kinematic viscosity, and τ_{ij} the stress tensor.

Common aspect: All these equations contain derivatives.

Approximation of a derivative

- We want to compute the derivative of a function $y(t)$ in the point t_i .



See video on Learn: “Lecture 4 - Differential Equations, Euler and Runge-Kutta methods - Video 1”
in Course Material - Lecture Slides

- The derivative (the slope of the function) can be approximated as:

$$\frac{dy}{dt} \approx \frac{\Delta y}{\Delta t} = \frac{y(t_{i+1}) - y(t_i)}{t_{i+1} - t_i}$$

Approximation of a derivative

- The approximate expression for the slope

$$\frac{dy}{dt} \approx \frac{\Delta y}{\Delta t} = \frac{y(t_{i+1}) - y(t_i)}{t_{i+1} - t_i}$$

is called a **Finite Divided Difference**

- The expression

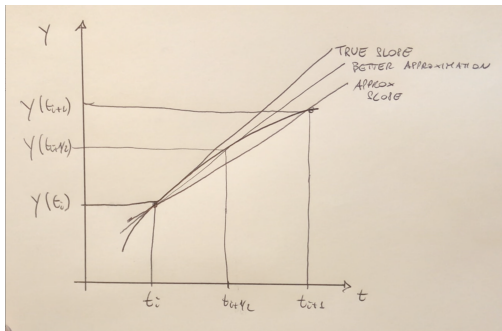
$$\frac{dy}{dt} \approx \frac{\Delta y}{\Delta t}$$

is approximate because Δ is finite.

- From calculus:

$$\frac{dy}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\Delta y}{\Delta t}$$

So if $\Delta t \rightarrow 0$ the approximate slope converges to the real one



See video on Learn: "Lecture 4 - Differential Equations, Euler and Runge-Kutta methods - Video 2"
in Course Material - Lecture Slides

- Taylor's theorem: If the function $f(x)$ of the independent variable x and its $n + 1$ derivatives are continuous in an interval containing the two points x_i and $x_{i+1} = x_i + h$, $f(x)$ can be expanded in the following series:

$$\begin{aligned} f(x_{i+1}) = & f(x_i) + \frac{f'(x_i)}{1!}(x_{i+1} - x_i) + \\ & + \frac{f''(x_i)}{2!}(x_{i+1} - x_i)^2 + \frac{f'''(x_i)}{3!}(x_{i+1} - x_i)^3 + \\ & + \dots + \\ & + \frac{f^{(n)}(x_i)}{n!}(x_{i+1} - x_i)^n + R_n \end{aligned}$$

where

$$R_n = \int_{x_i}^{x_{i+1}} \frac{(x_i - t)^n}{n!} f^{(n+1)}(t) dt$$

- R_n can be also expressed in the **Lagrangian form** (the derivation of R_n is not important for our purpose).

$$R_n = \frac{f^{(n+1)}(\xi)}{(n+1)!} (x_{i+1} - x_i)^{n+1}$$

where $x_{i+1} < \xi < x_i$

- R_n is often called **Truncation Error**

- ▶ Defining $h = x_{i+1} - x_i$ the Taylor series can be written as:

$$f(x_{i+1}) = f(x_i) + \frac{f'(x_i)}{1!}h + \frac{f''(x_i)}{2!}h^2 + \frac{f'''(x_i)}{3!}h^3 + \dots + \frac{f^{(n)}(x_i)}{n!}h^n + R_n$$

- ▶ It can be used to approximate a function in a point in terms of the value of the function and its derivatives in another point
- ▶ Depending on the number of terms we keep in the series, we have different levels of approximation:

- ▶ *zero-order approximation:*

$$f(x_{i+1}) = f(x_i)$$

If $f(x)$ is constant, this is a perfect estimate

- ▶ *first-order approximation*

$$f(x_{i+1}) = f(x_i) + f'(x_i)h$$

This can predict a change in the function, but it is exact only if the function is linear

- ▶ *order n approximation*

$$f(x_{i+1}) = f(x_i) + f'(x_i)h + \frac{f''(x_i)}{2!}h^2 + \dots + R_n \quad \text{with} \quad R_n = \frac{f^{(n+1)}(\xi)}{(n+1)!}h^{n+1}$$

► The Truncation Error

$$R_n = \frac{f^{(n+1)}(\xi)}{(n+1)!} h^{n+1}$$

cannot be determined since ξ is not known. We only know that it lies between x_i and x_{i+1}

- However, we have control over h . For different orders (equivalently, different values of n), the error decreases in different ways if we decrease h .
- We often write $R_n = \mathcal{O}(h^{n+1})$.
The expression $\mathcal{O}(h^{n+1})$ states that the error is of order of h^{n+1} , which means that the error is proportional to h^{n+1} .

- ▶ Let's consider the Taylor series, truncated after the first derivative:

$$f(x_{i+1}) = f(x_i) + f'(x_i)(x_{i+1} - x_i) + R_1$$

- ▶ This equation can be solved for:

$$f'(x_i) = \frac{f(x_{i+1}) - f(x_i)}{x_{i+1} - x_i} - \frac{R_1}{x_{i+1} - x_i}$$

- ▶ Which is the same expression we wrote in the graphical exercise we did before, but now we have also an estimate of the error:

$$\frac{R_1}{x_{i+1} - x_i} = \frac{f''(\xi)(x_{i+1} - x_i)^2}{2!} \frac{1}{x_{i+1} - x_i} = \mathcal{O}(x_{i+1} - x_i)$$

- ▶ In the usual more compact form:

$$f'(x_i) = \frac{\Delta f_i}{h} + \mathcal{O}(h)$$

where Δf_i is called **first forward difference** and h is called the **step size**.
It is called forward since we used data in i and $i + 1$

- ▶ Then, $\Delta f_i/h$ is the **first forward divided difference**

The first forward divided difference is one of many ways to approximate the derivative using the Taylor series.

- We want to solve equations in the form

$$\frac{dy}{dx} = f(x, y)$$

with a given initial condition $y_0 = y(x_0)$

- To solve this equation with a numerical method on a set of discrete points, we need to be able to extrapolate from a value y_i to a new value y_{i+1} over a step h (usually starting from the initial condition y_0).
- In a general mathematical form, this translates to:

$$y_{i+1} = y_i + \phi h$$

where ϕ is an estimate of an appropriate slope of the function y over the step h .

- Then, this formula can be applied step-by-step to compute on an interval as large as we want.
- This formula is the general way to express all **one-step** methods.
- The difference for different methods could be in the way we specify the slope ϕ .
- The simplest approach is to estimate the slope from the differential equation itself as the first derivative of y at the point x_i , which is nothing else than $f(x_i, y_i)$.
- This approach is called the **Euler Method**.
- Alternative ways to estimate the slope ϕ can result in more accurate predictions, for example in **one-step Runge-Kutta methods**.

- For the equation:

$$\frac{dy}{dx} = f(x, y) \quad \text{with} \quad y_0 = y(x_0)$$

The formula

$$y_{i+1} = y_i + f(x_i, y_i)h$$

is referred to as **Euler Method** or sometimes the **Euler-Cauchy Method**.

- A new value of y is computed extrapolating **linearly** over the step h using a slope approximated with the derivative in the original point x_i , where the solution and its derivatives are known.

Euler method

```
# -----
# importing modules
import numpy as np
import matplotlib.pyplot as plt
import math
# -----

# -----
# inputs

# functions that returns dy/dx
# i.e. the equation we want to solve:  $dy/dx = -y$ 
def model(y,x):
    k = - 1
    dydx = k * y
    return dydx

# initial conditions
x0 = 0
y0 = 1
# total solution interval
x_final = 1
# step size
h = 0.2
# -----

# -----
# Euler method

# number of steps
n_step = math.ceil(x_final/h)

# Definition of arrays to store the solution
y_eul = np.zeros(n_step+1)
x_eul = np.zeros(n_step+1)

# Initialize first element of solution arrays
# with initial condition
y_eul[0] = y0
x_eul[0] = x0

# Populate the x array
for i in range(n_step):
    x_eul[i+1] = x_eul[i] + h

# Apply Euler method n_step times
for i in range(n_step):
    # compute the slope using the differential equation
    slope = model(y_eul[i],x_eul[i])
    # use the Euler method
    y_eul[i+1] = y_eul[i] + h * slope
# -----
```

Output

```
# -----
# super refined sampling of the exact solution  $c \cdot e^{-x}$ 
# n_exact linearly spaced numbers
# only needed for plotting reference solution

# Definition of array to store the exact solution
n_exact = 1000
x_exact = np.linspace(0,x_final,n_exact+1)
y_exact = np.zeros(n_exact+1)

# exact values of the solution
for i in range(n_exact+1):
    y_exact[i] = y0 * math.exp(-x_exact[i])
# -----

# -----
# print results on screen
print ('Solution: step x y-eul y-exact error%')
for i in range(n_step+1):
    print(i,x_eul[i],y_eul[i], y0 * math.exp(-x_eul[i]),
          (y_eul[i]- y0 * math.exp(-x_eul[i]))/
          (y0 * math.exp(-x_eul[i])) * 100)
# -----

# -----
# print results in a text file (for later use if needed)
file_name= 'output_h' + str(h) + '.dat'
f_io = open(file_name,'w')
for i in range(n_step+1):
    s1 = str(i)
    s2 = str(x_eul[i])
    s3 = str(y_eul[i])
    s4 = s1 + ' ' + s2 + ' ' + s3
    f_io.write(s4 + '\n')
f_io.close()
# -----

# -----
# plot results
plt.plot(x_eul, y_eul , 'b.-',x_exact, y_exact , 'r-')
plt.xlabel('x')
plt.ylabel('y(x)')
plt.show()
# -----
```

Solution obtained with the Euler method

- We consider the ODE:

$$\frac{dy}{dx} = -y \quad \text{with} \quad y(x=0) = 1$$

- which has the analytical (exact) solution

$$y(x) = e^{-x}$$

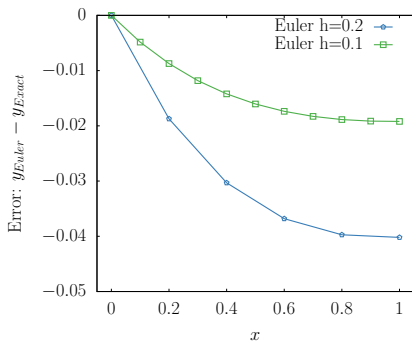
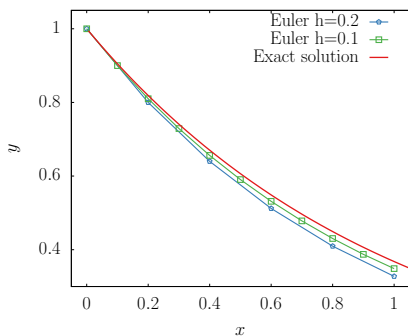
- We run the Euler methods to compute the solution in the interval $0 \leq x \leq 1$ for two different steps $h = 0.2$ and $h = 0.1$:

x	y exact	y Euler h = 0.1	error % h = 0.1	y Euler h = 0.2	error % h = 0.2
0.0	1.0	1.0	0.0	1.0	0.0
0.1	0.904	0.9	-0.534	NA	NA
0.2	0.818	0.81	-1.066	0.800	-2.287
0.3	0.740	0.729	-1.595	NA	NA
0.4	0.670	0.656	-2.121	0.640	-4.523
0.5	0.606	0.590	-2.644	NA	NA
0.6	0.548	0.531	-3.165	0.512	-6.707
0.7	0.496	0.478	-3.682	NA	NA
0.8	0.449	0.43	-4.197	0.409	-8.841
0.9	0.406	0.387	-4.709	NA	NA
1.0	0.367	0.348	-5.219	0.327	-10.92

Table: Solution and error for the Euler method with two two different steps (NA = not available)

Solution obtained with the Euler method

- Plotting the two solutions in the interval $0 \leq x \leq 1$ with different steps $h = 0.2$ and $h = 0.1$ we obtain:



- The error decreases by approximately a factor of 2 if the step h is halved
- Note that we used very large values of h in the example to highlight the numerical error

- **Runge-Kutta** methods achieve high accuracy without the use of higher order derivatives like in the case of the Taylor series

- Many versions exist, but all can be cast as:

$$y_{i+1} = y_i + \phi(x_i, y_i, h)h$$

where $\phi(x_i, y_i, h)$ is usually called an **increment function**.

- The increment function $\phi(x_i, y_i, h)$ is written, in general form as:

$$\phi = a_1 k_1 + a_2 k_2 + \dots + a_n k_n$$

where the coefficients a_j are constant and:

$$k_1 = f(x_i, y_i)$$

$$k_2 = f(x_i + p_{11}h, y_i + q_{11}k_1h)$$

$$k_3 = f(x_i + p_{21}h, y_i + q_{21}k_1h + q_{22}k_2h)$$

.

.

$$k_n = f(x_i + p_{n-1,1}h, y_i + q_{n-1,1}k_1h + q_{n-1,2}k_2h + \dots + q_{n-1,n-1}k_{n-1}h)$$

where also the coefficients p_j and q_j are constant.

- The $k_1 \dots k_n$ can be computed in a cascade from k_1 to k_2 to k_n (recurrence relations)
- Once n is selected, which is the order of the method, all the constants are computed by equating $y_{i+1} = y_i + \phi(x_i, y_i, h)h$ to the terms of the Taylor series.
- Finally, it is worth noting that the Runge-Kutta method with $n = 1$ is the Euler method.

Computation of the coefficients of the second-order Runge-Kutta method

- With $n = 2$, the Runge-Kutta method is written as:

$$y_{i+1} = y_i + (a_1 k_1 + a_2 k_2)h$$

where

$$k_1 = f(x_i, y_i)$$

$$k_2 = f(x_i + p_1 h, y_i + q_{11} k_1 h)$$

and we need to compute a_1 , a_2 , p_1 , and q_{11} .

- To do this, we start from the Taylor series:

$$y_{i+1} = y_i + f(x_i, y_i)h + \frac{f'(x_i, y_i)}{2}h^2$$

We use chain rule to compute:

$$f'(x_i, y_i) = \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} \frac{dy}{dx}$$

and substituting this in the Taylor series we get

$$y_{i+1} = y_i + f(x_i, y_i)h + \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} \frac{dy}{dx} \right) \frac{h^2}{2}$$

- Now, we have to equate this with $y_{i+1} = y_i + (a_1 k_1 + a_2 k_2)h$ to find the coefficients.

Computation of the coefficients of the second-order Runge-Kutta method

- ▶ Interpreting $k_2 = f(x_i + p_1 h, y_i + q_{11} k_1 h)$ as a Taylor series expansion with respect to both variables x and y , we can write:

$$f(x_i + p_1 h, y_i + q_{11} k_1 h) = f(x_i, y_i) + p_1 h \frac{\partial f}{\partial x} + q_{11} k_1 h \frac{\partial f}{\partial y} + \mathcal{O}(h^2)$$

- ▶ Inserting this and the expression $k_1 = f(x_i, y_i)$ in $y_{i+1} = y_i + (a_1 k_1 + a_2 k_2)h$:

$$y_{i+1} = y_i + \boxed{[a_1 f(x_i, y_i) + a_2 f(x_i, y_i)] h} + \left[a_2 p_1 \frac{\partial f}{\partial x} + a_2 q_{11} f(x_i, y_i) \frac{\partial f}{\partial y} \right] h^2 + \mathcal{O}(h^2)$$

- ▶ We equate term-by-term this equation with

$$y_{i+1} = y_i + \boxed{f(x_i, y_i) h} + \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} \frac{dy}{dx} \right) \frac{h^2}{2}$$

- ▶ In order for the two highlighted terms to be equal, we must have
- ▶ Equating also the other terms, we have

$$a_1 + a_2 = 1$$

$$a_2 p_1 = \frac{1}{2} \quad \text{and} \quad a_2 q_{11} = \frac{1}{2}$$

- So the second-order Runge-Kutta method is:

$$y_{i+1} = y_i + (a_1 k_1 + a_2 k_2)h$$

where

$$k_1 = f(x_i, y_i)$$

$$k_2 = f(x_i + p_1 h, y_i + q_{11} k_1 h)$$

with

$$a_1 + a_2 = 1$$

$$a_2 p_1 = \frac{1}{2}$$

$$a_2 q_{11} = \frac{1}{2}$$

- Since we have 3 equations and 4 unknowns, one of the coefficients needs to be specified arbitrarily. Therefore, there is an infinite number of second-order Runge-Kutta methods, which will give different results if the equation is more complicated than a simple linear or quadratic one.

Computation of the coefficients of the second-order Runge-Kutta method

- If we assume $a_2 = 1/2$, we can calculate the other coefficients to obtain $a_1 = 1/2$, $p_1 = 1$, and $q_{11} = 1$. Then:

$$y_{i+1} = y_i + \left(\frac{1}{2}k_1 + \frac{1}{2}k_2\right)h$$

with

$$k_1 = f(x_i, y_i)$$

$$k_2 = f(x_i + h, y_i + k_1 h)$$

This is the predictor-corrector method described before, since k_1 and k_2 are the derivatives in x_i and x_{i+1}

- If $a_2 = 1$, we obtain the midpoint method also described before
- If $a_2 = 2/3$, we obtain the **Ralston method**, which is the second-order Runge-Kutta method with the minimum truncation error:

$$y_{i+1} = y_i + \left(\frac{1}{3}k_1 + \frac{2}{3}k_2\right)h$$

with

$$k_1 = f(x_i, y_i)$$

$$k_2 = f\left(x_i + \frac{3}{4}h, y_i + \frac{3}{4}k_1 h\right)$$

- ▶ The most popular Runge-Kutta methods are the fourth order
- ▶ The most used version is:

$$y_{i+1} = y_i + \left(\frac{1}{6}k_1 + \frac{2}{6}k_2 + \frac{2}{6}k_3 + \frac{1}{6}k_4 \right) h$$

with

$$k_1 = f(x_i, y_i)$$

$$k_2 = f\left(x_i + \frac{1}{2}h, y_i + \frac{1}{2}k_1h\right)$$

$$k_3 = f\left(x_i + \frac{1}{2}h, y_i + \frac{1}{2}k_2h\right)$$

$$k_4 = f(x_i + h, y_i + k_3h)$$

- ▶ The slope

$$\frac{1}{6}k_1 + \frac{2}{6}k_2 + \frac{2}{6}k_3 + \frac{1}{6}k_4$$

is a weighted average of four slopes with

- ▶ K_1 computed in the first point x_i (like in the Euler method)
- ▶ K_2 and K_3 computed in the middle $x_{i+1/2}$ (like in the mid-point method)
- ▶ K_4 computed in the end-point x_{i+1} (like in the predictor-corrector method)

- As usual, we want to solve

$$\frac{dy}{dx} = f(x, y) \quad \text{with} \quad y_0 = y(x_0)$$

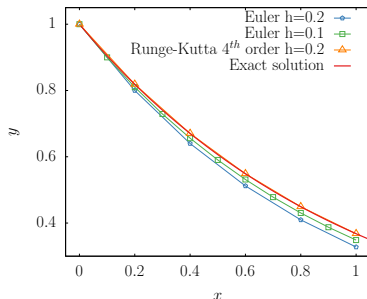
Euler method

```
# -----  
# Euler method  
  
# Apply Euler method n_step times  
for i in range(n_step):  
    # compute the slope using the differential equation  
    slope = model(y_eul[i], x_eul[i])  
    # use the Euler method  
    y_eul[i+1] = y_eul[i] + h * slope  
# -----
```

Runge-Kutta fourth order method

```
# -----  
# Fourth Order Runge-Kutta method  
  
# Apply RK method n_step times  
for i in range(n_step):  
    # Compute the four slopes  
    x_dummy = x_rk[i]  
    y_dummy = y_rk[i]  
    k1 = model(y_dummy, x_dummy)  
  
    x_dummy = x_rk[i] + h/2  
    y_dummy = y_rk[i] + k1 * h/2  
    k2 = model(y_dummy, x_dummy)  
  
    x_dummy = x_rk[i] + h/2  
    y_dummy = y_rk[i] + k2 * h/2  
    k3 = model(y_dummy, x_dummy)  
  
    x_dummy = x_rk[i] + h  
    y_dummy = y_rk[i] + k3 * h  
    k4 = model(y_dummy, x_dummy)  
  
    # compute the slope as weighted average of four slopes  
    slope = 1/6 * k1 + 2/6 * k2 + 2/6 * k3 + 1/6 * k4  
  
    # use the RK method  
    y_rk[i+1] = y_rk[i] + h * slope  
# -----
```

Solution obtained with the Runge-Kutta method



x	y exact	error Euler % $h = 0.1$	error R-K4 % $h = 0.2$
0.0	1.0	0.0	0.0
0.2	0.818	-1.066	3.2×10^{-4}
0.4	0.670	-2.121	6.3×10^{-4}
0.6	0.548	-3.165	9.4×10^{-4}
0.8	0.449	-4.197	1.2×10^{-3}
1.0	0.367	-5.219	1.5×10^{-4}

Table: Error for the Euler and Runge-Kutta methods

- The error for the Runge-Kutta method is remarkably smaller compared to the error obtained with the Euler method
- It is worth noting that the error is small even with the rather large step $h = 0.2$ used