

Numerical problem on Quantum
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Problem: 1: →

The given Hamiltonian: $H = -\frac{\hbar^2}{2I} \frac{\partial^2}{\partial \phi^2} + Mga(1 - \cos \phi)$

So T.I.S.E gives:

$$-\frac{\hbar^2}{2I} \frac{d^2 \psi(\phi)}{d\phi^2} + Mga(1 - \cos \phi) \cdot \psi(\phi) = E \psi(\phi) \dots (1)$$

Given: $M=1$; $a=1$; $g=1$; $I=300$.

Given assumed $\hbar = 1$. (in proper unit for all)

To solve eq (1); I've taken the method of discretisation.

The limit of ϕ be given by $\phi \in [0, 2\pi]$.

I've ~~break~~ broken the interval in N discrete spacial points of equal interval.

$$\text{i.e. } \{\phi_1, \phi_2, \dots, \phi_N\} ; \begin{cases} \phi_{i+1} - \phi_i = h \\ \phi_1 = 0; \phi_N = 2\pi; \phi_{N+1} = \phi_1 \end{cases}$$

$$\text{Let } \psi_i = \psi(\phi_i) ; V_i = V(\phi_i)$$

Now at small h ; the derivative of ψ be given (approximated) by:

$$\frac{d^2 \psi}{d\phi^2} \simeq \frac{\psi_{i+1} + \psi_{i-1} - 2\psi_i}{h^2} \quad | \quad h = \text{small.}$$

So: in discrete form; eq (1) can be written by:

$$-\frac{\hbar^2}{2m\Delta x^2} \frac{\psi_{i+1} + \psi_{i-1} - 2\psi_i}{\Delta x^2} + V_i \psi_i = E \psi_i$$

$$\Rightarrow \underbrace{\left(+ \frac{\hbar^2 \times 2}{2m\Delta x^2} + V_i \right)}_{\beta_i} \psi_i + \underbrace{\left(- \frac{\hbar^2}{2m\Delta x^2} \right)}_{\alpha} (\psi_{i+1} + \psi_{i-1}) = E \psi_i$$

$$\Rightarrow \alpha \psi_{i-1} + \beta_i \psi_i + \alpha \psi_{i+1} = E \psi_i$$

putting $i = 1, 2, 3, \dots, N$ we get; the set of eqn's, (using $\psi_{N+1} = \psi_1$)

$$\beta_1 \psi_1 + \alpha \psi_2 + \alpha \psi_N = E \psi_1$$

$$\alpha \psi_1 + \beta_2 \psi_2 + \alpha \psi_3 = E \psi_2$$

$$\alpha \psi_2 + \beta_3 \psi_3 + \alpha \psi_4 = E \psi_3$$

$$\alpha \psi_N = E \psi_1$$

$$= E \psi_2$$

$$= E \psi_3$$

$$\alpha \psi_{N-2} + \beta_{N-1} \psi_{N-1} + \alpha \psi_N = E \psi_{N-1}$$

$$\alpha \psi_{N-1} + \beta_N \psi_N = E \psi_N$$

$$\alpha \psi_1 +$$

So in matrix form we get:

$$\begin{pmatrix} \beta_1 & \alpha & 0 & 0 & \dots & 0 & \alpha \\ \alpha & \beta_2 & \alpha & 0 & 0 & 0 & 0 \\ 0 & \alpha & \beta_3 & \alpha & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & \dots & \alpha & \beta_i & \alpha & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & \dots & \alpha & \beta_{N-1} & \alpha & 0 \\ \alpha & 0 & \dots & 0 & \alpha & \beta_N & \alpha \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \\ \psi_N \end{pmatrix} = E \begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \\ \psi_N \end{pmatrix}$$

\hat{H} matrix.

In machine; I've given \hat{H} as input matrix and its eigenvalues will be possible energy values.

Obviously we must have to take at least $N \gg 300$ (i.e. 300 space interval) to get a $(N \times N)$ matrix to get N eigenvalues. Here the required no of eig-val is 300. Moreover larger N will give more & more precise eigenvalues.

With these eig-val; I've plotted the energy bands and the band gaps (not properly band gap but energy difference between states.)

(Spin interaction problem is a good checkpoint for this discretization method.)

As asked in question; first; 50th & 200th energy eigenvalues are given by:

$$\left. \begin{aligned} E_1 &= 0.02876 \text{ unit} \\ E_{50} &= 2.17278 \text{ unit.} \\ E_{200} &= 17.61328 \text{ unit.} \end{aligned} \right\} \underline{\text{Answer}}$$

To get a good precision; I ^{have} first broken the whole interval into 3000 small intervals. After that from ~~first~~ 3000 eigenvalues of H ; I have collected first 300 & worked on them.

After $n=43$ state the energy eigenvalue exceeds the maximum of potential.

$$\left. \begin{aligned} E(43) &\simeq 1.98 \text{ unit} \\ \& E(44) &\simeq 2.001 \text{ " } \end{aligned} \right\} \& V_{\max} = 2 \text{ unit.}$$

So after $n=43$; the energy levels come in pairs & the energy gap between them are very small (actually zero as they are degenerate levels.)

Ans we get (the difference is due to numerical error)

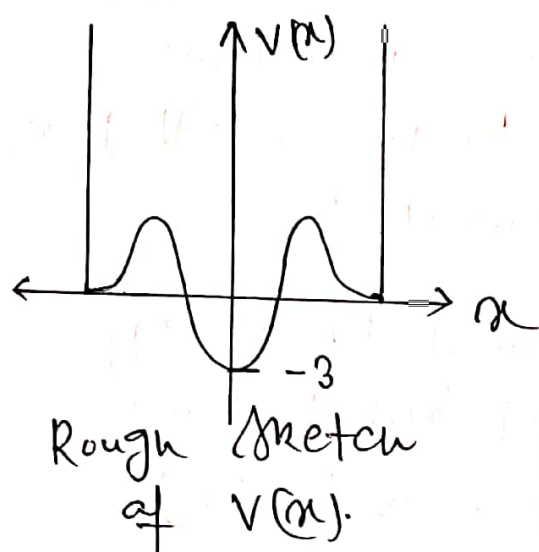
$$E_{45} \simeq E_{46}; E_{47} \simeq E_{48}; E_{49} \simeq E_{50}; \dots \text{ etc}$$

Answer

Problem: 2 →

Given potential $V(x) = (x^2 - 3)e^{-x^2/2}$

And it's in a box of length 10; ($x = -5$ to 5)



Clearly from fig we get:

$$E \geq -3 \text{ (must)}$$

As V is symmetric so we must get for all eig-state $\psi(x)$:

$$\psi(-x) = \pm \psi(x)$$

due to infinity potential at edge: $\psi(\pm 5) = 0$.

Now here I've used the difference eq method to solve T.I.S.E.

where I've taken, $\frac{\hbar^2}{2m} = 1$.

$$\text{i.e. } -\frac{d^2\psi}{dx^2} + V\psi = E\psi$$

discretizing space into small intervals like previous problem $x \rightarrow \{x_0, x_1, \dots, x_N\}$
the eq becomes:

$$-\frac{\psi_{i+1} + \psi_{i-1} - 2\psi_i}{h^2} + V_i\psi_i = E\psi_i$$

$$\Rightarrow \psi_{i+1} = -\left(h^2(E - V_i) + 2\right)\psi_i + \psi_{i-1}$$

By running the loop for i with different values of E with the B-C

$$\Psi_0 = 0 \text{ (obvious)}$$

$$\Psi_1 = 0.1 \text{ or } 0.01 \text{ or, } \dots$$

any number upto normalization will give a similar curve with different amplitude.

For proper choice of E ; Ψ will be close to zero at $x \rightarrow 5$

$$\text{i.e. } \Psi_N \approx 0.$$

o/w for wrong choice Ψ will blow up at edge. Using proper error bounds & shooting different E values in the loop; give collected the allowed energy values.

Lastly with proper Ψ ($\Psi_N \approx 0$); the only task is to normalize it & going to higher energy value.

This method can give upto any arbitrary no of states (E_1, E_2, \dots) as per we allow to vary E .

The checking was done by counting no of maxima, minima of Ψ .

ψ_{i+1}^{eig} excited will have one more extrema than ψ_i^{eig} state

And if ψ_i^{eig} is even/odd for ψ_{i+1} is odd/even

(Infinite well is a good checkpoint for this method, where it gives very good result.)

Comment:-

☒ I have done the same problem in python & mathematica. The mathematica code gives ~~more~~ better result. Here one important fact is as I've taken $m = \frac{1}{2}$; $\hbar = 1$ I got one '-ve' energy where as my friends have taken $m = 1$ & for them there are two states with '-ve' energy. Also the eig-values are different. However my mathematica & python file are producing same result I have checked my code for other potentials (infinite well) as well as $m = 1$ (other friends.)

One problem with the ~~py~~ python code is that, I was unable to find the ground state with the code. It can produce all states from 1st excited states and the plots are also consistent. I request the grader to kindly mention the possible change in my python code to give the complete result.

So the consecutive energy eigenvalues for the Schrodinger eq with $m = \frac{1}{2}$ & $\hbar = 1$ & $V(x) = (x^4 - 3)e^{-x^2/2}$ be given by:

i) From Python code:

0.676	unit.
1.298	"
1.464	"
2.738	"
3.772	"
4.879	"

ground state
is missing in python
code.

ii) From Mathematica:-

-1.7097	unit.
0.679	"
1.301	"
1.467	"
2.744	"

Answer

They are very close and the plots are also similar for different eigenstates.
(As sent both of them.)

2-2 . Band Gap near first Brillouin Zone = 2.046 unit.

☐ I am greatly thankful to my friend Krishnendu Maji; who has helped me for doing problem 2-2. I have used the same algorithm developed by Krishnendu for that problem.