

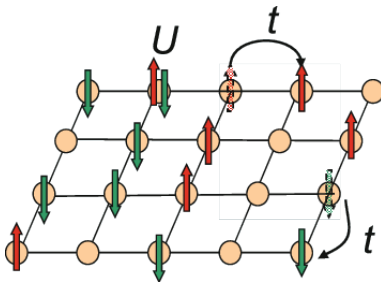
# Cuprates and Stripes

using Hubbard Model

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# Objective

- ▶ Looking for the effects of electron-phonon coupling on the physics of stripes seen in Hubbard model for square lattice in 2D designed for Cuprates.
- ▶ Modelling the coupling using bond-SSH mechanism.
- ▶ In future,
  - ▶ plan to execute more “physical” mechanism, namely, optical-SSH;
  - ▶ looking for the possible similarities and differences with the results obtained from Holstein mechanism.

# Model

- ▶ Hubbard model on 2D square lattice with nearest and next nearest neighbour hopping.
- ▶ Adding bond SSH coupling.

*equation*

- ▶ Relationship between  $\alpha$  and  $\lambda$ .
- ▶ Working in hole doped regime. For starters, going till  $p = 0.2$  and point of interest around  $p = 0.125$  due to ...

# Assesment of sign for different temperature

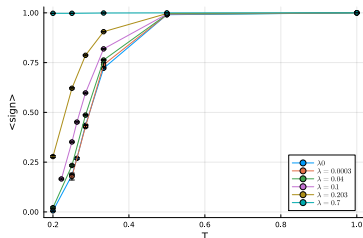


Figure: 1a.  $\langle \text{sign} \rangle$  vs  $T$

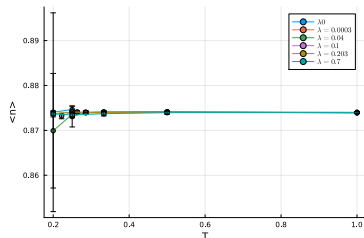


Figure: 1b.  $\langle n \rangle$  vs  $T$

NOTES: To check the limit on lower value of  $T$ , note: with varying  $T$ , important to keep  $\langle n \rangle$  constant.  $\langle n \rangle = 0.875$ ,  $U = 6$ ,  $t' = -0.25$ ,  $t = 1$ , varying  $T$  to check  $\langle \text{sign} \rangle$  for a set of  $\lambda$  values set by  $\alpha$

## Is Mott gap decreasing?

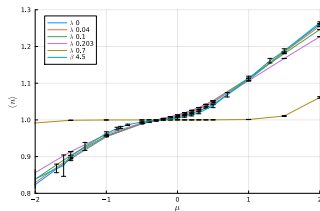
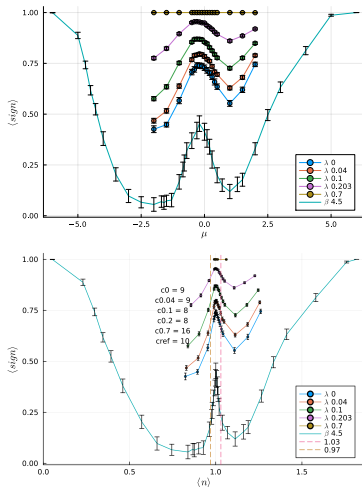


Figure:  $\langle n \rangle$  vs  $\mu$

Initial assessment of Mott gap by looking at  $\langle n \rangle$  vs  $\mu$  curve. The trend of slope increasing with increasing  $\lambda$  indicates that the mott gap is decreasing.

Parameters used:  $\beta = 3.5$  others are same as before.



- ▶ The sign gets worse very quickly as move away from half filling.
- ▶ Same distribution of chemical potential used for varying coupling strength to evaluate filling fraction.
- ▶ The denser the peak at half filling, more number of  $\mu$  points near half filling, meaning larger the gap.
- ▶  $c\{\lambda\}$  (e.g.,  $c_0$ ) counts such number within indicated upper and lower bounds and it again indicates decrease in Mott gap.
- ▶ This is not so evident from  $\mu$  vs sign scan.

# Structure factor

Direct look at the spin  $S(\mathbf{Q}, \omega = 0)$  and charge density  $N(\mathbf{Q}, \omega = 0)$  structure factor calculated for varying  $\lambda$ .  
SDW for  $q = (\pi/a, \pi/a)$  are found suppressed. As an example,

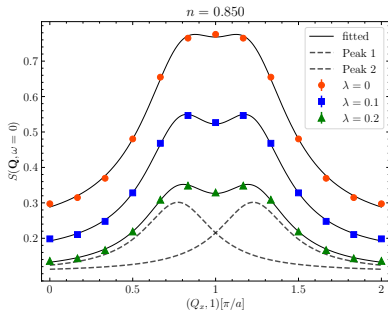
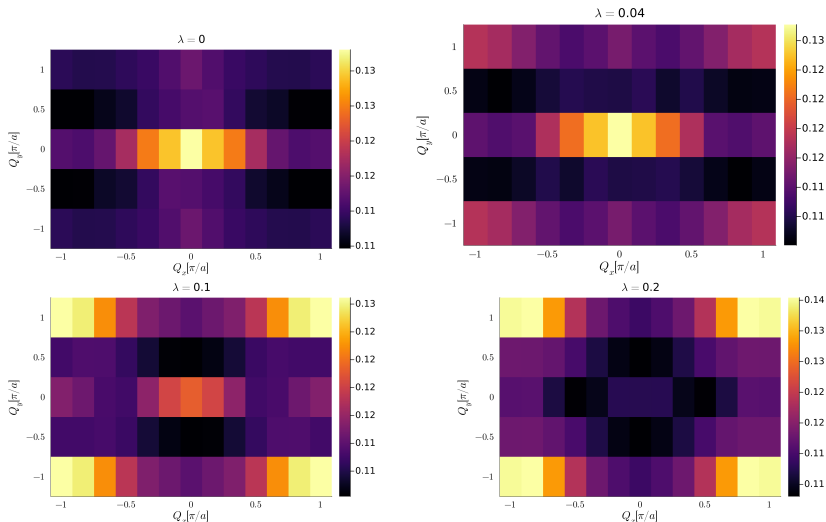


Figure: Spin structure factor along  $(q_x, q_y = \pi/a)$

For the charge density wave, however the location of maxima shifts from  $(0, 0)$  to  $(\pi/a, \pi/a)$ .

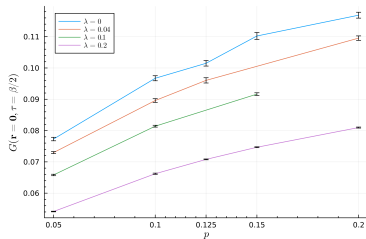
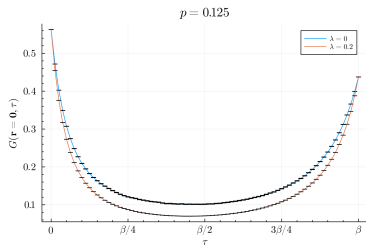


$$p = 0.1, \beta = 4$$

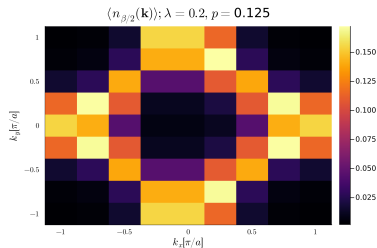
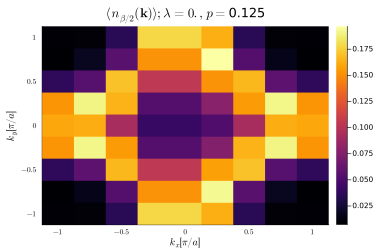
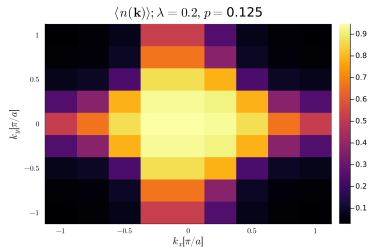
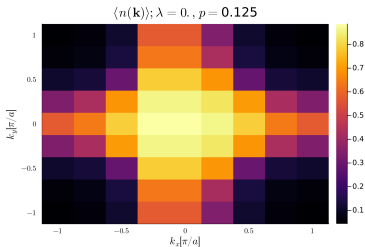


# Studying Mott Gap

- ▶  $G(r=0, \tau = \beta/2)$  gives the spectral weight near fermi energy ( $\mu$  at  $T \neq 0$ ).
- ▶  $G(r=0, \tau)$  records the cumulative spectral weights over an energy range depending on  $\tau$ .
- ▶  $G(k, \tau)$  projects this on  $k$ -space, and with the model being single band, each  $k$  hosts non-degenerate single particle eigenstate.
- ▶  $G(k, \tau = 0^-) = 1 - \langle n_k \rangle$ , recording  $\langle n_k \rangle$



Both tell us that the states near fermi energy are moving away.  
 The former graph also indicates increase in Mott gap, contrary to what we (think) we see earlier in  $\langle n \rangle$  vs  $\mu$  graph.  
 To see clearer on what is happening, looking at  $\langle n_k \rangle$



where,  $\langle n_{\mathbf{k}} \rangle = G(\mathbf{k}, \tau = \beta/2)$  captures the spectral weight near fermi energy, making it clearer to see the fermi surface(FS).

Fermi surface(FS) appears to sharpen.

# Hypothesis:

With Mott gap reducing and the FS sharpening,  $U$  might effectively reducing.

(This has been already seen in Hubb-Holstein model)

## Ideas to check this?

1. Looking for compressibility,  $\kappa(= \partial \langle n \rangle / \partial \mu)$ .
2. Direct comparison of ...(results) between low  $U$ ,  $\lambda = 0$  and high  $U$  with  $\lambda \neq 0$ .
3. Comparison with the weak-coupling physics, RPA limit results.

# Compressibility

With increasing  $\alpha$ , I find that  $\mu$  required to fix specific  $\langle n \rangle$  (let us say, 0.8) shifts away from 0, and  $\kappa$  reduces as well, implying the rate at which  $\langle n \rangle$  is changing w.r.t  $\mu$  is slowing down (slower the rate, lower the number of states).

Both of these lead to the conclusion that the Mott gap is increasing with e-ph coupling.

- ▶ Looking closely at the data of  $\langle n \rangle, \kappa$  with varying  $\mu$ , closer to half-filling, let us say  $\sim 1.004$  or  $\sim 0.997$  as the filling,  $\kappa$  increases with both fixed  $\mu$  or  $\langle n \rangle$  for increasing  $\alpha$ ; hinting decrease in Mott gap.
- ▶ On the other hand, the rough  $\mu$  range for  $\langle n \rangle$  to change from  $\sim 1.008$  to  $\sim 0.997$ , the interval shifts toward lower values of  $\mu$  but stay roughly same  $\sim 0.25 - 0.3$ .
- ▶ It should be noted that for this part of the analysis, better data is needed to conclude the statement with more confidence as the range interval has pattern of  $[0.3, 0.25, 0.25, 0.3]$  for  $[0, 0.12, 0.19, 0.27]$  as  $\alpha$  values. More number of alpha points and more number of  $\mu$  points are needed at  $\langle n \rangle \sim 1$  (there were mostly 3-4 points).
- ▶ IDEA : Fine tune  $\langle n \rangle - \mu$  map by running calc. for denser set of  $\mu$  points. Result?

# Result

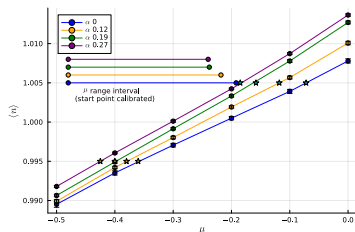


Figure:  $\langle n \rangle$  vs  $\mu$

This shows much clearly and conclusively that the gap is reducing.  $\mu$  range interval for  $0.995 < \langle n \rangle < 1.005$ .

As we move away from  $\langle n \rangle = 0.99$ ,  $\langle n \rangle$  drops drastically, hinting this is very close to Mott gap.

For further concrete comments, however, we need to look at DOS. >>> Analytic Continuation!