



5 电场力的功 电势

Work done by Electric Force

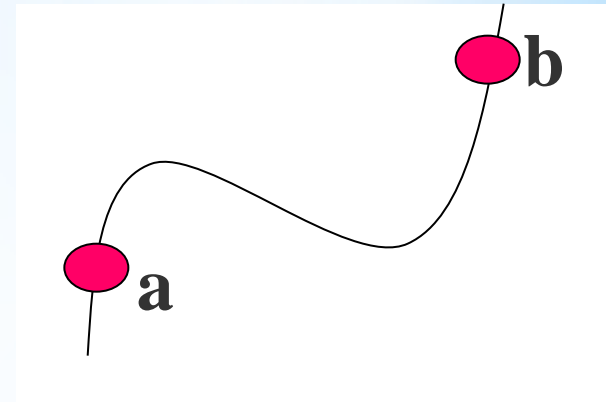
Electric Potential



1. Work done by Electric Force 静电力做功的特点

The work done by the force is given by

$$W = \int_a^b \vec{F} \cdot d\vec{r} = \int_a^b F \cos \theta ds$$



We will show that the force on a charged particle in a electric field produced by any combination of charges at rest is always a conservative force field(保守力场)

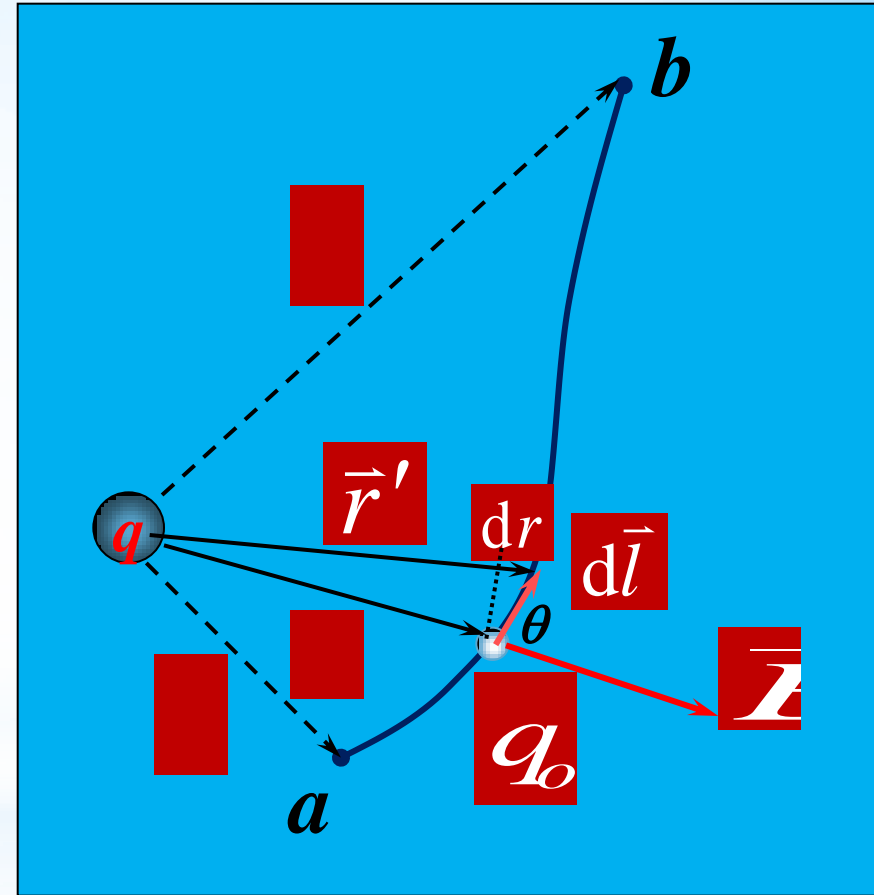
1) Consider the electric field of a point charge. A test charge q_0 moves from a to b in the electric field by the electric force. The work done by the electric force is given by

$$dW = q_0 \vec{E} \cdot d\vec{l} = q_0 E \cos \theta dl$$

$$dW = \frac{q_0 q}{4\pi\epsilon_0 r^2} \cos \theta dl = \frac{q_0 q}{4\pi\epsilon_0 r^2} dr$$

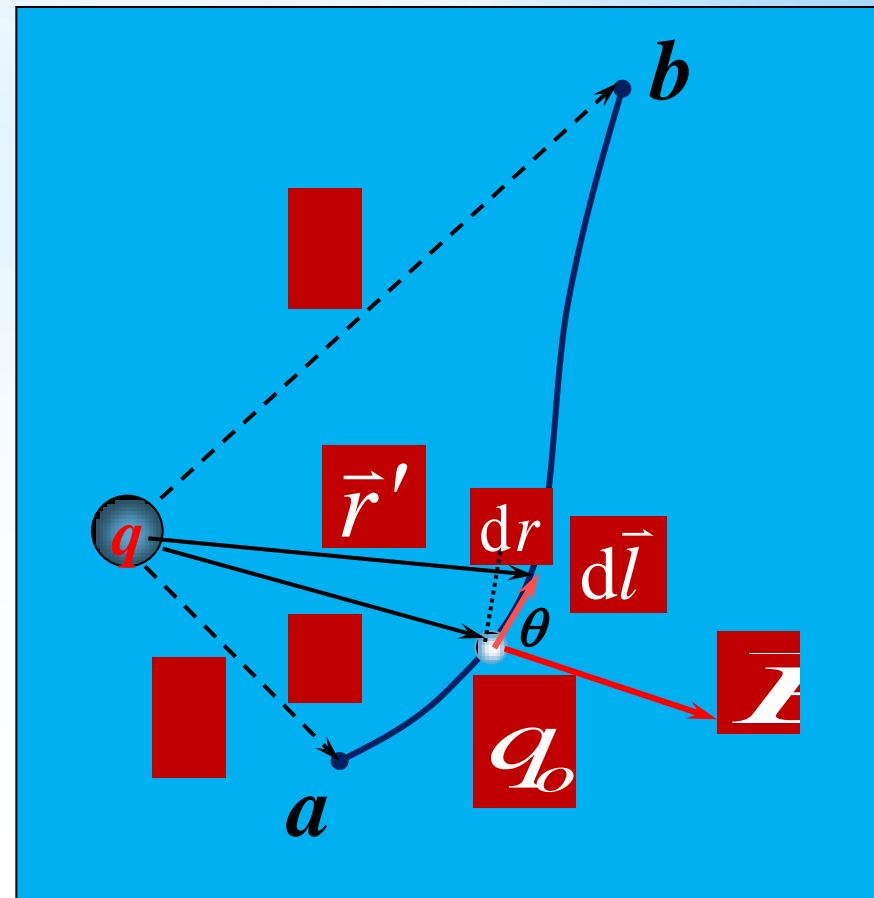
$$W_{ab} = \int_{r_a}^{r_b} \frac{q_0 q}{4\pi\epsilon_0 r^2} dr$$

$$= \frac{qq_0}{4\pi\epsilon_0} \left(\frac{1}{r_a} - \frac{1}{r_b} \right)$$



$$W = \frac{qq_0}{4\pi\epsilon_0} \left(\frac{1}{r_a} - \frac{1}{r_b} \right)$$

结论： 给定试验电荷在静电场中移动时，电场力所作的功只与试验电荷的起点和终点的位置有关，而与路径无关。



Conclusion: work done by electric field depends only on the initial & final positions, but not on the path.

2) In general, the electric field equals to

$$\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3 + \dots\dots$$

The work done on test charge by the electric field \vec{E} is given by

$$W = \int_a^b q_0 \vec{E} \cdot d\vec{r} = \int_a^b q_0 \vec{E}_1 \cdot d\vec{r} + \int_a^b q_0 \vec{E}_2 \cdot d\vec{r} + \dots\dots$$

Since the every term in the above formula (公式) depends only on the positions a and b, the sum depends only on the positions a and b. Therefore, the electrostatic force is conservative.

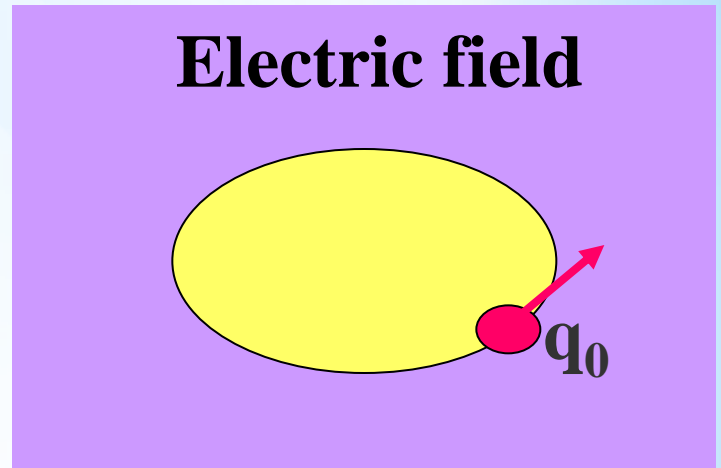
即电场力是保守力。静电场是保守场。

2. Circular theorem of the electrostatic field

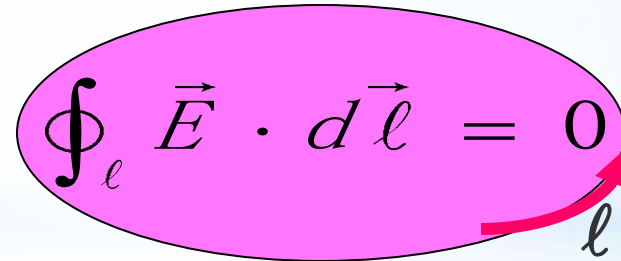
静电场的环路定理

If the test charge moves around a closed path, the work by the electric field is zero:

$$W = \oint_{\ell} q_0 \vec{E} \cdot d\vec{\ell} = 0$$



静电场的环路定理:



The diagram shows a pink oval representing a closed path. A red arrow points along the bottom edge of the oval, labeled ℓ , indicating the direction of integration. The equation $\oint_{\ell} \vec{E} \cdot d\vec{\ell} = 0$ is written inside the oval.

$$\oint_{\ell} \vec{E} \cdot d\vec{\ell} = 0$$

静电场中电场强度沿任意闭合路径线积分（环流）为零。

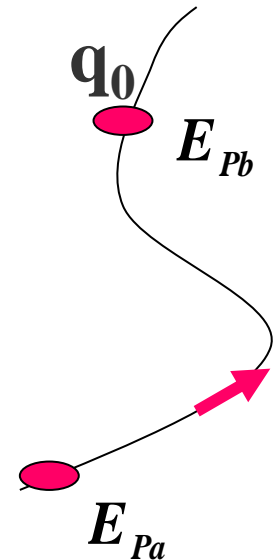
3. Electric potential energy and electrical potential 电势能和电势

- (1) 静电力是保守力，可引入电势能的概念。
- (2) 定义静电力由 a 点 $\rightarrow b$ 点作的功 为电势能增量的负值：

$$W_{ab} = q_o \int_a^b \vec{E} \cdot d\vec{l} = - (E_{Pb} - E_{Pa}) = E_{Pa} - E_{Pb}$$

E_{Pa} : the electric potential energy of q_0 at point a.

E_{Pb} : the electric potential energy of q_0 at point b.



令***b***点的势能为零 ($E_{pb}=0$)

$$W_{ab} = \int_a^b q_o \vec{E} \cdot d\vec{l} = E_{pa} - E_{pb}$$

*a*点的势能:

$$E_{pa} = \int_a^b q_o \vec{E} \cdot d\vec{l}$$

结论: 试验电荷 q_o 在空间某处的电势能在数值上就等于将 q_o 从该处移至势能的零点电场力所作的功。

电势能的零点可以任意选取，习惯上，当场源电荷为**有限带电体**时，通常把**电势能的零点选取在无穷远处**。

空间*a*点的电势能:

$$E_{pa} = \int_a^{\infty} q_o \vec{E} \cdot d\vec{l}$$

注意:

- 电势能为电场和位于电场中的电荷这个系统所共有。
- 电势能是标量，可正可负。

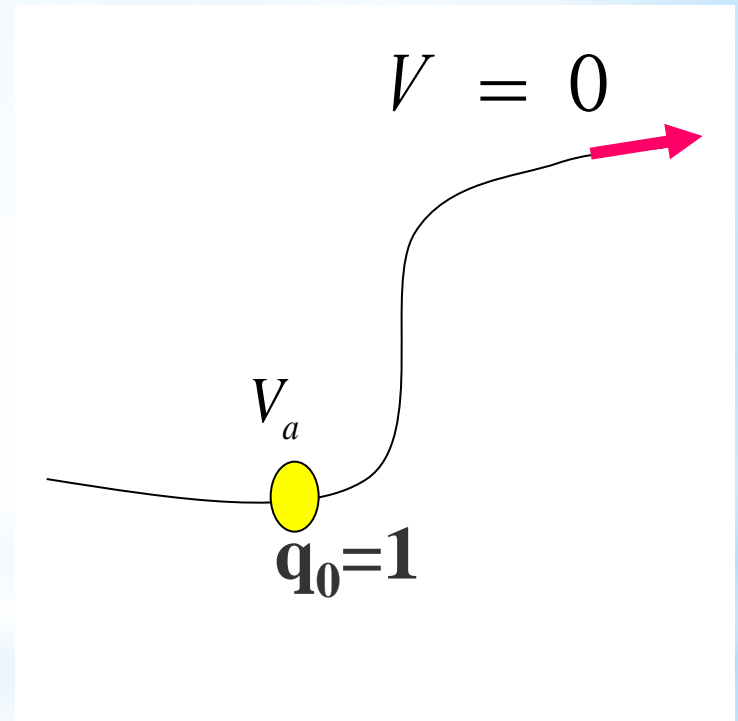
Electrical potential 电势:

$$V_a = \frac{E_{pa}}{q_0} = \int_a^{\text{势能零点}} \vec{E} \cdot d\vec{l}$$

单位: 伏特 ($\text{V} = \text{J} \cdot \text{C}^{-1}$)

结论: 电场中 a 点的电势, 在数值上等于把单位正电荷从 a 点移至势能的零点处电场力所作的功。

也等于在 a 点的单位正电荷所具有的电势能。



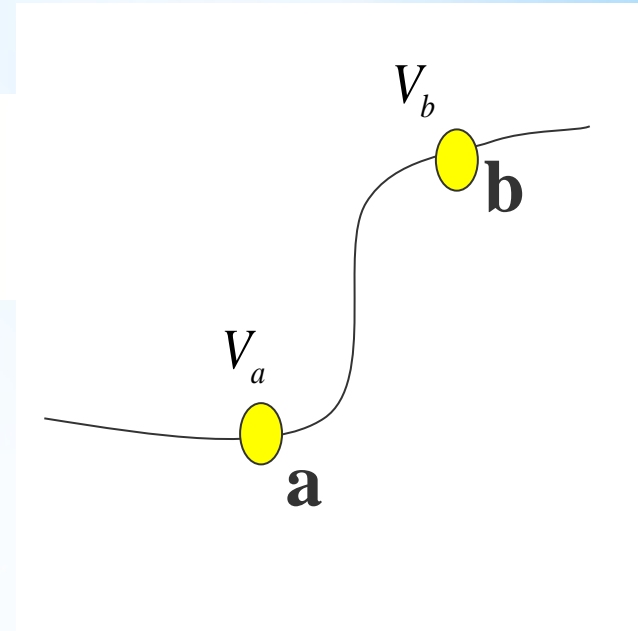
电场中一点的电势其数值与势能零点的选取有关。

The potential difference between point a and point b is

电势差：

$$V_a - V_b = \int_a^{\infty} \vec{E} \cdot d\vec{\ell} - \int_b^{\infty} \vec{E} \cdot d\vec{\ell}$$

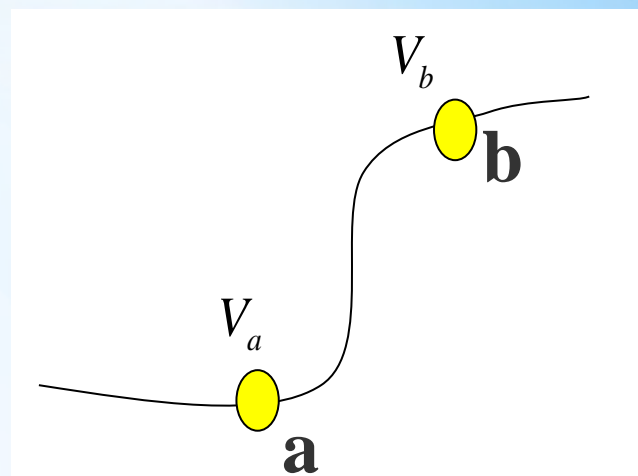
$$= \int_a^{\infty} \vec{E} \cdot d\vec{\ell} + \int_{\infty}^b \vec{E} \cdot d\vec{\ell} = \int_a^b \vec{E} \cdot d\vec{\ell}$$



结论： 静电场中 a, b 两点的电势差，等于将单位正电荷从 a 点移至 b 点电场力所作的功。是绝对的。

In summary: $E_{Pa} = qV_a$

$$\begin{aligned} W_{a \rightarrow b} &= -(E_{Pb} - E_{Pa}) = E_{Pa} - E_{Pb} \\ &= q(V_a - V_b) = q\Delta V \end{aligned}$$



- 电势是描述电场性质的物理量（能量）；
- 电势 V 是标量；而 \vec{E} 指向电势降低的方向；
 仅在电场力作用下, 正电荷的运动: 电势高→低;
 仅在电场力作用下, 负电荷的运动: 电势低→高.
- 电势的零参考点的选取是任意的。
 有限大带电体一般选无穷远为电势零点;
 对无限大带电体的电场, 通常选特殊点或线或面上的电势为零.

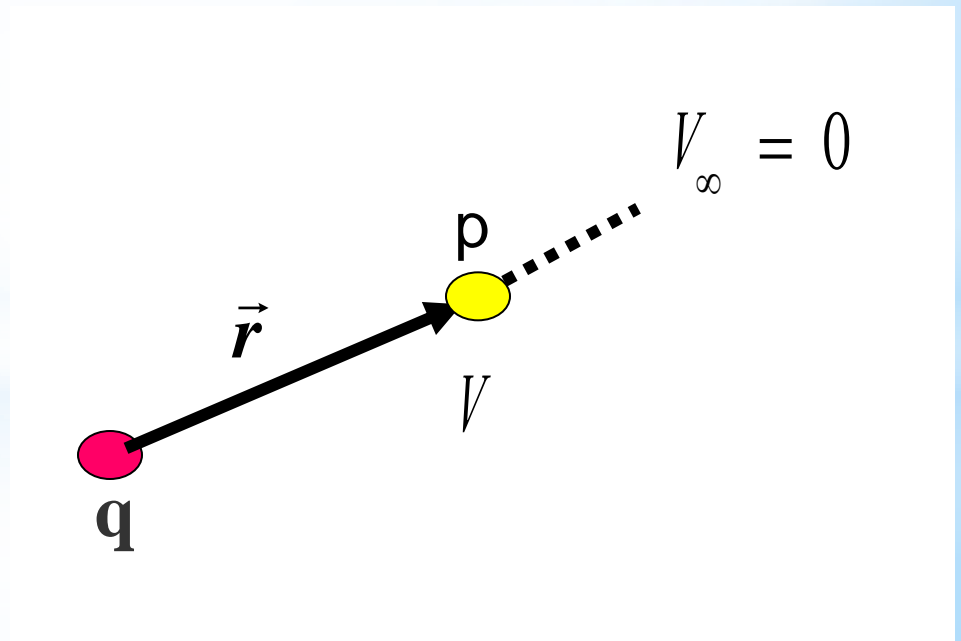
4. The calculation of the electrical potential 电势的计算

1) The electrical potential of field produced by a point charge q 点电荷的电势:

点电荷电场中的电势:

无穷远为电势零点!

$$\begin{aligned} V_p &= \int_p^\infty \vec{E} \cdot d\vec{l} = \int_r^\infty \frac{q}{4\pi\epsilon_0 r^2} dr \\ &= \frac{q}{4\pi\epsilon_0 r} \end{aligned}$$

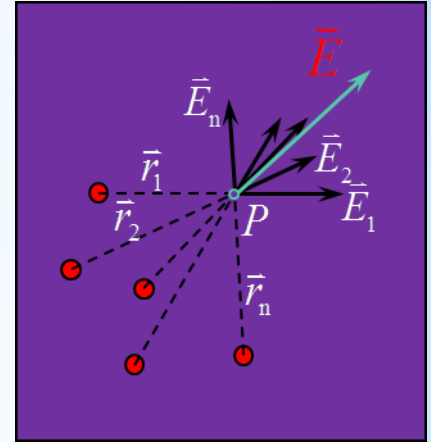


2) The electrical potential of field produced by the any charged system 点电荷系电场中的电势:

$$\vec{E} = \vec{E}_1 + \vec{E}_2 + \cdots + \vec{E}_n \quad V_p = \int_p^\infty \vec{E} \cdot d\vec{l}$$

$$V_p = \int_p^\infty (\vec{E}_1 + \vec{E}_2 + \cdots + \vec{E}_n) \cdot d\vec{l} = \sum_i V_i$$

$$\therefore V = \sum V_i = \sum_i \frac{q_i}{4\pi\epsilon_0 r_i}$$



the superposition principle of electrical potential.

电势叠加原理: 点电荷系电场中任一点的电势, 等于各个点电荷单独存在时在该点处的电势之代数和。

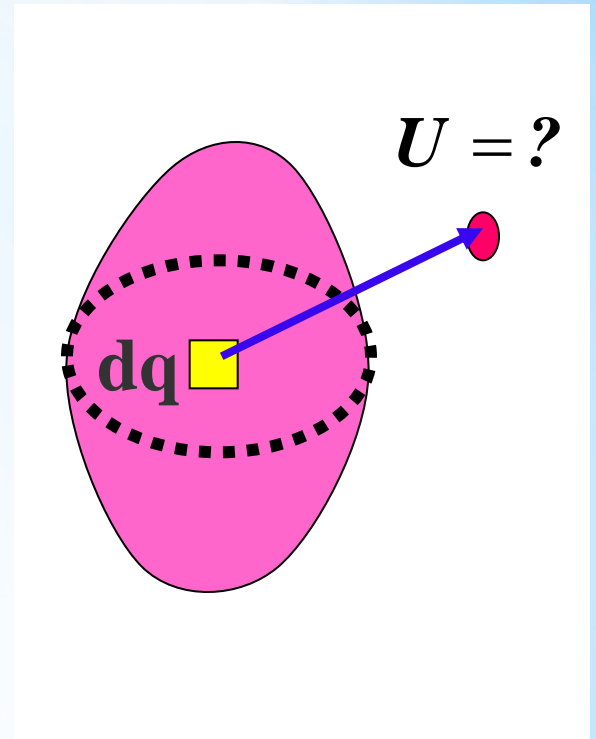
3) 连续分布电荷电场中的电势:

$$V = \int_V dV = \int_V \frac{dq}{4\pi\epsilon_0 r}$$

$$= \int_V \frac{\rho dV}{4\pi\epsilon_0 r} \quad \text{体电荷}$$

$$= \int_S \frac{\sigma dS}{4\pi\epsilon_0 r} \quad \text{面电荷}$$

$$= \int_\ell \frac{\lambda d\ell}{4\pi\epsilon_0 r} \quad \text{线电荷}$$



特点：标量积分！！！！

In summary:

The two methods calculating the electrical potential:

法一. 由电势定义式 $V_p = \int_p^{\text{零点}} \vec{E} \cdot d\vec{l}$ 计算P点电势。

法二. 根据点电荷的电势公式和电势迭加原理求电势。

例1. 半径为 R 的均匀带电球体, 带电量为 q 。求电势分布。

解: $\oiint_s \vec{E} \cdot d\vec{s} = \frac{1}{\epsilon_0} \sum q_i \begin{cases} \vec{E}_{\text{内}} = \frac{qr}{4\pi\epsilon_0 R^3} \hat{r} & (r < R) \\ \vec{E}_{\text{外}} = \frac{q}{4\pi\epsilon_0 r^2} \hat{r} & (r > R) \end{cases}$

$(r \geq R)$

$$V_{\text{外}} = \int_r^{\infty} \vec{E}_{\text{外}} \cdot d\vec{r}$$
$$= \int_r^{\infty} E_{\text{外}} dr = \int_r^{\infty} \frac{q}{4\pi\epsilon_0 r^2} dr = \frac{q}{4\pi\epsilon_0 r}$$

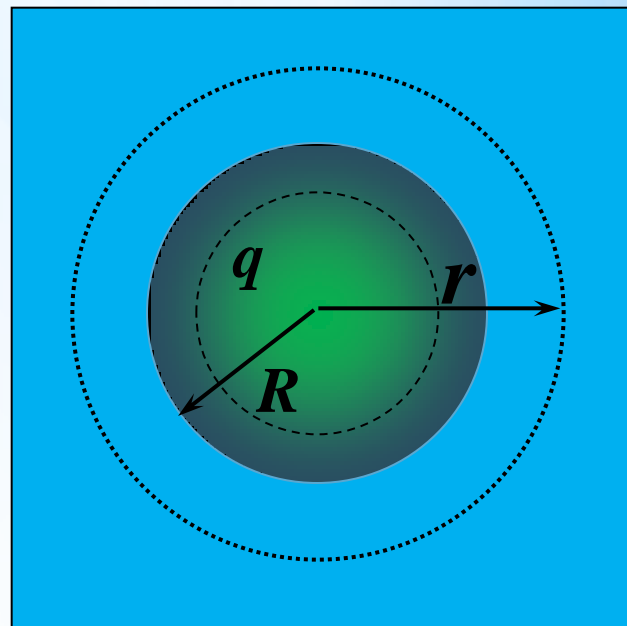
$(r < R)$

$$V_{\text{内}} = \int_r^{\infty} \vec{E} \cdot d\vec{r}$$

$$= \int_r^R \vec{E}_{\text{内}} \cdot d\vec{r} + \int_R^{\infty} \vec{E}_{\text{外}} \cdot d\vec{r} = \int_r^R E_{\text{内}} dr + \int_R^{\infty} E_{\text{外}} dr$$

$$= \int_r^R \frac{qr}{4\pi\epsilon_0 R^3} dr + \int_R^{\infty} \frac{q}{4\pi\epsilon_0 r^2} dr$$

$$= \frac{q}{8\pi\epsilon_0 R^3} (R^2 - r^2) + \frac{q}{4\pi\epsilon_0 R} = \frac{q(3R^2 - r^2)}{8\pi\epsilon_0 R^3}$$



例2. 均匀带电圆环，带电量为 q ，半径为 a ，求轴线上任意一点的 P 电势。

解： $dq = \lambda dl = \frac{q}{2\pi a} dl$

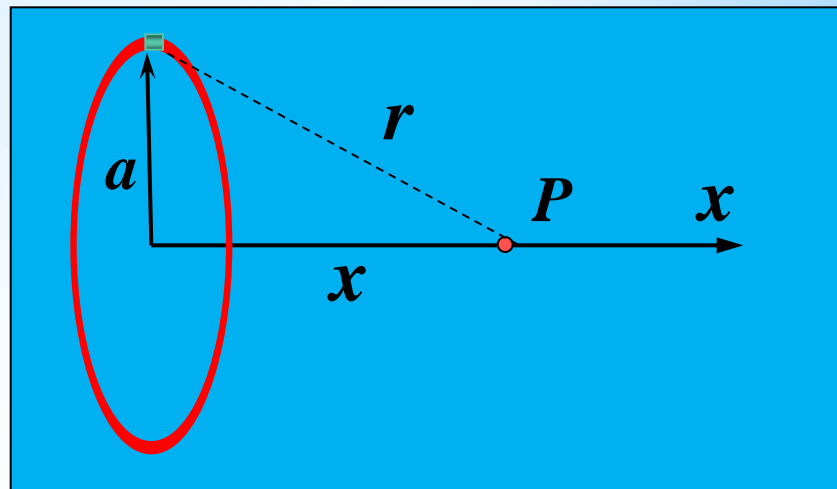
$$dV = \frac{dq}{4\pi\epsilon_0 r} = \frac{q dl}{8\pi^2 \epsilon_0 a r}$$

$$V = \int dV = \frac{q}{8\pi^2 \epsilon_0 a r} \int_L dl = \frac{q \cdot 2\pi a}{8\pi^2 \epsilon_0 a r}$$

$$\therefore V = \frac{q}{4\pi\epsilon_0 r} = \frac{q}{4\pi\epsilon_0 \sqrt{x^2 + a^2}}$$

法二： $E = \frac{1}{4\pi\epsilon_0} \frac{qx}{(x^2 + a^2)^{3/2}}$

$$V = \int_x^\infty \vec{E} \cdot d\vec{x} = \frac{q}{4\pi\epsilon_0} \int_x^\infty \frac{x dx}{(x^2 + a^2)^{3/2}} = \frac{q}{4\pi\epsilon_0 \sqrt{x^2 + a^2}}$$



例3. 求无限长均匀带电直线外任一点P的电势。
(电荷密度 λ)

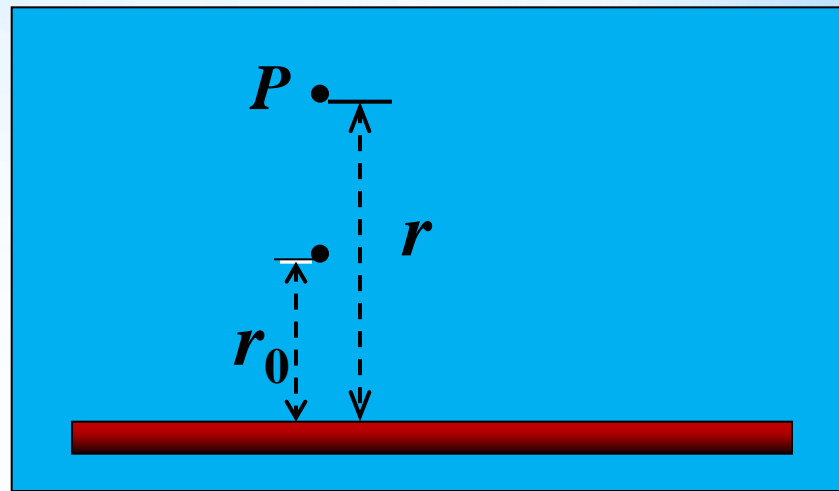
解: $E = \frac{\lambda}{2\pi\epsilon_0 r}$

$$V = \int_r^{r_0} \vec{E} \cdot d\vec{l} = \int_r^{r_0} \frac{\lambda}{2\pi\epsilon_0 r} dr$$

$$= \frac{\lambda}{2\pi\epsilon_0} \ln r \Big|_r^{r_0} = \frac{\lambda}{2\pi\epsilon_0} (\ln r_0 - \ln r) = \frac{\lambda}{2\pi\epsilon_0} \ln \frac{r_0}{r}$$

如果势能零点在 $r_0=1\text{m}$

$$V = \frac{-\lambda}{2\pi\epsilon_0} \ln r$$



例4:长为L的带电细杆，电荷线密度为 λ ，求其中垂线上距X轴 h 远处的电势。

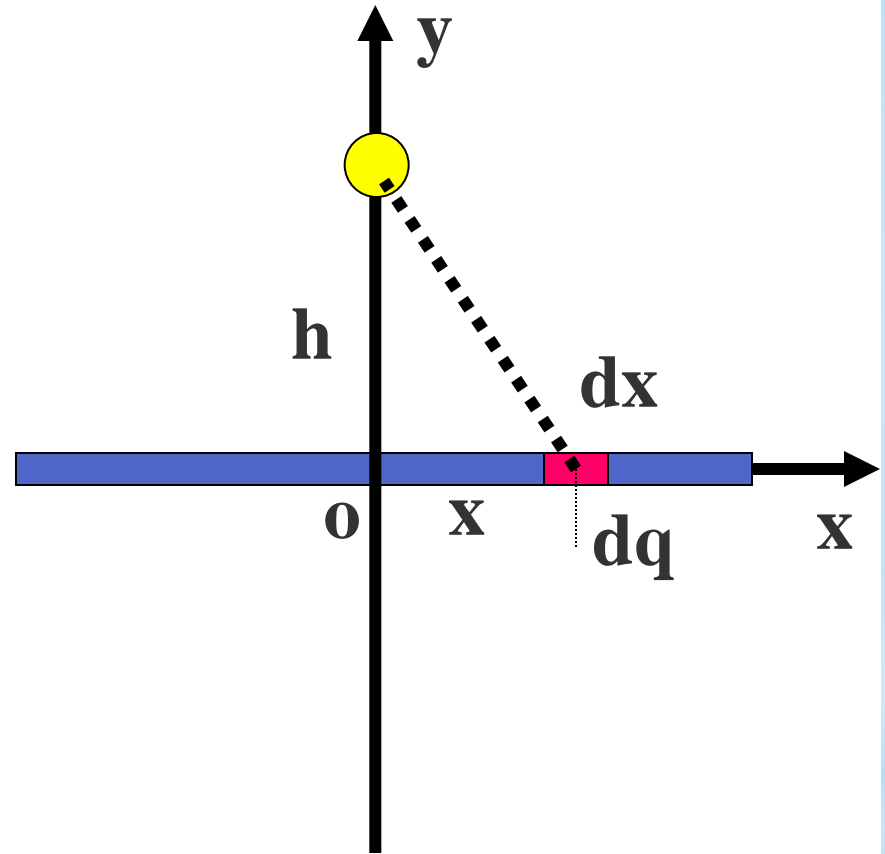
解:(1)选无限远处电势为零;

(2) 用电势叠加原理, 有

$$U = \int_{\ell} \frac{dq}{4\pi\epsilon_0 r}$$

$$= \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{\lambda dx}{4\pi\epsilon_0 (x^2 + h^2)^{1/2}}$$

$$= \frac{\lambda}{2\pi\epsilon_0} \ln \left(\frac{L}{2h} + \sqrt{1 + \frac{L^2}{4h^2}} \right)$$



例5: 一均匀电场, 电场强度 $\vec{E} = 4\vec{i} + 6\vec{j}$, 求XY平面内点a (3, 2) 和点b (1, 0) 间的电势差 (单位均取国际单位制)

解:

$$\Delta U = \int_a^b \vec{E} \cdot d\vec{\ell}$$

$$= \int_a^b E_x dx + E_y dy$$

$$= 4 \times (1 - 3) + 6 \times (0 - 2)$$

$$= -20 (V)$$

作业： 25, 26, 28