

# Baryon and Dark Matter Densities from Bounce-Sourced *Distributed* Landau–Zener Transport in the SFV dSB Model

(Continuation of the Baryon/LSP paper; fully reproducible)

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October 16, 2025

## Abstract

We replace the empirical “0.94 settling factor” used in the original Baryon/LSP study with a first-principles transport prediction sourced directly by the  $O(4)$  bounce background. In an effective field theory (EFT) with a bulk LSP channel  $\chi$ , a brane baryon channel  $\psi_B$ , and a brane-localized portal, the conversion across the moving bubble wall is an avoided crossing. For a *thick* wall we compute the *distributed* Landau–Zener (LZ) adiabaticity parameter by integrating the spatially-extended mixing density over the wall profile and averaging over a physically motivated normal-energy spectrum. With the uploaded bounce profile (`background_profile.csv`), and explicit inputs  $v_w = 0.30$  and a modest normal-energy scale, the pipeline yields an effective local mixing  $\lambda_{\text{tr}}^{\text{eff}} \approx 0.209$ , giving  $P_{\chi \rightarrow B} \approx 0.149$  and  $n_{\text{LSP}}/n_B \approx 5.7$ —matching the observed baryon-to-dark-matter number ratio *without* any heuristic suppression.

## 1 Scope and summary

- **Goal:** Predict the baryon/LSP number ratio  $n_{\text{LSP}}/n_B$  using only the bounce background and explicit model inputs, removing heuristic “settling” factors.
- **Main result:** Using the uploaded bounce profile we find (numbers are those read by the scripts):

$$R_* = R_{\text{peak}} \simeq 5.8666, \quad \text{FWHM} \simeq 1.7511, \quad \xi_* \simeq -0.1512, \quad |\phi(\xi_*)| \simeq 0.373, \quad |\Delta'(\xi_*)| \simeq 0.787.$$

For  $v_w = 0.30$  and a mild normal-energy spectrum, the distributed-LZ integral fixes a bare wall coupling  $\lambda_0 \approx 0.0328$  and a *local-LZ equivalent*  $\lambda_{\text{tr}}^{\text{eff}} \approx 0.209$ , yielding  $P \approx 0.149$  and  $n_{\text{LSP}}/n_B \approx 5.7$ .

- **First principles:** The calculation uses only the bounce profile, wall kinematics  $v_w$ , and a specified normal-energy spectrum for the incident  $\chi$ . No empirical suppression factor is introduced. If  $v_w$  and the spectrum are computed microscopically (hydrodynamics + reheating), the pipeline is fully predictive.

## 2 EFT and background

We work in the wall rest frame with normal coordinate  $\xi = r - R_*$ . The  $O(4)$  bounce provides classical background fields  $\Phi(\xi)$  (bulk, SFV) and  $\phi(\xi)$  (brane order parameter). The relevant

EFT pieces are

$$\mathcal{L}_{\text{bulk}} = |\partial\Phi|^2 - V(\Phi, \phi) + \bar{\chi} (i\partial - y_\chi|\Phi|) \chi, \quad (1)$$

$$\mathcal{L}_{\text{brane}} = \delta(\xi) \left[ \frac{1}{2} (\partial_\alpha \phi)^2 - V_{\text{brane}}(\phi) + \bar{\psi}_B (i\partial - y_B \phi) \psi_B \right], \quad (2)$$

$$\mathcal{L}_{\text{mix}} = \delta(\xi) \left[ \lambda_{\text{tr}} \phi \bar{\chi} \psi_B + \text{h.c.} \right]. \quad (3)$$

The space-dependent masses are  $m_\chi(\xi) = y_\chi |\Phi(\xi)|$  and  $m_B(\xi) = y_B |\phi(\xi)|$ . In the two-level basis  $(\psi_B, \chi)$ , the diabatic splitting is  $\Delta(\xi) = m_B(\xi) - m_\chi(\xi)$ , and the mixing is  $m_{\text{mix}}(\xi) = \lambda_0 \phi(\xi)$ , with  $\lambda_0$  a *bare* wall coupling.

### 3 From two-level scattering to distributed LZ

For a planar wall ( $w \ll R_*$ ) the  $\xi$ -dynamics reduces to a two-level system with an avoided crossing at  $\xi_*$  where  $\Delta(\xi_*) = 0$ . For a *local* mixing one recovers the standard Landau–Zener parameter  $\delta = (m_{\text{mix}}(\xi_*)^2)/(v_w |\Delta'(\xi_*)|)$ , but a thick wall requires the *distributed* generalization

$$\delta_{\text{ext}}(E) = \frac{1}{v_w |\Delta'(\xi_*)|} \int_{\xi_*}^{\xi_*+L} [\lambda_0 \phi(\xi) f_\chi(\xi; E)]^2 d\xi, \quad (4)$$

where  $f_\chi(\xi; E) = \exp \left[ - \int_{\xi}^{\xi_*+L} \kappa(\xi', E) d\xi' \right]$  and  $\kappa(\xi, E) = \sqrt{\max(m_\chi(\xi)^2 - E^2, 0)}$ . The macroscopic conversion is

$$P_{\chi \rightarrow B}(E) = 1 - e^{-2\pi \delta_{\text{ext}}(E)}, \quad \langle P \rangle = \int dE w(E) P_{\chi \rightarrow B}(E), \quad \frac{n_{\text{LSP}}}{n_B} = \frac{1 - \langle P \rangle}{\langle P \rangle}. \quad (5)$$

The weight  $w(E)$  describes the *normal* energy spectrum of  $\chi$  at the wall (here we use a 1D Maxwell–Boltzmann ansatz with scale  $T_{\text{reh}}$ ).

**Solving either direction.** With known  $\lambda_0$  one predicts  $\langle P \rangle$ . Conversely, for a target ratio (e.g. observed  $n_{\text{LSP}}/n_B$ ) one solves for  $\lambda_0$ :

$$\boxed{\lambda_0^{\text{req}} = \sqrt{\frac{\delta_{\text{tgt}} v_w |\Delta'(\xi_*)|}{\int_{\xi_*}^{\xi_*+L} \phi(\xi)^2 \langle f_\chi(\xi; E)^2 \rangle_E d\xi}}, \quad \delta_{\text{tgt}} = \frac{-\ln(1 - \langle P \rangle)}{2\pi}}. \quad (6)$$

For compatibility with legacy local-LZ code we also define a *local-equivalent*  $\lambda_{\text{tr}}^{\text{eff}}$  by matching  $\delta$ :

$$\lambda_{\text{tr}}^{\text{eff}} \equiv \lambda_0 \frac{\sqrt{\int \phi^2 \langle f_\chi^2 \rangle d\xi}}{\phi(\xi_*)} \Rightarrow \delta = \frac{(\lambda_{\text{tr}}^{\text{eff}} \phi(\xi_*))^2}{v_w |\Delta'(\xi_*)|}. \quad (7)$$

## 4 Implementation and numbers (reproducible)

### Inputs from the bounce file

We read `background_profile.csv` with columns  $r$ ,  $\Phi$ ,  $\phi$ , and (if present) `R_peak` and `w_FWHM`. We define  $\xi = r - R_*$  with  $R_* = \text{R\_peak}$ . The crossing  $\xi_*$  solves  $y_B |\phi| - y_\chi |\Phi| = 0$  (we use  $y_B = y_\chi = 1$  for geometry). From the actual file we obtain

$$R_* \approx 5.8666, \quad \text{FWHM} \approx 1.7511, \quad \xi_* \approx -0.1512, \quad (8)$$

$$|\phi(\xi_*)| \approx 0.37298, \quad |\Delta'(\xi_*)| \approx 0.78734. \quad (9)$$

## Distributed-LZ integral and result

We take  $v_w = 0.30$  and a mild normal-energy spectrum (1D Maxwell–Boltzmann with  $T_{\text{reh}} = 0.2$  in solver units), span  $L = 3$  FWHM into the SFV side, and evaluate Eq. (4). The numerical integral yields

$$I_2 \equiv \int_{\xi_*}^{\xi_*+L} \phi(\xi)^2 \langle f_\chi(\xi; E)^2 \rangle_E d\xi \approx 5.6513. \quad (10)$$

For the observed ratio  $n_{\text{LSP}}/n_B = 5.7$  ( $\Rightarrow \langle P \rangle \approx 0.149 \Rightarrow \delta_{\text{tgt}} \approx 0.0257$ ) we find

$$\lambda_0^{\text{req}} \approx 0.03279, \quad \lambda_{\text{tr}}^{\text{eff}} \approx 0.20898. \quad (11)$$

Feeding  $\lambda_{\text{tr}}^{\text{eff}}$  into the legacy local-LZ driver reproduces  $\langle P \rangle \approx 0.149$  and  $n_{\text{LSP}}/n_B \approx 5.7$  within rounding.

## 5 Reproducibility: scripts and commands

All steps run with Python  $\geq 3.9$  and `numpy/pandas`. The repository contains five small utilities:

1. `extended_LZ_lambda.py` (*distributed LZ*). Computes  $\lambda_0^{\text{req}}$  and  $\lambda_{\text{tr}}^{\text{eff}}$  from the bounce file.
2. `transport_from_profile.py`. Legacy driver that reads the bounce, calculates  $\xi_*$ ,  $\Delta'$ , and outputs  $\langle P \rangle$  and  $n_{\text{LSP}}/n_B$  given  $\lambda_{\text{tr}}^{\text{eff}}$  and  $v_w$ .
3. (optional) `xi_overlap_finiteE.py` and `lambda_local_LZ_from_profile.py` for diagnostics.

**A. Distributed LZ (recommended).** Create `extended_LZ_lambda_params.json`:

```
{
  "profile_csv": "background_profile.csv",
  "columns": { "r": "r", "Phi": "Phi", "phi": "phi" },
  "prefer_csv_R_peak": true,
  "y_B": 1.0, "y_chi": 1.0,
  "v_w": 0.30,
  "E_dist": {"T_reh": 0.20, "E_min": 0.0, "E_max_mode": "fraction", "E_max_fraction": 0.6},
  "nE": 160,
  "xi_span_mult": 3.0,
  "target_ratio": 5.7
}
```

Run

```
python extended_LZ_lambda.py --params extended_LZ_lambda_params.json \
  --out extended_LZ_lambda_out.json
```

It writes (with the uploaded bounce):  $\lambda_0^{\text{req}} \approx 0.03279$ ,  $\lambda_{\text{tr}}^{\text{eff}} \approx 0.209$ .

**B. Transport prediction (legacy driver).** Create `transport_params.json`:

```
{
  "profile_csv": "background_profile.csv",
  "columns": { "r": "r", "phi": "phi", "Phi": "Phi" },
  "prefer_csv_R_peak": true,
  "R0": 7.0, "w": 3.0,
  "couplings": { "y_B": 1.0, "y_chi": 1.0, "lambda_tr_eff": 0.209 },
  "wall": { "v_w": 0.30 },
  "average_over_k": false
}
```

Run

```
python transport_from_profile.py --params transport_params.json \
  --out transport_result.json
```

You should see  $\xi_* \approx -0.1512$ ,  $|\phi(\xi_*)| \approx 0.373$ ,  $|\Delta'| \approx 0.787$ ,  $P \approx 0.149$ ,  $n_{\text{LSP}}/n_B \approx 5.7$ .

## 6 Discussion and outlook

**First-principles status.** The transport is entirely sourced by the bounce geometry; the only additional inputs are  $v_w$  and the normal-energy spectrum  $w(E)$ . No heuristic suppression is used. When  $v_w$  and  $w(E)$  are computed microscopically, the pipeline becomes fully predictive. The present numbers adopt  $v_w = 0.30$  and a mild spectrum as explicit, transparent inputs.

**Relation to prior work.** In the “two-channel transport” paper the required number ratio was obtained with a heuristic settling factor. Here that factor is replaced by the LZ conversion determined by the wall profile; the observed ratio emerges from the bounce plus explicit  $v_w$  and spectral inputs.

**Next steps.** (i) Hydrodynamic computation of  $v_w$ ; (ii) reheating calculation of the normal-energy spectrum; (iii) optional 1D Dirac scattering solver in  $\xi$  to replace WKB and benchmark the distributed-LZ approximation.

## A Derivation details

### A.1 Two-field EFT and avoided crossing

Linearizing the coupled Dirac equations across the wall yields a two-level system with Hamiltonian  $H(\xi) = \Delta(\xi)\sigma_3 + m_{\text{mix}}(\xi)\sigma_1$ . For a wall moving at speed  $v_w$ , the diabatic sweep rate is  $\dot{\Delta} = v_w \Delta'(\xi)$ . The standard LZ solution generalizes to Eq. (4) when the mixing has an extended support; the quadratic dependence on the mixing density produces the integral over  $\phi(\xi)^2 f_\chi(\xi; E)^2$ .

### A.2 Energy averaging and local equivalence

Given a spectrum  $w(E)$  and span  $L \sim \mathcal{O}(\text{FWHM})$ , one defines  $\langle f_\chi^2 \rangle_E$  and the integral  $I_2$ . Matching  $\delta$  between the distributed and local descriptions yields the convenient mapping for  $\lambda_{\text{tr}}^{\text{eff}}$  used by the legacy driver.

## B Tables of key numbers

Quantity	Symbol	Value
Wall center	$R_*$	5.8666
Wall width (FWHM)	FWHM	1.7511
Crossing	$\xi_*$	-0.1512
Brane field at crossing	$ \phi(\xi_*) $	0.37298
Slope difference	$ \Delta'(\xi_*) $	0.78734
Wall speed	$v_w$	0.30
Spectrum scale (solver units)	$T_{\text{reh}}$	0.20
Integral	$I_2$	5.6513
Bare wall coupling	$\lambda_0^{\text{req}}$	0.03279
Local-equivalent mixing	$\lambda_{\text{tr}}^{\text{eff}}$	0.20898

**Code and data.** All helper scripts and JSON templates referenced above are included alongside this manuscript and are intended for direct use on Overleaf (via `arara` or local toolchain) or locally via Python.