

# Computational Electromagnetics

## FDTD (Finite Difference Time Domain)

Maxwell's Equations :-

Effect = Cause

- ①  $\nabla \cdot \vec{D}(t) = \rho_v(t)$  — Gauss' Law
- ②  $\nabla \cdot \vec{B}(t) = 0$  — No Magnetic Charge
- ③  $\nabla \times \vec{E}(t) = -\frac{\partial \vec{B}(t)}{\partial t}$  — Faraday's Law
- ④  $\nabla \times \vec{H}(t) = \vec{J}(t) + \frac{\partial \vec{D}(t)}{\partial t}$  — Ampere's Circuital Law

Current Continuity:-

$$\nabla \cdot \vec{J} = -\frac{\partial \rho_v}{\partial t}$$

Constitutive Relations:-

$$\vec{D}(t) = [\epsilon(t)] * \vec{E}(t)$$
 (Electric Response)

$$\vec{B}(t) = [\mu(t)] * \vec{H}(t)$$
 (Magnetic Response)

Lorentz Force Law:-

$$\vec{F} = q\vec{E} + q(\vec{v} \times \vec{B})$$

D: Electric flux density  
 B: Magnetic flux density  
 E: Electric field intensity  
 H: Magnetic field intensity

- ① Linear, isotropic, non-dispersive materials:-

$$D(t) = \epsilon \vec{E}(t)$$
 (linear relation)

- ② Dispersive materials:-

$$\vec{D}(t) = \epsilon(t) * \vec{E}(t)$$
 (convolution)

- ③ Anisotropic materials:-

$$D(t) = [\epsilon] \vec{E}(t)$$
 (tensor)

- ④ Non-linear materials:-

$$D(t) = \epsilon_0 \chi_e^{(1)} E(t) + \epsilon_0 \chi_e^{(2)} E^2(t) + \dots$$

Material refractive index :

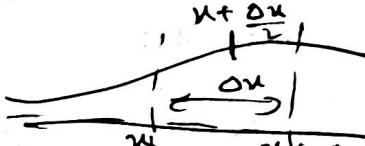
$$n = \sqrt{\mu_r \epsilon_r}$$

Material impedance :

$$\eta \approx \frac{|E|}{|H|}$$

$$\eta = \eta_0 \sqrt{\frac{\mu_r}{\epsilon_r}}$$

Finite Difference Approximations :-

$$\boxed{\frac{df_{1.s}}{du} \approx \frac{f_2 - f_1}{\Delta u}} \quad (\text{first order derivative})$$


- Each term in a finite-difference equation must exist at the same point in time and space.

$$\frac{df(u)}{du} + f(u) = 0$$

$$\rightarrow \underbrace{\frac{f(u + \Delta u) - f(u)}{\Delta u}}_{\text{exists at } u + \frac{\Delta u}{2}} + \underbrace{f(u)}_{\text{exists at } u} = 0$$

↳ Simulation Unstable

Solution :-

$$\frac{f(u + \Delta u) - f(u)}{\Delta u} + f\left(u + \frac{\Delta u}{2}\right) = 0$$

$$\frac{f(u + \Delta u) - f(u)}{\Delta u} + \frac{f(u + \Delta u) + f(u)}{2} = 0$$

— (Averaging)  
/Interpolation)

## Approximation of Time Derivatives in Maxwell's eqns:-

$$\nabla \times \vec{E}(t) = -\mu \frac{\partial \vec{H}(t)}{\partial t}$$

$$\Rightarrow \nabla \times \vec{E}(t) = -\mu \frac{\vec{H}(t+\Delta t) - \vec{H}(t)}{\Delta t} \hookrightarrow \text{Unstable}$$

$$\nabla \times \vec{H}(t) = \epsilon \frac{\vec{E}(t+\Delta t) - \vec{E}(t)}{\Delta t} \hookrightarrow \text{Unstable}$$

Solution :- (Staggering  $\vec{E}$  and  $\vec{H}$  in time)

$$\nabla \times \vec{E}(t) = -\mu \frac{\vec{H}(t+\Delta t/2) - \vec{H}(t-\Delta t/2)}{\Delta t}$$

$$\nabla \times \vec{H}(t) \rightarrow \epsilon \vec{E}$$

$$\nabla \times \vec{H}(t+\Delta t/2) = \epsilon \frac{\vec{E}(t+\Delta t) - \vec{E}(t)}{\Delta t}$$

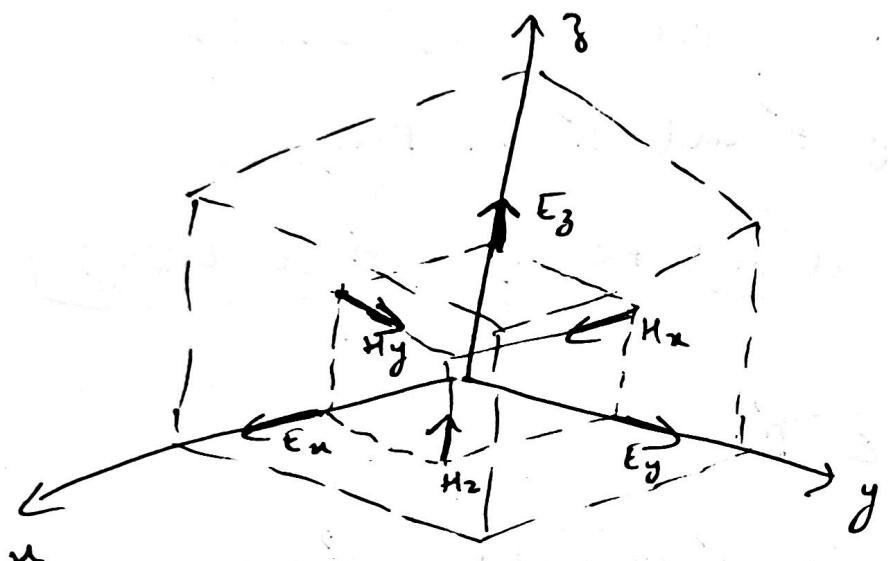
Update Equations :-

$$\vec{E}|_{t+\Delta t} = \vec{E}|_t + \frac{\Delta t}{\epsilon} (\nabla \times \vec{H}|_{t+\frac{\Delta t}{2}})$$

$$\vec{H}|_{t+\frac{\Delta t}{2}} = \vec{H}|_{t-\frac{\Delta t}{2}} - \frac{\Delta t}{\mu} (\nabla \times \vec{E}|_t)$$

## Yee Grid Scheme

- Assigning a function value to a specific point within the grid.
- Stagger the position of each field component within the grid cells.



results in divergence free grid

$$\nabla \cdot (\epsilon \vec{E}) = 0$$

$$\nabla \cdot (\mu \vec{H}) = 0$$

Maxwell's curl equations are approximated.

### Normalization:

$$\tilde{\vec{H}} = \sqrt{\frac{\mu_0}{\epsilon_0}} \vec{H}$$

$$\Rightarrow \nabla \times \vec{E} = - \frac{[\epsilon_r]}{c_0} \frac{\partial \tilde{\vec{H}}}{\partial t}$$

$$\nabla \times \tilde{\vec{H}} = \frac{[\epsilon_r]}{c_0} \frac{\partial \vec{E}}{\partial t}$$

Curl Eqs :-

(assume anisotropic materials and ignore off-diagonal tensor elements)

$$\nabla \times \vec{E} = - \frac{[\mu_{xx}]}{c_0} \frac{\partial \tilde{H}_x}{\partial t}$$

$$\Rightarrow \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = - \frac{\mu_{xx}}{c_0} \frac{\partial \tilde{H}_x}{\partial t}$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = - \frac{\mu_{yy}}{c_0} \frac{\partial \tilde{H}_y}{\partial t}$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = - \frac{\mu_{zz}}{c_0} \frac{\partial \tilde{H}_z}{\partial t}$$

Finite Difference Eqs:-

$$\textcircled{2} \quad \frac{E_z^{i,j+1,k}|_t - E_z^{i,j,k}|_t}{\Delta y} - \frac{E_y^{i,j,k+1}|_t - E_y^{i,j,k}|_t}{\Delta z}$$

$$= - \frac{\mu_{xx}}{c_0} \frac{\tilde{H}_x^{i,j,k}|_{t+\frac{\Delta t}{2}} - \tilde{H}_x^{i,j,k}|_{t-\frac{\Delta t}{2}}}{\Delta t}$$

$$\frac{\partial \tilde{H}_z}{\partial y} - \frac{\partial \tilde{H}_y}{\partial z} = \frac{E_{xx}}{c_0} \frac{\partial E_x}{\partial t}$$

$$\frac{\tilde{H}_z^{i,j,k}|_{t+\frac{\Delta t}{2}} - \tilde{H}_z^{i,j-1,k}|_{t+\frac{\Delta t}{2}}}{\Delta y} - \frac{\tilde{H}_y^{i,j,k}|_{t+\frac{\Delta t}{2}} - \tilde{H}_y^{i,j,k-1}|_{t+\frac{\Delta t}{2}}}{\Delta z}$$

$$= \frac{E_{xx}}{c_0} \frac{E_x^{i,j,k}|_{t+\Delta t} - E_x^{i,j,k}|_t}{\Delta t}$$

Reduction to 1-dimension :-

\* Cancel out  $\frac{\partial}{\partial u}$  and  $\frac{\partial}{\partial y}$  terms.

\* longitudinal terms  $\vec{E}_z$  and  $\vec{H}_z$  are 0.

\* ignore i and j indices  
(positions in u & y directions)

Eqn

$$\frac{E_y^{k+1}|_t - E_y^k|_t}{\Delta y} = -\frac{\mu_{uu}^k}{c_0} \frac{\tilde{H}_u^k|_{t+\frac{\Delta t}{2}} - \tilde{H}_u^k|_{t-\frac{\Delta t}{2}}}{\Delta t}$$

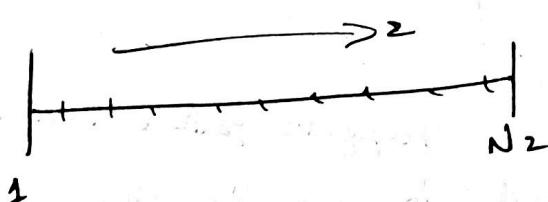
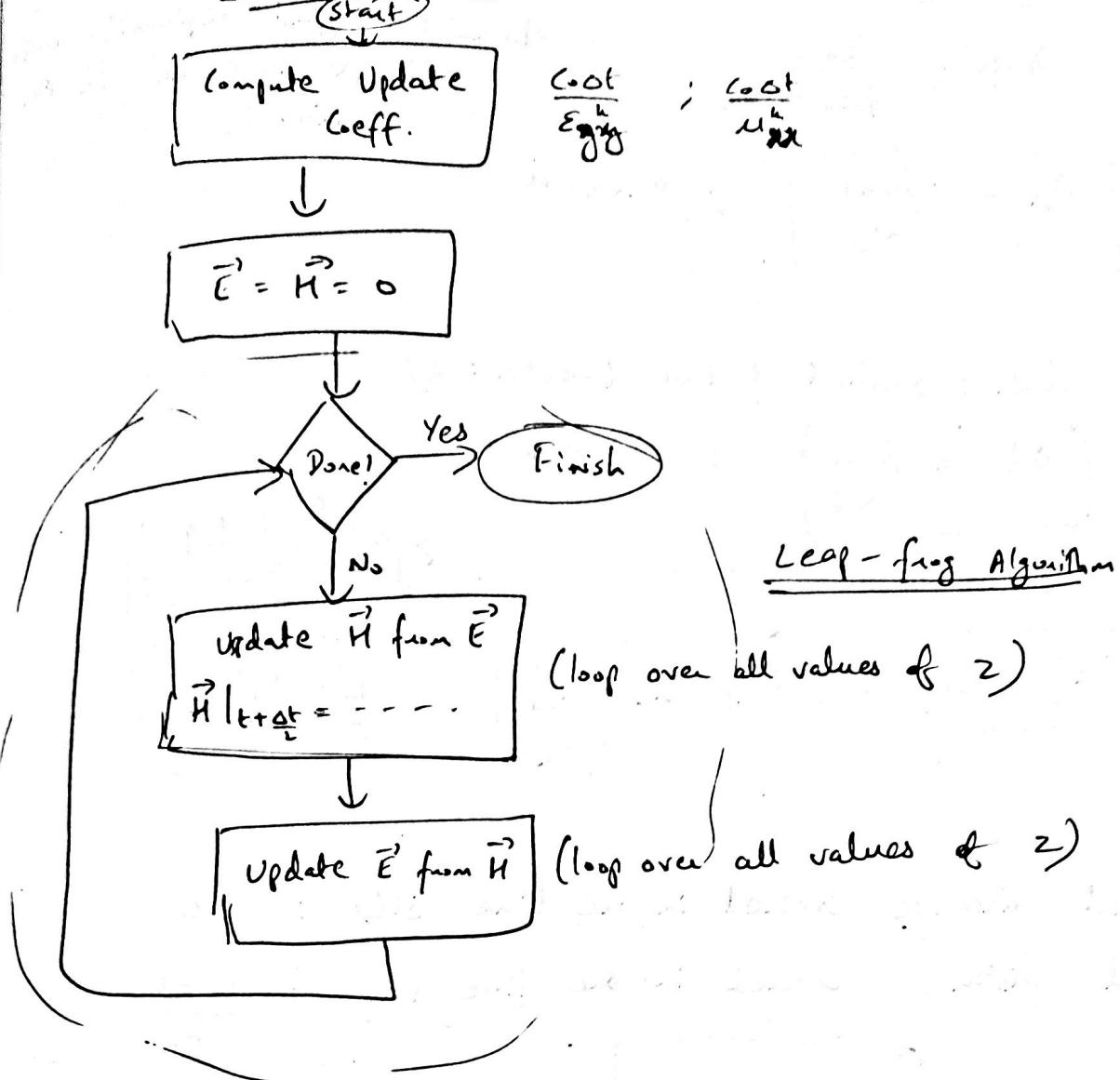
$$\frac{\tilde{H}_u^k|_{t+\frac{\Delta t}{2}} - \tilde{H}_u^{k-1}|_{t+\frac{\Delta t}{2}}}{\Delta y} = \frac{\epsilon_{yy}^k}{c_0} \frac{E_y^k|_{t+\Delta t} - E_y^k|_t}{\Delta t}$$

Update Equation :-

$$E_y^k|_{t+\Delta t} = E_y^k|_t + \left( \frac{c_0 \Delta t}{\epsilon_{yy}^k} \right) \left[ \frac{\tilde{H}_u^k|_{t+\frac{\Delta t}{2}} - \tilde{H}_u^{k-1}|_{t+\frac{\Delta t}{2}}}{\Delta y} \right]$$

$$\tilde{H}_u^k|_{t+\frac{\Delta t}{2}} = \tilde{H}_u^k|_{t-\frac{\Delta t}{2}} + \left( \frac{c_0 \Delta t}{\mu_{uu}^k} \right) \left[ \frac{E_y^{k+1}|_t - E_y^k|_t}{\Delta y} \right]$$

# FDTD Algo :



## Dirichlet Boundary Condition :-

Assume fields outside the grid (boundary) are 0.

$$\tilde{H}_x^k|_{t+\frac{\Delta t}{2}} = \tilde{H}_x^k|_{t-\frac{\Delta t}{2}} + M_{H_x}^k \left( \frac{\tilde{E}_y^{k+1}|_t - \tilde{E}_y^k|_t}{\Delta z} \right) ; k < \frac{N_2}{2}$$

$$\tilde{H}_x^k|_{t+\frac{\Delta t}{2}} = \tilde{H}_x^k|_{t-\frac{\Delta t}{2}} + M_{H_x}^k \left( 0 - \frac{\tilde{E}_y^{N_2}|_t}{\Delta z} \right) ; k = N_2$$

$$\tilde{E}_y^k|_{t+\frac{\Delta t}{2}} = \tilde{E}_y^k|_t + M_{E_y}^k \left( \frac{\tilde{H}_x^{k+1}|_t - \tilde{H}_x^k|_t}{\Delta z} \right) ; k = 1$$

### Grid resolution:-

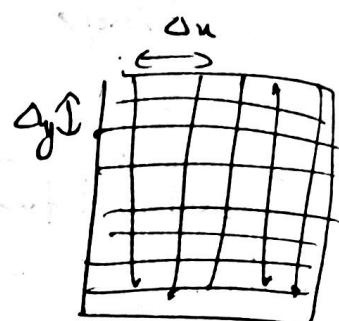
$$\lambda_{\min} = \frac{c_0}{\text{max } n_{\max}}$$

$n_{\max}$ : largest refractive index found anywhere in grid.

$$\boxed{\Delta z = \frac{\lambda_{\min}}{n_{\lambda}}} ; \quad N_{\lambda} \geq 10$$

$d_{\min}$ : smallest feature (mechanical)

$$\boxed{\delta d \approx \frac{d_{\min}}{N_d}} ; \quad N_d \geq 10$$



$$\boxed{\Delta u = \Delta y = \min [\Delta z, \Delta d]}$$

Numerical distance covered in one time step :  $\Delta z$

Physical distance covered in one time step :  $\frac{c_0 \Delta t}{n}$

$$\boxed{\frac{c_0 \Delta t}{n} < \Delta z}$$

↳ wave cannot propagate farther than a single unit cell. in one time step.

$$\Rightarrow \boxed{\Delta t < \frac{n \Delta z}{c_0}} \quad n \approx 1 \text{ (smallest refractive index)}$$

$$\rightarrow \boxed{t < \frac{\Delta z}{c_0}} \rightarrow \text{for 1D}$$

Constant stability condition:-

$$\boxed{\Delta t < \frac{1}{\frac{c_0}{(\Delta x)^2} + \frac{1}{(\Delta y)^2} + \frac{1}{(\Delta z)^2}}}$$

→ for 2D / 3D

$$\frac{c_0}{(\Delta x)^2} + \frac{1}{(\Delta y)^2} + \frac{1}{(\Delta z)^2}$$

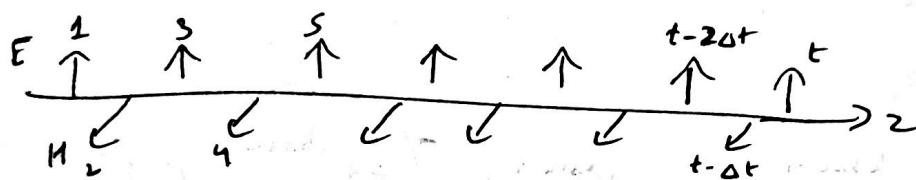
$$\Delta t = \frac{n_{bc} \Delta z}{2c_0}$$

$n_{bc}$  : boundary refractive index

0.5 factor for safety margin.

Wave travels the distance of one grid in exactly two time steps.

$$E_y^{Nz+1} |_t = E_y^{Nz} |_{t-2\Delta t}$$



Source : Short pulse  $g(t) \propto e^{-\left(\frac{t-t_0}{\tau}\right)^2}$

$\hookrightarrow$  Gaussian pulse

(1) max freq. =  $f_{max}$

$$B = f_{max} = \frac{1}{\pi \tau} \Rightarrow \tau \leq \frac{1}{\pi f_{max}}$$

$$\tau \approx \frac{0.5}{f_{max}}$$

$$\Delta t = \frac{\tau}{N_t}$$

;  $N_t \geq 10$

(3) minimum of  $\Delta t$  obtained by constant condition  
and max. freq. condition.

Delay for introducing source :- (No step function)

$$t_0 = 6\tau$$

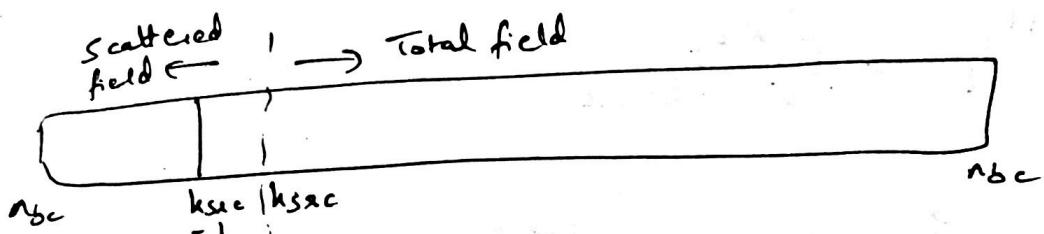


Soft Source:

$$\tilde{H}_u|_{t+\frac{\Delta t}{2}} = \tilde{H}_u|_{t-\frac{\Delta t}{2}} + g_{H_u}|_k$$

$$\tilde{E}_y|_{t+\frac{\Delta t}{2}} = \tilde{E}_y|_{t-\frac{\Delta t}{2}} + g_E|_k$$

### Total Field / Scattered Field



$$\tilde{H}_u|_{t+\frac{\Delta t}{2}}^{hsac-1} = \tilde{H}_u|_{t-\frac{\Delta t}{2}}^{hsac-1} + (m_{H_u})^{hsac-1} \left[ \frac{\tilde{E}_y|_t^{hsac} - \tilde{E}_y|_t^{hsac-1}}{\Delta z} \right]$$

total field term in scattered field region

### Correction:-

Subtract source  ~~$E_y$~~  from  $E_y^{hsac}$  to make it look like a scattered field.

$$\tilde{H}_u|_{t+\frac{\Delta t}{2}}^{hsac-1} = \tilde{H}_u|_{t-\frac{\Delta t}{2}}^{hsac-1} + (m_{H_u})^{hsac-1} \left[ \frac{\tilde{E}_y|_t^{hsac} - \tilde{E}_y|_t^{hsac-1}}{\Delta z} - \frac{(m_{H_u}) \tilde{E}_y|_t^{hsac}}{\Delta z} \right]$$

$$\Rightarrow \tilde{H}_u|_{t+\frac{\Delta t}{2}}^{hsac-1} = \tilde{H}_u|_{t-\frac{\Delta t}{2}}^{hsac-1} + (m_{H_u}) \left[ \frac{\tilde{E}_y|_t^{hsac} - \tilde{E}_y|_t^{hsac-1}}{\Delta z} \right] - \frac{(m_{H_u}) \tilde{E}_y|_t^{hsac}}{\Delta z}$$

Original update eqn.

Similarly,

$$\tilde{E}_y|_{t+\frac{\Delta t}{2}}^{hsac} = \tilde{E}_y|_t^{hsac} + (m_E) \left[ \frac{\tilde{H}_u|_{t+\frac{\Delta t}{2}}^{hsac} - \tilde{H}_u|_{t+\frac{\Delta t}{2}}^{hsac-1}}{\Delta z} \right]$$

Scattered field term in total field region.

Correction :-

$$E_y|_{t+\Delta t}^{ksec} = E_y|_t^{ksec} + (m_{E_y}) \left[ \frac{\tilde{H}_u|_{t+\frac{\Delta t}{2}}^{ksec} - \tilde{H}_u|_{t-\frac{\Delta t}{2}}^{ksec-1}}{\Delta \tau} \right] - \boxed{(m_{E_y}) \frac{\tilde{H}_u|_{t+\frac{\Delta t}{2}}^{ksec-1}}{\Delta \tau}}$$

original update eqn.

↓  
correction term

$$E_y = \sin(\omega t - \beta_j)$$

$$H_u = A \sin(\omega t - \beta_j - \phi)$$

$$A_2 = \sqrt{\frac{E_x^{sec}}{m_e^{sec}}}$$

$$\boxed{E_y|_t^{ksec} = g(t)}$$

$$\tilde{H}_u|_{t+\frac{\Delta t}{2}}^{ksec-1} = -\sqrt{\frac{E_x}{m_e}} g\left(t + \frac{ksec \Delta \tau}{2c_0} + \frac{\Delta t}{2}\right)$$

↓  
delay through  
half of grid cell  
(staggered)

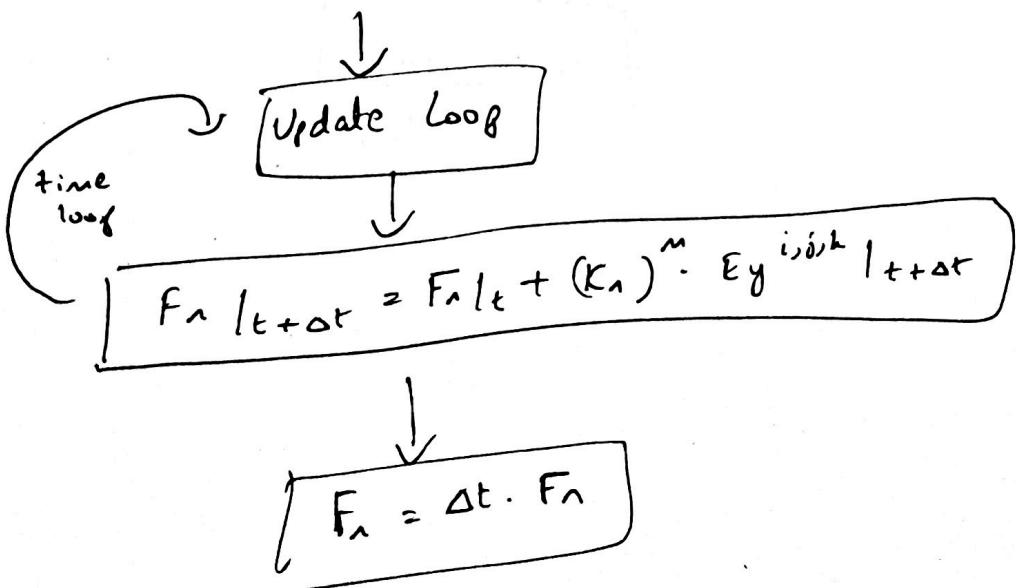
↳ staggered  
in time

Fourier Transform :-

$$F(f) \approx \Delta t \sum_{m=1}^M (e^{-j2\pi f \Delta t})^m \cdot f(m)$$

Compute Kernel (for each freq.) :-

$$K_n = e^{-j2\pi f_n \Delta t}$$



Adding Material to Grid :-

$$n_{\text{max}} = \sqrt{\mu_r \epsilon_r}$$

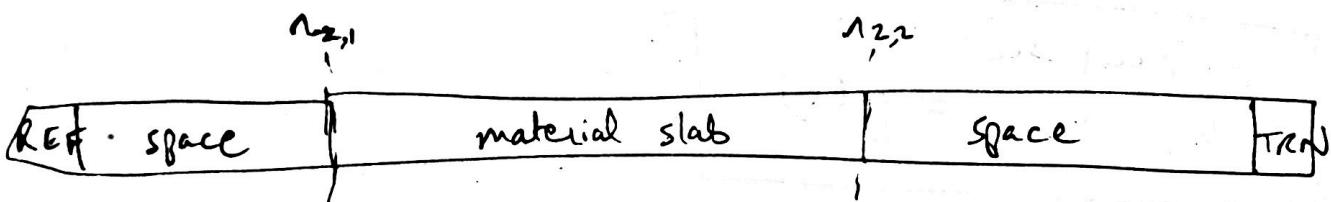
find  $\lambda_{\text{min}}$ .

$$\text{find } \Delta_\lambda = \frac{\lambda_{\text{min}}}{N_2}, \Delta_d = \frac{d}{N_d}$$

$$d_2 = \min(\Delta_\lambda, \Delta_d)$$

$$N' = \frac{d}{d_2}$$

$N = \text{ceil}(N')$  = Thickness (no. of cells) of the material.

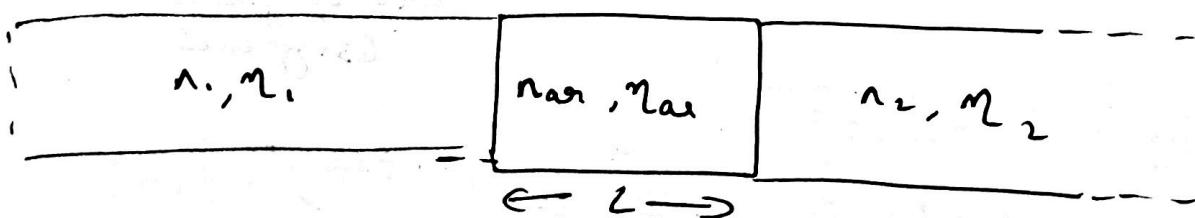


$$N_2 = N_{\text{material}} + 2 (\text{Space} + TRW/REF)$$

$$UR(n_{z,1}; n_{z,2}) = U_R;$$

$$ER(n_{z,1}; n_{z,2}) = E_R;$$

Anti-Reflection layer :-



$$m_{\text{air}} = \sqrt{n_1 n_2}$$

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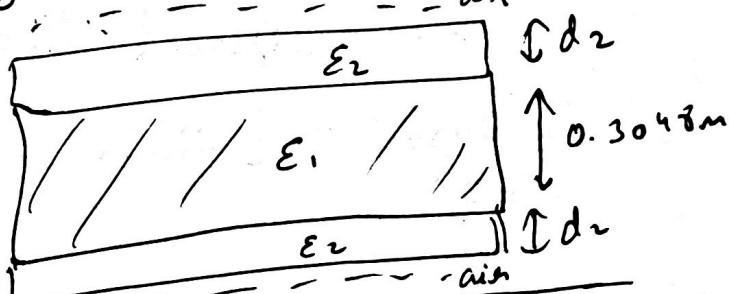
$$L = \frac{\lambda_0}{4 m_{\text{air}}}$$

problem :-

A radome is to be designed to protect an antenna operating at 2.4 GHz. For mechanical reasons, it must be constructed from 1 ft thick plastic with dielectric constant 12. Modify the design to minimize transmission through radome. Simulate the design using 1D FDTD.

$$1 \text{ feet} = 0.3048 \text{ m}$$

Using Anti-Reflection layer:-



$$\epsilon_2 = \sqrt{\epsilon_1 \cdot \epsilon_{\text{air}}} = \sqrt{12(1)} = 3.46$$

$$n_2 = \sqrt{\epsilon_2} = \sqrt{3.46} = 1.86$$

$$\begin{aligned} \lambda_0 &= \frac{c_0}{f_0} = \frac{3 \times 10^8}{2.4 \times 10^9} \\ &= 12.49 \text{ cm} \\ &= 0.1249 \text{ m} \end{aligned}$$

$$d_2 = \frac{\lambda_0}{4n_2} = \frac{0.1249}{4(1.86)} = 0.0167 \text{ m}$$