

Spectral flow demonstrated with N unit cells - Changing K1 and K2 simultaneously; adding a K2 spring at the end of the system (needed to get edge mode)

```
restart;
with(LinearAlgebra) : with(Physics) : with(Statistics) : with(plots) :
```

Nominal values for mass m, stiffnesses K1 and K2, and the number of unit cells (comes from paper with Nehemiah and Sai)

```
m := 0.1; k := 6000; N := 20;
```

$$\begin{aligned} m &:= 0.1 \\ k &:= 6000 \\ N &:= 20 \end{aligned} \tag{1.1}$$

Set this to 1 if the final K2 spring to ground should be added at the end of the chain. Otherwise, set to 0. The system of equations are only chiral when edge spring is present (see Theocharis's paper, Phys Rev. B, 024106, 2021) for acoustic system. Without the edge spring, we do not get the topological mode (unless use quantum system, at which point diagonalModifier should be set to zero)

```
endSpring := 1;
```

$$endSpring := 1 \tag{1.2}$$

Note: if I set diagonalModifier to zero, I get the QM tight-binding system where the diagonals of the "stiffness" matrix are zero. Otherwise, set to 1.

```
diagonalModifier := 1;
```

$$diagonalModifier := 1 \tag{1.3}$$

```
M := Matrix(2·N, 2·N, [ ]) : K := Matrix(2·N, 2·N, [ ]) :
Mass matrix for system
for i from 1 to 2·N do
    for j from 1 to 2·N do
        M(i, j) := m·KroneckerDelta[i, j];
    end do;
end do;
```

This helps in the stiffness rows to assign either K1 or K2 to the entry to the left/right of the diagonal depending on whether the row is even numbered or odd numbered

```
mod(1, 2);mod(2, 2);mod(3, 2);mod(4, 2);mod(5, 2);
```

$$\begin{array}{c} 1 \\ 0 \\ 1 \\ 0 \\ 1 \end{array} \tag{1.4}$$

Here, we let K1 and K2 be functions of lambda
 $K1 := k \cdot (1 - \lambda)$; $K2 := k \cdot (1 + \lambda)$;

$$K1 := 6000 - 6000 \lambda$$

$$K2 := 6000 + 6000 \lambda \quad (1.5)$$

Stiffness matrix for system.

```

for i from 1 to 2·N do
  if (i ≠ 1 and i ≠ 2·N) then K(i, i) := diagonalModifier·(K1 + K2); K(i, i - 1) := -K1·
    mod(i - 1, 2) - K2·mod(i, 2); K(i, i + 1) := -K1·mod(i, 2) - K2·mod(i + 1, 2); end if;
  if (i = 1) then K(1, 1) := diagonalModifier·K1; K(1, 2) := -K1; end if;
  if (i = 2·N) then K(2·N, 2·N) := diagonalModifier·(K1 + endSpring·K2); K(2·N, 2·N
    - 1) := -K1; end if;
end do;
```

The eigenvalues from the generalized EVP, $K^*x=\lambda M^*x$, are found from $\omega^2=\lambda$. Note that the usage of lambda here is different from the spectral flow lambda used above

$$\begin{aligned}
eigs &:= x \rightarrow \text{Eigenvalues}(\text{subs}(\lambda=x, K), M); \\
eigs &:= x \mapsto \text{Eigenvalues}(\text{subs}(\lambda=x, K), M)
\end{aligned} \quad (1.6)$$

#eigs(0.1);

Create a sequence for various values of lambda from -1 to 1. Do this by specifying the number of lambda values to use.

numLambda := 300;

$$\text{numLambda} := 300 \quad (1.7)$$

$$\begin{aligned}
data &:= \text{seq}\left(\left[\left(\frac{2 \cdot x}{\text{numLambda}} - 1\right), \text{eigs}\left(\frac{2 \cdot x}{\text{numLambda}} - 1\right)\right], x=0..\text{numLambda}\right) : \\
&\quad \text{data[1]; data[1][1]; data[1][2];} \\
&\quad \left[\begin{array}{c} 1..40 \text{ Vector}_{\text{column}} \\ \text{Data Type: complex}_8 \\ \text{Storage: rectangular} \\ \text{Order: Fortran_order} \end{array} \right] \\
&\quad -1 \\
&\quad \left[\begin{array}{c} 1..40 \text{ Vector}_{\text{column}} \\ \text{Data Type: complex}_8 \\ \text{Storage: rectangular} \\ \text{Order: Fortran_order} \end{array} \right]
\end{aligned} \quad (1.8)$$

$$\begin{aligned}
\text{lambdaVals} &:= j \rightarrow \text{data}[j][1]; \\
\text{lambdaVals} &:= j \rightarrow \text{data}_{j_1}
\end{aligned} \quad (1.9)$$

$\text{lambdaVals}(2);$

$$-\frac{149}{150} \quad (1.10)$$

$\text{eigVals} := j \rightarrow \text{data}[j][2];$

$$\text{eigVals} := j \rightarrow \text{data}_{j_2} \quad (1.11)$$

eigVals(1);

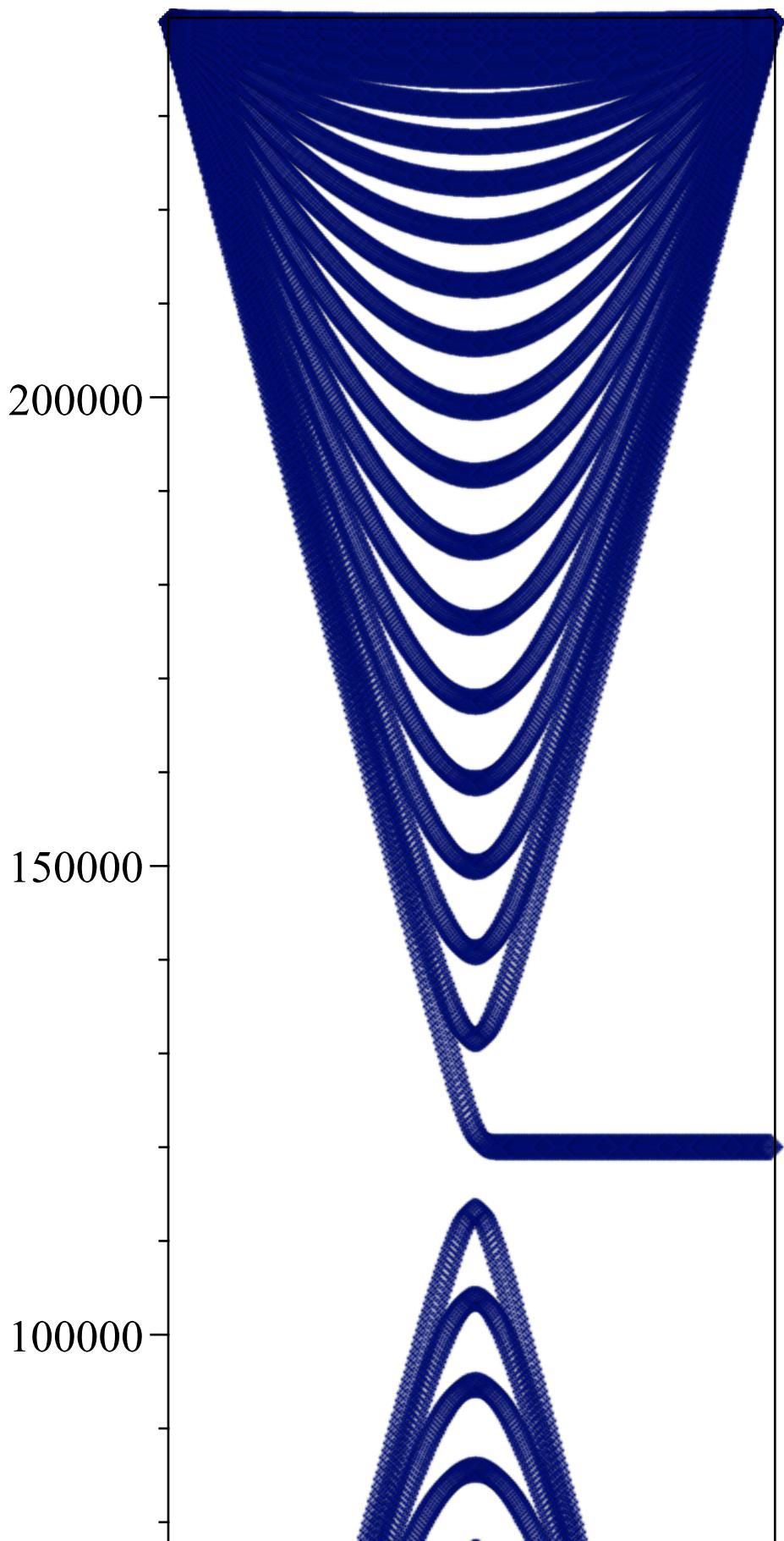
<i>I .. 40 Vector_{column}</i>
<i>Data Type: complex₈</i>
<i>Storage: rectangular</i>
<i>Order: Fortran_order</i>

(1.12)

I plot either the eigenvalue or the sqrt(eigenvalue). When sqrt, we recover omega.

```
#ScatterPlot( <seq(seq(lambdaVals(l),j=1..2·N),l=1..numLambda+1)>, <seq(sqrt
~(eigVals(j)),j=1..numLambda+1)>, labels = [ 'lambda','ω' ] );
ScatterPlot( <seq(seq(lambdaVals(l),j=1..2·N),l=1..numLambda+1)>, <seq(eigVals(j),j=1
..numLambda+1)>, labels = [ 'lambda','ω ω' ] );
```

$\psi^*(\omega, \omega)$



Have to be careful to take Transpose at times since eigenvectors are columns (see further below)
 $eigs := x \rightarrow Eigenvectors(subs(\lambda=x, K), M);$

$eigs := x \mapsto Eigenvectors(subs(\lambda=x, K), M)$ (1.13)
 $eigs(0);$

$$\begin{bmatrix} 1 .. 40 \text{ Vector}_{\text{column}} \\ \text{Data Type: complex}_8 \\ \text{Storage: rectangular} \\ \text{Order: Fortran_order} \end{bmatrix}, \begin{bmatrix} 40 \times 40 \text{ Matrix} \\ \text{Data Type: complex}_8 \\ \text{Storage: rectangular} \\ \text{Order: Fortran_order} \end{bmatrix} \quad (1.14)$$

$eigs(1.0)[1][3];$

$$240000.000000000 + 0. \text{ I} \quad (1.15)$$

This is the mode on the right end of the chain

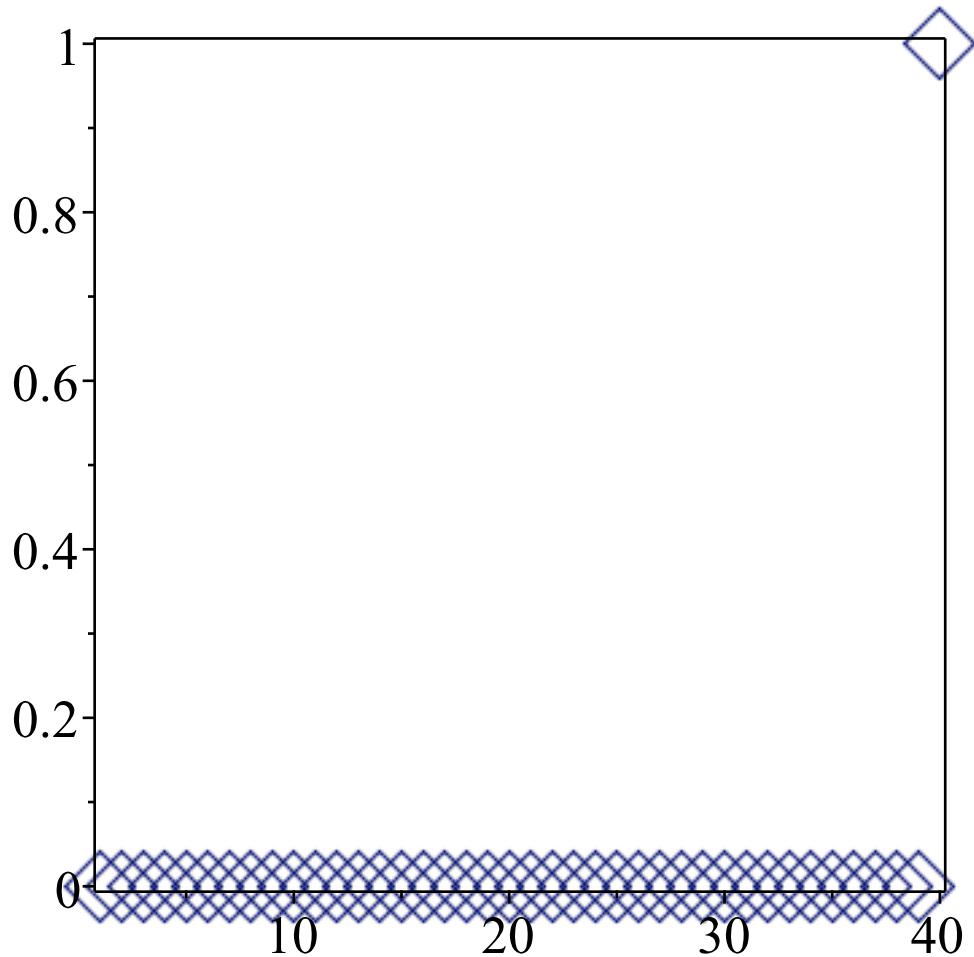
$eigs(1.0)[1][2 \cdot N]; \text{Transpose}(eigs(1.0)[2])[2 \cdot N];$

$$120000. + 0. \text{ I}$$

$$\begin{bmatrix} 1 .. 40 \text{ Vector}_{\text{row}} \\ \text{Data Type: complex}_8 \\ \text{Storage: rectangular} \\ \text{Order: Fortran_order} \end{bmatrix} \quad (1.16)$$

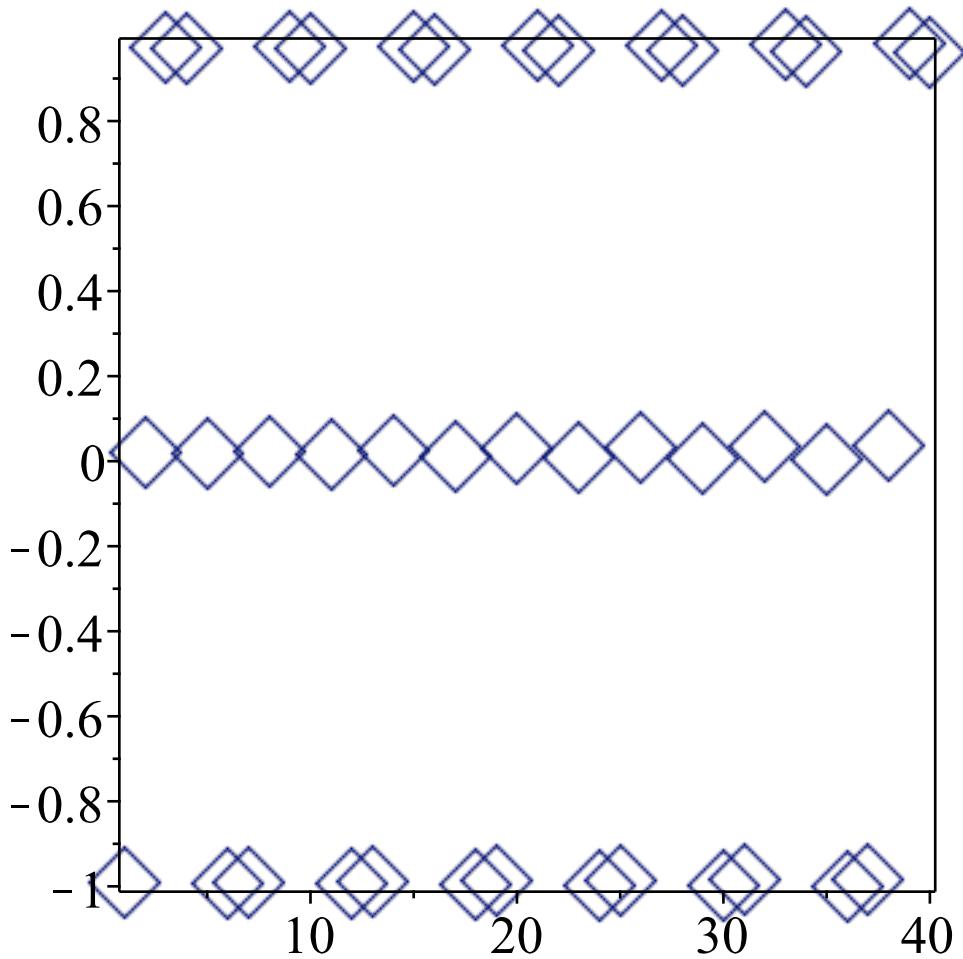
Note that the mode shape when $\lambda=1.0$ has a 1 at mass $2N$ and the rest of the masses have very little response

$\text{ScatterPlot}(\langle \text{seq}(j, j = 1 .. 2 \cdot N) \rangle, \text{Transpose}(eigs(1.0)[2])[2 \cdot N], \text{symbolsize} = 50);$



Note that the mode shape when lambda=0.02 the localized mode becomes distributed
eigs(0.02)[1][2·N]; ScatterPlot(seq(j,j=1..2·N)), Transpose(eigs(0.02)[2])[2·N],
symbolsize=50);

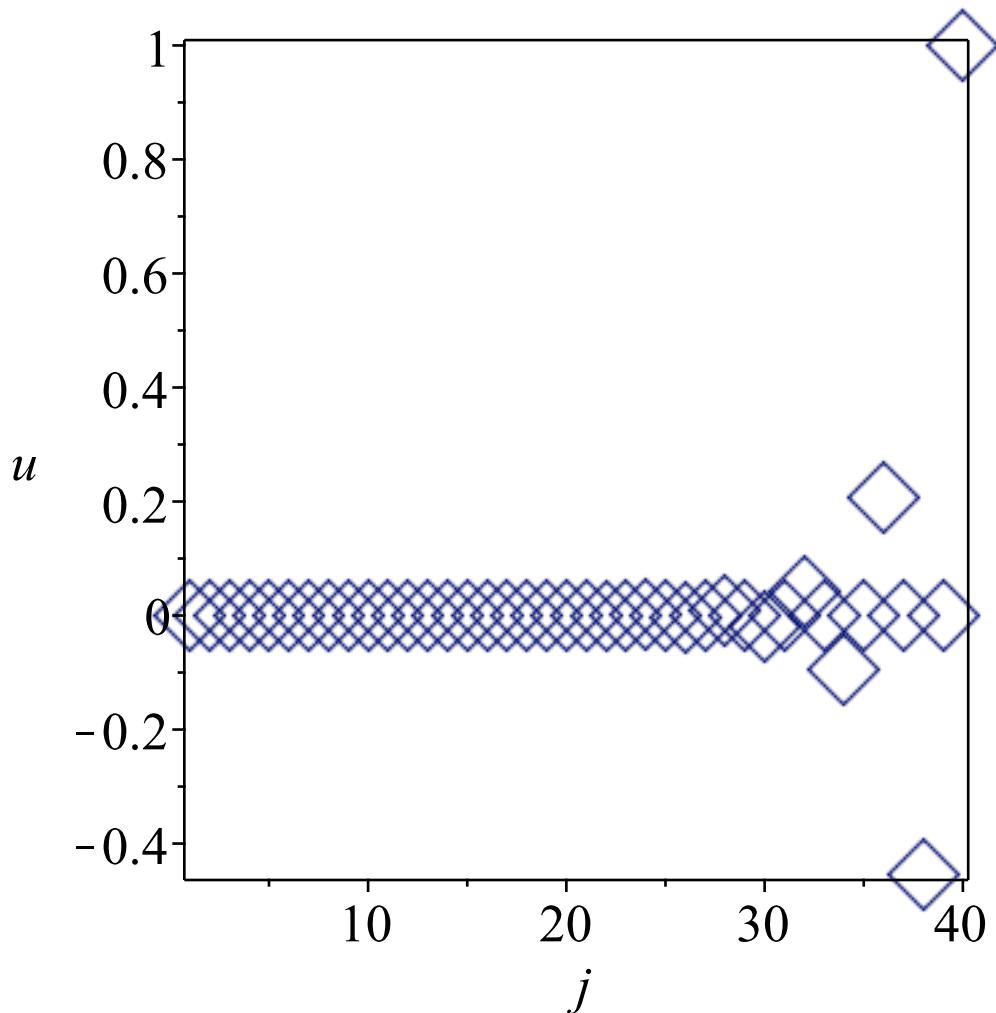
$$60008.4423987034 + 0. \mathrm{i}$$



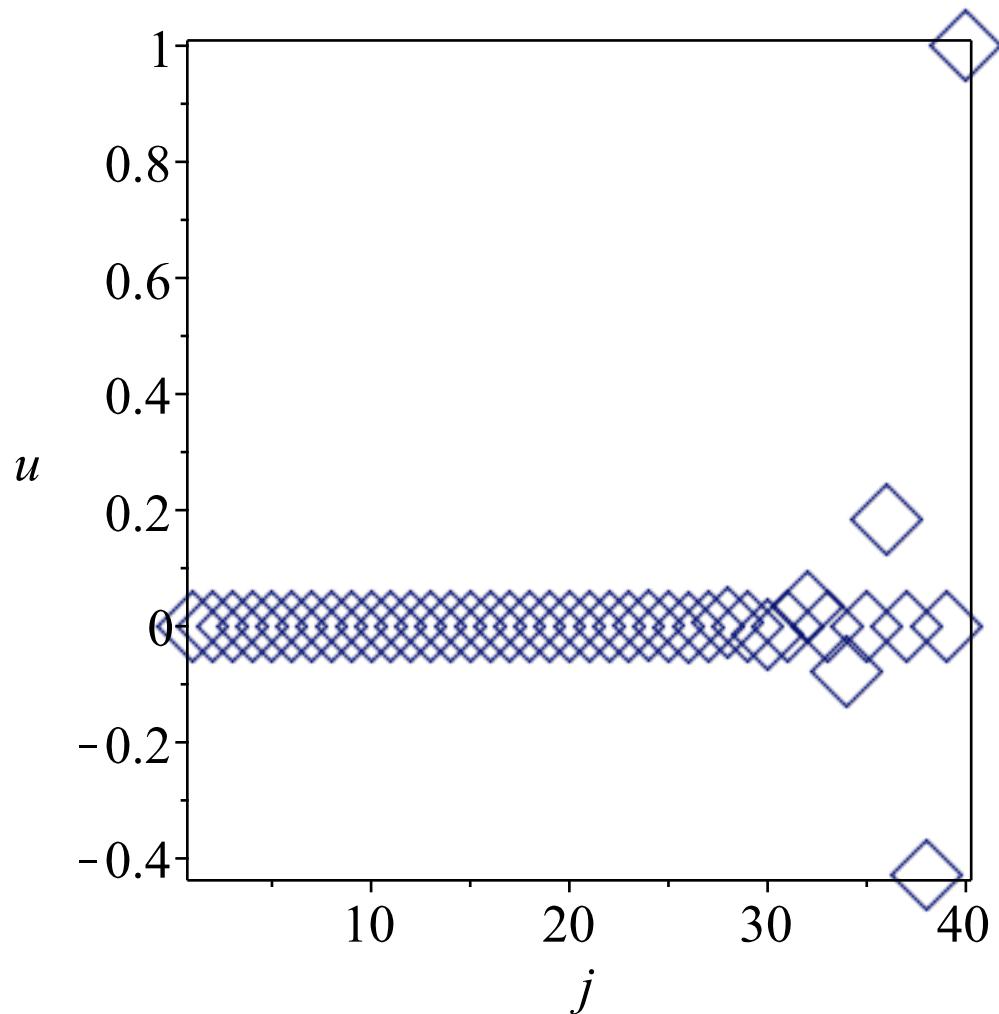
Here is the topological mode at lambda = 0.375, 0.4, 0.5, 0.6, ..., 1.0 (turns out to be the last eigenvector returned)

```
eigs( $\frac{3}{8}$ )[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs( $\frac{3}{8}$ )[2])[2·N],  
    symbolsize = 50, labels = ['j', 'u']);  
  
eigs(0.4)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(0.4)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);  
eigs(0.5)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(0.5)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);  
eigs(0.6)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(0.6)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);  
eigs(0.7)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(0.7)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);  
eigs(0.8)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(0.8)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);  
eigs(0.9)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(0.9)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);  
eigs(1.0)[1][2·N]; ScatterPlot(⟨seq(j, j = 1 .. 2·N)⟩, Transpose(eigs(1.0)[2])[2·N], symbolsize  
    = 50, labels = ['j', 'u']);
```

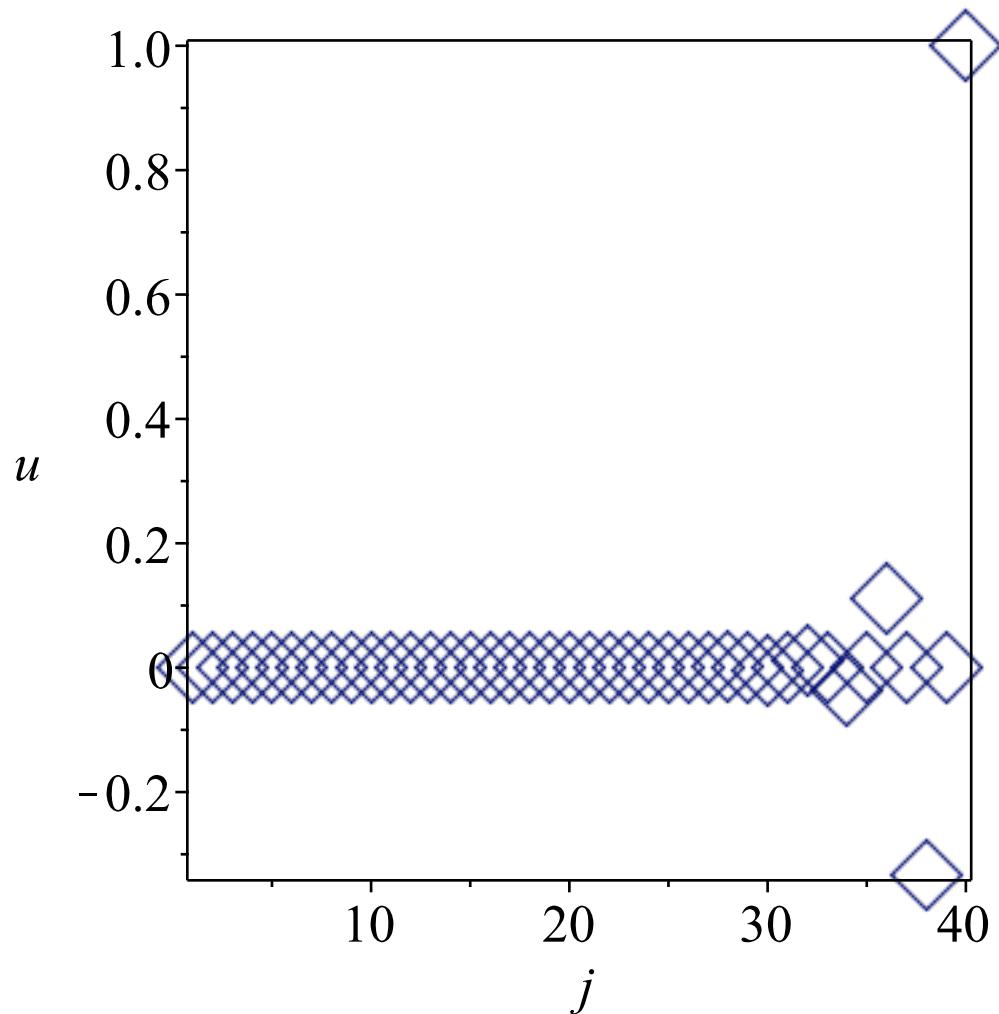
$120000.000000000 + 0. \mathrm{i}$



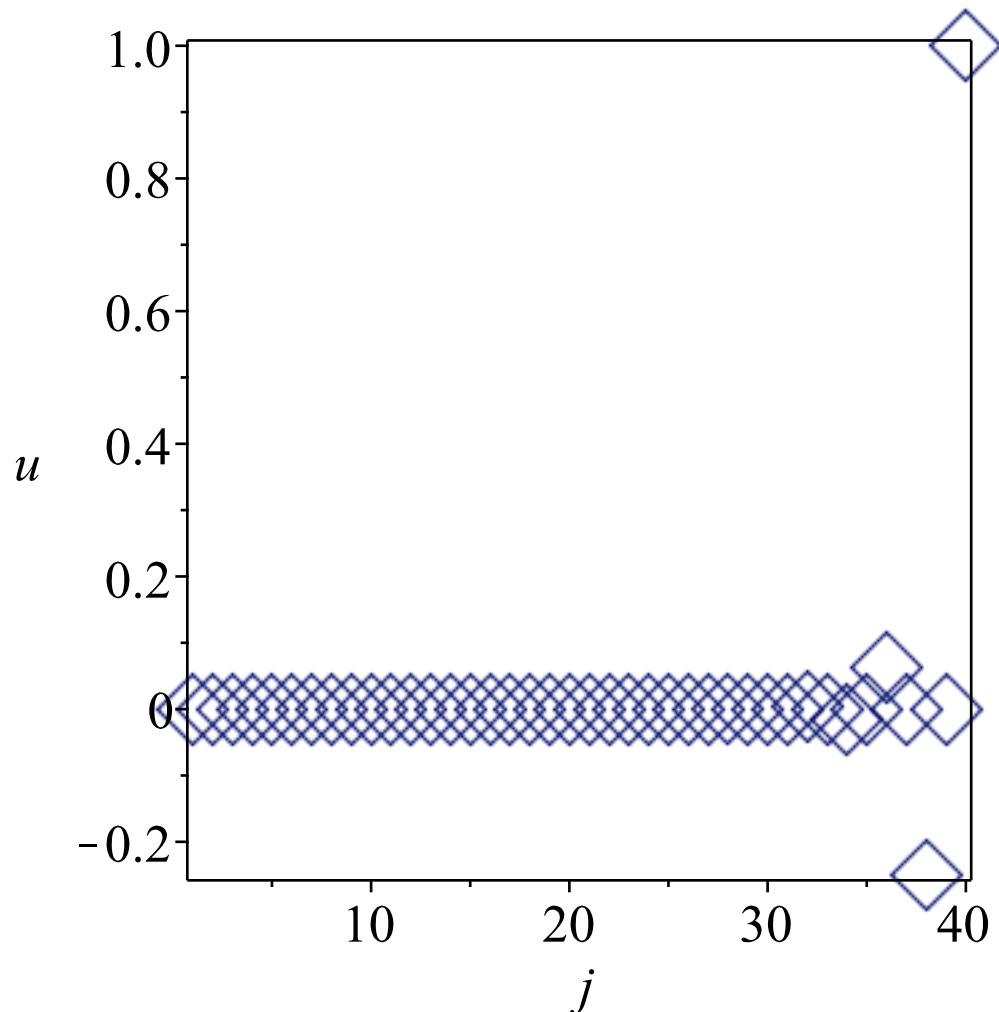
$120000.000000000 + 0. \mathrm{i}$



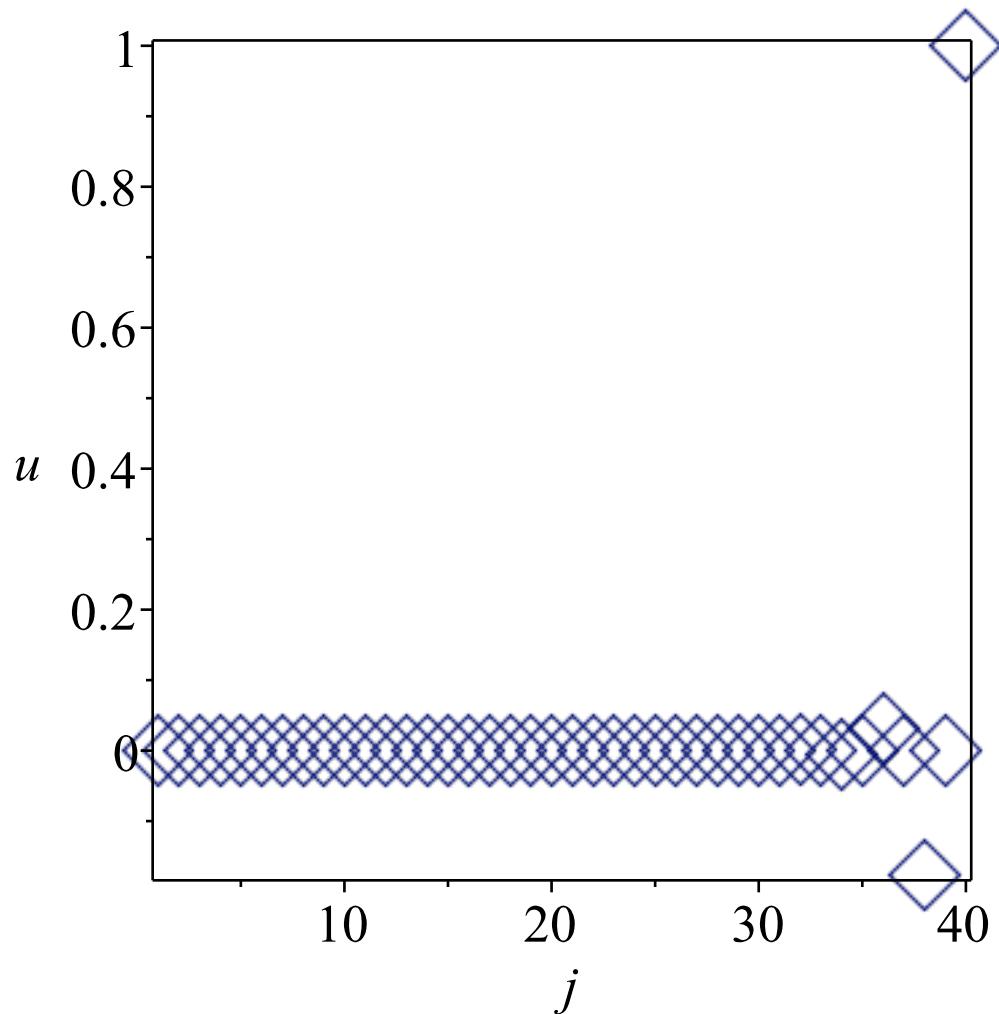
120000. + 0. I



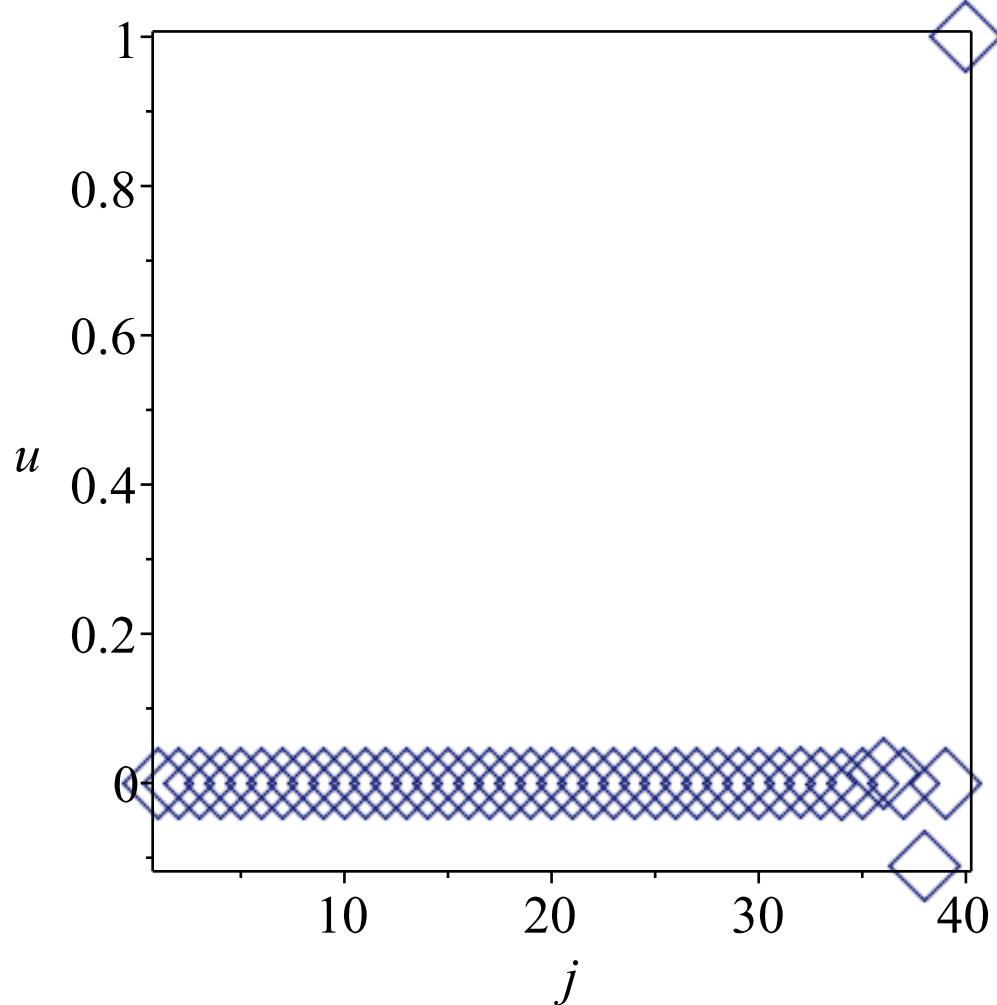
120000.000000000 + 0. I



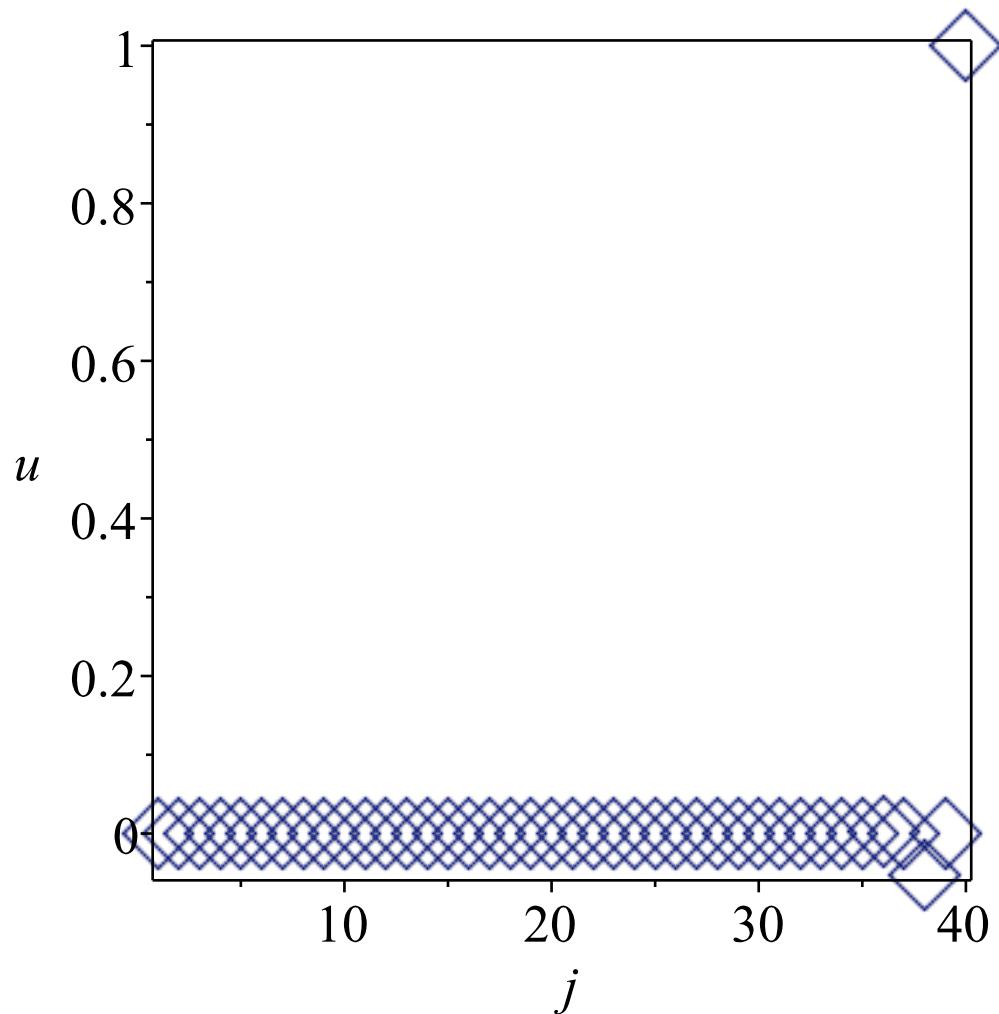
120000.000000000 + 0. I



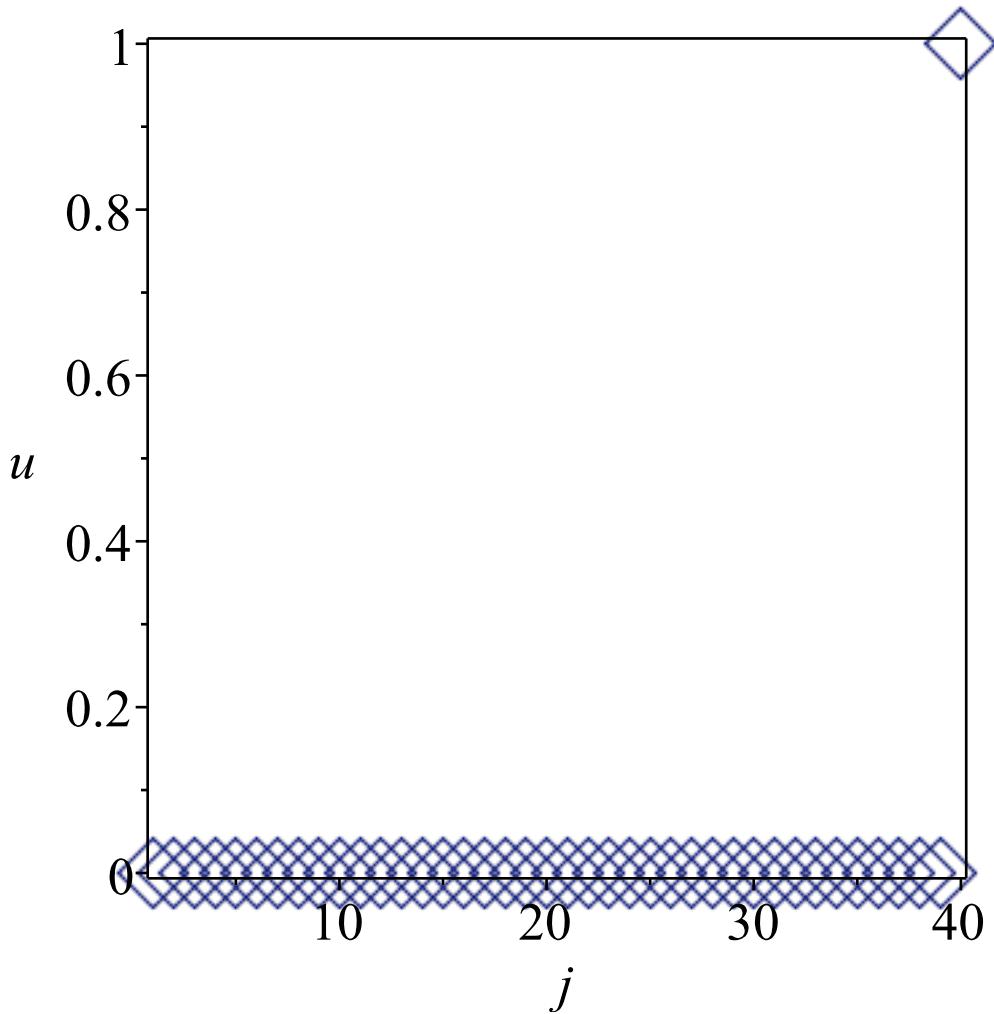
120000.000000000 + 0. I



120000. + 0. I



120000. + 0. I



Springs needed in the experiment to see a localized edge mode

$$\text{subs}\left(\lambda = \frac{3}{8}, K1\right); \text{subs}\left(\lambda = \frac{3}{8}, K2\right);$$

$$3750$$

$$8250 \quad (1.17)$$

Expected edge mode frequency in Hz

$$\text{evalf}\left(\frac{\sqrt{\text{eigs}\left(\frac{3}{8}\right)[1][2 \cdot N]}}{2 \cdot \text{Pi}}\right);$$

$$55.1328895276665 + 0. \text{i} \quad (1.18)$$