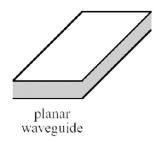
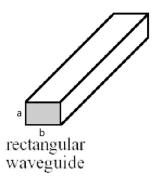
**Q**:

a) A flat metal waveguide is assumed, find its modes for that waveform.



b) A metal waveguide channel with a rectangular cross-section of length a and width b is assumed. Find its modes for that waveform.

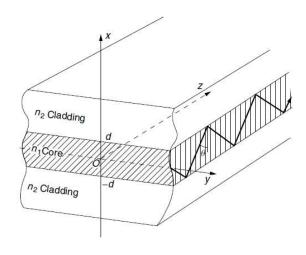


Sol:

a) In TM mode, we check the release:

$$E = E_0(x, y)e^{i(\beta z - \ )}$$

$$H = H_0(x, y)e^{i(\beta z - \omega)}$$



Because we are in TM mode:

$$H_z = H_x = 0$$
,  $E_v = 0$ 

Based on the Helmholtz equation:

$$\nabla^{2}H + (n_{1,2}k)^{2}H = 0 \rightarrow \nabla^{2}H + (n_{1,2}^{2}k^{2} - \beta^{2})H = 0$$

$$n_{1}k > \beta > n_{2}k$$

$$n_{1}^{2} - \beta^{2} = K^{2} \qquad 0 < x < d$$

Let's solve the above equation only inside the core:

$$H_{y} = Acos(Kx) + Bsin(Kx)$$

$$H_{y}(x) = 0 \quad x = 0, x = d \rightarrow H_{y}(x) = Bsin\left(\frac{m\pi}{d}x\right)$$

$$H_{y}(x, y, z, t) = Bsin\left(\frac{m\pi}{d}x\right)e^{i(\beta z - \omega t)}$$

Now, based on Maxwell's equations, we can get the rest of the equations:

$$\nabla \times H = \frac{\partial D}{\partial t}$$
 
$$\nabla \times E = -\frac{\partial B}{\partial t}$$
 
$$\hat{y} \quad \hat{z}$$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & H_{y} & 0 \end{vmatrix} = -\epsilon * i\omega(\mathbf{E}_{x}\hat{x} + \mathbf{E}_{z}\hat{z}) \rightarrow \begin{cases} -\frac{\partial H_{y}}{\partial z} = -\epsilon * i\omega\mathbf{E}_{x} \\ \frac{\partial H_{y}}{\partial x} = -\epsilon * i\omega\mathbf{E}_{z} \end{cases}$$

$$\begin{cases} E_{x}(x, y, z, t) = \frac{i\beta H_{y}}{\epsilon * i\omega} = \frac{\beta}{\epsilon * \omega} H_{y} = \frac{\beta}{\epsilon * \omega} \sin\left(\frac{m\pi}{d}x\right) e^{i(\beta z - \omega t)} \\ E_{z}(x, y, z, t) = \frac{i}{\epsilon * \omega} \frac{\partial H_{y}}{\partial x} = \frac{i}{\epsilon * \omega} \frac{m\pi}{d} \cos\left(\frac{m\pi}{d}x\right) e^{i(\beta z - \omega t)} \end{cases}$$

b) In TM mode, we check the release:

$$E = E_0(x, y)e^{i(\beta z - \omega t)}$$

$$H = H_0(x, y)e^{i(\beta z - \omega t)}$$

$$H(x, y) = H(x)H(y)$$

Because we are in TM mode, Of course, it should be noted that since the two states x and y are independent of each other, we must solve two equations:

$$H_y(x,0) = H_y(x,b) = H_y(0,y) = H_y(a,y) = 0$$
  
 $H_y(x,0) = H_y(x,b) = H_y(0,y) = H_y(a,y) = 0$ 

Based on the Helmholtz equation:

$$\nabla^2 H + (k)^2 H = 0 \rightarrow \nabla^2 H + (K_x^2 - K_y^2) H = 0$$
  
 $K_x^2 - K_y^2 = K^2$   $0 < x < a, 0 < y < b$ 

Let's solve the above equation only inside the core:

$$K_{x} = \frac{m\pi}{a}, K_{y} = \frac{n\pi}{b}$$

$$H_{y} = A\cos(Kx) + B\sin(Kx)$$

$$H_{y}(x,y) = 0 \quad x = 0, x = a, y = 0, y = b \to H_{y}(x,y) = B sin\left(\frac{m\pi}{a}x\right) sin\left(\frac{n\pi}{b}y\right)$$

$$H_{y}(x,y,z,t) = B sin(k_{y}x) sin(k_{y}y) e^{i(\beta z - \omega t)}$$

Now, based on Maxwell's equations, we can get the rest of the equations:

$$\nabla \times H = \frac{\partial D}{\partial t}$$

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix} = \mu * i\omega \left( \mathbf{H}_x \hat{x} + \mathbf{H}_y \hat{y} \right) \rightarrow \begin{cases} \mu * i\omega \mathbf{H}_x = \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \\ \mu * i\omega \mathbf{H}_y = \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \\ 0 = \frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \end{cases}$$

$$\rightarrow \begin{cases} \mu * i\omega \mathbf{H}_x = \frac{\partial E_z}{\partial y} - i\beta E_y \\ \mu * i\omega \mathbf{H}_y = i\beta E_x - \frac{\partial E_z}{\partial x} \\ 0 = \frac{\partial E_y}{\partial x} - \frac{\partial E_z}{\partial y} \end{cases}$$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ H_x & H_y & 0 \end{vmatrix} = -\epsilon * i\omega (\mathbf{E}_x \hat{x} + \mathbf{E}_y \hat{y} + \mathbf{E}_z \hat{z})$$

$$\rightarrow \begin{cases} -\epsilon * i\omega \mathbf{E}_x = -\frac{\partial H_y}{\partial z} \\ -\epsilon * i\omega \mathbf{E}_y = \frac{\partial H_x}{\partial z} \end{cases} \rightarrow \begin{cases} -\epsilon * i\omega \mathbf{E}_x = -i\beta H_y \\ -\epsilon * i\omega \mathbf{E}_y = i\beta H_x \end{cases}$$

$$-\epsilon * i\omega \mathbf{E}_z = \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y}$$

$$-\epsilon * i\omega E_x = -\frac{\partial H_y}{\partial z} \rightarrow E_x(x, y, z, t) = \frac{\beta}{\epsilon * \omega} Bsin(k_x x) sin(k_y y) e^{i(\beta z - \omega t)}$$

$$0 = \frac{\partial E_{y}}{\partial x} - \frac{\partial E_{x}}{\partial y} \to \mathbf{E}_{y}(\mathbf{x}, \mathbf{y}, \mathbf{z}, \mathbf{t}) = \int \frac{\beta}{\epsilon * \omega} k_{y} B \sin(k_{x}x) \cos(k_{y}y) e^{i(\beta z - \omega t)} dx$$

$$= \frac{\beta}{\epsilon * \omega} k_{y} * \frac{-1}{k_{x}} B \cos(k_{x}x) \cos(k_{y}y) e^{i(\beta z - \omega t)}$$

$$= -\frac{\beta}{\epsilon * \omega} \frac{k_{y}}{k_{x}} B \cos(k_{x}x) \cos(k_{y}y) e^{i(\beta z - \omega t)}$$

$$-\epsilon * i\omega E_y = \frac{\partial H_x}{\partial z} \rightarrow$$

$$H_x(x, y, z, t) = \beta \frac{k_y}{k_x} B\cos(k_x x) \cos(k_y y) e^{i(\beta z - \omega t)}$$

$$-\epsilon * i\omega \mathbf{E}_{z} = \frac{\partial H_{y}}{\partial x} - \frac{\partial H_{x}}{\partial y} \rightarrow$$

$$\mathbf{E}_{z}(x, y, z, t) = (k_{x}B\cos(k_{x}x)\sin(k_{y}y)e^{i(\beta z - \omega t)} + \beta \frac{k_{y}^{2}}{k_{x}}B\cos(k_{x}x)\sin(k_{y}y)e^{i(\beta z - \omega t)}) \frac{1}{-\epsilon * i\omega}$$

$$= \frac{-\mathbf{B}}{\epsilon * i\omega}\cos\left(\frac{m\pi}{a}x\right)\sin\left(\frac{n\pi}{b}y\right)e^{i(\beta z - \omega t)}\left(k_{x} + \beta \frac{k_{y}^{2}}{k_{x}}\right)$$