

## Note on oblique parameters, JH, April 15, 2021

I have used 2HDMC [1] in the IDM mode as well as the formulas from [2] and [3] to compute the  $S$  and  $T$  parameters (oblique or Peskin-Takeuchi parameter). For  $T$  the expressions from [2] [eq. (27)] and [3] [eq. (5.26)] are the same while for the  $S$  parameter they are different. The simple expression in [2] [eq. (65)] can just be integrated numerically<sup>1</sup> while the one in [3] [eq. (5.25)] requires the evaluation of Passarino-Veltman integrals using *e.g.* LoopTools. For the  $T$  parameter I found agreement on the per-mille level between 2HDMC and the analytic expression. For the  $S$  parameter the deviations are a bit larger, see table 1.

$m_{H^0}$	$m_{A^0}$	$m_{H^\pm}$	$S(2\text{HDMC})$	$S(\text{Barbieri})$	$S(\text{Haber})$
106.2	1040.4	873.9	-0.0119	-0.01233	-0.01235
411.8	2362.8	2655.8	-0.0259	-0.02687	-0.02685
754.4	815.9	860.8	-0.00468	-0.00486	-0.00487
187.6	277.3	261.2	-0.00584	-0.00561	-0.00569
246.9	1112.8	1287.7	-0.02667	-0.02763	-0.02766
1255.2	1687.2	1860.8	-0.01166	-0.01212	-0.01211
349.5	1483.7	1171.2	-0.0068	-0.00708	-0.0071
175.8	1658.	1275.8	-0.00734	-0.00764	-0.00773
351.7	834.7	996.3	-0.02405	-0.0249	-0.02494
783.3	2749.8	3248.5	-0.02651	-0.0275	-0.02731
624.3	2405.2	2018.5	-0.00948	-0.00987	-0.01013
1357.4	2005.9	2459.5	-0.01886	-0.01958	-0.01958
307.4	1646.9	1352.4	-0.00965	-0.01003	-0.00999
294.	915.5	892.3	-0.01595	-0.01652	-0.01654
380.1	2241.3	1779.4	-0.00816	-0.00849	-0.00846
305.3	786.2	607.6	-0.00258	-0.00265	-0.00266
1140.9	1225.7	1335.	-0.00612	-0.00638	-0.00638
180.9	699.1	661.4	-0.01574	-0.01624	-0.01629
859.7	2116.6	2616.8	-0.02617	-0.02714	-0.02716
1089.2	1197.8	1385.4	-0.00976	-0.01014	-0.01015

Table 1: Comparison of the  $S$  oblique parameter from using 2HDMC [1] and the analytic expressions from Barbieri *et al.* [2] and Haber *et al.* [3] for 20 random parameter points.

In our study, we want to use the analytic expressions [2, 4] (see footnote 1):

$$S = \frac{1}{2\pi} \left( \frac{1}{6} \ln \frac{M_{H^0}^2}{M_{H^\pm}^2} - \frac{5}{36} + \frac{M_{H^0}^2 M_{A^0}^2}{3(M_{A^0}^2 - M_{H^0}^2)^2} + \frac{M_{A^0}^4 (M_{A^0}^2 - 3M_{H^0}^2)}{6(M_{A^0}^2 - M_{H^0}^2)^3} \ln \frac{M_{A^0}^2}{M_{H^0}^2} \right) \quad (1)$$

$$T = \frac{1}{16\pi^2 \alpha_{\text{ew}} v^2} \left( F(M_{H^\pm}, M_{H^0}) + F(M_{H^\pm}, M_{A^0}) - F(M_{A^0}, M_{H^0}) \right), \quad (2)$$

where  $F(m_1, m_2) = (m_1^2 + m_2^2)/2 - m_1^2 m_2^2 / (m_1^2 - m_2^2) \ln(m_1^2 / m_2^2)$ . I used  $\alpha_{\text{ew}} = 1/127.9$ . Furthermore, eq. (2) is written for  $v \simeq 246 \text{ GeV}$  (note the difference in the pre-factor with respect to [4]).

## References

- [1] D. Eriksson, J. Rathsman, and O. Stal, “2HDMC: Two-Higgs-Doublet Model Calculator Physics and Manual,” *Comput. Phys. Commun.* **181** (2010) 189–205, [arXiv:0902.0851 \[hep-ph\]](#).

<sup>1</sup>BTW, it seems it can actually be integrated numerically, as [4] give a fully algebraic expression that numerically resembles the one from [2] to machine precision.

- [2] R. Barbieri, L. J. Hall, and V. S. Rychkov, “Improved naturalness with a heavy Higgs: An Alternative road to LHC physics,” *Phys. Rev. D* **74** (2006) 015007, [arXiv:hep-ph/0603188](#).
- [3] H. E. Haber and D. O’Neil, “Basis-independent methods for the two-Higgs-doublet model III: The CP-conserving limit, custodial symmetry, and the oblique parameters S, T, U,” *Phys. Rev. D* **83** (2011) 055017, [arXiv:1011.6188 \[hep-ph\]](#).
- [4] M. Baak, M. Goebel, J. Haller, A. Hoecker, D. Ludwig, K. Moenig, M. Schott, and J. Stelzer, “Updated Status of the Global Electroweak Fit and Constraints on New Physics,” *Eur. Phys. J. C* **72** (2012) 2003, [arXiv:1107.0975 \[hep-ph\]](#).