

# GRAVITATIONAL LENSING

## LECTURE 6

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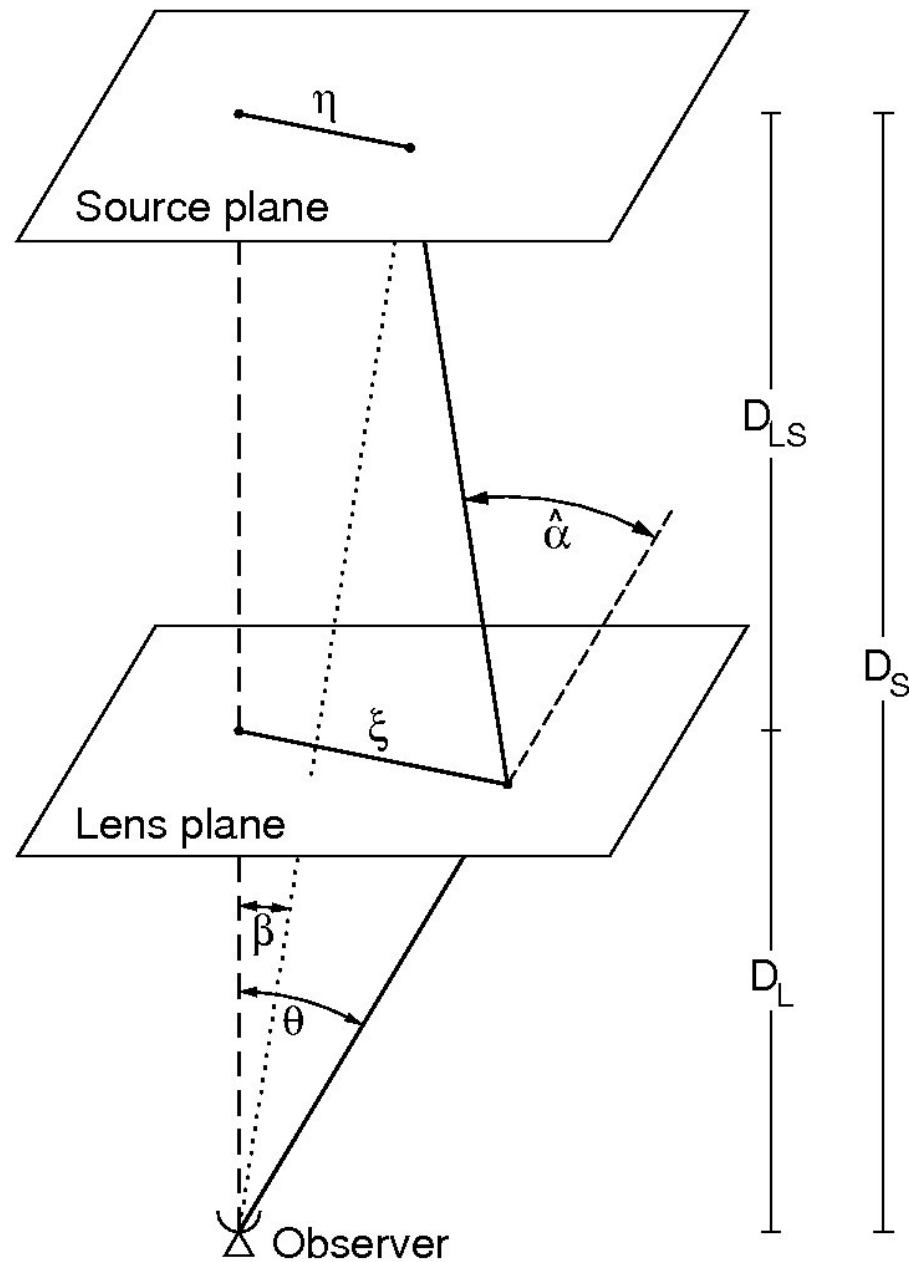
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*AA 2016-2017*

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- lens mapping (first order)

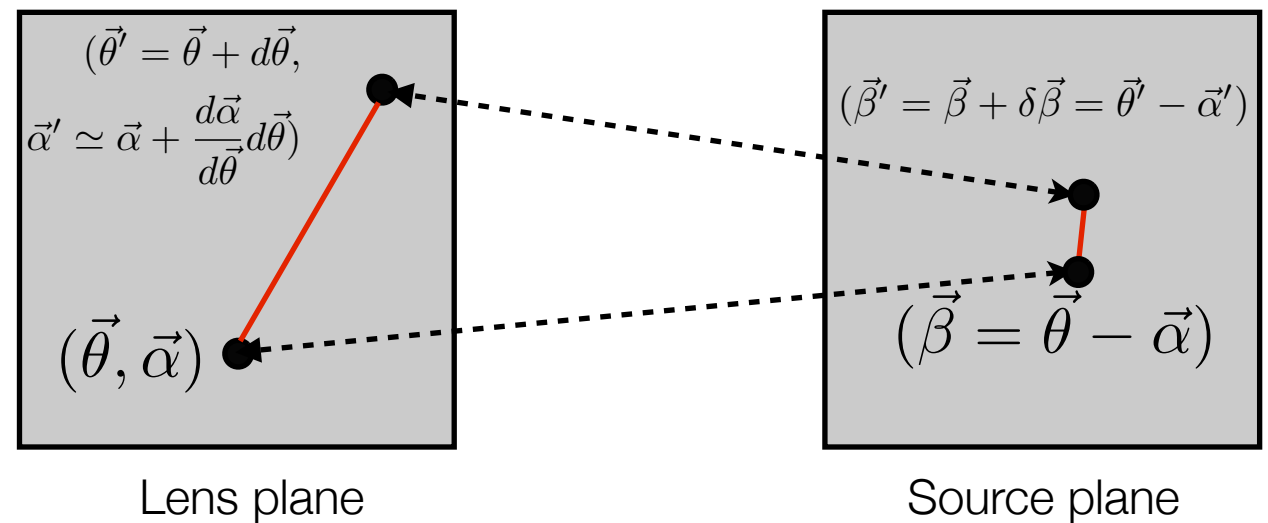
# LENS MAPPING (FIRST ORDER)



- we derived the lens equation

$$\vec{\beta} = \vec{\theta} - \frac{D_{LS}}{D_S} \hat{\alpha}(\vec{\theta}) = \vec{\theta} - \vec{\alpha}(\vec{\theta})$$

- Assuming that the d.a. does not vary significantly over the scale  $d\theta$ :



$$(\vec{\beta}' - \vec{\beta}) = \left( I - \frac{d\vec{\alpha}}{d\vec{\theta}} \right) (\vec{\theta}' - \vec{\theta}) = A(\vec{\theta}' - \vec{\theta})$$

# LENS MAPPING (FIRST ORDER)

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$$A = \frac{\partial \beta}{\partial \theta} = \left( \delta_{ij} - \frac{\partial \alpha_i(\vec{\theta})}{\partial \theta_j} \right) = \left( \delta_{ij} - \frac{\partial^2 \psi(\vec{\theta})}{\partial \theta_i \partial \theta_j} \right)$$

*A is called “the lensing Jacobian”: it is a symmetric second rank tensor describing the first order mapping between lens and source planes.*

*This tensor can be written as the sum of an isotropic part, proportional to its trace, and an anisotropic traceless part.*

$$A_{iso,i,j} = \frac{1}{2} \text{Tr} A \delta_{i,j}$$

$$A_{aniso,i,j} = A_{i,j} - \frac{1}{2} \text{Tr} A \delta_{i,j}$$

# ANISOTROPIC PART

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$$A_{aniso,i,j} = A_{i,j} - \frac{1}{2} \text{Tr} A \delta_{i,j}$$

$$\begin{aligned} \left( A - \frac{1}{2} \text{tr} A \cdot I \right)_{ij} &= \delta_{ij} - \Psi_{ij} - \frac{1}{2} (1 - \Psi_{11} + 1 - \Psi_{22}) \delta_{ij} \\ &= -\Psi_{ij} + \frac{1}{2} (\Psi_{11} + \Psi_{22}) \delta_{ij} \\ &= \begin{pmatrix} -\frac{1}{2}(\Psi_{11} - \Psi_{22}) & -\Psi_{12} \\ -\Psi_{12} & \frac{1}{2}(\Psi_{11} - \Psi_{22}) \end{pmatrix} \end{aligned}$$

Introducing the *shear*:

$$\gamma_1 = \frac{1}{2} (\Psi_{11} - \Psi_{22})$$

$$\gamma_2 = -\Psi_{12} = -\Psi_{21}$$

$$\begin{pmatrix} \gamma_1 & \gamma_2 \\ \gamma_2 & -\gamma_1 \end{pmatrix}$$

*Symmetric, trace-less tensor*

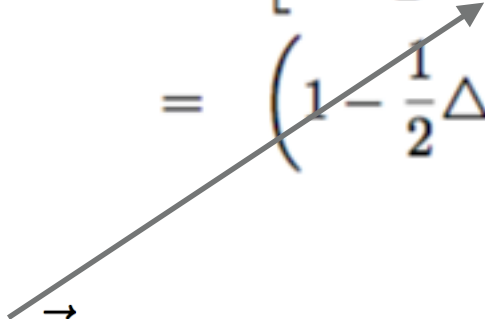
# ISOTROPIC PART

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$$A_{iso,i,j} = \frac{1}{2} \text{Tr} A \delta_{i,j}$$

$$\begin{aligned} \frac{1}{2} \text{tr} A \cdot I &= \left[ 1 - \frac{1}{2} (\Psi_{11} + \Psi_{22}) \right] \delta_{ij} \\ &= \left( 1 - \frac{1}{2} \Delta \Psi \right) \delta_{ij} = (1 - \kappa) \delta_{ij} \end{aligned}$$

*Remember:*  $\Delta_{\theta} \Psi(\vec{\theta}) = 2\kappa(\vec{\theta})$



# LENSING JACOBIAN

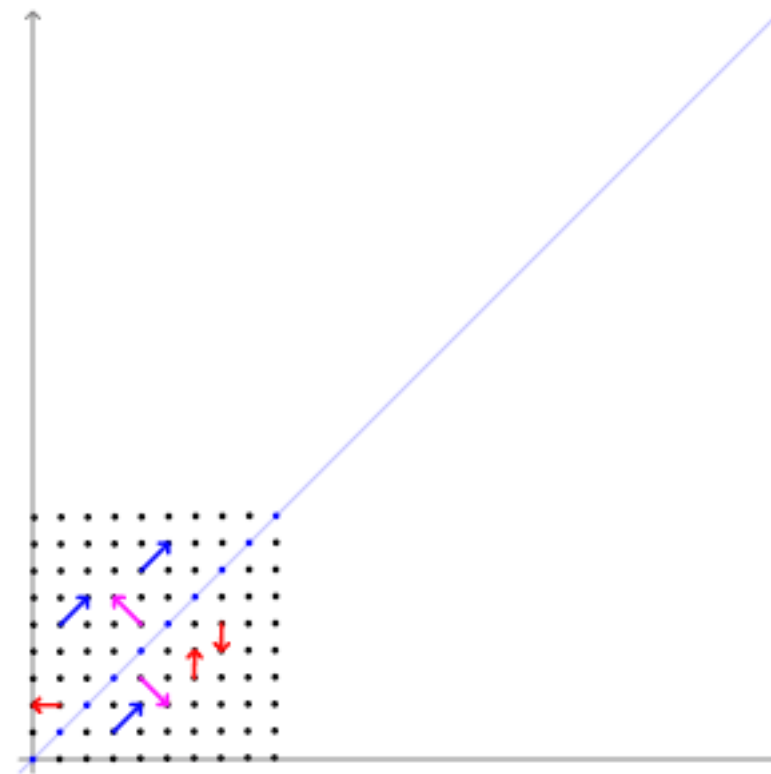
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$$A = (1 - \kappa) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \begin{pmatrix} \gamma_1 & \gamma_2 \\ \gamma_2 & -\gamma_1 \end{pmatrix}$$

*Lens mapping at first order is a linear application, distorting areas.*

*Distortion directions are given by the **eigenvectors** of  $A$ .*

*Distortion amplitudes in these directions are given by the **eigenvalues**.*



# LENSING JACOBIAN

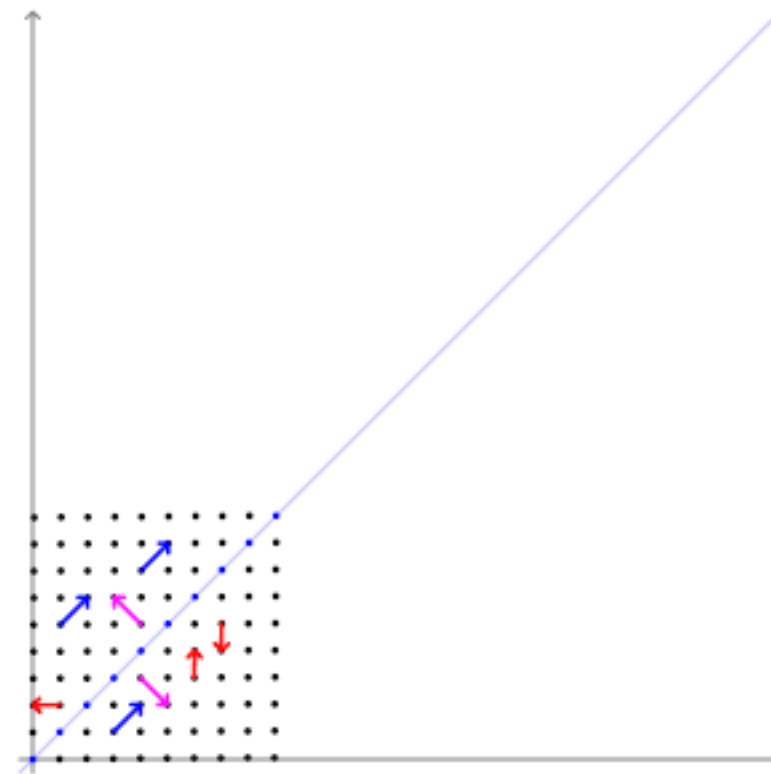
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$$A = (1 - \kappa) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \begin{pmatrix} \gamma_1 & \gamma_2 \\ \gamma_2 & -\gamma_1 \end{pmatrix}$$

*Lens mapping at first order is a linear application, distorting areas.*

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*Distortion amplitudes in these directions are given by the **eigenvalues**.*





# EIGENVALUES

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$$\det A = (1 - \kappa - \gamma)(1 - \kappa + \gamma)$$
$$\gamma = \sqrt{\gamma_1^2 + \gamma_2^2}$$

*There is thus an orthogonal coordinate transformation  $R(\varphi)$ , a rotation by an angle  $\varphi$ , which brings the Jacobian matrix into diagonal form.*

*Generally, the Jacobian matrix transforms as*

$$A \rightarrow A' = R(\varphi)^T A R(\varphi)$$

*This shows that the shear components transform under coordinate rotations as*

$$\gamma_1 \rightarrow \gamma'_1 = \gamma_1 \cos(2\varphi) + \gamma_2 \sin(2\varphi)$$
$$\gamma_2 \rightarrow \gamma'_2 = -\gamma_1 \sin(2\varphi) + \gamma_2 \cos(2\varphi)$$

*i.e. unlike a vector! Since the shear components are mapped onto each other after rotations of  $\varphi = \pi$  rather than  $\varphi = 2\pi$ , they form a so-called spin-2 field.*

# EXAMPLE: FIRST ORDER DISTORTION OF A CIRCULAR SOURCE

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$$\beta_1^2 + \beta_2^2 = \beta^2$$

*In the reference frame where  $A$  is diagonal:*

$$\begin{pmatrix} \beta_1 \\ \beta_2 \end{pmatrix} = \begin{pmatrix} 1 - \kappa - \gamma & 0 \\ 0 & 1 - \kappa + \gamma \end{pmatrix} \begin{pmatrix} \theta_1 \\ \theta_2 \end{pmatrix}$$

$$\beta_1 = (1 - \kappa - \gamma)\theta_1$$

$$\beta_2 = (1 - \kappa + \gamma)\theta_2$$

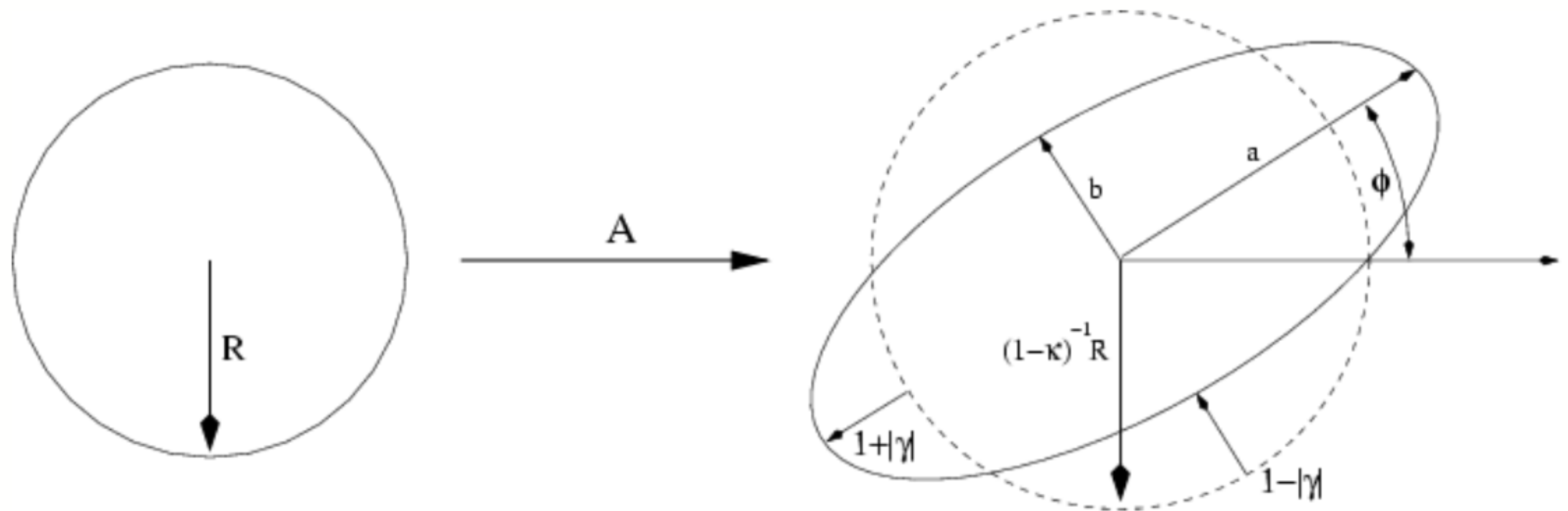
$$\beta^2 = \beta_1^2 + \beta_2^2 = (1 - \kappa - \gamma)^2 \theta_1^2 + (1 - \kappa + \gamma)^2 \theta_2^2$$

*This is the equation of an ellipse with semi-axes:*

$$a = \frac{\beta}{1 - \kappa - \gamma} \qquad b = \frac{\beta}{1 - \kappa + \gamma}$$

# EXAMPLE: FIRST ORDER DISTORTION OF A CIRCULAR SOURCE

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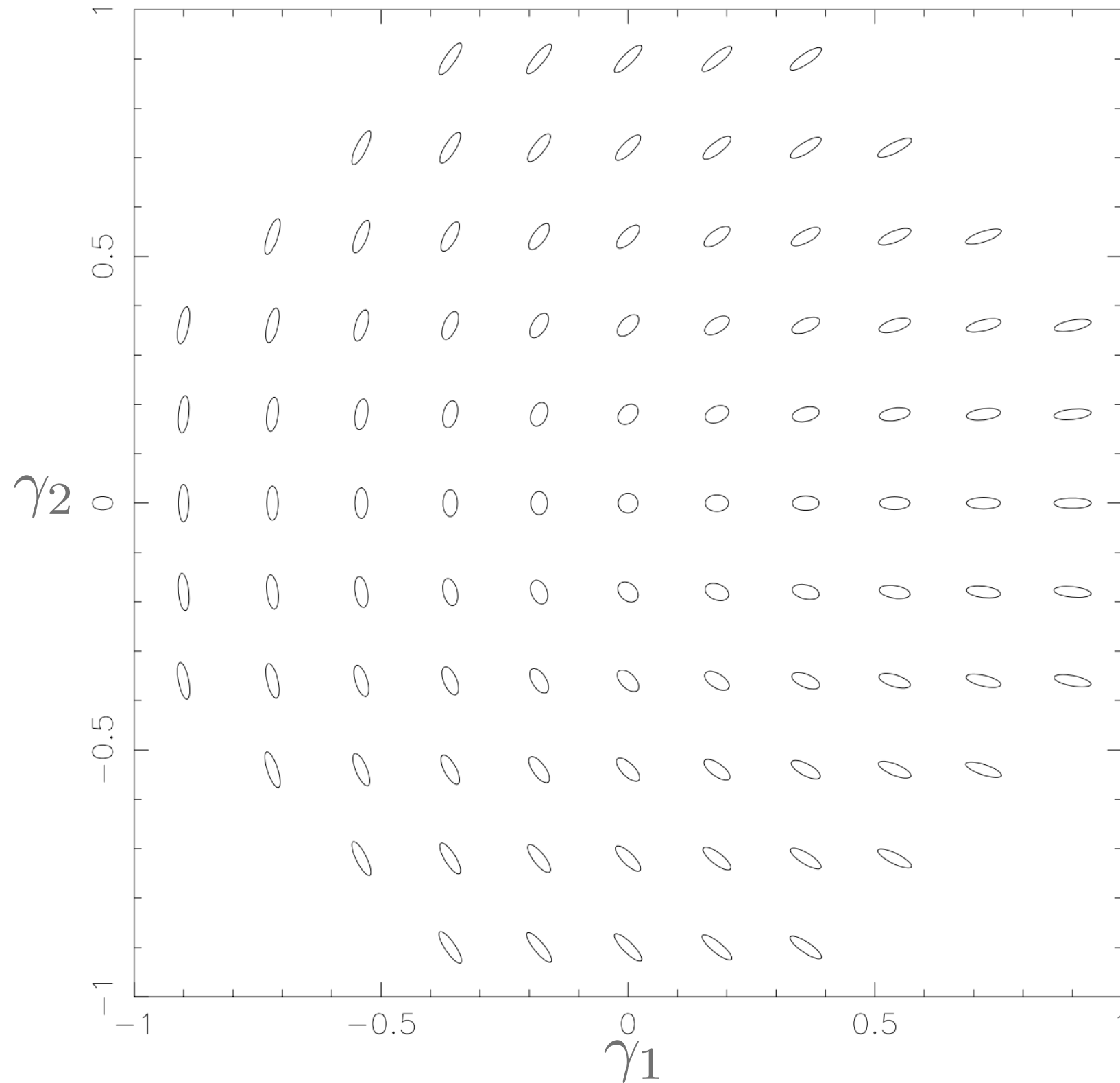
*convergence: responsible for isotropic expansion or contraction*

*shear: responsible for anisotropic distortion*

*Ellipticity:* 
$$e = \frac{a - b}{a + b} = \frac{\gamma}{1 - \kappa} = g$$

# SHEAR DISTORTIONS

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# CONSERVATION OF SURFACE BRIGHTNESS

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*The source surface  
brightness is*

$$I_\nu = \frac{dE}{dt dA d\Omega d\nu}$$

*In phase space, the radiation emitted is characterized by the density*

$$f(\vec{x}, \vec{p}, t) = \frac{dN}{d^3x d^3p}$$

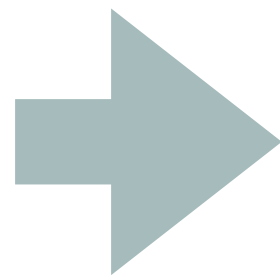
*In absence of photon creations or absorptions,  $f$  is conserved (Liouville theorem)*

$$dN = \frac{dE}{h\nu} = \frac{dE}{cp}$$

$$f(\vec{x}, \vec{p}, t) = \frac{dN}{d^3x d^3p} = \frac{dE}{h c p^3 dA dt d\nu d\Omega} = \frac{I_\nu}{h c p^3}$$

$$d^3x = c dt dA$$

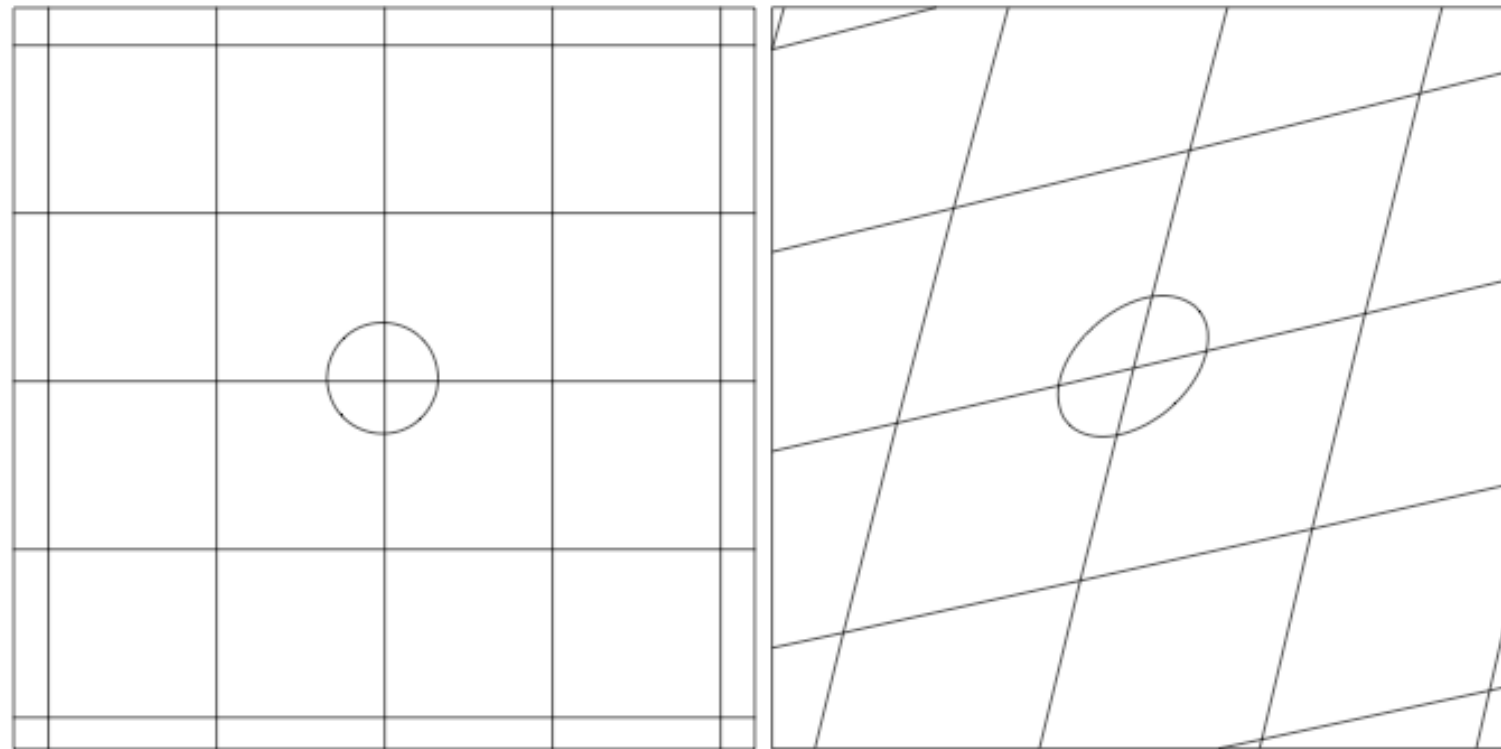
$$d^3\vec{p} = p^2 dp d\Omega$$



*Since GL does not involve creation or absorption of photons, neither it changes the photon momenta (achromatic!), surface brightness is conserved!*

# MAGNIFICATION

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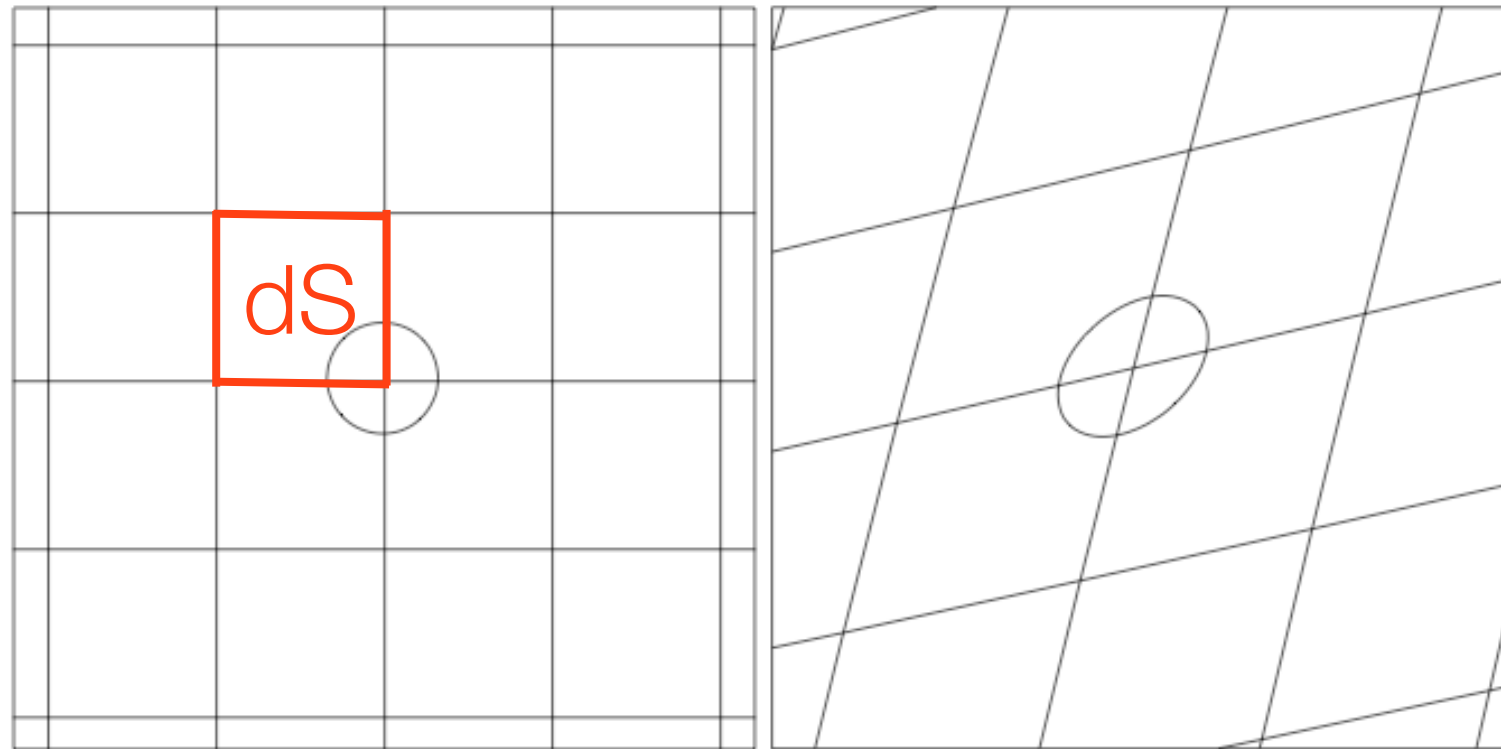
Kneib & Natarajan (2012)

$$F_{\nu} = \int_I I_{\nu}(\vec{\theta}) d^2\theta = \int_S I_{\nu}^S[\vec{\beta}(\vec{\theta})] \mu d^2\beta$$

*Lensing changes the amount of photons (flux) we receive from the source by changing the solid angle the source subtends*

# MAGNIFICATION

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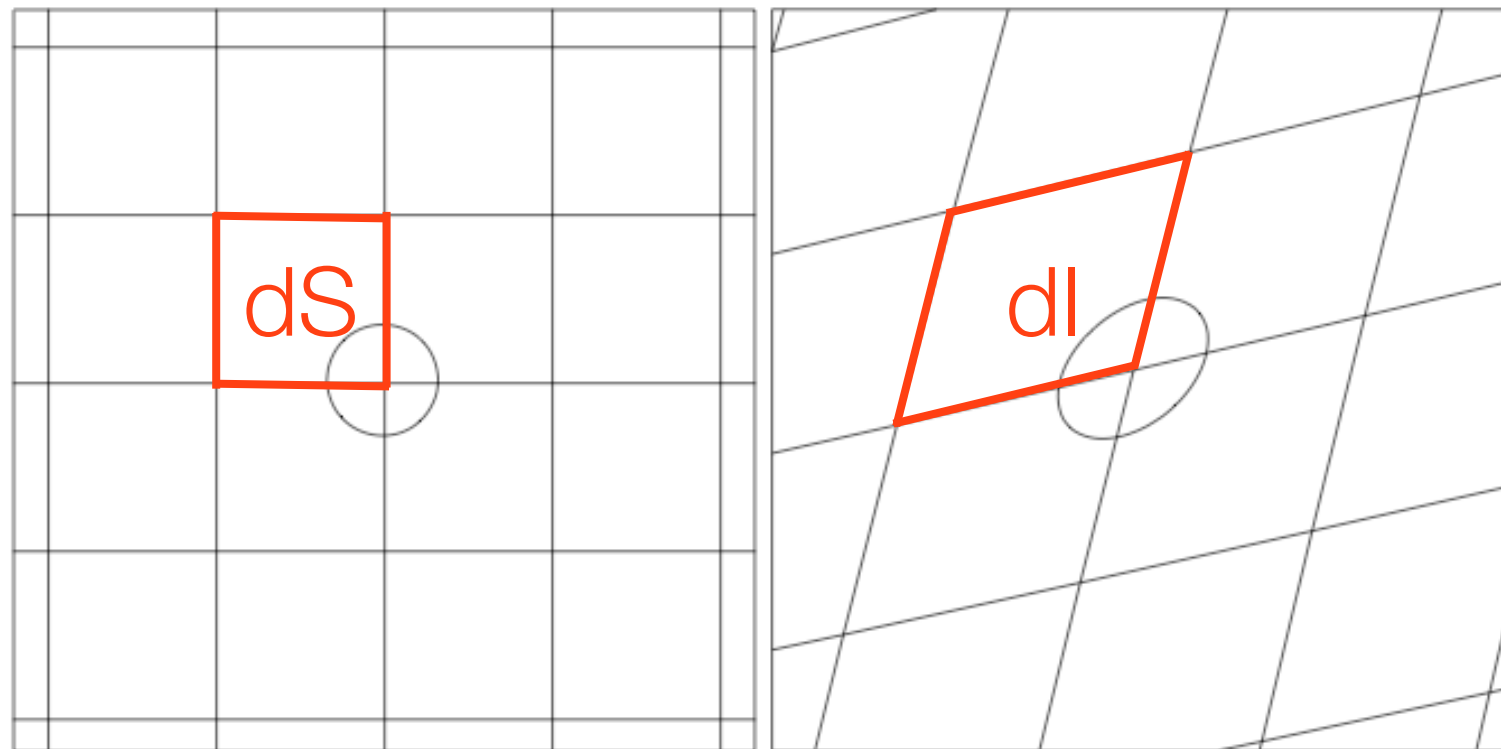
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# MAGNIFICATION

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Kneib & Natarajan (2012)

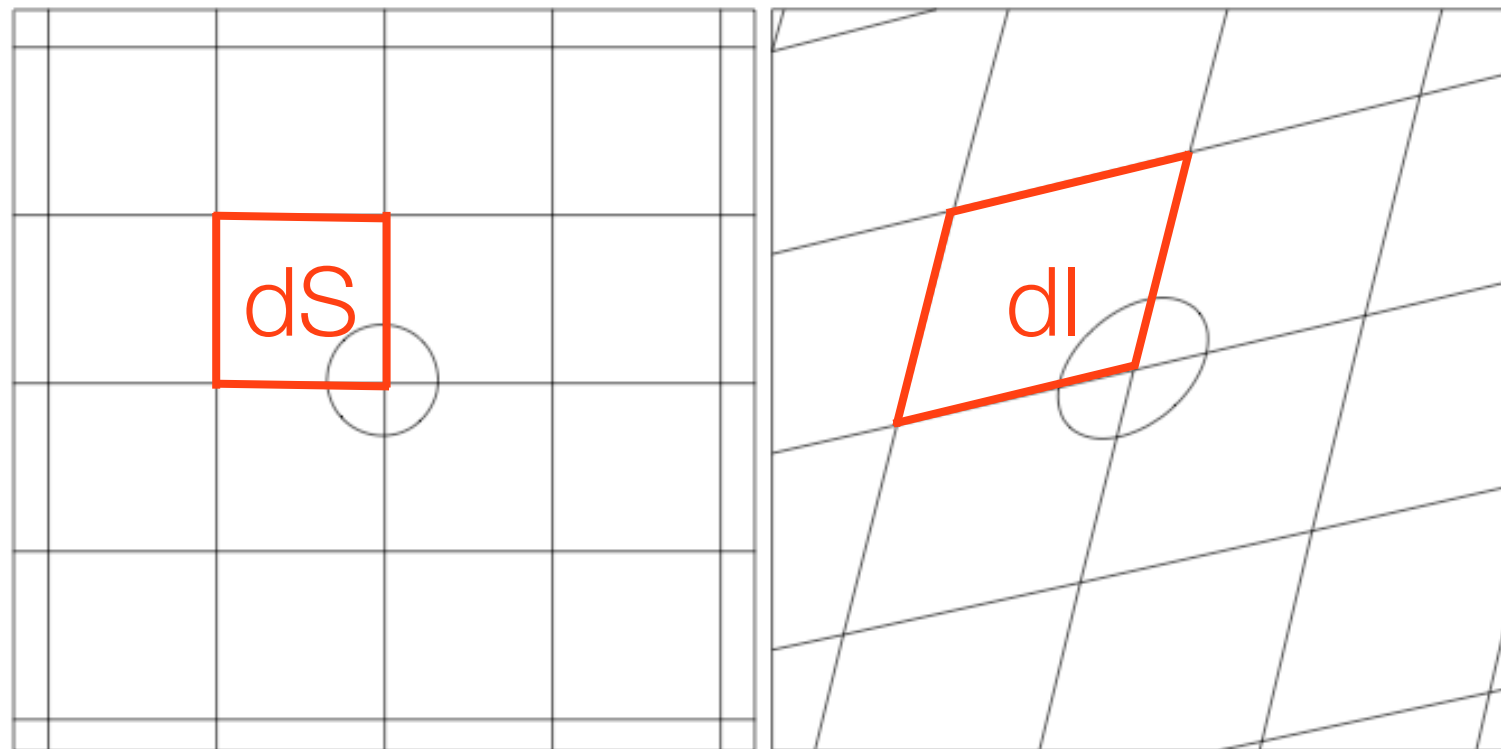
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# MAGNIFICATION

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Kneib & Natarajan (2012)

$$\mu = \frac{dI}{dS} = \frac{\delta\theta^2}{\delta\beta^2} = \det A^{-1}$$

$$F_\nu = \int_I I_\nu(\vec{\theta}) d^2\theta = \int_S I_\nu^S[\vec{\beta}(\vec{\theta})] \mu d^2\beta$$

*Lensing changes the amount of photons (flux) we receive from the source by changing the solid angle the source subtends*

# CRITICAL LINES AND CAUSTICS

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*Both convergence and shear are functions of position on the lens plane:*

$$\kappa = \kappa(\vec{\theta})$$

$$\gamma = \gamma(\vec{\theta})$$

*The determinant of the lensing Jacobian is*

$$\det A = (1 - \kappa - \gamma)(1 - \kappa + \gamma) = \mu^{-1}$$

*The **critical lines** are the lines where the eigenvalues of the Jacobian are zero:*

$$(1 - \kappa - \gamma) = 0 \quad \text{tangential critical line}$$

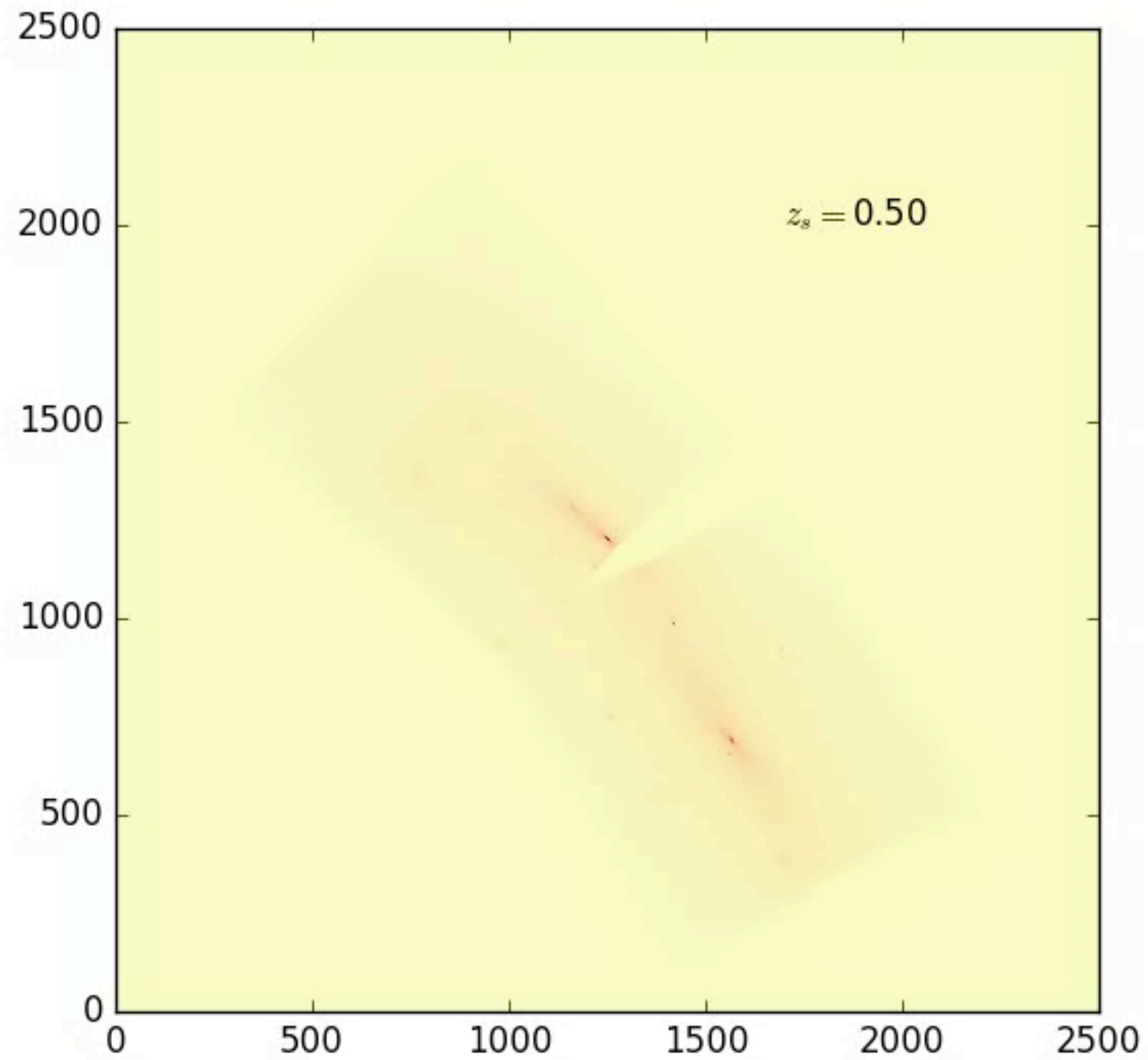
$$(1 - \kappa + \gamma) = 0 \quad \text{radial critical line}$$

*Along these lines the magnification diverges!*

*Via the lens equations, they are mapped into the **caustics**...*

# VISUALIZING THE CAUSTICS

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# VISUALIZING THE CAUSTICS

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