

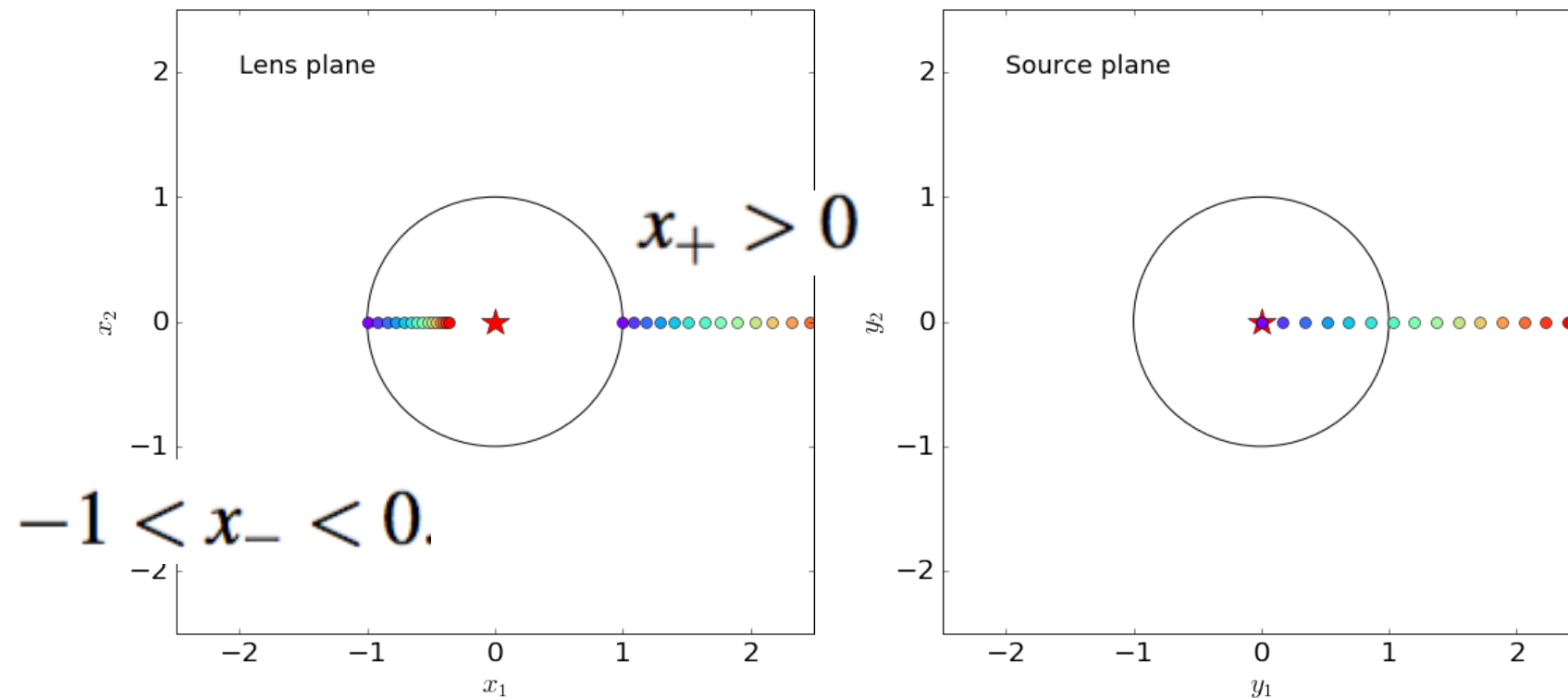
GRAVITATIONAL LENSING

9 – MICROLENSING LIGHT CURVES

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PROPERTIES OF THE IMAGES

From the last lesson



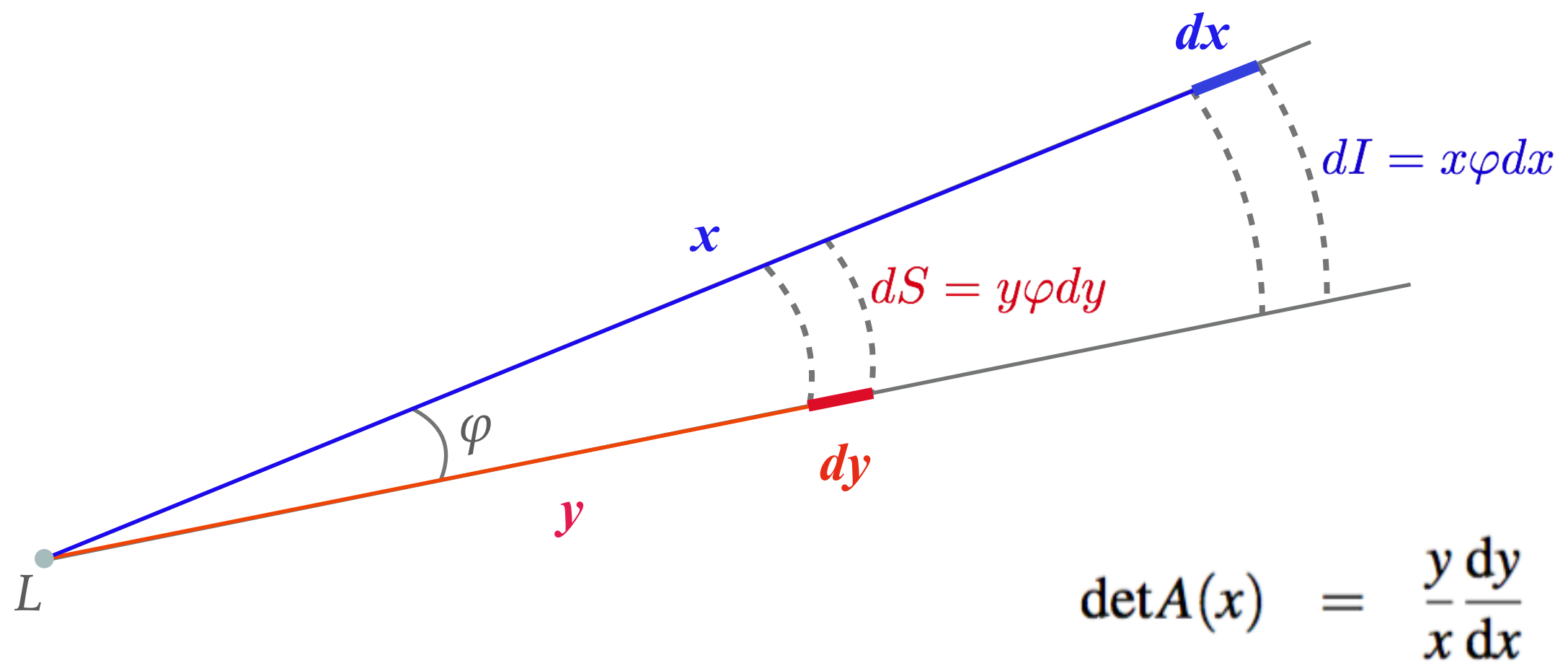
$$x_{\pm} = \frac{1}{2} \left[y \pm \sqrt{y^2 + 4} \right]$$

One of the images is internal to the Einstein radius, the other is external

For $y=0$, the image is a full ring: $x_{\pm} = \pm 1$

*This is the **Einstein ring***

MAGNIFICATION



CRITICAL LINES AND CAUSTICS

From the lens equation, it follows that:

$$\begin{aligned}\lambda_t(x) &= \frac{y}{x} = \left(1 - \frac{1}{x^2}\right) \\ \lambda_r(x) &= \frac{dy}{dx} = \left(1 + \frac{1}{x^2}\right) .\end{aligned}$$

The second eigenvalue is always positive (no critical line). The first is zero on the circle

$$x^2 = 1$$

Thus, the Einstein ring is the tangential critical line! The corresponding caustic is a point at $y=0$

IMAGE MAGNIFICATION

Clearly,

$$\det A(x) = \frac{y}{x} \frac{dy}{dx}$$

$$\begin{aligned}\lambda_t(x) &= \frac{y}{x} = \left(1 - \frac{1}{x^2}\right) \\ \lambda_r(x) &= \frac{dy}{dx} = \left(1 + \frac{1}{x^2}\right) .\end{aligned}$$



$$\mu(x) = \left(1 - \frac{1}{x^4}\right)^{-1}$$

IMAGE PARITY

Note that:

$$y > 0 \quad \Rightarrow \quad \begin{aligned} x_+ &> 0 \\ x_- &< 0 \end{aligned}$$

$$\mu_t = \frac{x}{y} \quad \Rightarrow \quad \begin{aligned} \mu_t(x_+) &> 0 \\ \mu_t(x_-) &< 0 \end{aligned}$$

$$\mu_r = \frac{dx}{dy} > 0$$

Thus the parity of the images is different!

SOURCE MAGNIFICATION

Let's compute now the source magnification. This is the sum of the magnifications of the two images

$$x_{\pm} = \frac{1}{2} \left[y \pm \sqrt{y^2 + 4} \right] \quad \rightarrow \quad \begin{aligned} \frac{x}{y} &= \frac{1}{2} \left(1 \pm \frac{\sqrt{y^2 + 4}}{y} \right) \\ \frac{dx}{dy} &= \frac{1}{2} \left(1 \pm \frac{y}{\sqrt{y^2 + 4}} \right) . \end{aligned}$$

Thus the magnifications at the two image positions are

$$\begin{aligned} \mu_{\pm}(y) &= \frac{1}{4} \left(1 \pm \frac{\sqrt{y^2 + 4}}{y} \right) \left(1 \pm \frac{y}{\sqrt{y^2 + 4}} \right) \\ &= \frac{1}{4} \left(1 \pm \frac{\sqrt{y^2 + 4}}{y} \pm \frac{y}{\sqrt{y^2 + 4}} + 1 \right) \\ &= \frac{1}{4} \left(2 \pm \frac{2y^2 + 4}{y\sqrt{y^2 + 4}} \right) = \frac{1}{2} \left(1 \pm \frac{y^2 + 2}{y\sqrt{y^2 + 4}} \right) \end{aligned}$$

SOURCE MAGNIFICATION

The total magnification is obtained by summing the magnifications of the images:

$$\begin{aligned}\mu_{\pm}(y) &= \frac{1}{4} \left(1 \pm \frac{\sqrt{y^2 + 4}}{y} \right) \left(1 \pm \frac{y}{\sqrt{y^2 + 4}} \right) \\ &= \frac{1}{4} \left(1 \pm \frac{\sqrt{y^2 + 4}}{y} \pm \frac{y}{\sqrt{y^2 + 4}} + 1 \right) \\ &= \frac{1}{4} \left(2 \pm \frac{2y^2 + 4}{y\sqrt{y^2 + 4}} \right) = \frac{1}{2} \left(1 \pm \frac{y^2 + 2}{y\sqrt{y^2 + 4}} \right)\end{aligned} \quad \Rightarrow \quad \mu(y) = \mu_+(y) + |\mu_-(y)| = \frac{y^2 + 2}{y\sqrt{y^2 + 4}}$$

The sum of the signed magnification is one!

We can take a power series of the magnification to see that $\mu \propto 1 + 2/y^4$ for $y \rightarrow \infty$.

Thus, the magnification drops quickly as the source moves away from the lens!

SOURCE MAGNIFICATION

In addition:

$$\begin{aligned}\left|\frac{\mu_+}{\mu_-}\right| &= \frac{1 + \frac{y^2+2}{y\sqrt{y^2+4}}}{\frac{y^2+2}{y\sqrt{y^2+4}} - 1} \\ &= \frac{y^2 + 2 + y\sqrt{y^2+4}}{y^2 + 2 - y\sqrt{y^2+4}}\end{aligned}$$



$$\begin{aligned}\left|\frac{\mu_+}{\mu_-}\right| &= \left(\frac{y + \sqrt{y^2+4}}{y - \sqrt{y^2+4}}\right)^2 \\ &= \left(\frac{x_+}{x_-}\right)^2.\end{aligned}$$

Power series at infinity:

$$\begin{aligned}\frac{1}{2} \left(y + \sqrt{y^2+4}\right)^2 &= y^2 + 2 + y\sqrt{y^2+4} \\ \frac{1}{2} \left(y - \sqrt{y^2+4}\right)^2 &= y^2 + 2 - y\sqrt{y^2+4}\end{aligned}$$

$$\left|\frac{\mu_+}{\mu_-}\right| \propto y^4$$

As we move the source away from the lens, the image in x_+ dominates the flux budget very soon.

$$\lim_{y \rightarrow \infty} \mu_- = 0$$

$$\lim_{y \rightarrow \infty} \mu_+ = 1$$

A SOURCE ON THE EINSTEIN RING

For a source on the Einstein ring:

$$x_{\pm} = \frac{1 \pm \sqrt{5}}{2}$$

$$\mu_{\pm} = \left[1 - \left(\frac{2}{1 \pm \sqrt{5}} \right)^4 \right]^{-1}$$

Therefore: $\mu = |\mu_+| + |\mu_-| = 1.17 + 0.17 = 1.34$

$$\Delta m = -2.5 \log \mu \sim 0.3$$

Given how quickly the magnification drops by moving the source away from the lens, we can assume that only sources within the Einstein radius are magnified in a significant way.

For this reason, the circle within the Einstein radius is assumed to be the cross section for microlensing.

SIZE OF THE EINSTEIN RADIUS

From the last lesson

$$\theta_E \equiv \sqrt{\frac{4GM}{c^2} \frac{D_{LS}}{D_L D_S}}$$

$$D \equiv \frac{D_L D_S}{D_{LS}}$$

$$\begin{aligned}\theta_E &\approx (10^{-3})'' \left(\frac{M}{M_\odot} \right)^{1/2} \left(\frac{D}{10 \text{ kpc}} \right)^{-1/2}, \\ &\approx 1'' \left(\frac{M}{10^{12} M_\odot} \right)^{1/2} \left(\frac{D}{\text{Gpc}} \right)^{-1/2},\end{aligned}$$

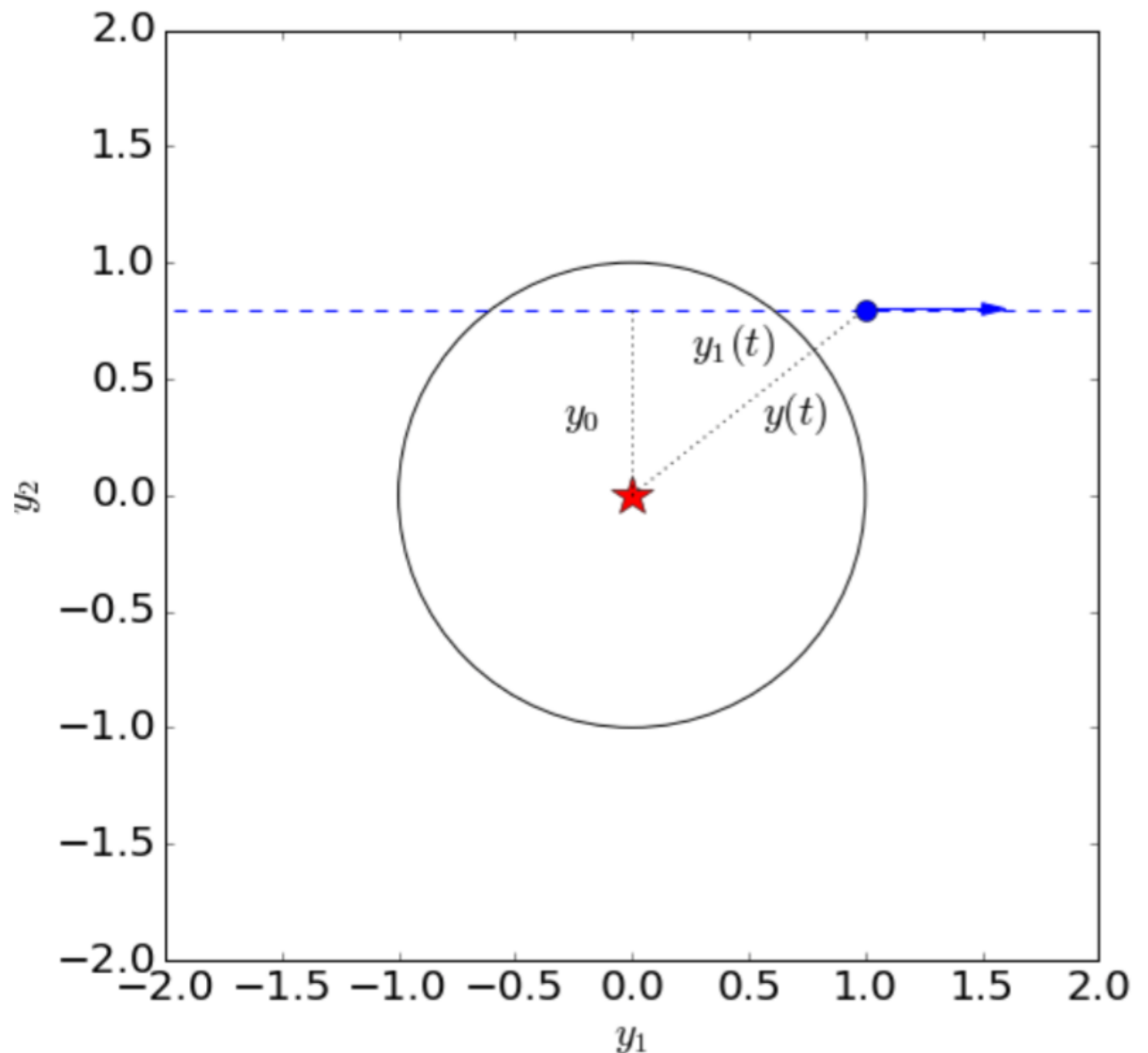
For a star like the sun within the MW, the Einstein radius is of the order of milli-arcseconds!

MICROLENSING LIGHT CURVE

- typical Einstein radii for lenses in the MW are ~ 1 mas
- thus, the image separation is too small to resolve the images
- magnification is small also for relatively close pairs of lenses and sources
- how to detect a microlensing event?

MICROLENSING LIGHT CURVE

- stars (including the sun) rotate around the galactic center
- rotation is differential (i.e. speed depends on distance)
- this introduces a relative velocity between the lenses and the sources (either in the bulge or in the MCs)
- this causes the relative distance between the sources and the lenses to vary over time...



MICROLENSING LIGHT CURVE

Assume a linear trajectory of the source relative to the lens, with impact parameter y_0

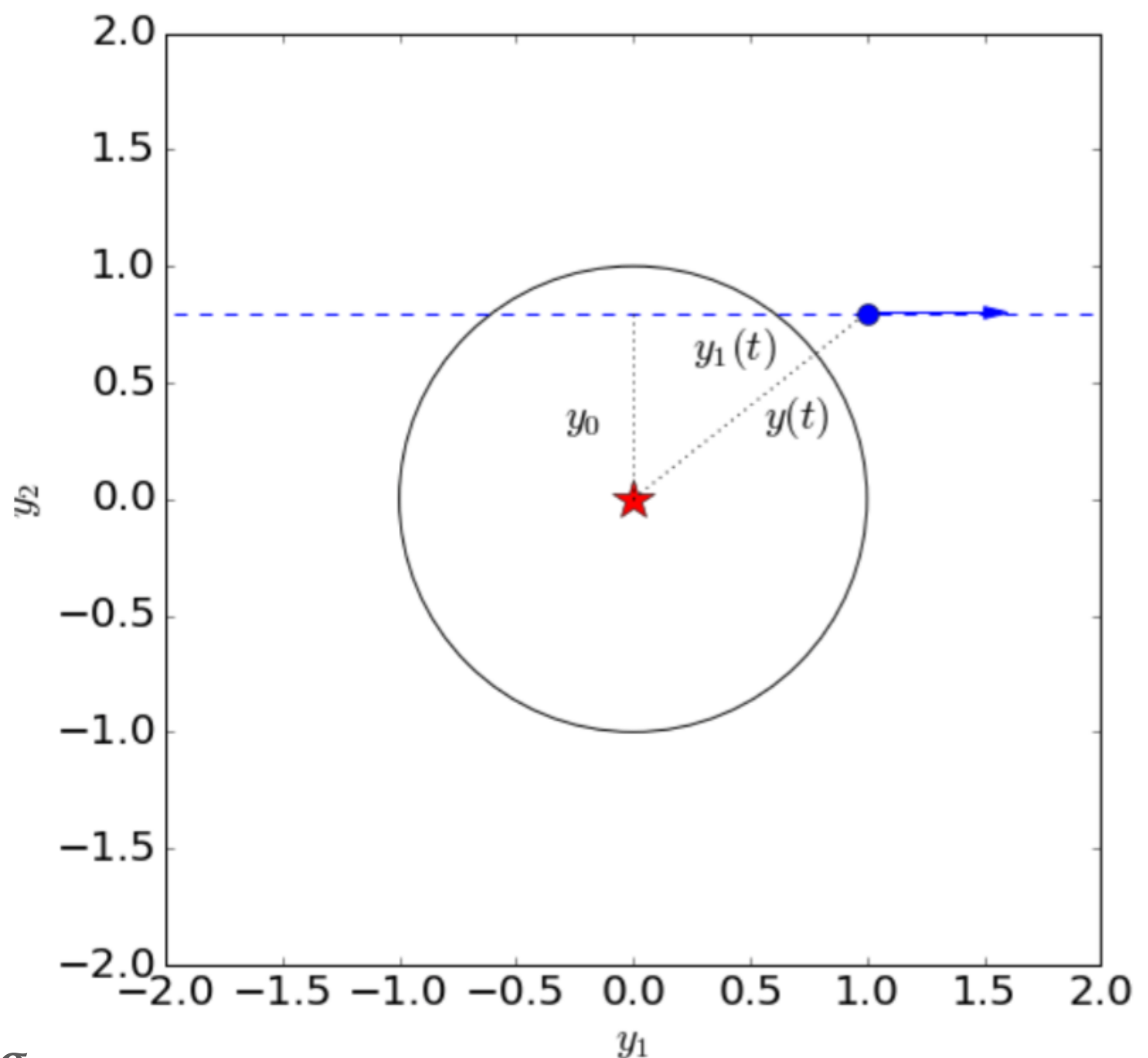
Assume also constant transverse velocity v :

$$y_1(t) = \frac{v(t - t_0)}{D_L \theta_E}$$

We can define a characteristic time of the event:

$$t_E = \frac{D_L \theta_E}{v} = \frac{\theta_E}{\mu_{rel}}$$

This is the Einstein radius crossing time



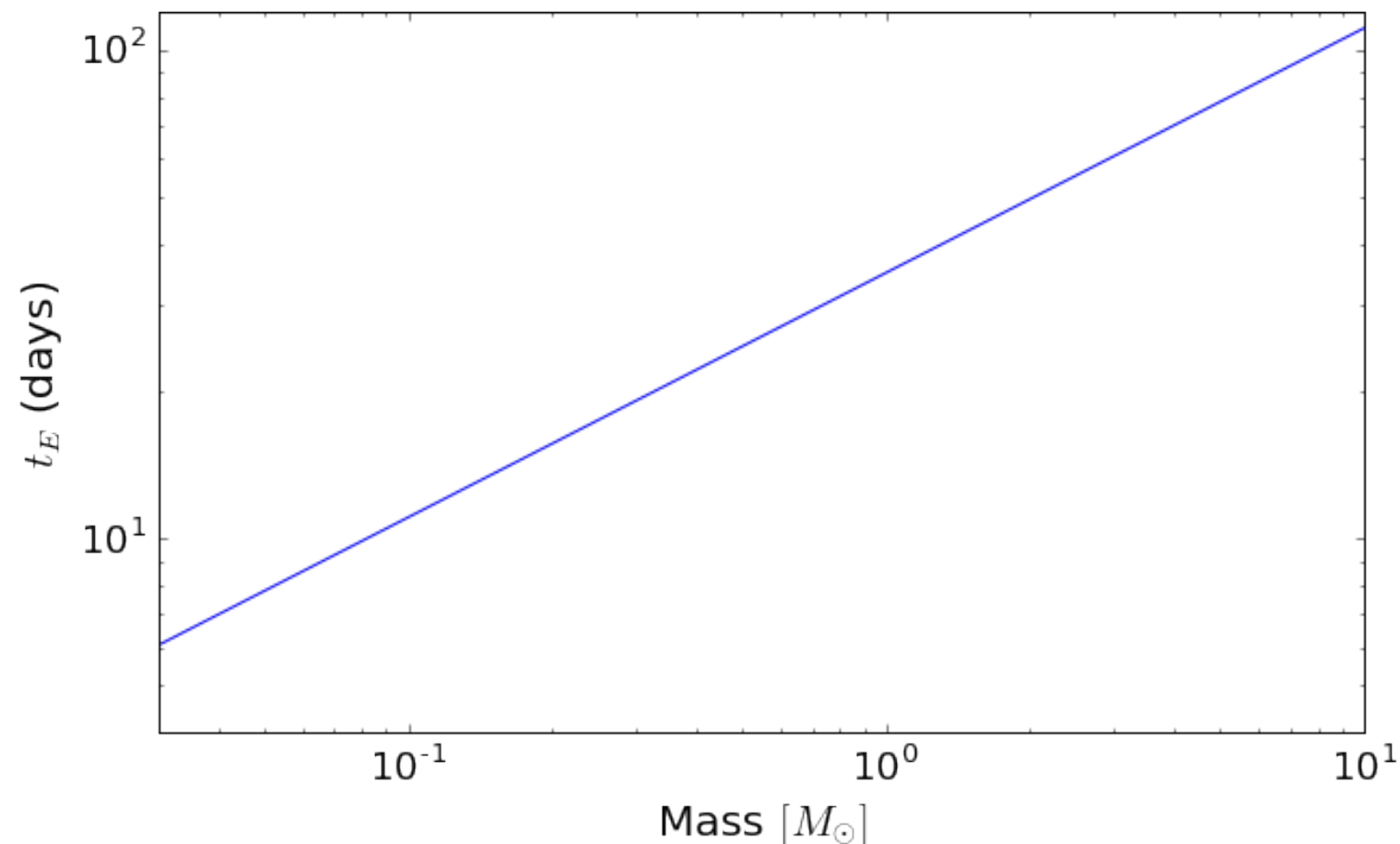
MICROLENSING LIGHT CURVE

Given the definition of Einstein radius

$$\theta_E \equiv \sqrt{\frac{4GM}{c^2} \frac{D_{LS}}{D_L D_S}}$$

The order of magnitude of the t_E is

$$t_E \approx 19 \text{ days} \sqrt{4 \frac{D_L}{D_S} \left(1 - \frac{D_L}{D_S}\right)} \left(\frac{D_S}{8 \text{ kpc}}\right)^{1/2} \left(\frac{M}{0.3 M_\odot}\right)^{1/2} \left(\frac{v}{200 \text{ km/s}}\right)^{-1}$$



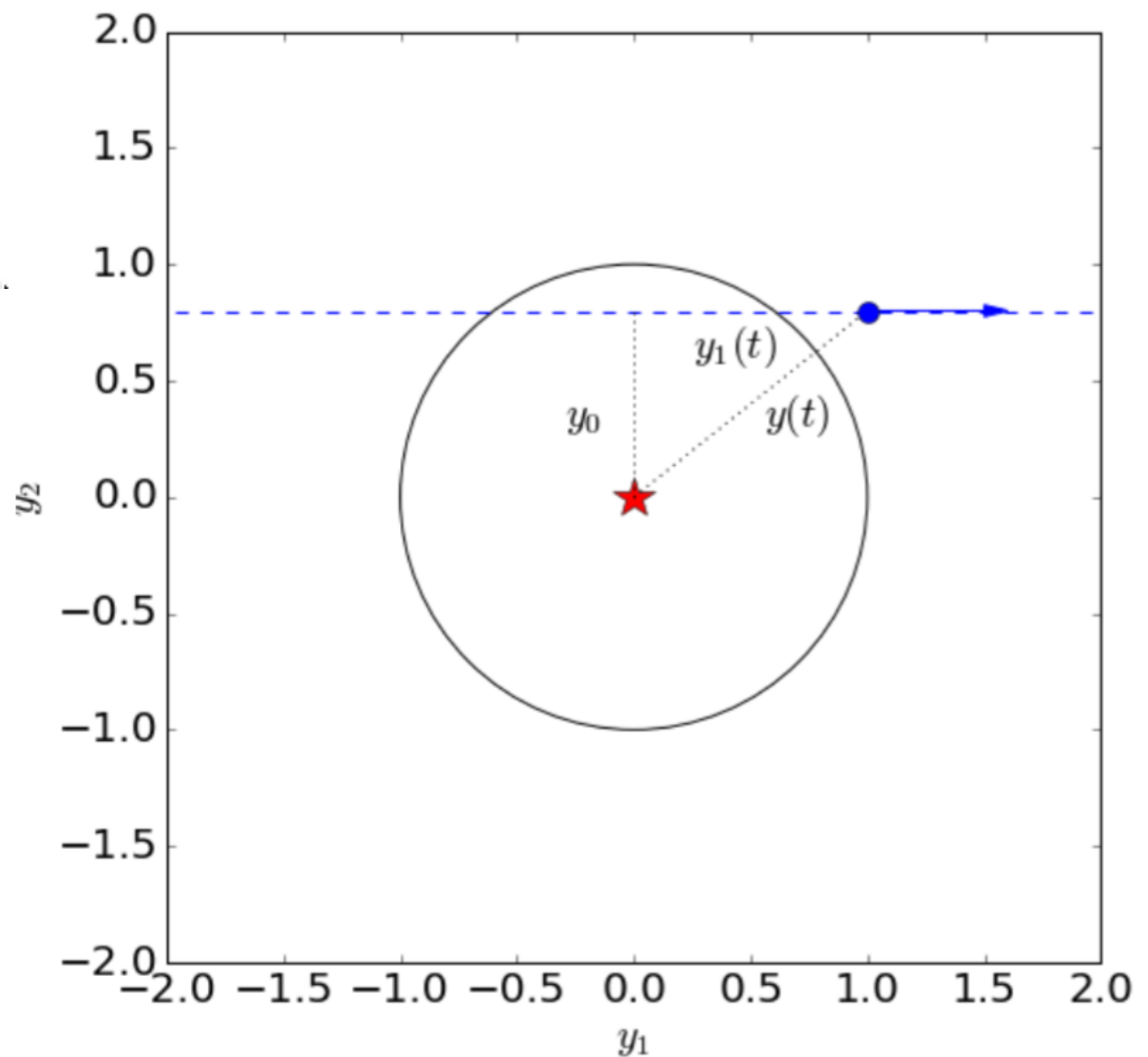
MICROLENSING LIGHT CURVE

We obtain

$$y_1(t) = \frac{v(t - t_0)}{D_L \theta_E} \quad \rightarrow \quad y_1(t) = \frac{(t - t_0)}{t_E}$$

Thus:

$$y(t) = \sqrt{y_0^2 + y_1^2(t)} = \sqrt{y_0^2 + \frac{(t - t_0)^2}{t_E^2}}$$



EXAMPLE OF STANDARD LIGHT CURVE

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