

PERIODIC MOTION II

PHYS1112

Lecture 15

Intended Learning Outcomes

- After this lecture you will learn:
 - 1) various types of harmonic oscillators
 - 2) damping effect on a harmonic oscillator
 - 3) forced oscillator and the phenomenon of resonance

Recap from last lecture

A mechanical system with equation of motion

$$d^2x/dt^2 = -\omega^2x$$

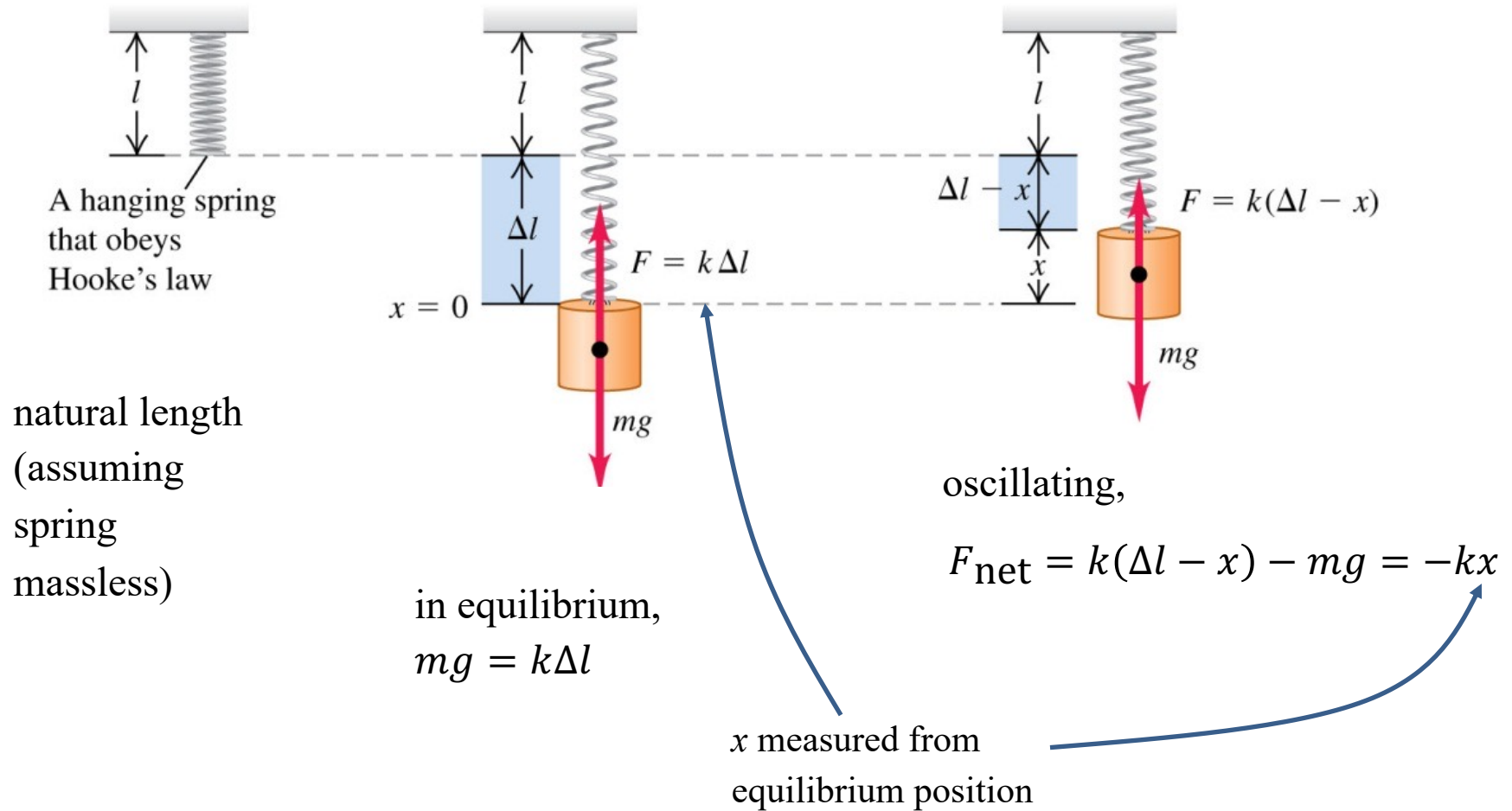
is said to execute simple harmonic motion in the coordinate x , with frequency

$$f = \omega/2\pi$$

and period

$$T = 1/f = 2\pi/\omega$$

Vertical Spring and Mass System



Implication: same as horizontal spring and mass system *if* x measured from equilibrium position, *not* from natural (unextended) position.

Example

Suppose a car's shock absorbers are worn out so that it provides no damping to oscillations. Its mass is 1000 kg. A 100 kg person sits in it and its center of gravity lowers by 2.8 cm. It then hits a bump and start oscillating.

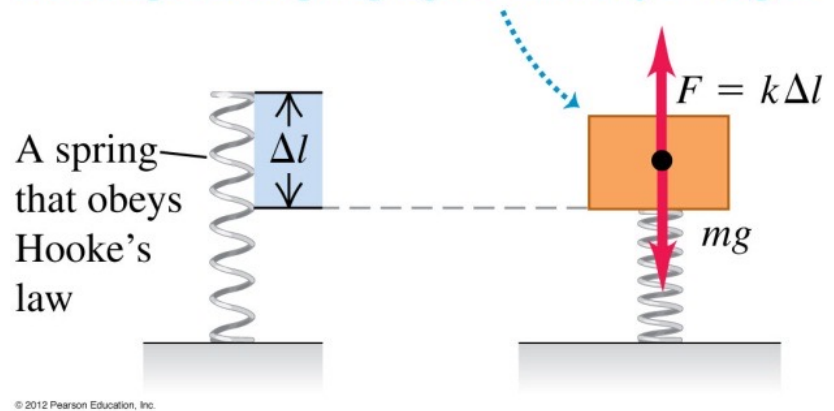
Spring constant

$$k = -\frac{F}{x} = -\frac{980 \text{ N}}{-0.028 \text{ m}} \\ = 3.5 \times 10^4 \text{ kg/s}^2$$

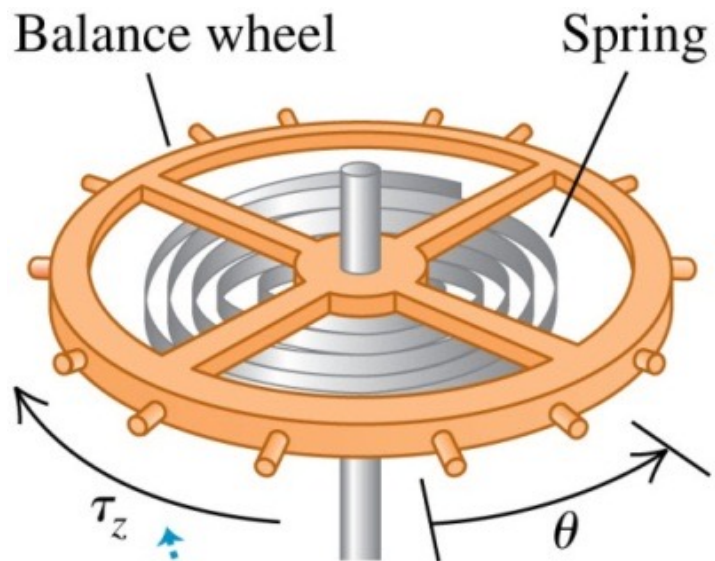
Period of the oscillation

$$T = 2\pi \sqrt{\frac{m}{k}} = 2\pi \sqrt{\frac{1100 \text{ kg}}{3.5 \times 10^4 \text{ kg/s}^2}} \\ = 1.11 \text{ s}$$

A body is placed atop the spring. It is in equilibrium when the upward force exerted by the compressed spring equals the body's weight.



Angular Oscillation



The spring torque τ_z opposes the angular displacement θ .

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Assumption: restoring *torque* proportional to angular displacement

$$\tau = -\kappa\theta = I\alpha \quad \Rightarrow \quad \boxed{\frac{d^2\theta}{dt^2} = -\frac{\kappa}{I}\theta}$$

$$\omega = \sqrt{\frac{\kappa}{I}}, \quad f = \frac{\omega}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{\kappa}{I}},$$

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{I}{\kappa}}$$

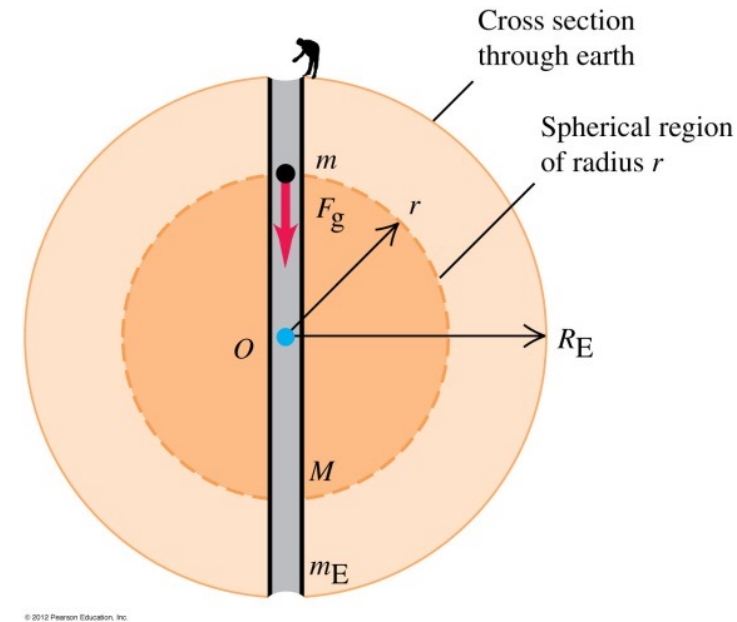
An imaginary case – A Tunnel Through the Earth

From lecture 13, a mass in the tunnel experience a force

$$F_g = \underbrace{-\frac{Gm_E m}{R_E^3} r}_{F_g \text{ opposite to } r} \Rightarrow \frac{d^2 r}{dt^2} = \frac{F_g}{m} = -\underbrace{\frac{Gm_E}{R_E^3}}_{\omega^2} r$$

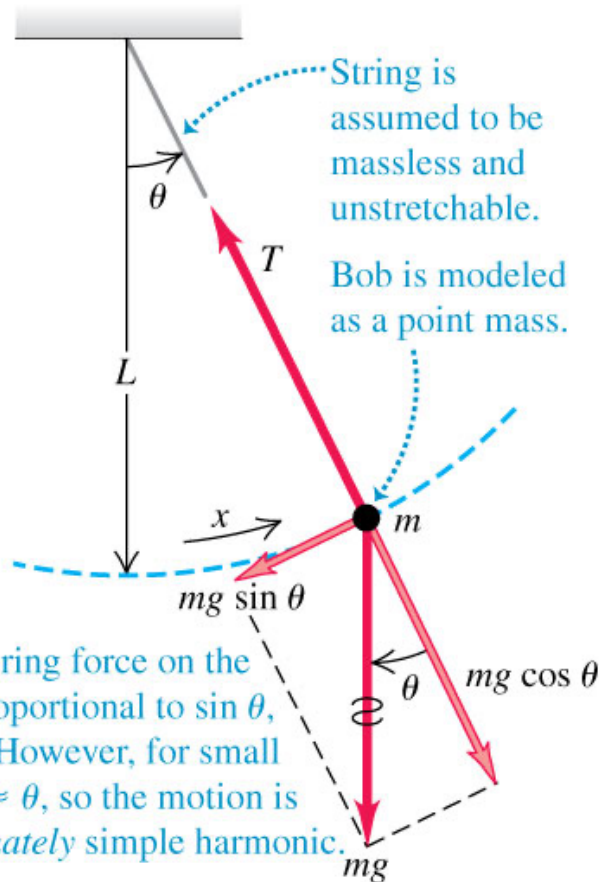
Execute SHM with period with

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{R_E^3}{Gm_E}}$$



Simple Pendulum

Galilei observed that a lamp hung from the ceiling of a church swung with constant period



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x is the arc length

If θ small enough, $\sin \theta \approx \theta = x/L$

In the tangential direction, the restoring force

$$F_{\theta} = -mg \sin \theta \approx -mg \frac{x}{L} = ma_{\text{tan}}$$

$$\Rightarrow \boxed{\frac{d^2 x}{dt^2} = -\frac{g}{L} x}$$

$$\omega = \sqrt{\frac{g}{L}}, \quad f = \frac{\omega}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{g}{L}},$$

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{L}{g}}$$

Q14.9

A simple pendulum consists of a point mass suspended by a massless, unstretchable string. If the mass is doubled while the length of the string remains the same, the period of the pendulum

- A. becomes four times greater.
- B. becomes twice as great.
- C. becomes greater by a factor of $\sqrt{2}$.
- D. remains unchanged.
- E. decreases.

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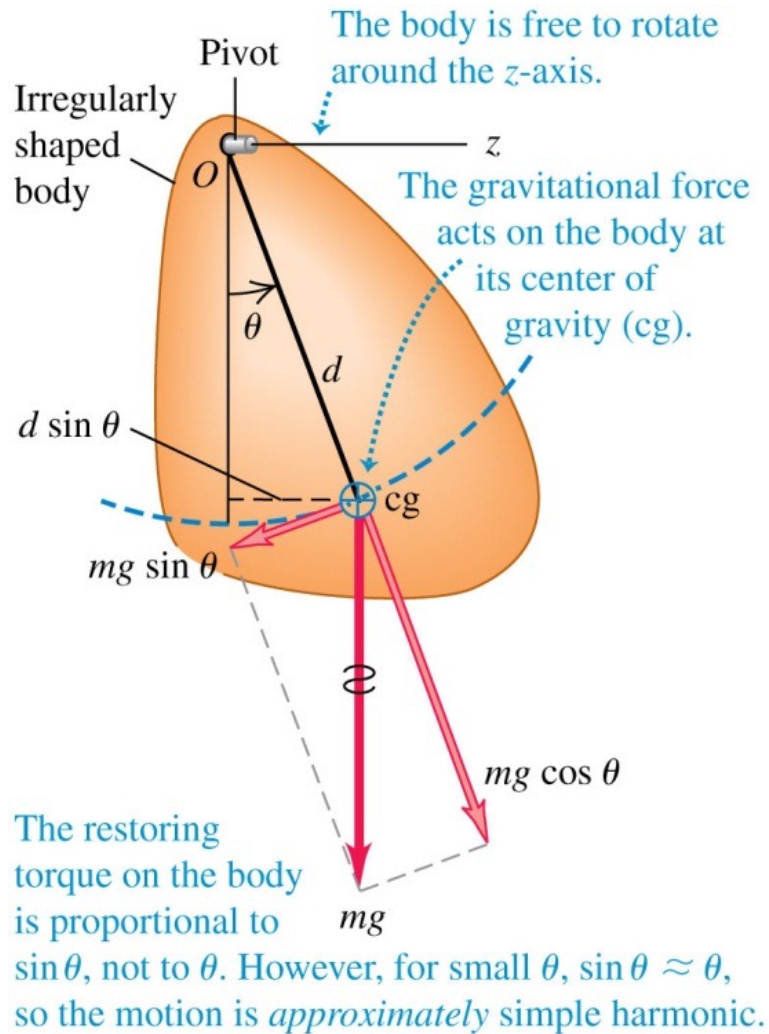
C. becomes greater by a factor of $\sqrt{2}$.



D. remains unchanged.

E. decreases.

Physical Pendulum



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In small θ limit:

$$\tau = -(mg)d \sin \theta \approx -(mgd)\theta = I\alpha$$

$$\Rightarrow \boxed{\frac{d^2 \theta}{dt^2} = -\frac{mgd}{I} \theta}$$

$$\omega = \sqrt{\frac{mgd}{I}},$$

$$f = \frac{\omega}{2\pi} = \frac{1}{2\pi} \sqrt{\frac{mgd}{I}},$$

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{I}{mgd}}$$

Demonstration:

Simple Pendulum vs. Meter Stick (same center of mass, different periods)



Question:

The CG of a simple pendulum of mass m and length L is at a distance L from its pivot. The CG of a uniform rod of mass m and length $2L$, pivoted at one end, is also at a distance L from its pivot. The period of the rod is (longer / shorter / the same) as the period of the pendulum.

Damped Harmonic Oscillator

Suppose oscillator experience fluid resistance at low speed, $f = -bv$ (oppose to v)

$$ma = -kx - bv,$$

or

$$m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

2nd order *differential equation*

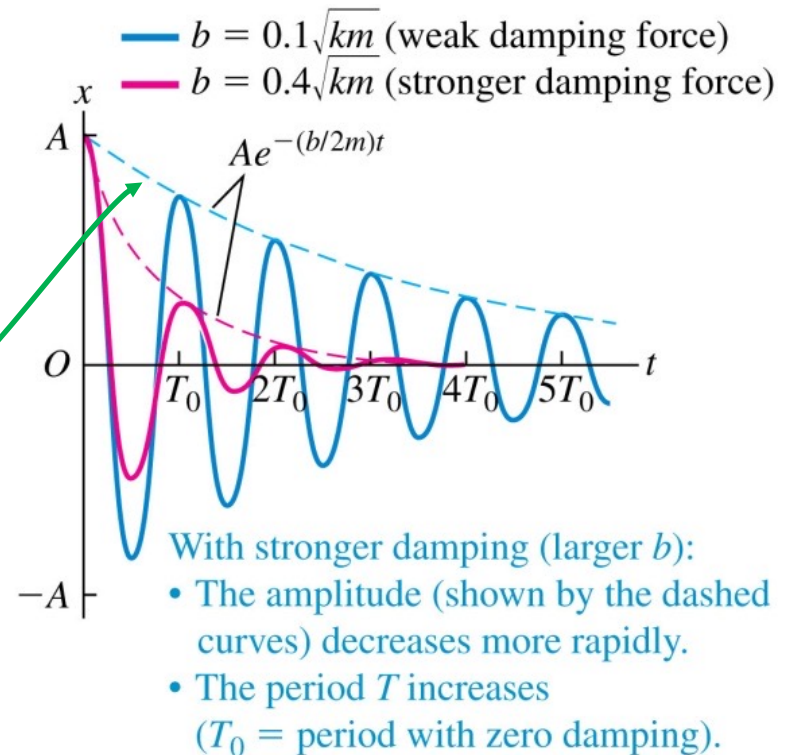
Solution:

$$x(t) = \underbrace{Ae^{-(\frac{b}{2m})t}}_{\text{exponentially decaying amplitude, or envelop, damp out the oscillation}} \underbrace{\cos(\omega' t + \phi)}_{\text{periodic oscillation with angular frequency } \omega'}$$

exponentially decaying
amplitude, or envelop,
damp out the oscillation

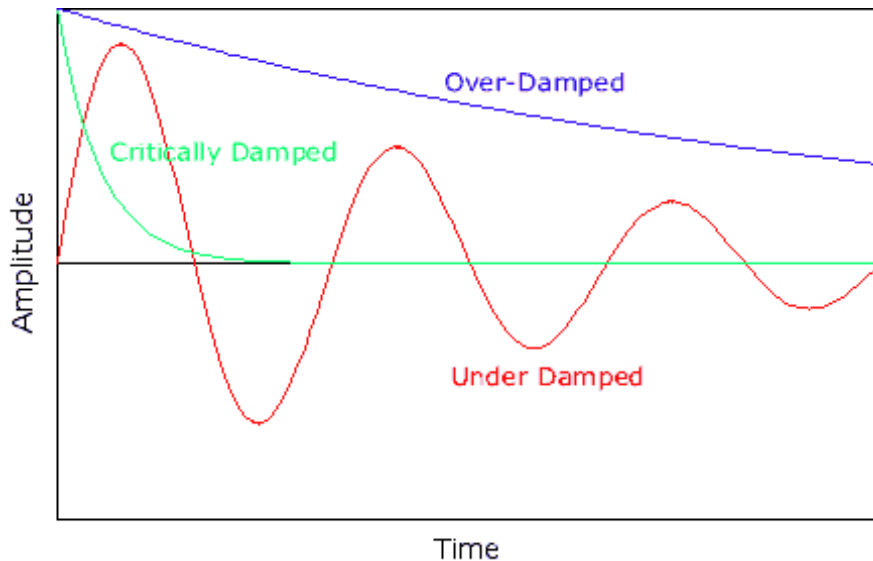
periodic oscillation with
angular frequency ω'

$$\omega' = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$$



$$x(t) = Ae^{-\left(\frac{b}{2m}\right)t} \cos(\omega' t + \phi)$$

$$\omega' = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$$



When $b < 2\sqrt{km}$, $\omega' > 0$, called **underdamping**.

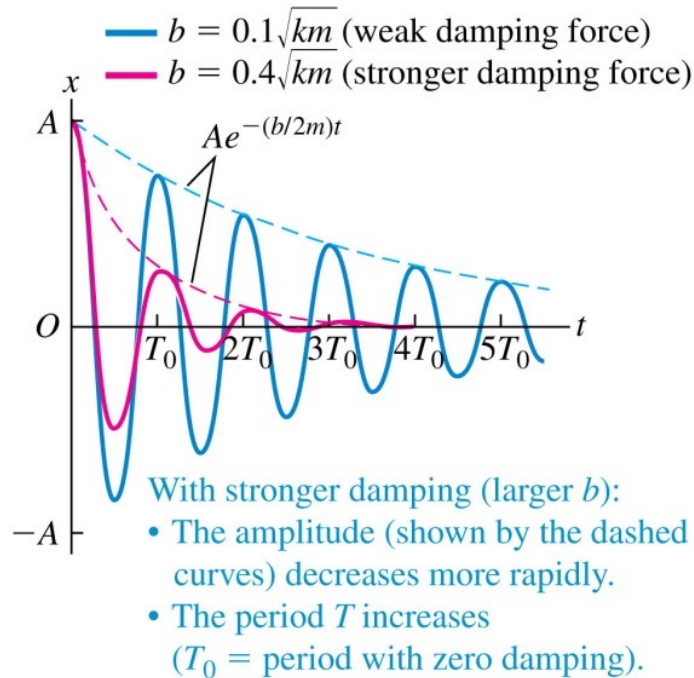
Stronger damping (larger b), oscillation dies off faster

When $b > 2\sqrt{km}$, ω' imaginary (only if you know complex numbers), no oscillation, solution becomes

$$x(t) = C_1 e^{-a_1 t} + C_2 e^{-a_2 t}$$

When $b = 2\sqrt{km}$, $\omega' = 0$, no oscillation, return to equilibrium position in shortest time, called **critical damping**

Rate of energy dissipation



$$ma = -kx - bv$$

$$\frac{dE}{dt} = \frac{d}{dt} \left(\frac{1}{2}mv^2 + \frac{1}{2}kx^2 \right) = mv \underbrace{\frac{dv}{dt}}_a + kx \underbrace{\frac{dx}{dt}}_v = v \underbrace{(ma + kx)}_{-bv} = -bv^2$$

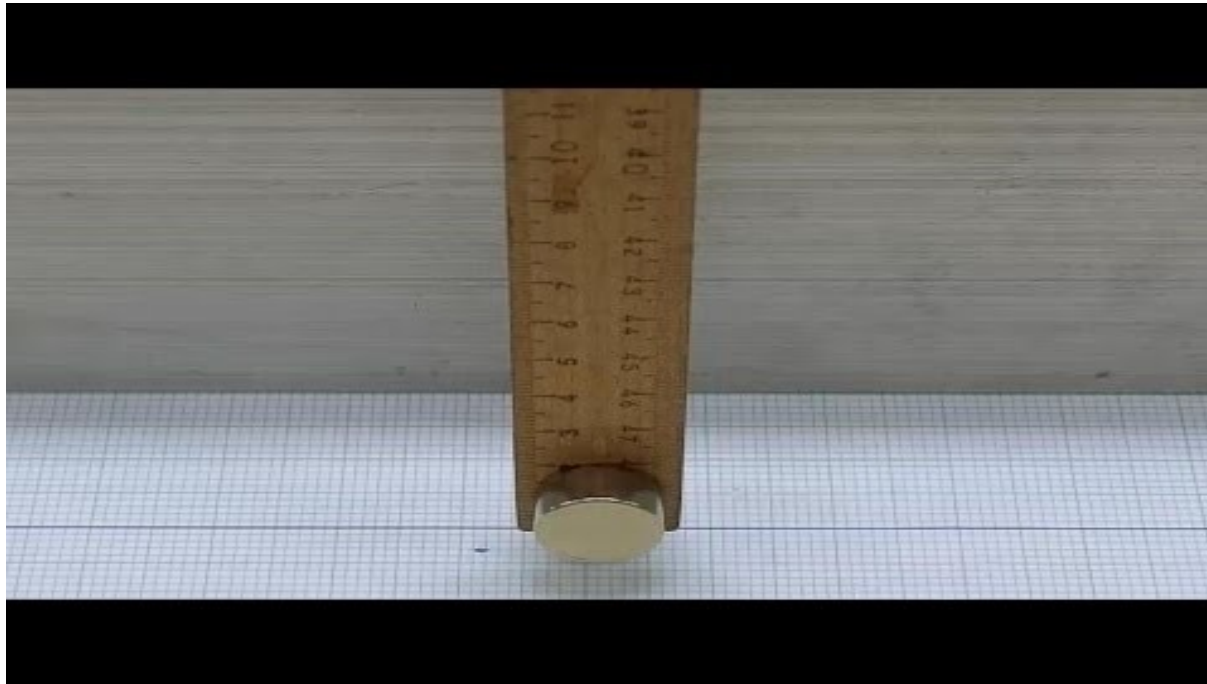
⚠ $dE/dt < 0$ means energy decreases, i.e., dissipation

⚠ consistent with $dE/dt = Fv = (-bv)v$

Demonstration

Under, critical, and overdamping

<https://www.youtube.com/watch?v=99ZE2RGwgSM>



Forced Oscillations

Suppose an external periodic driving force $F(t) = F_{\max} \cos \omega_d t$

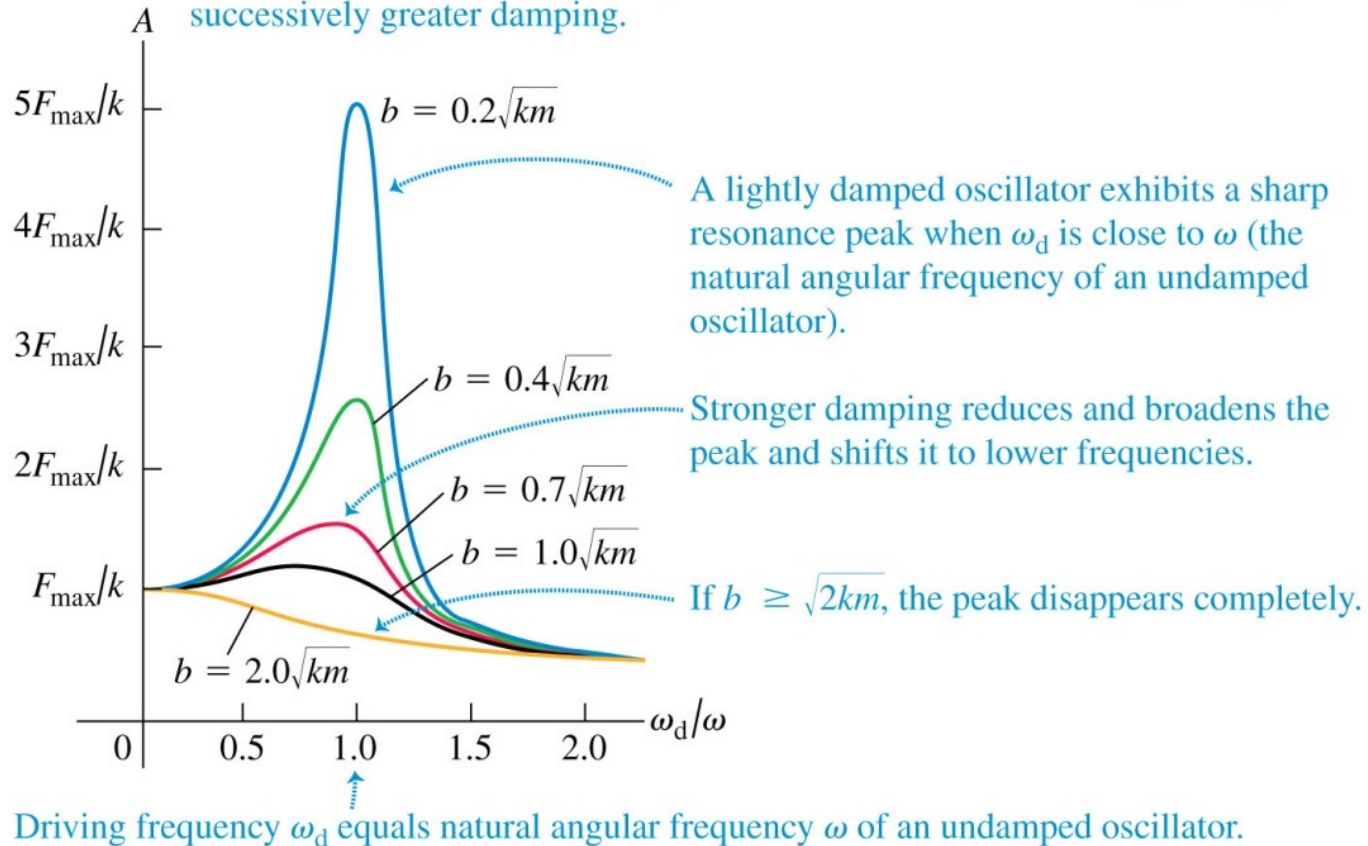
$$ma = -kx - bv + F(t)$$

$$\Rightarrow \boxed{m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx = F_{\max} \cos \omega_d t}$$

Oscillator vibrates with ω_d . What if $\omega_d \rightarrow \omega = \sqrt{k/m}$, the **natural frequency** of the free (undamped) oscillator?

Driving force deposit energy into the “natural mode” of vibration, expect amplitude to increase

Each curve shows the amplitude A for an oscillator subjected to a driving force at various angular frequencies ω_d . Successive curves from blue to gold represent successively greater damping.



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$$A = \frac{F_{\max}}{\sqrt{(k - m\omega_d^2)^2 + b^2\omega_d^2}}$$

A reaches maximum when $\omega_d \approx \sqrt{k/m}$, called **resonance**

- ⚠ when damping b small, peak is higher, sharper, and closer to the **natural frequency** of the free (undamped) oscillator
- ⚠ resonance vibration can be strong enough to make bridges and buildings collapse

Demonstration

forced oscillation and resonance of a wine glass

<http://www.youtube.com/watch?v=EjVq96h4zD4>

