

## Thermal Properties of matter II

Heat capacity of ideal gas

Recall: Molar heat capacity  $= C = \frac{1}{n} \frac{dQ}{dT}$

Consider a gas with fixed  $V$ .

Heat transfer to the gas increases the kinetic energy of the gas.

$$\Delta Q = \Delta K \quad \text{from kinetic theory } K = \frac{3}{2} nRT.$$

$$\Rightarrow \Delta Q = \frac{3}{2} nR \Delta T$$

$$\text{Molar heat capacity} = C_v = \frac{1}{n} \frac{\Delta Q}{\Delta T} = \frac{3}{2} R = 12.47 \text{ J/mol}\cdot\text{K}.$$

at fixed  $V$

True for all kinds of ideal gas?

From experiment

Gas

$C_v$  J/mol·K

Ne

12.47

$\sim \frac{3}{2} R = 12.47$

Ar

12.47

H<sub>2</sub>

20.42

O<sub>2</sub>

20.85

N<sub>2</sub>

20.76

CO

20.85

$\sim \frac{5}{2} R = 20.78$

Seems to have some patterns

Recall how  $K = \frac{3}{2} nRT$  is obtained.

for gas consisting of non-interacting point particles,

$$E_{\text{tot}} = K_{\text{tr}} = \sum_i \frac{1}{2} m \vec{v}_i^2$$

$$= \underbrace{\sum_i \frac{1}{2} m v_{ix}^2}_{\frac{1}{2} pV} + \underbrace{\sum_i \frac{1}{2} m v_{iy}^2}_{\frac{1}{2} pV} + \underbrace{\sum_i \frac{1}{2} m v_{iz}^2}_{\frac{1}{2} pV}$$

$$= \frac{3}{2} pV$$

$$= \frac{3}{2} nRT$$

$\therefore$  ideal gas law  
 $pV = nRT$

$$E_{\text{tot}} = \frac{3}{2} N k_B T$$

$$\text{every energy per particle/molecule} = \frac{E_{\text{tot}}}{N} = \frac{3}{2} k_B T = \underbrace{\frac{1}{2} k_B T}_{v_x} + \underbrace{\frac{1}{2} k_B T}_{v_y} + \underbrace{\frac{1}{2} k_B T}_{v_z}$$

- each particle/molecule needs 3 parameters to specify its energy:  
( $v_x, v_y, v_z$ )

We refer this number of parameters as degrees of freedom

- each degree of freedom carries  $\frac{1}{2} k_B T$  amount of energy per molecule. This principle is called

Principle of equipartition.

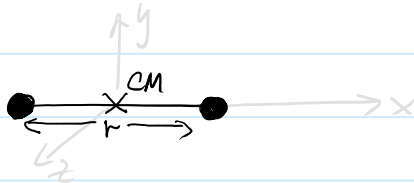
## Applications of equipartition principle.

For an ideal gas consisting of molecules with  $f$  degrees of freedom,

$$\begin{aligned}\text{Total energy of the gas} &= E_{\text{tot}} = f \times \frac{1}{2} k_B T \times N \\ &= \frac{f}{2} N k_B T \approx \frac{f}{2} n R T.\end{aligned}$$

$$\Rightarrow \boxed{C_v = \frac{f}{2} R}$$

Example. Diatomic ideal gas. e.g.  $\text{H}_2$ ,  $\text{N}_2$ ,  $\text{O}_2$ ,  $\text{CO}$ , ...



Possible

Translation of CM (x 3)  
 $V_x, V_y, V_z$ .

Rotation about  $y$  and  $z$  axis. (x 2)  
 $\omega_y, \omega_z$

Vibration of the two atoms (x 2)  
(each vibration needs 2 parameters to specify the potential & kinetic energy).

$$\text{Degree of freedom} = 3 + 2 + 2 = 7.$$

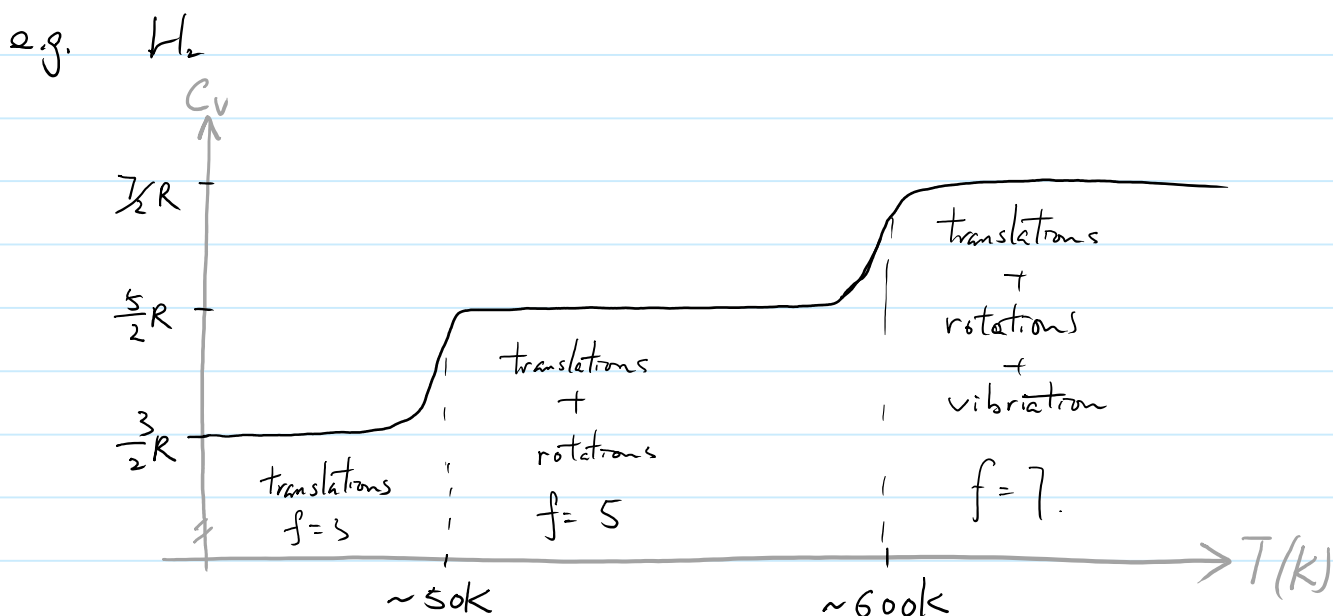
for monoatomic gas,  $f = 3$ ,  $C_v = \frac{3}{2}R = \underline{12.465 \text{ J/mol}\cdot\text{K}}$

(agree well with data  
see table on p.2)

for diatomic gas,  $f = 7$ ,  $C_v = \frac{7}{2}R = 29 \text{ J/mol}\cdot\text{K}$ .

but most of them on the table have  $C_v \approx 20.7 \text{ J/mol}\cdot\text{K}$   
what's wrong?

### Temperature dependence of $C_v$



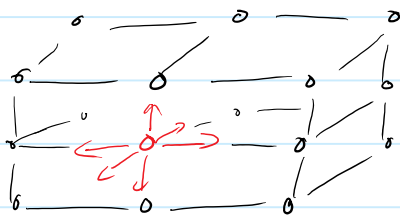
At room temperature, the vibration of diatomic molecules is not probable. So, degrees of freedom equals 5.

$$\Rightarrow C_v = \frac{5}{2}R \approx \underline{20.78 \text{ J/mol}\cdot\text{K}}$$

agree with experiments done in  
room temperature.

In solid, molecules are located at a lattice site.

They can vibrate in 3 directions but not translate.

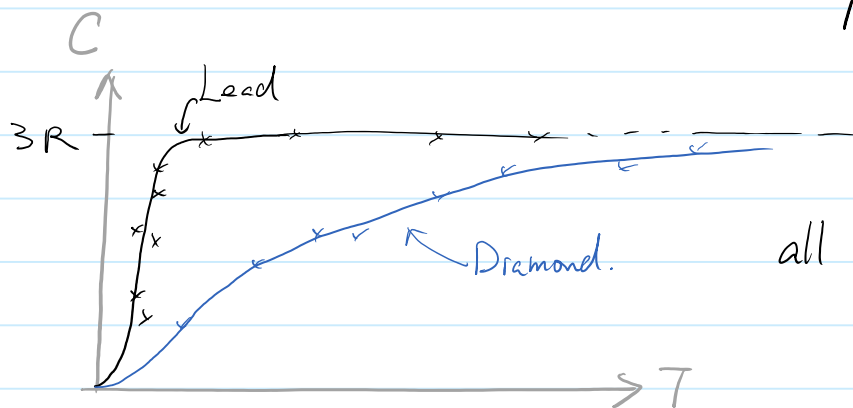


3 vibrations

$$\Rightarrow f = 6$$

$$\Rightarrow C = \frac{6}{2}R = 3R.$$

when all 3 vibrations are probable.



all saturate at  $C = 3R$ .  
at high  $T$ .

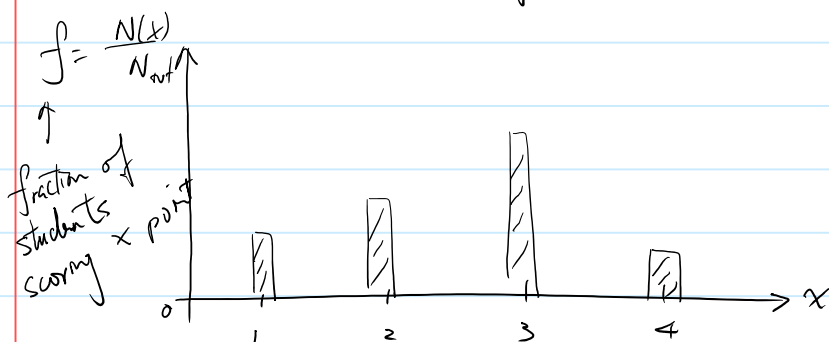
Molecular speed of gas

Speed distribution in a gas. (histogram of a continuous value)

Analogy: score distribution in class

$N(x)$  = # of students score  $x$  point

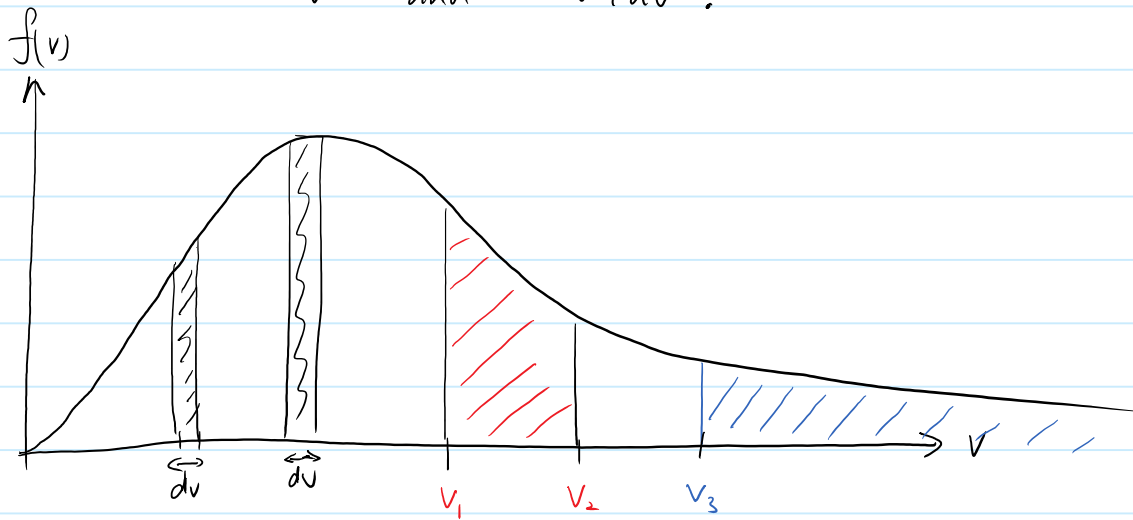
$N_{tot}$  = total # of students in the class



what if the  $x$  is continuous?

$v$  = speed of gas molecule

$f(v)dv$  = fraction of gas molecules having a speed between  $v$  and  $v+dv$ .



Probability / fraction of molecules having speed between  $v_1$  &  $v_2$

$$= \int_{v_1}^{v_2} f(v) dv = \text{Area under the curve from } v_1 \text{ to } v_2.$$

Probability / fraction of molecules having speed larger than  $v_3$

$$= \int_{v_3}^{\infty} f(v) dv.$$

By definition, probability / fraction of molecules having speed

from 0 to  $\infty$  = 100% = 1

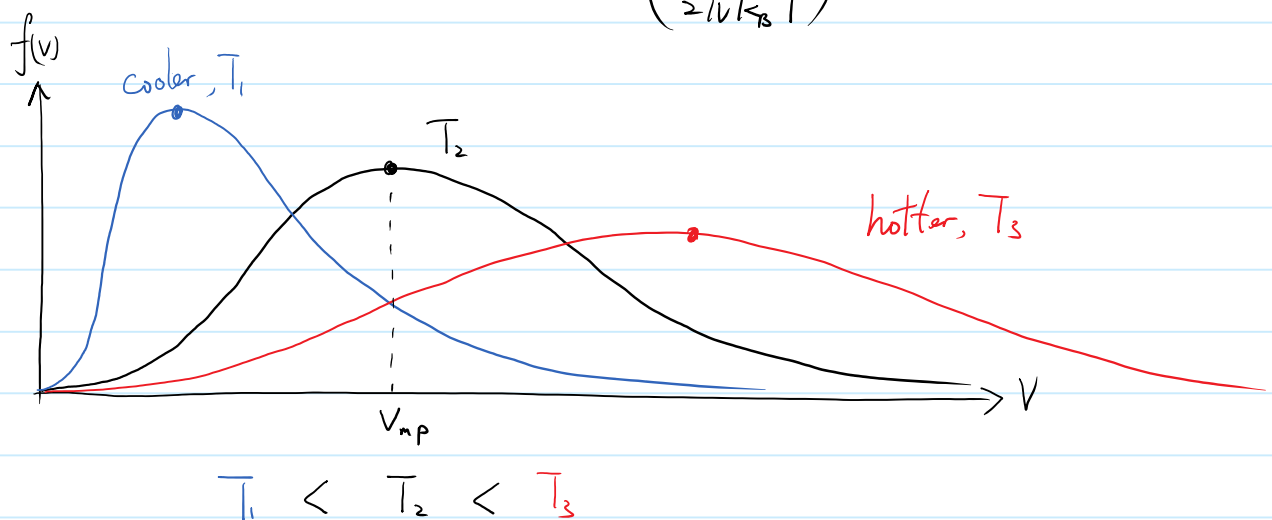
$$\Rightarrow \int_0^{\infty} f(v) dv = 1$$

$$\text{Average speed} = \int_0^{\infty} v \cdot f(v) \cdot dv$$

$$V_{\text{rms}} = \left[ \overline{(V^2)}_{\text{avg}} \right]^{\frac{1}{2}} = \left[ \int_0^{\infty} v^2 f(v) dv \right]^{\frac{1}{2}}$$

## Maxwell-Boltzmann Distribution.

for ideal gas:  $f(v) = 4\pi \left( \frac{m}{2\pi k_B T} \right)^{\frac{3}{2}} v^2 e^{-\frac{1}{2} \frac{mv^2}{k_B T}}$



The area under the curve must be 1.

$\Rightarrow$  longer the tail, lower the peak.

Most populated speed =  $V_{\text{mp}}$

To find  $V_{\text{mp}}$ :  $\left. \frac{df}{dv} \right|_{v=V_{\text{mp}}} = 0$

$$\Rightarrow 2v e^{-\frac{1}{2} \frac{mv^2}{k_B T}} + v^2 e^{-\frac{1}{2} \frac{mv^2}{k_B T}} \cdot \left( -\frac{mv}{k_B T} \right) = 0$$

$$\Rightarrow 2v = \frac{mv^3}{k_B T} \Rightarrow V_{\text{mp}} = \sqrt{\frac{2k_B T}{m}}$$

The fact that at any temperature there are always some fraction of molecules which have very high speed is the reason why evaporation could occur.

Imagine a cup of water at room temperature.

The water molecules could have a wide range of speeds.

Some of the molecules will have enough speed to escape

from the inter-molecular force holding them together. Those

molecules will fly away from the water. In principle,

the water becomes cooler due to the loss of

the high energy molecules. But the temperature will

be equilibrate by the environment and, practically

remains at room temperature. So, there will be

the same fraction of molecules which could evaporate.

The process will continue until, in principle, only

infinitesimal amount of water left in the cup given

long enough times.