

# **DYNAMICS OF RIGID BODIES II**

PHYS1112

Lecture 9

# Intended Learning Outcomes

- After this lecture you will learn:
  - 1) how to calculate the moment of inertia of simple symmetric rigid bodies
  - 2) the parallel axis theorem to find the moment of inertia about different rotation axis
  - 3) Vector product
  - 4) torque, and the Newton's second law in rotational dynamics

## Parallel axis theorem

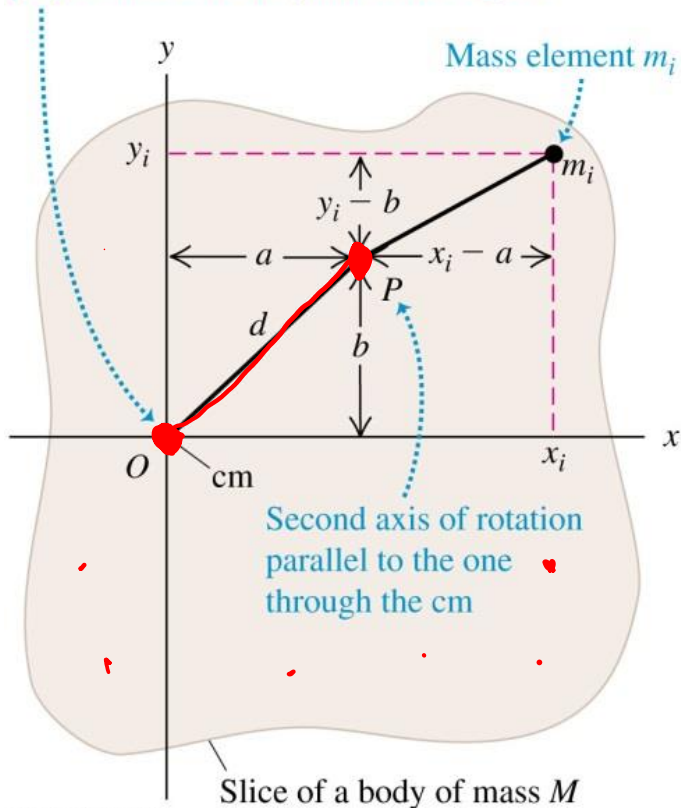
$I_{\text{cm}}$ : moment of inertia about an axis through its CM

$I_p$ : moment of inertia about another axis  $\parallel$  to the original one and at  $\perp$  distance  $d$

$$I_p = I_{\text{cm}} + Md^2$$

Proof: take CM as the origin, rotation axis as the  $z$  axis. A point mass  $m_i$  in the solid has coordinates  $(x_i, y_i, z_i)$

Axis of rotation passing through cm and perpendicular to the plane of the figure



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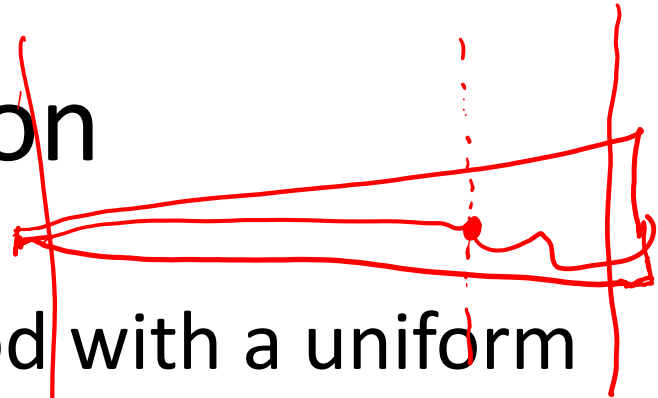
square of  $\perp$  distance of  $m_i$  to rotation axis

$$I_{\text{cm}} = \sum m_i \overbrace{(x_i^2 + y_i^2)}^{\text{not } \sum m_i(x_i^2 + y_i^2 + z_i^2)}$$

$$\begin{aligned} I_p &= \sum m_i [(x_i - a)^2 + (y_i - b)^2] \\ &= \underbrace{\sum m_i (x_i^2 + y_i^2)}_{I_{\text{cm}}} - 2a \underbrace{\sum m_i x_i}_{Mx_{\text{cm}} = 0} - 2b \underbrace{\sum m_i y_i}_{My_{\text{cm}} = 0} \end{aligned}$$

$$+ \underbrace{(a^2 + b^2)}_{d^2} \underbrace{\sum m_i}_M = \frac{\sum m_i x_i}{M} = x_{\text{cm}} = 0$$

# Question



- A pool cue is a wooden rod with a uniform composition and tapered with a larger diameter at one end than at the other end.

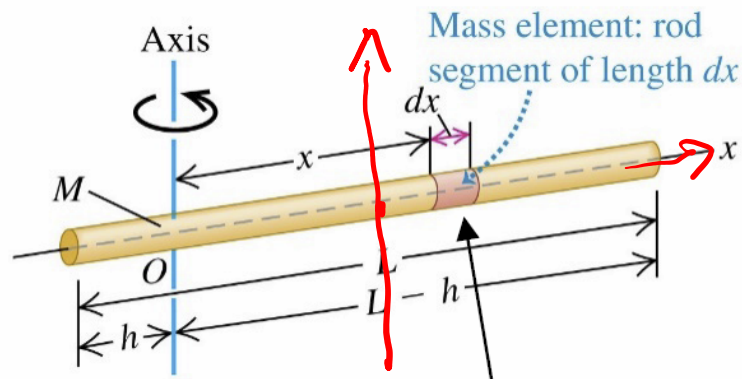
$$I = \sum_i m_i x_i^2 + \underline{M d^2}$$

Does it have a larger moment of inertia

- ① for an axis through the thicker end of the rod and perpendicular to the length of the rod, or
- ② for an axis through the thinner end of the rod and perpendicular to the length of the rod?

**Significance of the  
parallel axis theorem:**  
need formula for  $I_{cm}$  only

Example A thin rod with uniform linear density  $\rho = M/L$



⚠ Before calculating moment of inertia, must specify rotation axis

$$I = \sum m_i r_i^2 \rightarrow \int \underbrace{r^2}_{\perp \text{ distance of } m_i \text{ to rotation axis}} dm$$

$\perp$  distance of  $m_i$  to rotation axis

Choose a convenient mass element  $dm$  – a segment of length  $dx$  at a  $\perp$  distance  $x$  from the axis, and mass  $dm = \rho dx$

$$I_O = \int_{-h}^{L-h} x^2 (\rho dx) = \frac{\rho}{3} [(L-h)^3 + \underbrace{h^3}] = \frac{M}{3} (L^2 - 3Lh + 3h^2)$$

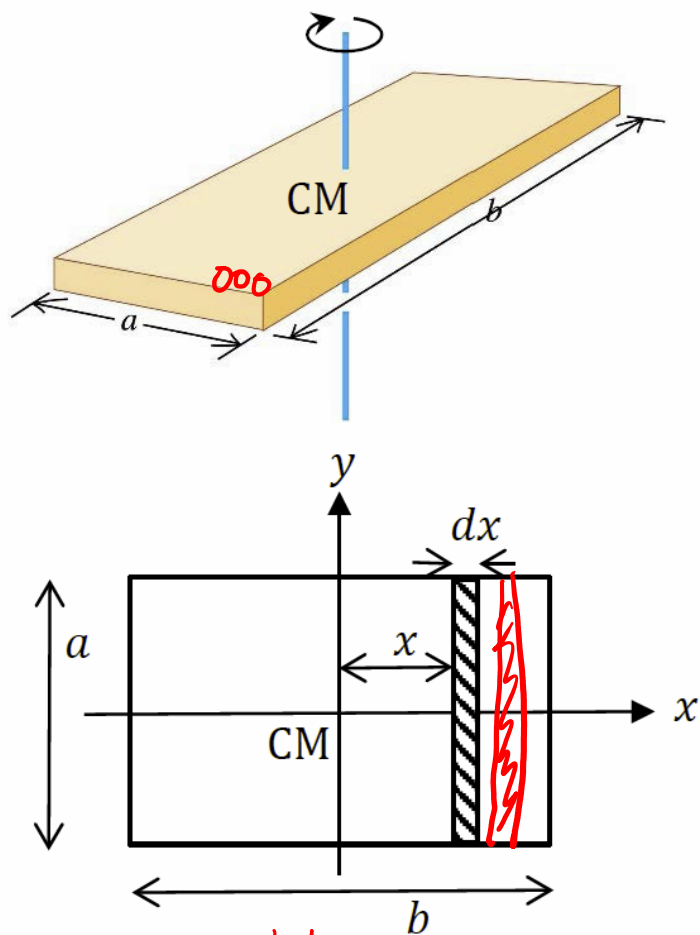
Handwritten red notes:

$$\int_{-h}^{L-h} x^2 dx = \frac{x^3}{3} \Big|_{-h}^{L-h} \quad (L-h)^3 = \frac{(L-h)(L-h)^2}{1} = L^3 - 3Lh(L-h) + h^3$$

⚠ Put  $h = L/2$ , we get  $I_{\text{cm}} = ML^2/12$ .

⚠ Check the parallel axis theorem  $I_O = I_{\text{cm}} + M(\quad)^2$

**Example** A rectangular plate



Choose the mass element  $dm$  to be a rod at  $\perp$  distance  $x$  from the axis. *Why? Because you know its moment of inertia!*

$$dI = \underbrace{\frac{(dm)a^2}{12}}_{\text{about CM of the rod, not of the plate}} + \underbrace{(dm)x^2}_{\text{parallel axis theorem}}$$

about CM of the rod,  
not of the plate

parallel axis  
theorem

Since  $dm = \left(\frac{M}{b}\right) dx$   $\frac{M}{b} \frac{a^2}{12} \cdot b$

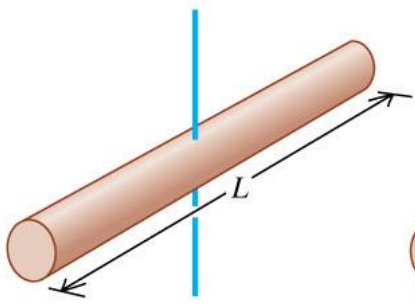
$$I = \int dI = \frac{M}{b} \int_{-b/2}^{b/2} \left[ \frac{a^2}{12} + x^2 \right] dx = \frac{1}{12} M(a^2 + b^2)$$

$\frac{M}{b} \int_{-b/2}^{b/2} x^2 dx = \frac{M}{b} \frac{x^3}{3} \Big|_{-b/2}^{b/2} = \frac{M}{b} \cdot 2 \frac{b^3}{3 \cdot 8}$

**Table 9.2 Moments of Inertia of Various Bodies**

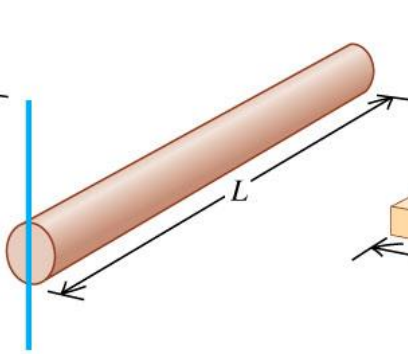
(a) Slender rod,  
axis through center

$$I = \frac{1}{12} ML^2$$



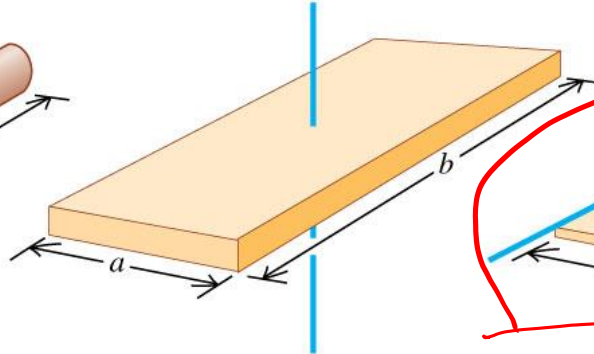
(b) Slender rod,  
axis through one end

$$I = \frac{1}{3} ML^2$$



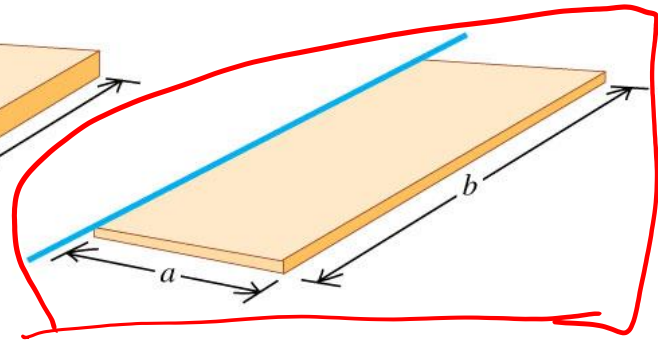
(c) Rectangular plate,  
axis through center

$$I = \frac{1}{12} M(a^2 + b^2)$$



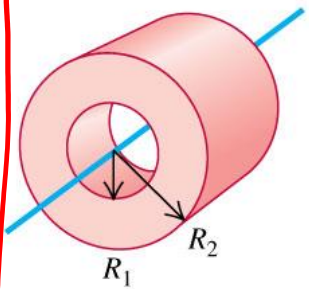
(d) Thin rectangular plate,  
axis along edge

$$I = \frac{1}{3} Ma^2$$



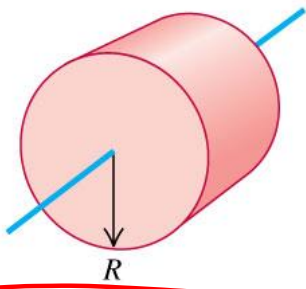
(e) Hollow cylinder

$$I = \frac{1}{2} M(R_1^2 + R_2^2)$$



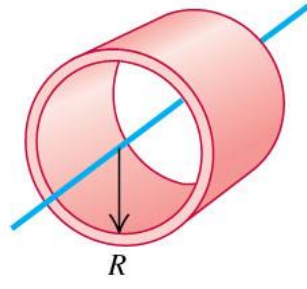
(f) Solid cylinder

$$I = \frac{1}{2} MR^2$$



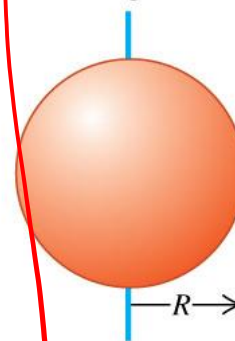
(g) Thin-walled hollow  
cylinder

$$I = MR^2$$



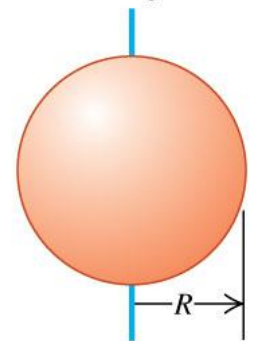
(h) Solid sphere

$$I = \frac{2}{5} MR^2$$



(i) Thin-walled hollow  
sphere

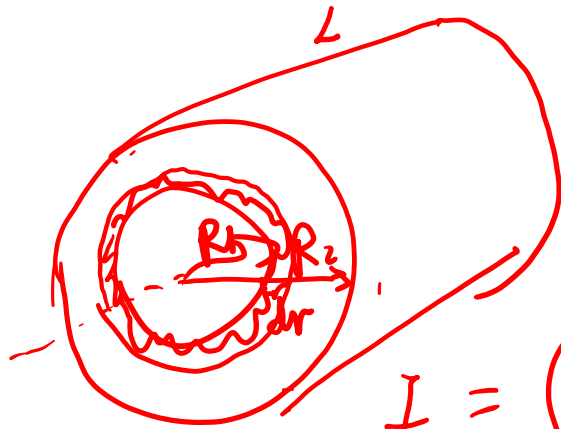
$$I = \frac{2}{3} MR^2$$



$R_1 = 0$

$R_1 = R_2 = R$





$$I = \int r^2 dm$$

$$dm = \rho dV = \rho 2\pi r dr \cdot L$$

$$I = \int_{R_1}^{R_2} 2\pi \rho L r^3 dr = \frac{2\pi \rho L}{4} r^4 \bigg|_{R_1}^{R_2}$$

$$= \frac{\pi \rho L}{2} (R_2^4 - R_1^4)$$

$$\rho = \frac{M}{V} = \frac{M}{\pi (R_2^2 - R_1^2) \cdot L} = \frac{\pi \cancel{L} \frac{M}{\pi (R_2^2 - R_1^2) \cancel{L}}}{2} (R_2^4 - R_1^4)$$

$$I = \frac{1}{2} M (R_2^2 + R_1^2)$$

torque

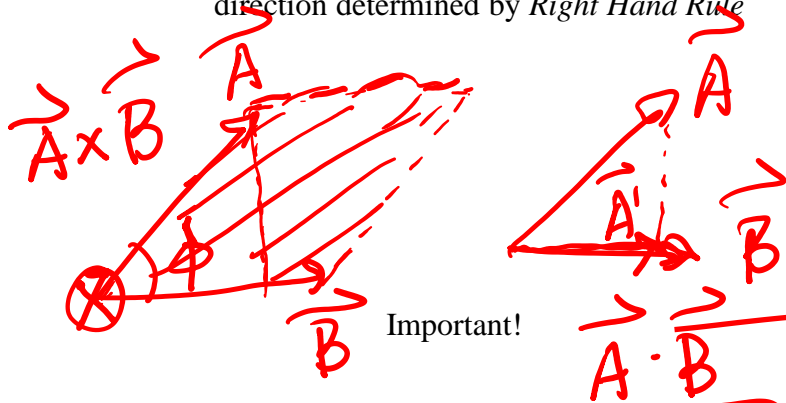
# Vector (Cross) Product

$$\vec{A} \cdot \vec{B} = AB \cos \phi = \vec{B} \cdot \vec{A}$$

$$\vec{C} = \vec{A} \times \vec{B}$$

Magnitude:  $C = AB \sin \phi$

direction determined by *Right Hand Rule*



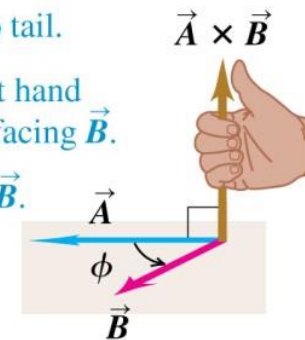
(a) Using the right-hand rule to find the direction of  $\vec{A} \times \vec{B}$

① Place  $\vec{A}$  and  $\vec{B}$  tail to tail.

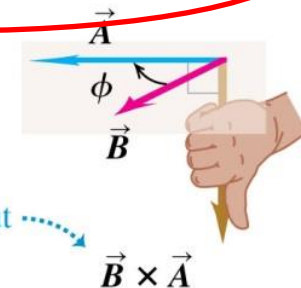
② Point fingers of right hand along  $\vec{A}$ , with palm facing  $\vec{B}$ .

③ Curl fingers toward  $\vec{B}$ .

④ Thumb points in direction of  $\vec{A} \times \vec{B}$ .



(b)  $\vec{B} \times \vec{A} = -\vec{A} \times \vec{B}$  (the vector product is anticommutative)



Same magnitude but opposite direction

Special cases:

(i) if  $\vec{A} \parallel \vec{B}$ ,  $|\vec{A} \times \vec{B}| = 0$ ,  
in particular,  $\hat{i} \times \hat{i} = \hat{j} \times \hat{j} = \hat{k} \times \hat{k} = 0$

(ii) if  $\vec{A} \perp \vec{B}$ ,  $|\vec{A} \times \vec{B}| = AB$   
in particular,

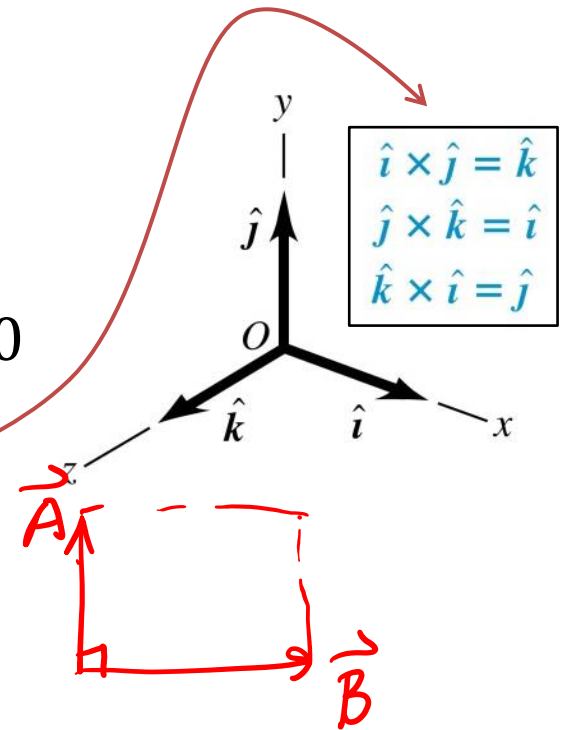
In analytical form (no need to memorize)

$$\vec{A} \times \vec{B}$$

$$= (A_y B_z - A_z B_y) \hat{i} + (A_z B_x - A_x B_z) \hat{j} \\ + (A_x B_y - A_y B_x) \hat{k}$$

$$= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix}$$

3x3

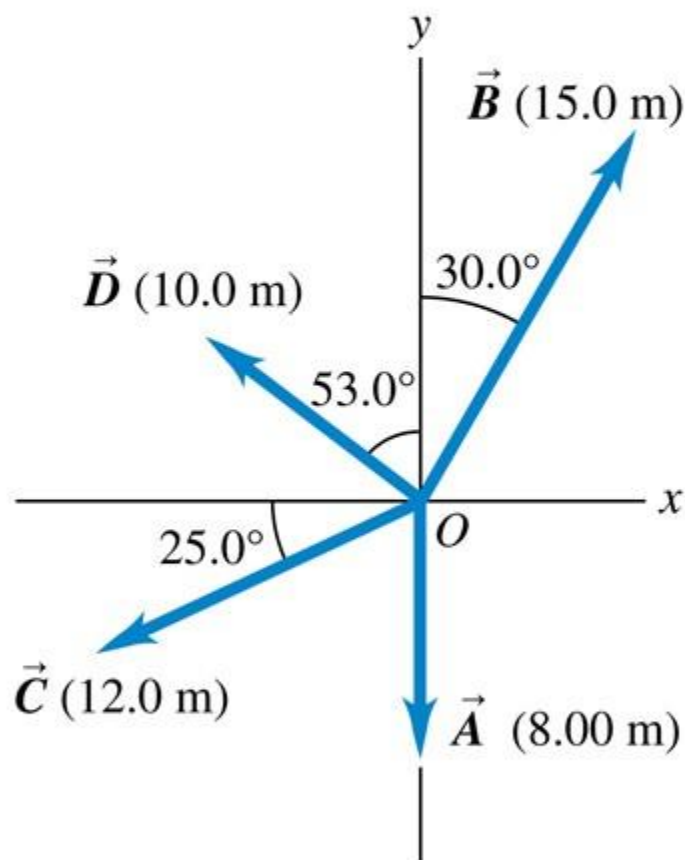


linear algebra  
don't worry if you  
have not learnt  
determinants in  
high school

big data  
n x n

## Q1.14

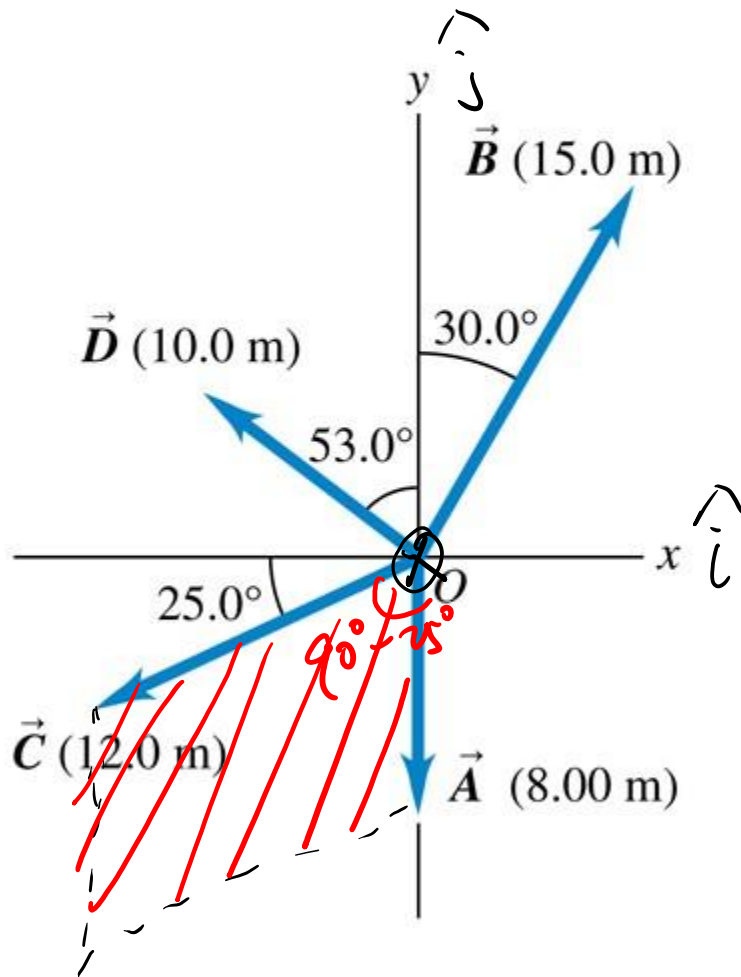
Consider the vectors shown. What is the cross product  $\vec{A} \times \vec{C}$ ?



- A.  $(96.0 \text{ m}^2) \sin 25.0^\circ \hat{k}$
- B.  $(96.0 \text{ m}^2) \cos 25.0^\circ \hat{k}$
- C.  $-(96.0 \text{ m}^2) \sin 25.0^\circ \hat{k}$
- D.  $-(96.0 \text{ m}^2) \cos 25.0^\circ \hat{k}$
- E. none of these

A1.14

Consider the vectors shown. What is the cross product  $\vec{A} \times \vec{C}$ ?



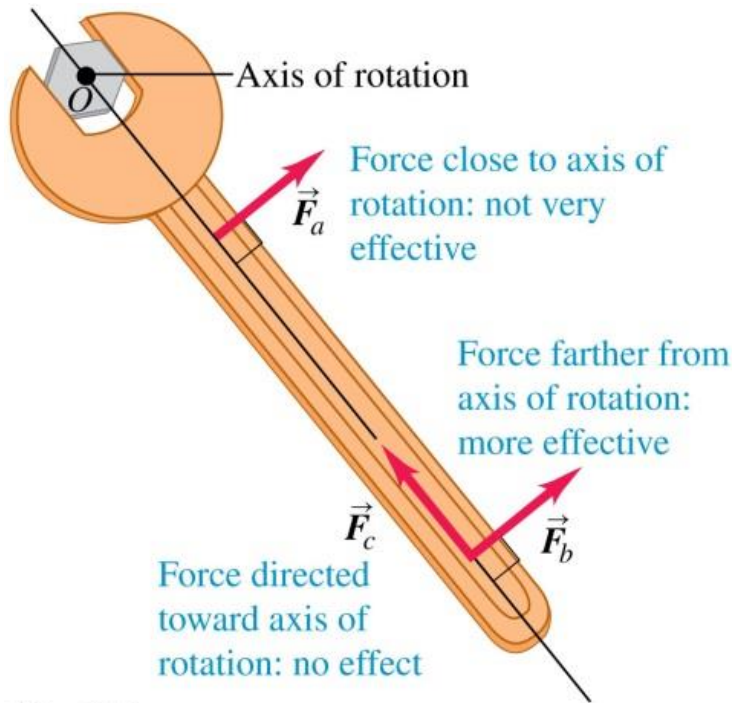
- A.  $(96.0 \text{ m}^2) \sin 25.0^\circ \hat{k}$
- B.  $(96.0 \text{ m}^2) \cos 25.0^\circ \hat{k}$
- C.  $-(96.0 \text{ m}^2) \sin 25.0^\circ \hat{k}$
- ✓ D.  $-(96.0 \text{ m}^2) \cos 25.0^\circ \hat{k}$
- E. none of these

$$\underline{|A| \cdot |C| \cdot \sin(90^\circ - 25^\circ)}$$

$\cos 25^\circ$

# Torque

Besides magnitude and direction, the **line of action** of a force is important because it produces rotation effect.



$\vec{F}_a$  and  $\vec{F}_b$  have the same magnitudes and directions, but different line of action: they produce different physical effects – which force would you apply if you were to tighten/loosen the screw?

Define **torque** about a point  $O$  as  
a vector

$\vec{r}$  : "lever arm vector"

$$\vec{\tau} = \vec{r} \times \vec{F}$$

$\vec{\tau} = -\vec{F} \times \vec{r}$



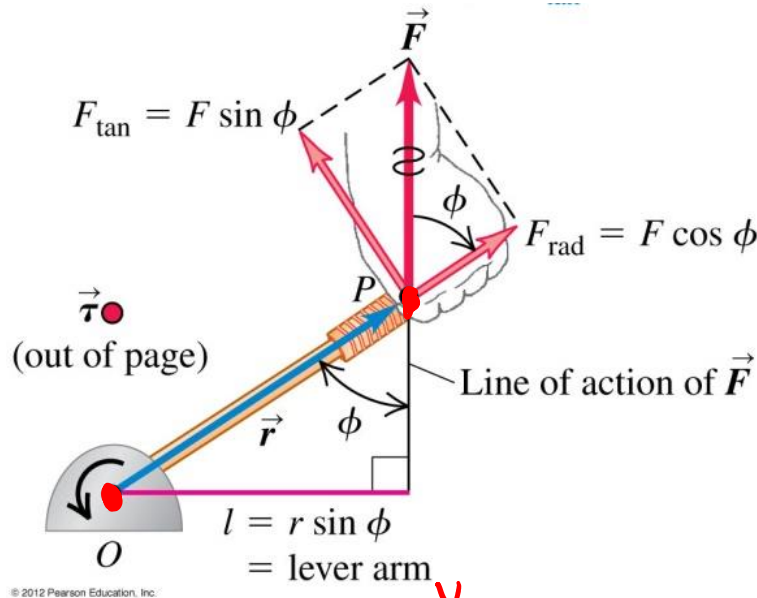
$\vec{\tau}$  is  $\perp$  to both  $\vec{r}$  and  $\vec{F}$

Magnitude:

$$\tau = r(F \sin \phi) = (r \sin \phi)F$$

component  
of  $\vec{F} \perp$  to  $\vec{r}$

$\perp$  distance  
from  $O$  to  
line of  
actions of  $\vec{F}$



" $\vec{r} \perp$  axis"

Direction gives the sense of rotation about  $O$  through the right-hand-rule.

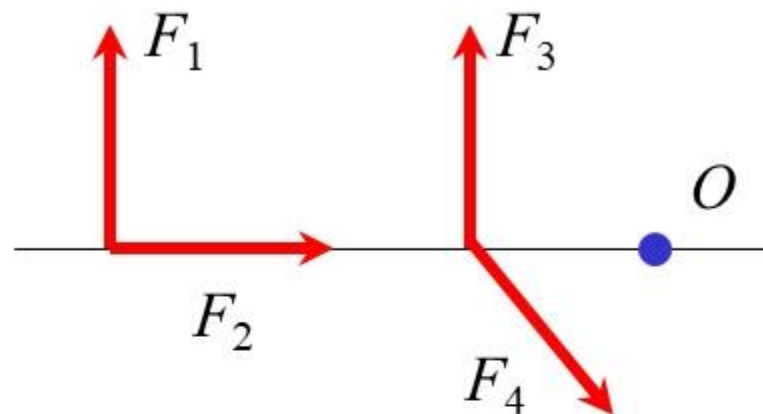
Notation:  $\odot$  out of the plane

$\otimes$  into the plane

SI unit for torque: Nm (just like work done)

## Q10.2

Which of the four forces shown here produces a torque about  $O$  that is directed *out of* the plane of the drawing?

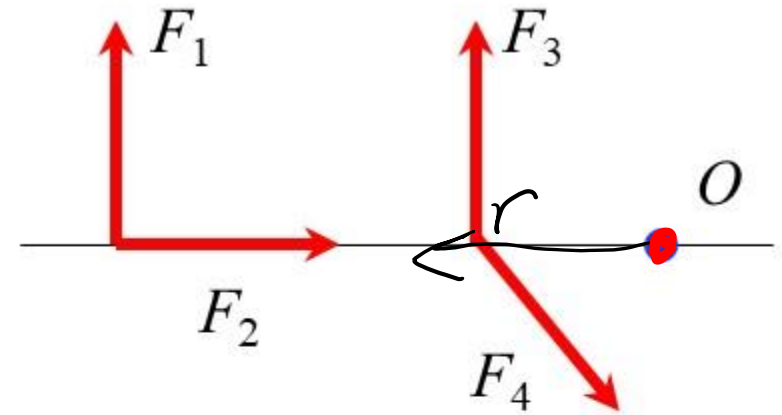


- A.  $F_1$
- B.  $F_2$
- C.  $F_3$
- D.  $F_4$
- E. more than one of these



## A10.2

Which of the four forces shown here produces a torque about  $O$  that is directed *out of* the plane of the drawing?



A.  $F_1$

B.  $F_2$

C.  $F_3$

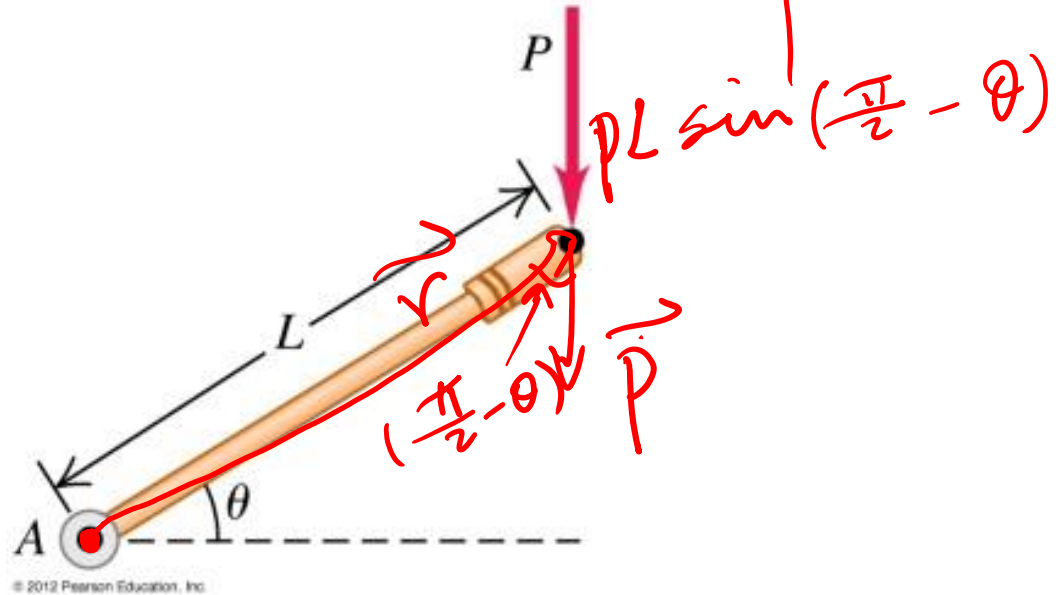
✓ D.  $F_4$

E. more than one of these

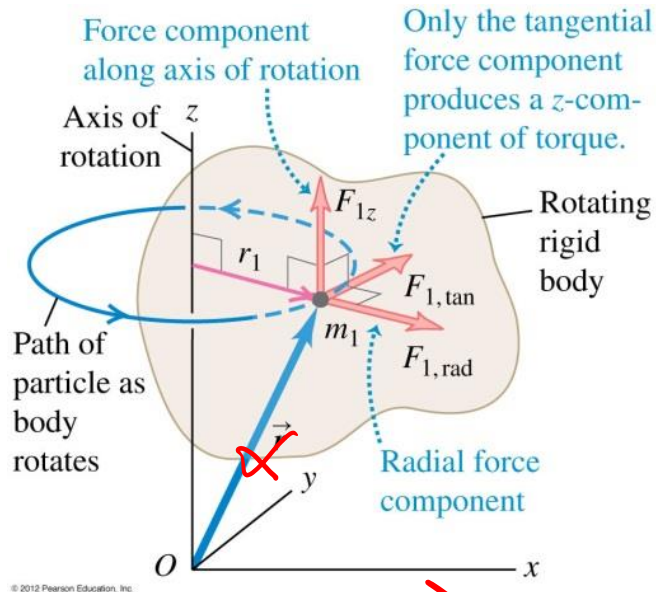
⊙  $\vec{r} \times \vec{F}$

# Question

A force  $P$  is applied to one end of a lever of length  $L$ . The magnitude of the torque of this force about point  $A$  is ( $PL \sin \theta$  /  $PL \cos \theta$  /  $PL \tan \theta$ )



Suppose a rigid body is rotating about a fixed axis which we arbitrarily call the  $z$  axis.  
 $m_1$  is a small part of the total mass.



$F_{1,\text{rad}}$ ,  $F_{1,\text{tan}}$ , and  $F_{1,z}$  are the 3 components of the total force acting on  $m_1$

Only  $F_{1,\text{tan}}$  produces the desired rotation,  $F_{1,\text{rad}}$  and  $F_{1,z}$  produce some other effects which are irrelevant to the rotation about the  $z$  axis.

$$F_{1,\text{tan}} = m_1 a_{1,\text{tan}} = m_1 (r_1 \alpha_z)$$

$$\underbrace{F_{1,\text{tan}} r_1}_{\text{torque on } m_1 \text{ about } z, \tau_{1z}} = m_1 r_1^2 \alpha_z$$

torque on  $m_1$  about  $z$ ,  $\tau_{1z}$

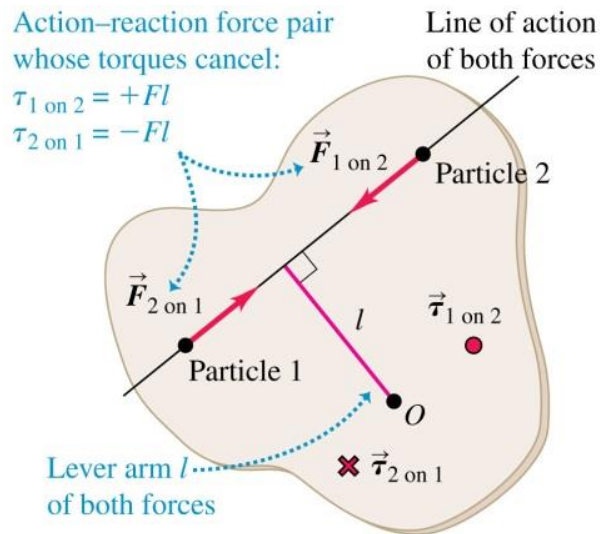
$$\tau = \vec{r}_1 \times \vec{F} = \vec{r}_1 \times (F_x \hat{i} + F_y \hat{j} + F_z \hat{k}) =$$

Sum over all mass in the body, since they all have the same  $\alpha_z$

$$\sum \tau_{iz} = \left( \sum m_r r_i^2 \right) \alpha_z = I \alpha_z$$

$$\sum \vec{F}_i = M \vec{a}$$

$$\begin{aligned} & \vec{r}_1 \times F_x \hat{i} \\ & + \vec{r}_1 \times F_y \hat{j} \\ & + \vec{r}_1 \times F_z \hat{k} \end{aligned}$$



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Need to consider torque due to external forces only. Internal forces (action and reaction pairs) produce equal and opposite torques which have no net rotational effect.

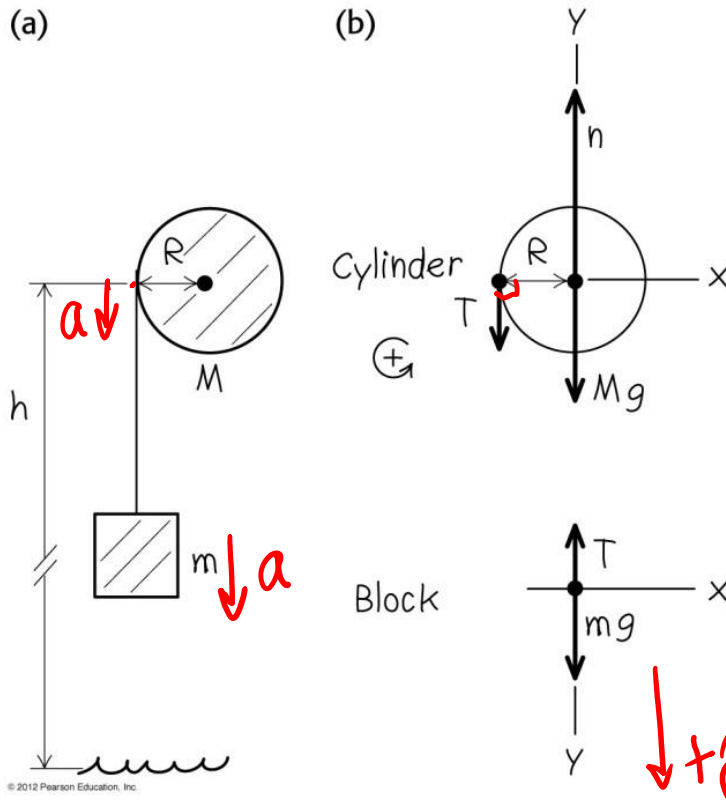
Conclusion: for rigid body rotation about a fixed axis,

$$\sum \tau_{\text{ext}} = I\alpha$$

*c.f.* Newton's second law  $\sum \vec{F}_{\text{ext}} = M\vec{a}$

## Example

Pulley rotates about a fixed axis. What is the acceleration  $a$  of the block?



For the cylinder

$$\underbrace{TR}_{\text{torque due to } T} = \underbrace{\left(\frac{1}{2}MR^2\right)}_{\text{moment of inertia of cylinder}} \underbrace{\left(\frac{a}{R}\right)}_{\text{angular acceleration}}$$

i.e.  $T = \frac{1}{2}Ma$

For the block

$$mg - T = ma$$

Therefore

$$a = \frac{g}{1 + M/2m}$$

Suppose the block is initially at rest at height  $h$ . At the moment it hits the floor:

$$v^2 = 0 + 2 \left( \frac{g}{1 + M/2m} \right) h \Rightarrow v = \sqrt{\frac{2gh}{1 + M/2m}}$$

c.f. lecture 8 in which we get the same result using energy conservation.

# Question

Mass  $m_1$  slides on a frictionless track. The pulley has moment of inertia  $I$  about its rotation axis, and the string does not slip nor stretch. When the hanging mass  $m_2$  is released, arrange the forces  $T_1$ ,  $T_2$ , and  $m_2g$  in increasing order of magnitude.

