

## Chapter 5

# DOUBLE AND TRIPLE INTEGRALS

### 5.1 Multiple-Integral Notation

Previously ordinary integrals of the form

$$\int_J f(x) \, dx = \int_a^b f(x) \, dx \quad (5.1)$$

where  $J = [a, b]$  is an interval on the real line, have been studied. Here we study *double integrals*

$$\iint_{\Omega} f(x, y) \, dx \, dy \quad (5.2)$$

where  $\Omega$  is some **region** in the  $xy$ -plane, and a little later we will study *triple integrals*

$$\iiint_T f(x, y, z) \, dx \, dy \, dz \quad (5.3)$$

where  $T$  is a solid (**volume**) in the  $xyz$ -space.

### 5.2 Double Integrals

#### 5.2.1 Properties

##### (1) Area property

$$\iint_{\Omega} dx \, dy = \text{Area of } \Omega.$$

In particular if  $\Omega$  is the rectangle  $\Omega = [a, b] \times [c, d]$  then  $\iint_{\Omega} dx \, dy = (b - a)(d - c)$ .

##### (2) Linearity

$$\iint_{\Omega} [\alpha f(x, y) + \beta g(x, y)] \, dx \, dy = \alpha \iint_{\Omega} f(x, y) \, dx \, dy + \beta \iint_{\Omega} g(x, y) \, dx \, dy \quad (5.4)$$

where  $\alpha$  and  $\beta$  are constants.

**(3) Additivity**

If  $\Omega$  is broken up into a finite number of nonoverlapping basic regions  $\Omega_1, \dots, \Omega_n$ , then

$$\int \int_{\Omega} f(x, y) \, dx \, dy = \int \int_{\Omega_1} f(x, y) \, dx \, dy + \dots + \int \int_{\Omega_n} f(x, y) \, dx \, dy. \quad (5.5)$$

**5.2.2 Geometric Interpretation**

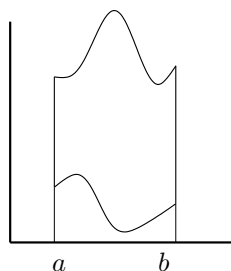
The double integral over  $\Omega$  gives the volume of the solid  $T$  whose upper boundary is the surface  $z = f(x, y)$  and whose lower boundary is the region  $\Omega$  in the  $xy$ -plane:

$$\int \int_{\Omega} f(x, y) \, dx \, dy = \text{volume of } T. \quad (5.6)$$

**5.2.3 The Evaluation of Double Integrals by Repeated Integrals**

If an ordinary integral  $\int_a^b f(x) \, dx$  proves difficult to evaluate, it is not because of the interval  $[a, b]$  but because of the integrand  $f$ . Difficulty in evaluating a double integral  $\int \int_{\Omega} f(x, y) \, dx \, dy$  can come from two sources: from the integrand  $f$  or from the domain  $\Omega$ . Even such a simple looking integral as  $\int \int_{\Omega} 1 \, dx \, dy$  is difficult to evaluate if  $\Omega$  is complicated.

In this section we introduce a technique for evaluating double integrals over domains that have special shapes. The key idea is that double integrals over such special domains can be reduced to a pair of ordinary integrals.

**Horizontally simple domain**

The **projection** of the domain  $\Omega$  onto the  $x$ -axis is a closed interval  $[a, b]$  and  $\Omega$  consists of all points  $(x, y)$  with

$$a \leq x \leq b, \text{ and } \phi_1(x) \leq y \leq \phi_2(x). \quad (5.7)$$

Then

$$\int \int_{\Omega} f(x, y) \, dx \, dy = \int_a^b \left( \int_{\phi_1(x)}^{\phi_2(x)} f(x, y) \, dy \right) dx. \quad (5.8)$$

Here we first calculate  $\int_{\phi_1(x)}^{\phi_2(x)} f(x, y) \, dy$  by integrating  $f(x, y)$  with respect to  $y$  from  $y = \phi_1(x)$  to  $y = \phi_2(x)$ . The resulting expression is a function of  $x$  alone, which we then integrate with respect to  $x$  from  $x = a$  to  $x = b$ .

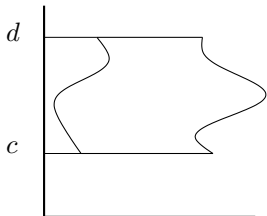
**Example**

Evaluate  $\int \int_{\Omega} (x^4 - 2y) \, dx \, dy$ , where the domain  $\Omega$  consists of all points  $(x, y)$  with  $-1 \leq x \leq 1$  and  $-x^2 \leq y \leq x^2$ .

**Solution**

$$\begin{aligned} \int \int_{\Omega} (x^4 - 2y) \, dx \, dy &= \int_{x=-1}^{x=1} \int_{y=-x^2}^{y=x^2} (x^4 - 2y) \, dy \, dx = \int_{x=-1}^{x=1} [x^4 y - y^2]_{y=-x^2}^{y=x^2} dx \\ &= \int_{x=-1}^{x=1} 2x^6 \, dx = [2x^7/7]_{x=-1}^{x=1} = 4/7 \end{aligned}$$

### Vertically simple domain



The *projection* of the domain  $\Omega$  onto the  $y$ -axis is a closed interval  $[c, d]$  and  $\Omega$  consists of all points  $(x, y)$  with

$$c \leq y \leq d, \text{ and } \psi_1(y) \leq x \leq \psi_2(y). \quad (5.9)$$

Then

$$\int \int_{\Omega} f(x, y) \, dx \, dy = \int_c^d \left( \int_{\psi_1(y)}^{\psi_2(y)} f(x, y) \, dx \right) dy. \quad (5.10)$$

Here we first calculate  $\int_{\psi_1(y)}^{\psi_2(y)} f(x, y) \, dx$  by integrating  $f(x, y)$  with respect to  $x$  from  $x = \psi_1(y)$  to  $x = \psi_2(y)$ . The resulting expression is a function of  $y$  alone, which we then integrate with respect to  $y$  from  $y = c$  to  $y = d$ .

The integrals in the right-hand sides of formulae (5.8) and (5.10) are called *repeated integrals*.

#### Remark 1

Sometimes a domain can be expressed both as a horizontally simple domain:  $a \leq x \leq b$ ,  $\phi_1(x) \leq y \leq \phi_2(x)$ , and as a vertically simple domain:  $c \leq y \leq d$ ,  $\psi_1(y) \leq x \leq \psi_2(y)$ . Then

$$\int \int_{\Omega} f(x, y) \, dx \, dy = \int_a^b \left( \int_{\phi_1(x)}^{\phi_2(x)} f(x, y) \, dy \right) dx = \int_c^d \left( \int_{\psi_1(y)}^{\psi_2(y)} f(x, y) \, dx \right) dy. \quad (5.11)$$

Therefore we can, at least in theory, perform the integration in either order. However, there are situations where one order is preferable over the other.

#### Remark 2

Finally, if  $\Omega$ , the domain of integration, is neither horizontally nor vertically simple, then it is usually possible to break it up into a finite number of domains, say  $\Omega_1, \dots, \Omega_n$ , each of which is either horizontally or vertically simple. Then we can use the additivity property given by eq. (5.5).

### 5.2.4 Evaluating Double Integrals Using Polar Coordinates

Let  $\Omega$  be a domain formed with all points  $(x, y)$  that have polar coordinates  $(r, \theta)$  in the set

$$\Gamma : \alpha \leq \theta \leq \beta, \, \rho_1(\theta) \leq r \leq \rho_2(\theta) \quad (5.12)$$

where  $\beta \leq \alpha + 2\pi$ . Then

$$\begin{aligned} \int \int_{\Omega} f(x, y) \, dx \, dy &= \int \int_{\Gamma} f(r \cos(\theta), r \sin(\theta)) r \, dr \, d\theta \\ &= \int_{\alpha}^{\beta} \int_{\rho_1(\theta)}^{\rho_2(\theta)} f(r \cos(\theta), r \sin(\theta)) r \, dr \, d\theta. \end{aligned} \quad (5.13)$$

## 5.3 Triple Integrals

### 5.3.1 Properties

#### (1) Volume property

$$\int \int \int_T dx \, dy \, dz = \text{Volume of } T.$$

In particular if  $T$  is the box  $T = [a, b] \times [c, d] \times [e, f]$  then  $\int \int \int_T dx \, dy \, dz = (b-a)(d-c)(f-e)$ .

#### (2) Linearity

$$\begin{aligned} \int \int \int_T [\alpha f(x, y, z) + \beta g(x, y, z)] \, dx \, dy \, dz \\ = \alpha \int \int \int_T f(x, y, z) \, dx \, dy \, dz + \beta \int \int \int_T g(x, y, z) \, dx \, dy \, dz \end{aligned} \quad (5.14)$$

where  $\alpha$  and  $\beta$  are constants.

#### (3) Additivity If $T$ is broken up into a finite number of nonoverlapping basic regions $T_1, \dots, T_n$ , then

$$\int \int \int_T f(x, y, z) \, dx \, dy \, dz = \int \int \int_{T_1} f(x, y, z) \, dx \, dy \, dz + \dots + \int \int \int_{T_n} f(x, y, z) \, dx \, dy \, dz. \quad (5.15)$$

### 5.3.2 The Evaluation of Triple Integrals by Repeated Integrals

Let  $T$  be a solid whose projection onto the  $xy$ -plane is labelled  $\Omega_{xy}$ . Then the solid  $T$  is the set of all points  $(x, y, z)$  satisfying

$$(x, y) \in \Omega_{xy}, \chi_1(x, y) \leq z \leq \chi_2(x, y). \quad (5.16)$$

The triple integral over  $T$  can be evaluated by setting

$$\int \int \int_T f(x, y, z) \, dx \, dy \, dz = \int \int_{\Omega_{xy}} \left( \int_{\chi_1(x, y)}^{\chi_2(x, y)} f(x, y, z) \, dz \right) \, dx \, dy. \quad (5.17)$$

In eq. (5.17) we can evaluate the integration with respect to  $z$  first and then evaluate the double integral over the domain  $\Omega_{xy}$  as studied for double integrals. In particular if  $\Omega_{xy}$  is horizontally simple, say

$$a \leq x \leq b, \quad \phi_1(x) \leq y \leq \phi_2(x). \quad (5.18)$$

then the solid  $T$  itself is the set of all points  $(x, y, z)$  such that

$$a \leq x \leq b, \quad \phi_1(x) \leq y \leq \phi_2(x), \quad \chi_1(x, y) \leq z \leq \chi_2(x, y) \quad (5.19)$$

and the triple integral over  $T$  can be expressed by three ordinary integrals as:

$$\int \int \int_T f(x, y, z) \, dx \, dy \, dz = \int_a^b \left[ \int_{\phi_1(x)}^{\phi_2(x)} \left( \int_{\chi_1(x, y)}^{\chi_2(x, y)} f(x, y, z) \, dz \right) \, dy \right] \, dx. \quad (5.20)$$

Here we first integrate with  $z$  [from  $z = \chi_1(x, y)$  to  $z = \chi_2(x, y)$ ], then with respect to  $y$  [from  $y = \phi_1(x)$  to  $y = \phi_2(x)$ ], and finally with respect to  $x$  [from  $x = a$  to  $x = b$ ].

There is nothing special about this order of integration. Other orders of integration are possible and in some cases more convenient. Suppose for example that the projection of  $T$  onto the  $xz$ -plane is a domain  $\Omega_{xz}$  of the form

$$z_1 \leq z \leq z_2, \quad \phi_1(z) \leq x \leq \phi_2(z). \quad (5.21)$$

If  $T$  is the set of all  $(x, y, z)$  with

$$z_1 \leq z \leq z_2, \quad \phi_1(z) \leq x \leq \phi_2(z), \quad \psi_1(x, z) \leq y \leq \psi_2(x, z) \quad (5.22)$$

then

$$\int \int \int_T f(x, y, z) \, dx \, dy \, dz = \int_{z_1}^{z_2} \left[ \int_{\phi_1(z)}^{\phi_2(z)} \left( \int_{\psi_1(x, z)}^{\psi_2(x, z)} f(x, y, z) \, dy \right) \, dx \right] \, dz. \quad (5.23)$$

In this case we integrate first with respect to  $y$ , then with respect to  $x$ , and finally with respect to  $z$ . Still four other orders of integration are possible.

### 5.3.3 Evaluating Triple Integrals Using Cylindrical Coordinates

Let  $T$  be a solid whose projection onto the  $xy$ -plane is labelled  $\Omega_{xy}$ . Then the solid  $T$  is the set of all points  $(x, y, z)$  satisfying

$$(x, y) \in \Omega_{xy}, \quad \chi_1(x, y) \leq z \leq \chi_2(x, y). \quad (5.24)$$

The domain  $\Omega_{xy}$  has polar coordinates in some set  $\Omega_{r\theta}$  and then the solid  $T$  in cylindrical coordinates is some solid  $S$  satisfying

$$(r, \theta) \in \Omega_{r\theta}, \quad \chi_1(r \cos(\theta), r \sin(\theta)) \leq z \leq \chi_2(r \cos(\theta), r \sin(\theta)). \quad (5.25)$$

Then

$$\begin{aligned} \int \int \int_T f(x, y, z) \, dx \, dy \, dz &= \int \int_{\Omega_{xy}} \left( \int_{\chi_1(x, y)}^{\chi_2(x, y)} f(x, y, z) \, dz \right) \, dx \, dy \\ &= \int \int_{\Omega_{r\theta}} \left( \int_{\chi_1(r \cos(\theta), r \sin(\theta))}^{\chi_2(r \cos(\theta), r \sin(\theta))} f(r \cos(\theta), r \sin(\theta), z) \, dz \right) r \, dr \, d\theta = \\ &\quad \int \int \int_S f(r \cos(\theta), r \sin(\theta), z) r \, dr \, d\theta \, dz. \end{aligned} \quad (5.26)$$

### 5.3.4 Evaluating Triple Integrals Using Spherical Coordinates

Let  $T$  be a solid in  $xyz$ -space with spherical coordinates in the solid  $S$  of  $\rho\theta\phi$ -space. Then

$$\int \int \int_T f(x, y, z) \, dx \, dy \, dz = \int \int \int_S f(\rho \sin \theta \cos \phi, \rho \sin \theta \sin \phi, \rho \cos \theta) \rho^2 \sin \theta \, d\rho \, d\theta \, d\phi. \quad (5.27)$$

## 5.4 Jacobians and changing variables in multiple integration

During the course of the last few sections you have met several formulae for changing variables in multiple integration: to polar coordinates, to cylindrical coordinates, to spherical coordinates. The purpose of this section is to bring some unity to that material and provide a general description for other changes of variable.

### 5.4.1 Change of variables for double integrals

Consider the change of variables  $x = x(u, v)$  and  $y = y(u, v)$ , which maps the points  $(u, v)$  of some domain  $\Gamma$  into the points  $(x, y)$  of some other domain  $\Omega$ . Then

$$\text{The area of } \Omega = \int \int_{\Gamma} \left| \frac{\partial(x, y)}{\partial(u, v)} \right| du dv. \quad (5.28)$$

Suppose now that we want to integrate some function  $f(x, y)$  over  $\Omega$ . If this proves difficult to do directly, then we can change variables  $(x, y)$  to  $(u, v)$  and try to integrate over  $\Gamma$  instead. Then

$$\int \int_{\Omega} f(x, y) dx dy = \int \int_{\Gamma} f(x(u, v), y(u, v)) \left| \frac{\partial(x, y)}{\partial(u, v)} \right| du dv. \quad (5.29)$$

### 5.4.2 Change of variables for triple integrals

Consider the change of variables  $x = x(u, v, w)$ ,  $y = y(u, v, w)$ ,  $z = z(u, v, w)$  which maps the points  $(u, v, w)$  of some solid  $S$  into the points  $(x, y, z)$  of some other solid  $T$ . Then

$$\text{The volume of } T = \int \int \int_S \left| \frac{\partial(x, y, z)}{\partial(u, v, w)} \right| du dv dw. \quad (5.30)$$

Suppose now that we want to integrate some function  $f(x, y, z)$  over  $T$ . If this proves difficult to do directly, then we can change variables  $(x, y, z)$  to  $(u, v, w)$  and try to integrate over  $S$  instead. Then

$$\int \int \int_T f(x, y, z) dx dy dz = \int \int \int_S f(x(u, v, w), y(u, v, w), z(u, v, w)) \left| \frac{\partial(x, y, z)}{\partial(u, v, w)} \right| du dv dw. \quad (5.31)$$

Referring back to equations (5.26) and (5.27), and the Jacobians given at the end of §4.5, we can verify that this formula is correct for a change from Cartesian to cylindrical coordinates (Jacobian is  $r$ ) and for a change from Cartesian to spherical coordinates (Jacobian is  $\rho^2 \sin \theta$ ).

## Chapter 6

# LINE INTEGRALS AND SURFACE INTEGRALS

In this chapter we will study integration along curves and integration along surfaces. At the heart of this subject lie three great theorems: *Green's theorem*, *Gauss's theorem* (commonly known as the *divergence theorem*) and *Stokes's theorem*. All of these are ultimately based on the *fundamental theorem of integral calculus*, and all can be cast in the same general form: *An integral over a region  $S$  = An integral over the boundary of  $S$ .*

### 6.1 Line integrals

Let  $\underline{h}(x, y, z) = (h_1(x, y, z), h_2(x, y, z), h_3(x, y, z))$  be a vector function that is continuous over a smooth curve  $C$  parametrised by  $C : \underline{r}(u) = (x(u), y(u), z(u))$  with  $u \in [a, b]$ . The *line integral* of  $\underline{h}$  over  $C$  is the number

$$\int_C \underline{h}(\underline{r}) \cdot d\underline{r} = \int_a^b [\underline{h}(\underline{r}(u)) \cdot \underline{r}'(u)] \, du. \quad (6.1)$$

Although we stated this definition in terms of three-dimensional vectorial functions  $\underline{h}(x, y, z)$  and curves in space  $\underline{r}(u) = (x(u), y(u), z(u))$ , it also includes the two-dimensional case:  $\underline{h}(x, y)$  and plane curves  $\underline{r}(u) = (x(u), y(u))$ .

If the curve  $C$  is not smooth but is made up of a finite number of adjoining smooth pieces  $C_1, \dots, C_n$ , i.e. it is *piecewise smooth*, then we define the integral over  $C$  as the sum of the integrals over  $C_i$  for  $i = 1, \dots, n$ , that is  $\int_C = \int_{C_1} + \dots + \int_{C_n}$ . All polygonal paths are piecewise smooth.

When we integrate over a parametrised curve, we integrate in the direction determined by the parametrisation. If we integrate in the opposite direction, our answer is altered by a factor of  $-1$ , that is  $\int_{-C} = -\int_C$ .

#### 6.1.1 Another notation for line integrals

If  $\underline{h}(x, y, z) = (h_1(x, y, z), h_2(x, y, z), h_3(x, y, z))$  then the line integral over a curve  $C$  can be written as

$$\int_C \underline{h}(\underline{r}) \cdot d\underline{r} = \int_C \{h_1(x, y, z) \, dx + h_2(x, y, z) \, dy + h_3(x, y, z) \, dz\}. \quad (6.2)$$

## 6.2 The Fundamental Theorem for Line Integrals

In general, if we integrate a vector function  $\underline{h}$  from one point to another, the value of the line integral depends on the path chosen. There is, however, an important exception. If the vector function  $\underline{h}$  is a *gradient*, i.e. there exists a scalar function  $f$  such that  $\underline{h} = \underline{\nabla}f$ , then the value of the line integral depends only on the endpoints of the path and not on the path itself. The details are spelled out in the following theorem.

### Theorem

Let  $C$ , parametrised by  $\underline{r} = \underline{r}(u)$  with  $u \in [a, b]$ , be a piecewise smooth curve that begins at  $\underline{\alpha} = \underline{r}(a)$  and ends at  $\underline{\beta} = \underline{r}(b)$ . Then if the vector function  $\underline{h}$  is a gradient, i.e.  $\underline{h} = \underline{\nabla}f$ , we have

$$\int_C \underline{h}(\underline{r}) \cdot d\underline{r} = \int_C \underline{\nabla}f(\underline{r}) \cdot d\underline{r} = f(\underline{\beta}) - f(\underline{\alpha}). \quad (6.3)$$

**NOTE:** It is important to see that this result is an extension of the fundamental theorem of integral calculus:  $\int_a^b f'(x) dx = f(b) - f(a)$ .

### Corollary

If the curve  $C$  is closed, i.e.  $\underline{\alpha} = \underline{\beta}$ , then  $f(\underline{\alpha}) = f(\underline{\beta})$  and  $\int_C \underline{\nabla}f(\underline{r}) \cdot d\underline{r} = 0$ .

## 6.3 Line integrals with respect to arc length

Suppose that  $f$  is a scalar function continuous on a piecewise smooth curve  $C$  parametrised by  $\underline{r} = \underline{r}(u)$  with  $u \in [a, b]$ . If  $s(u)$  is the length of the curve from the tip of  $\underline{r}(a)$  to the tip of  $\underline{r}(u)$ , then, as we have seen in section 2.3,  $s'(u) = \|\underline{r}'(u)\|$ . The integral of  $f$  over  $C$  with respect to arc length  $s$  is defined by setting

$$\int_C f(\underline{r}) ds = \int_a^b f(\underline{r}(u)) s'(u) du. \quad (6.4)$$

## 6.4 Green's Theorem

If  $P(x, y)$  and  $Q(x, y)$  are scalar functions defined over a domain  $\Omega$  with piecewise smooth closed boundary  $C$ , then

$$\int \int_{\Omega} \left[ \frac{\partial Q}{\partial x}(x, y) - \frac{\partial P}{\partial y}(x, y) \right] dx dy = \oint_C P(x, y) dx + Q(x, y) dy \quad (6.5)$$

where the integral on the right is a line integral over  $C$  taken in the anticlockwise direction.

**Remark** As indicated, the symbol  $\oint$  is used to denote the line integral over a simple closed curve  $C$  taken in the anticlockwise direction.

## 6.5 Parametrised Surfaces; Surface Area

We have seen that a space curve  $C$  can be parametrised by a vector function  $\underline{r} = \underline{r}(u)$  where  $u$  ranges over some interval  $I$  of the  $u$ -axis. In an analogous manner, we can parametrise a surface  $S$  in space by a vector function  $\underline{r} = \underline{r}(u, v)$  where  $(u, v)$  ranges over some domain  $\Omega$  of the  $uv$ -plane.



**Example (The graph of a function)**

The graph of a function  $y = f(x)$ ,  $x \in [a, b]$  can be parametrised by setting  $\underline{r}(u) = (u, f(u))$ ,  $u \in [a, b]$ .

Similarly, the graph of a function  $z = f(x, y)$ ,  $(x, y) \in \Omega$  can be parametrised by setting  $\underline{r}(u, v) = (u, v, f(u, v))$ ,  $(u, v) \in \Omega$ .

**Example (A plane)**

If two vectors  $\underline{a}$  and  $\underline{b}$  are not parallel, then the set of all combinations  $u\underline{a} + v\underline{b}$  generates a plane  $P_0$  that passes through the origin. We can parametrise this plane by setting  $\underline{r}(u, v) = u\underline{a} + v\underline{b}$ ,  $u, v$  real numbers.

The plane  $P$  that is parallel to  $P_0$  and passes through the tip of a vector  $\underline{c}$  can be parametrised by setting  $\underline{r}(u, v) = u\underline{a} + v\underline{b} + \underline{c}$ ,  $u, v$  real numbers.

**Example (A sphere)**

The sphere of radius  $a$  centred at the origin can be parametrised by setting

$$\underline{r}(u, v) = (a \sin(u) \cos(v), a \sin(u) \sin(v), a \cos(u)), \quad (u, v) \in [0, \pi] \times [0, 2\pi]. \quad (6.6)$$

**6.5.1 The fundamental vector product**

Let  $S$  be a surface parametrised by  $\underline{r} = \underline{r}(u, v)$ ,  $(u, v) \in \Omega$ . The cross product

$$\underline{N} = \underline{r}'_u \times \underline{r}'_v \quad (6.7)$$

is called the *fundamental vector product* of the surface  $S$ .

The vector  $\underline{N}(u, v)$  is perpendicular to the surface  $S$  at the point with position vector  $\underline{r}(u, v)$  and, if different from zero, can be taken as the normal to the surface  $S$  at that point.

**Example**

For the plane  $\underline{r}(u, v) = u\underline{a} + v\underline{b} + \underline{c}$ , the vector  $\underline{a} \times \underline{b}$  is normal to the plane.

**Example**

The fundamental vector product for a sphere is parallel to the radius vector  $\underline{r}(u, v)$ . (Using the parametrisation given above,  $\underline{N} = a \sin(u)\underline{r}$ .)

**6.5.2 The area of a parametrised surface**

The area of a surface  $S$  parametrised by  $\underline{r} = \underline{r}(u, v)$ ,  $(u, v) \in \Omega$ , is given by

$$\text{Area of } S = \int \int_S \|\underline{N}(u, v)\| \, du \, dv. \quad (6.8)$$

**Example (The surface area of a sphere)**

Using the parametrisation given by equation (6.6), we had  $\underline{N} = a \sin(u)\underline{r}$  so  $\|\underline{N}\| = a^2 \sin(u)$  and the area is

$$\int \int_S \|\underline{N}(u, v)\| \, du \, dv = a^2 \int_{v=0}^{2\pi} \int_{u=0}^{\pi} \sin(u) \, du \, dv = 2\pi a^2 [-\cos(u)]_0^{\pi} = 4\pi a^2.$$

**Example (The area of a plane domain)**

A plane domain may be parametrised as  $\underline{r} = (u, v, 0)$  for  $(u, v) \in \Omega$ . Then  $\underline{r}'_u = (1, 0, 0)$  and  $\underline{r}'_v = (0, 1, 0)$  and so the fundamental vector product is  $\underline{N} = (0, 0, 1)$  which has magnitude 1.

$$\int \int_{\Omega} 1 \, du \, dv = \text{Area of } \Omega.$$

### 6.5.3 The area of a surface $z = f(x, y)$

Let the surface  $S$  be the graph of the function  $z = f(x, y)$  with  $(x, y) \in \Omega$ . Then

$$\text{Area of } S = \iint_{\Omega} \left[ \left( \frac{\partial f}{\partial x} \right)^2 + \left( \frac{\partial f}{\partial y} \right)^2 + 1 \right]^{1/2} dx dy. \quad (6.9)$$

In this case the parametrisation of  $S$  is  $\underline{r}(u, v) = (u, v, f(u, v))$ ,  $(u, v) \in \Omega$  and so  $\underline{N} = (-f_x, -f_y, 1)$ . The unit vector  $\underline{n} = \underline{N}/\|\underline{N}\|$  is called the *upper unit normal*.

## 6.6 Surface Integrals

Let  $H(x, y, z)$  be a scalar function, continuous over a surface  $S$  parametrised by  $\underline{r} = \underline{r}(u, v)$ ,  $(u, v) \in \Omega$ . The *surface integral* of  $H$  over  $S$  is the number

$$\iint_S H(x, y, z) d\sigma = \iint_{\Omega} H(\underline{r}(u, v)) \|\underline{N}(u, v)\| du dv. \quad (6.10)$$

Taking  $H \equiv 1$  and referring back to eq. (6.8) we get

$$\iint_S d\sigma = \text{Area of } S. \quad (6.11)$$

### 6.6.1 Flux of a vector function

Let  $\underline{q}(x, y, z)$  be a vector function that is continuous over a smooth surface  $S$  parametrised by  $\underline{r} = \underline{r}(u, v)$ ,  $(u, v) \in \Omega$ . The *flux* of  $\underline{q}$  across  $S$  in the direction of the unit normal  $\underline{n}$  to the surface  $S$  is the number

$$\iint_S \underline{q} \cdot \underline{n} d\sigma \quad (6.12)$$

which can be calculated as

$$\iint_S \underline{q} \cdot \underline{n} d\sigma = \iint_{\Omega} \underline{q}(\underline{r}(u, v)) \cdot \underline{n} \|\underline{N}\| du dv = \iint_{\Omega} \underline{q}(\underline{r}(u, v)) \cdot \underline{N} du dv. \quad (6.13)$$

#### Proposition

If  $S$  is the graph of a function  $z = f(x, y)$  with  $(x, y) \in \Omega$  and  $\underline{n}$  is the upper unit normal, then the flux of the vector function  $\underline{q} = (q_1(x, y, z), q_2(x, y, z), q_3(x, y, z))$  across  $S$  in the direction of  $\underline{n}$  is

$$\iint_S \underline{q} \cdot \underline{n} d\sigma = \iint_{\Omega} (-q_1 f_x - q_2 f_y + q_3) dx dy. \quad (6.14)$$

#### Proof

We can parametrise the surface by  $\underline{r} = (u, v, f(u, v))$  with  $(u, v) \in \Omega$ . Then the fundamental vector product is

$$\underline{N} = \underline{r}'_u \times \underline{r}'_v = (1, 0, f_u) \times (0, 1, f_v) = (-f_u, -f_v, 1)$$

and we have

$$\begin{aligned} \iint_S \underline{q} \cdot \underline{n} d\sigma &= \iint_{\Omega} (\underline{q} \cdot \underline{N}) du dv \\ &= \iint_{\Omega} (-q_1 f_u - q_2 f_v + q_3) du dv = \iint_{\Omega} (-q_1 f_x - q_2 f_y + q_3) dx dy. \end{aligned}$$

where we have simply changed the names of the variables at the end.

## 6.7 The Divergence (Gauss) Theorem

Recall that if  $P(x, y)$  and  $Q(x, y)$  are scalar functions defined over a domain  $\Omega$  with piecewise smooth closed boundary  $C$ , then Green's theorem (section 6.4) allowed us to express a double integral over  $\Omega$  as a line integral over  $C$ :

$$\int \int_{\Omega} \left[ \frac{\partial Q}{\partial x}(x, y) - \frac{\partial P}{\partial y}(x, y) \right] dx dy = \oint_C P(x, y) dx + Q(x, y) dy. \quad (6.15)$$

This formula can be rewritten in vector terms (using  $\underline{q} = (Q, -P)$ ) to give the *divergence theorem in two dimensions* as follows:

### The divergence theorem in two dimensions

Let  $\Omega$  be a two-dimensional domain bounded by a piecewise smooth closed curve  $C$ . Then for any (continuously differentiable) vector function  $\underline{q}(x, y)$  we have that

$$\int \int_{\Omega} (\nabla \cdot \underline{q}) dx dy = \oint_C (\underline{q} \cdot \underline{n}) ds \quad (6.16)$$

where  $\underline{n}$  is the outer unit normal and the integral on the right is taken with respect to arc length.

We can now give the three-dimensional analogue of the divergence (Gauss) theorem.

### The divergence theorem in three dimensions

Let  $T$  be a three-dimensional solid bounded by a piecewise smooth closed surface  $S$ . Then for any (continuously differentiable) vector function  $\underline{q}(x, y, z)$  we have that

$$\int \int \int_T (\nabla \cdot \underline{q}) dx dy dz = \int \int_S (\underline{q} \cdot \underline{n}) d\sigma \quad (6.17)$$

where  $\underline{n}$  is the outer unit normal.

#### 6.7.1 Divergence as outward flux per unit volume

In eq. (6.17), the right-hand side  $\int \int_S (\underline{q} \cdot \underline{n}) d\sigma$  represents the  $\underline{q}$  across  $S$  in the direction of  $\underline{n}$ . In this sense, from eq. (6.17) we can say that *the divergence is the outward flux per unit volume*, as we discussed in section 4.4.1.

Points  $(x, y, z) \in T$  for which

- $\nabla \cdot \underline{q}(x, y, z) < 0$  are called *sinks*.
- $\nabla \cdot \underline{q}(x, y, z) > 0$  are called *sources*.
- If  $\nabla \cdot \underline{q}(x, y, z) \equiv 0$  then  $\underline{q}$  is called *solenoidal*.

## 6.8 Stokes's Theorem

We return to Green's theorem (section 6.4):

$$\int \int_{\Omega} \left[ \frac{\partial Q}{\partial x}(x, y) - \frac{\partial P}{\partial y}(x, y) \right] dx dy = \oint_C P(x, y) dx + Q(x, y) dy. \quad (6.18)$$

and this time setting  $\underline{q} = (P, Q, R)$  a vector function, we have

$$(\underline{\nabla} \times \underline{q}) \cdot \underline{k} = \det \begin{bmatrix} \frac{i}{\partial x} & \frac{j}{\partial y} & \frac{k}{\partial z} \\ P & Q & R \end{bmatrix} \cdot \underline{k} = \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y}. \quad (6.19)$$

Thus in vector terms Green's theorem can be written as

$$\int \int_{\Omega} [(\underline{\nabla} \times \underline{q}) \cdot \underline{k}] \, dx \, dy = \oint_C \underline{q}(\underline{r}) \cdot \underline{dr}. \quad (6.20)$$

Since any plane can be coordinatised as the  $xy$ -plane, this result can be phrased in the following theorem

### Stokes's theorem

Let  $S$  be a smooth surface with smooth bounding curve  $C$ . Then for any (continuously differentiable) vectorial function  $\underline{q}(x, y, z)$  we have

$$\int \int_S [(\underline{\nabla} \times \underline{q}) \cdot \underline{n}] \, d\sigma = \oint_C \underline{q}(\underline{r}) \cdot \underline{dr} \quad (6.21)$$

where  $\underline{n}$  is a unit normal that varies continuously on  $S$ , and the line integral  $\oint_C$  is taken in the positive sense with respect to  $\underline{n}$ .

#### 6.8.1 The normal component of $\underline{\nabla} \times \underline{q}$ as circulation per unit area; Irrotational flow

Interpret the vector function  $\underline{q}(x, y, z)$  as the velocity of a fluid. In eq. (6.21), the right-hand side line integral  $\oint_C \underline{q}(\underline{r}) \cdot \underline{dr}$  is called the *circulation* of  $\underline{q}$  around the curve  $C$ . In this sense, from eq. (6.21), we can say that  $\underline{\nabla} \times \underline{q}$  in the direction  $\underline{n}$  is the circulation of  $\underline{q}$  per unit area, which relates to the rotation of the fluid as discussed in section 4.4.2.

If  $\underline{\nabla} \times \underline{q} \equiv \underline{0}$  then there is no circulation and  $\underline{q}$  is called *irrotational*, i.e. the fluid has no rotational tendency.