

HYDROSTATIC EQUILIBRIUM

$$M(R)$$

Let there be a gas planet whose mass is

The acceleration due to gravity for a non rotating planet be, $g(r) = \frac{Gm(r)}{r^2} = \frac{g(R)}{R}r$

G is the Universal gravitational constant.

The surface gravity is defined as,

$$g(R) = \frac{GM(R)}{R^2} = g_s \text{-----3}$$

M(R) is the total mass of the body. R is the mean radius of the large celestial body.

Now if we consider the rotation of the planet with time period $T = \frac{2\pi}{\omega}$

The angular velocity of rotation be ω

$$g(r)$$

The expression for $g(r)$ becomes

$$g(r, \phi) = g(r) - r\omega^2 \cos^2 \phi$$

$$g(r, \phi) = \left[\frac{g_s}{R} - \omega^2 \cos^2 \phi \right] r = A_\phi r$$

Where

$$A_\phi = \left[\frac{g_s}{R} - \omega^2 \cos^2 \phi \right]$$

The equation of state for an ideal gas be,

$$p(r) = \frac{kT(r)}{Am_H} \rho(r) \text{-----5}$$

P(r) is the pressure at r

$\rho(r)$ is the density at r

T(r) is the temperature at r

k is the Boltzmann constant

A is the atomic number of the elements involved in the process.

m_H is the rest mass of an Hydrogen atom

$$\rho(r) = \frac{Am_H}{kT(r)} p(r) \text{-----6}$$

Again we have,

$$\frac{\rho_c}{p_c} = \frac{Am_H}{kT_c} \text{-----7}$$

ρ_c is the mass density at center

p_c is the central pressure

$$p_c = \frac{2}{3} \pi G \bar{\rho}^2 R^2 \approx \frac{3GM^2(R)}{4\pi R^4} \text{-----8}$$

Equation 8 represents the expression for central pressure.

R is the radius of the body

$$\rho_c = \frac{Am_H}{kT_c} \cdot \frac{3GM^2(R)}{4\pi R^4} \text{-----9}$$

The temperature is very badly assumed to be in the form,

$$T(r) = \frac{a}{r} \text{-----10}$$

The temperature profile is evaluated by the formula,

$$v_{orbit} = v_{rms}$$

$$R \sqrt{\frac{g_s}{r}} = \sqrt{\frac{3kT(r)}{Am_H}} \text{-----10}$$

$$T(r) = \frac{a}{r}$$

Where a is a constant, $a = \frac{Am_H g_s R^2}{3k}$

$$T(0) = T_c$$

The central temperature be,

The surface temperature be,

$$T(R) = T_s = \frac{a}{R} = \frac{Am_H g_s R}{3k}$$

The modified First Stellar equation may be written as,

$$dp = -g(r, \phi) \rho(r) dr = -\left[\frac{g_s}{R} - \omega^2 \cos^2 \phi\right] r \cdot \frac{Am_H}{k \cdot \frac{Am_H g_s R^2}{3kr}} \cdot p(r) dr$$

$$\frac{dp}{p} = -\frac{3}{g_s R^2} \left[\frac{g_s}{R} - \omega^2 \cos^2 \phi\right] \cdot r^2 dr = -3 \left[\frac{1}{R^3} - \frac{\omega^2 \cos^2 \phi}{g_s R^2}\right] r^2 dr$$

$$\int_{p_c}^p \frac{dx}{x} = -3 \left[\frac{1}{R^3} - \frac{\omega^2 \cos^2 \phi}{g_s R^2}\right] \int_0^r y^2 dy$$

$$\ln \frac{p}{p_c} = -\left[\frac{1}{R^3} - \frac{\omega^2 \cos^2 \phi}{g_s R^2}\right] \cdot r^3 = -\alpha_\phi r^3$$

$$p(r) = p_c \exp[-\alpha_\phi r^3]$$

Where, $\alpha_\phi = \frac{1}{R^3} - \frac{\omega^2 \cos^2 \phi}{g_s R^2}$

As the pressure and density profile is same, So the density profile may be written as,

$$\rho(r) = p_c \exp[-\alpha_\phi r^3]$$

The second stellar equation may be written as,

$$\frac{d}{dr} M(r) = 4\pi \cdot r^2 \rho(r) \text{-----18}$$

The mass included with in the radius r is,

$$M(r) = 4\pi \rho_c \int_0^r r^2 \exp[-\alpha_\phi r^3] dr$$

$$M(r) = \frac{4\pi \rho_c}{3\alpha_\phi} (1 - \exp[-\alpha_\phi r^3])$$

The expression for the total mass may be written as r=R,

$$M(R) = \frac{4\pi\rho_c}{\frac{3}{R^2}\left[\frac{1}{R} - \frac{R\omega^2 \cos^2 \phi}{g_s}\right]} (1 - \exp[-\alpha_\phi R^3]) = \frac{4\pi\rho_c}{3\alpha_\phi} \left\{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s R^2}\right)\right]\right\}$$

$$M(R) = \frac{4\pi R^3}{3} \cdot \frac{\rho_c}{1 - \frac{R\omega^2 \cos^2 \phi}{g_s}} \left\{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)\right]\right\}$$

$$\bar{\rho}(R) = \frac{M(R)}{V(R)} = \frac{\rho_c}{1 - \frac{R\omega^2 \cos^2 \phi}{g_s}} \left\{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)\right]\right\}$$

$$\frac{\bar{\rho}}{\rho_c} = \frac{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)\right]}{1 - \frac{R\omega^2 \cos^2 \phi}{g_s}}$$

The expression for the total mass may be rearranged

$$\bar{\rho} = \frac{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)\right]}{1 - \frac{R\omega^2 \cos^2 \phi}{g_s}} \rho_c$$

$$\frac{M(r)}{V(R)} = \frac{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)\right]}{1 - \frac{R\omega^2 \cos^2 \phi}{g_s}} \frac{Am_H}{kT_c} \frac{3GM^2(R)}{4\pi R^4}$$

$$M(R) = \frac{kT_c}{GAm_H} \frac{\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)R}{1 - \exp\left[-\left(1 - \frac{R\omega^2 \cos^2 \phi}{g_s}\right)\right]}$$

The expression for central temperature may be written as ,

$$M(R) = \frac{kRT_c}{AGm_H} \frac{(1 - \frac{R\omega^2 \cos^2 \phi}{g_s})}{1 - \exp[-(1 - \frac{R\omega^2 \cos^2 \phi}{g_s})]}$$

$$T_c \frac{(1 - \frac{R\omega^2 \cos^2 \phi}{g_s})}{1 - \exp[-(1 - \frac{R\omega^2 \cos^2 \phi}{g_s})]} = \frac{AGm_H M(R)}{kR} = \lambda = \text{const}$$

$$T_c(\omega) = T_c(\frac{1}{T_{day}}) = \frac{AGm_H M(R)}{kR} \frac{1 - \exp[-(1 - \frac{R\omega^2 \cos^2 \phi}{g_s})]}{(1 - \frac{R\omega^2 \cos^2 \phi}{g_s})}$$