

EAE 298 Aeroacoustics  
Fall Quarter 2016  
Homework #3

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**Problem 1. [50 points]**

The acoustic wave equation without considering the source is expressed as follows:

$$\frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} - \nabla^2 p = 0$$

We can define a new function  $\tilde{p}$  using the imbedding technique as follows:

$$\begin{aligned} \tilde{p} &= p, & f > 0 \\ \tilde{p} &= 0, & f < 0 \end{aligned}$$

where  $f = 0$  describes the arbitrary moving body. Show that the wave equation whose sound is generated by an arbitrary moving body ( $f = 0$ ) can be expressed as follows:

$$\frac{1}{c^2} \frac{\partial^2 \tilde{p}}{\partial t^2} - \nabla^2 \tilde{p} = - \left[ \frac{M_n}{c} \frac{\partial p}{\partial t} + p_n \right] \delta(f) - \frac{1}{c} \frac{\partial}{\partial t} [M_n p \delta(f)] - \nabla \cdot [p \vec{n} \delta(f)]$$

where  $\vec{n}$  is the unit normal vector on the surface and  $p_n = \nabla p \cdot \vec{n}$ . Now we can use the Greens function of the wave equation in the unbounded space, the so-called free-space Greens function, to find the unknown function  $p(\vec{x}, t)$  everywhere in space. The result is the Kirchhoff formula for moving surfaces.

**Solution**

Using the rules of generalized differentiation, we obtain

$$\begin{aligned} \frac{d\tilde{p}}{dx} &= \frac{d\tilde{p}}{dx} + \Delta\tilde{p}\delta(x-c) \\ &= \frac{d\tilde{p}}{dx} + \Delta\tilde{p}\delta(f) \\ \bar{\nabla}\tilde{p} &= \frac{\bar{\partial}\tilde{p}}{\partial x_i} = \frac{\bar{\partial}\tilde{p}}{\partial f} \frac{\partial f}{\partial x_i} \\ &= \left( \frac{\partial\tilde{p}}{\partial f} + \Delta\tilde{p}\delta(f) \right) \frac{\partial f}{\partial x_i} \\ &= \frac{\partial\tilde{p}}{\partial f} \frac{\partial f}{\partial x_i} + \Delta\tilde{p} \frac{\partial f}{\partial x_i} \delta(f) \\ &= \frac{\partial\tilde{p}}{\partial x_i} + \frac{\partial f}{\partial x_i} \Delta\tilde{p}\delta(f) \\ &= \nabla\tilde{p} + \nabla f \Delta\tilde{p}\delta(f) \\ &= \nabla\tilde{p} + \hat{n}p\delta(f) \end{aligned}$$

Applying the generalized differentiation to the acoustic wave equation

$$\begin{aligned} \frac{1}{c} \frac{\bar{\partial}\tilde{p}}{\partial t} &= \frac{1}{c} \frac{\partial\tilde{p}}{\partial t} - M_n p \delta(f) \\ \frac{1}{c} \frac{\bar{\partial}^2 \tilde{p}}{\partial t^2} &= \frac{1}{c} \frac{\partial^2 \tilde{p}}{\partial t^2} - \frac{M_n}{c} \frac{\partial p}{\partial t} \delta(f) - \frac{1}{c} \frac{\partial}{\partial t} [M_n p \delta(f)] \end{aligned}$$

where  $M_n = v_n/c$  is the local normal Mach number of the surface  $f = 0$ . Similarly, we find

$$\begin{aligned}\bar{\nabla}\tilde{p} &= \nabla\tilde{p} + p\vec{n}\delta(f) \\ \bar{\nabla}^2\tilde{p} &= \nabla^2\tilde{p} + p_n\delta(f) + \nabla \cdot [p\vec{n}\delta(f)]\end{aligned}$$

Combining these equations yields

$$\begin{aligned}\frac{1}{c^2}\frac{\partial^2\tilde{p}}{\partial t^2} - \bar{\nabla}^2\tilde{p} &= \frac{1}{c}\frac{\partial^2\tilde{p}}{\partial t^2} - \frac{M_n}{c}\frac{\partial p}{\partial t}\delta(f) - \frac{1}{c}\frac{\partial}{\partial t}[M_n p\delta(f)] - (\nabla^2\tilde{p} + p_n\delta(f) + \nabla \cdot [p\vec{n}\delta(f)]) \\ &= \left[\frac{1}{c}\frac{\partial^2\tilde{p}}{\partial t^2} - \nabla^2\tilde{p}\right] - \left[\frac{M_n}{c}\frac{\partial p}{\partial t} + p_n\delta(f)\right]\delta(f) - \frac{1}{c}\frac{\partial}{\partial t}[M_n p\delta(f)] - \nabla \cdot [p\vec{n}\delta(f)] \\ &= -\left[\frac{M_n}{c}\frac{\partial p}{\partial t} + p_n\delta(f)\right]\delta(f) - \frac{1}{c}\frac{\partial}{\partial t}[M_n p\delta(f)] - \nabla \cdot [p\vec{n}\delta(f)]\end{aligned}$$

where we have used the property that the original acoustic wave equation is equal to zero to remove the first two terms. Upon use of the Green's Function of the wave operator in unbounded space, this equation gives

$$\begin{aligned}4\pi\tilde{p}(\vec{x}, t) &= -\int \frac{1}{r}\left(p_n + \frac{M_n}{c}p_{\tau(y)}\right)\delta(f)\delta(g)d\vec{y}d\tau - \frac{1}{c}\frac{\partial}{\partial t}\int \frac{1}{r}M_n p\delta(f)\delta(g)d\vec{y}d\tau \\ &\quad - \bar{\nabla} \cdot \int \frac{1}{r}p\vec{n}\delta(f)\delta(g)d\vec{y}d\tau\end{aligned}$$

where  $g = \tau - t + r/c$ , and  $p_{\tau(y)} = \partial p(\vec{y}, \tau)/\partial \tau$ .

## Problem 2. [50 points]

Farassats formulation 1 for the loading noise is given as

$$4\pi p'_L(\vec{x}, t) = \frac{1}{c}\frac{\partial}{\partial t}\int_{f=0}\left[\frac{L_r}{r(1-M_r)}\right]_{ret}dS + \int_{f=0}\left[\frac{L_r}{r^2(1-M_r)}\right]_{ret}dS \quad (1)$$

where  $L_r = \Delta P \vec{n} \cdot \hat{r} = \Delta P \cos \theta$ . This formulation 1 is difficult to compute since the observer time differentiation is outside the integrals. A much more efficient and practical formulation can be derived by carrying the observer time derivate inside the integrals (formulation 1A). Show that formulation 1A for the loading noise becomes

$$\begin{aligned}4\pi p'_L(\vec{x}, t) &= \frac{1}{c}\int_{f=0}\left[\frac{\dot{L}_r}{r(1-M_r)^2}\right]_{ret}dS \\ &\quad + \int_{f=0}\left[\frac{L_r - L_M}{r^2(1-M_r)^2}\right]_{ret}dS + \frac{1}{c}\int_{f=0}\left[\frac{L_r(r\dot{M}_r + c(M_r - M^2))}{r^2(1-M_r)^3}\right]_{ret}dS\end{aligned}$$

where  $L_M = \vec{L} \cdot \vec{M}$ .

## Solution

From the in class notes, we know that

$$\left. \frac{\partial}{\partial t} \right|_x = \left[ \frac{1}{1 - M_r} \frac{\partial}{\partial \tau} \right]_{ret}$$

where  $\frac{\partial}{\partial t}$  is evaluated along surface  $x$ . Applying this rule to Equation (1)

$$\frac{\partial}{\partial t} \int_{f=0} \left[ \frac{L_r}{r(1 - M_r)} \right]_{ret} ds = \int_{f=0} \left[ \frac{1}{1 - M_r} \frac{\partial}{\partial \tau} \left( \frac{L_r}{r(1 - M_r)} \right) \right]_{ret} ds \quad (2)$$

The produce rule requires us to calculate the derivative of each part of the equation. Evaluating each part of this derivative before taking the actual derivative yields

$$\frac{\partial}{\partial \tau} (L_r) = \dot{L}_r$$

From the class notes we compute the time derivative of  $\frac{1}{r}$  using the chain rule:

$$\begin{aligned} \frac{\partial}{\partial \tau} \left( \frac{1}{r} \right) &= \frac{\partial}{\partial r} \left( \frac{1}{r} \right) \frac{\partial r}{\partial \tau} \\ &= -\frac{1}{r^2} \frac{\partial r}{\partial y_i} \frac{\partial y_i}{\partial \tau} \\ &= -\frac{1}{r^2} (-1)(V_r) \\ &= \frac{V_r}{r^2} \end{aligned}$$

Computing the time derivative of  $\frac{1}{1 - M_r}$

$$\begin{aligned} \frac{\partial}{\partial \tau} \left( \frac{1}{1 - M_r} \right) &= \frac{1}{(1 - M_r)^2} \frac{\partial}{\partial \tau} (M_r) \\ &= \frac{1}{(1 - M_r)^2} \frac{\partial}{\partial \tau} (\vec{M} \cdot \hat{r}) \\ &= \frac{1}{(1 - M_r)^2} \left\{ \frac{\partial \vec{M}}{\partial \tau} \cdot \hat{r} + \vec{M} \frac{\partial \hat{r}}{\partial \tau} \right\} \\ &= \frac{1}{(1 - M_r)^2} \left\{ \dot{\vec{M}} \cdot \hat{r} + \vec{M} \frac{\partial \hat{r}}{\partial \tau} \right\} \end{aligned} \quad (3)$$

Solving for  $\frac{\partial \hat{r}}{\partial \tau}$  yields

$$\begin{aligned} \frac{\partial}{\partial \tau} (\hat{r}) &= \frac{\partial}{\partial \tau} \left( \frac{\vec{r}}{r} \right) = \frac{1}{r} \frac{\partial \vec{r}}{\partial \tau} + \vec{r} \frac{\partial}{\partial \tau} \left( \frac{1}{r} \right) \\ &= \frac{-\vec{V}}{r} + \vec{r} \cdot \frac{V_r}{r^2} = \frac{-\vec{V}}{r} + \frac{1}{r} \frac{\vec{r}}{r} \cdot V_r \\ &= -\frac{\vec{V}}{r} + \frac{1}{r} \hat{r} \cdot V_r \\ &= -\frac{c}{r} (\vec{M} - \vec{r} \cdot M_r) \end{aligned} \quad (4)$$

Plugging the solution for Equation 4 into Equation 3

$$\begin{aligned}
\frac{\partial}{\partial \tau} \left( \frac{1}{1 - M_r} \right) &= \frac{1}{(1 - M_r)^2} \left\{ \dot{\vec{M}} \cdot \hat{r} + \vec{M} \frac{\partial \hat{r}}{\partial \tau} \right\} \\
&= \frac{1}{(1 - M_r)^2} \left\{ \dot{\vec{M}} \cdot \hat{r} + \vec{M} \frac{c}{r} \left( -\vec{M} + \vec{r} \cdot M_r \right) \right\} \\
&= \frac{1}{(1 - M_r)^2} \left\{ \dot{\vec{M}}_r + \frac{c}{r} \left( M_r^2 - |\vec{M}|^2 \right) \right\} \\
&= \frac{\dot{\vec{M}}_r + \frac{c}{r} \left( M_r^2 - |\vec{M}|^2 \right)}{(1 - M_r)^2} \\
&= \frac{r \dot{\vec{M}}_r + c \left( M_r^2 - |\vec{M}|^2 \right)}{r(1 - M_r)^2}
\end{aligned}$$

Using these results, we can take the chain rule and find the derivative

$$\begin{aligned}
\frac{\partial}{\partial \tau} \left[ \frac{L_r}{r(1 - M_r)} \right] &= \frac{\partial L_r}{\partial \tau} \frac{1}{r(1 - M_r)} + L_r \frac{\partial}{\partial \tau} \left[ \frac{1}{r(1 - M_r)} \right] \\
&= \frac{\dot{L}_r}{r(1 - M_r)} + L_r \left[ \frac{\partial}{\partial \tau} \left( \frac{1}{r} \right) \frac{1}{1 - M_r} + \frac{1}{r} \frac{\partial}{\partial \tau} \left( \frac{1}{1 - M_r} \right) \right] \\
&= \frac{\dot{L}_r}{r(1 - M_r)} + L_r \left[ \frac{V_r}{r^2(1 - M_r)} + \frac{r \dot{\vec{M}}_r + c \left( M_r^2 - |\vec{M}|^2 \right)}{r^2(1 - M_r)^2} \right] \\
&= \frac{\dot{L}_r}{r(1 - M_r)} + \frac{L_r V_r}{r^2(1 - M_r)} + \frac{L_r \left[ r \dot{\vec{M}}_r + c \left( M_r^2 - |\vec{M}|^2 \right) \right]}{r^2(1 - M_r)^2} \\
&= \frac{\dot{L}_r}{r(1 - M_r)} + \frac{(\vec{L} \cdot \hat{r})(\vec{V} \cdot \hat{r})}{r^2(1 - M_r)} + \frac{L_r \left[ r \dot{\vec{M}}_r + c \left( M_r^2 - |\vec{M}|^2 \right) \right]}{r^2(1 - M_r)^2} \\
&= \frac{\dot{L}_r}{r(1 - M_r)} + \frac{(\vec{L} \cdot \vec{V})(\hat{r} \cdot \hat{r})}{r^2(1 - M_r)} + \frac{L_r \left[ r \dot{\vec{M}}_r + c \left( M_r^2 - |\vec{M}|^2 \right) \right]}{r^2(1 - M_r)^2} \\
&= \frac{\dot{L}_r}{r(1 - M_r)} + \frac{\vec{L} \cdot \vec{V}}{r^2(1 - M_r)} + \frac{L_r \left[ r \dot{\vec{M}}_r + c \left( M_r^2 - |\vec{M}|^2 \right) \right]}{r^2(1 - M_r)^2} \tag{5}
\end{aligned}$$

where the second term has made use of the commutative property of the dot product. Finally, plugging Equation (5) back into Equation (2) yields

$$\begin{aligned}
\frac{\partial}{\partial t} \int_{f=0} \left[ \frac{L_r}{r(1-M_r)} \right]_{ret} ds &= \int_{f=0} \left[ \frac{1}{1-M_r} \frac{\partial}{\partial \tau} \left( \frac{L_r}{r(1-M_r)} \right) \right]_{ret} ds \\
&= \int_{f=0} \left\{ \frac{1}{1-M_r} \left[ \frac{\dot{L}_r}{r(1-M_r)} + \frac{\vec{L} \cdot \vec{V}}{r^2(1-M_r)} + \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^2} \right] \right\}_{ret} ds \\
&= \int_{f=0} \left[ \frac{\dot{L}_r}{r(1-M_r)^2} + \frac{\vec{L} \cdot \vec{V}}{r^2(1-M_r)^2} + \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^3} \right]_{ret} ds
\end{aligned}$$

Plugging this result back into Equation (1)

$$\begin{aligned}
4\pi p'_L &= \frac{1}{c} \frac{\partial}{\partial t} \int_{f=0} \left[ \frac{L_r}{r(1-M_r)} \right]_{ret} ds + \int_{f=0} \left[ \frac{L_r}{r^2(1-M_r)} \right]_{ret} ds \\
&= \frac{1}{c} \int_{f=0} \left[ \frac{\dot{L}_r}{r(1-M_r)^2} + \frac{\vec{L} \cdot \vec{V}}{r^2(1-M_r)^2} + \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^3} \right]_{ret} ds + \int_{f=0} \left[ \frac{L_r}{r^2(1-M_r)} \right]_{ret} ds \\
&= \frac{1}{c} \int_{f=0} \left[ \frac{\dot{L}_r}{r(1-M_r)^2} \right]_{ret} ds + \int_{f=0} \left[ \frac{L_r}{r^2(1-M_r)} \right]_{ret} ds + \frac{1}{c} \int_{f=0} \left[ \frac{\vec{L} \cdot \vec{V}}{r^2(1-M_r)^2} \right]_{ret} ds \\
&\quad + \frac{1}{c} \int_{f=0} \left[ \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^3} \right]_{ret} ds \\
&= \frac{1}{c} \int_{f=0} \left[ \frac{\dot{L}_r}{r(1-M_r)^2} \right]_{ret} ds + \int_{f=0} \left[ \frac{L_r}{r^2(1-M_r)} + \frac{\frac{1}{c}(\vec{L} \cdot \vec{V})}{r^2(1-M_r)^2} \right]_{ret} ds \\
&\quad + \frac{1}{c} \int_{f=0} \left[ \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^3} \right]_{ret} ds \\
&= \frac{1}{c} \int_{f=0} \left[ \frac{\dot{L}_r}{r(1-M_r)^2} \right]_{ret} ds + \int_{f=0} \left[ \frac{L_r + \vec{L} \cdot \vec{M}}{r^2(1-M_r)^2} \right]_{ret} ds \\
&\quad + \frac{1}{c} \int_{f=0} \left[ \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^3} \right]_{ret} ds \\
&= \frac{1}{c} \int_{f=0} \left[ \frac{\dot{L}_r}{r(1-M_r)^2} \right]_{ret} ds + \int_{f=0} \left[ \frac{L_r + L_M}{r^2(1-M_r)^2} \right]_{ret} ds \\
&\quad + \frac{1}{c} \int_{f=0} \left[ \frac{L_r \left[ r\dot{M}_r + c(M_r^2 - |\vec{M}|^2) \right]}{r^2(1-M_r)^3} \right]_{ret} ds
\end{aligned}$$