MAE 298 Aeroacoustics – Homework #2
Lilley's Equation Solution and Application to Jet Noise

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Nomenclature

R_j	Axisymmetric jet radius	E	Exponential term: $kz + n\theta - \omega t$
W_j	Jet exit mean velocity	P	
$\overline{\rho_j}$	Jet exit mean density	k	
W_0	Ambient mean velocity	ω	
$\overline{\rho_0}$	Ambient mean density	n	
$W(r)$	Radial distribution of axial velocity	r	
$\overline{a^2}(r)$	Radial distribution of speed of sound squared	θ	
p'	Perturbation pressure	z	
i	Imaginary number $\sqrt{-1}$		

I. Background

The following analysis will derive solutions to Lilley's equation for parallel axisymmetric flow (Eqn 1):

$$\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)^3 p' - \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)\left(\overline{a^2}\nabla^2 p'\right) - \frac{d\overline{a^2}}{dr}\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)\frac{dp'}{dr} + 2\overline{a^2}\frac{dW}{dr}\frac{\partial^2 p'}{\partial z \partial r} = S(\vec{x}, t) \tag{1}$$

$$\text{where } \nabla^2 \equiv \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2}$$

$W(r)$ is the radial distribution of axial velocity and $\overline{a^2}(r)$ is the radial distribution of speed of sound squared.

After a general solution is derived, it will be applied to find homogeneous solutions to the far-field and potential core regions of an axial jet. Finally, the conditions for matching the solutions for this regions will be discussed.

II. Problem 1 – Solution to Lilley's Equation

Seek solutions of Lilley's equation in the form:

$$p'(r, \theta, z, t) \sim P(r) \exp\left[i(kz + n\theta - \omega t)\right] \tag{2}$$

Assume:

1. $\overline{a^2} = \dfrac{\gamma\overline{p}}{\overline{\rho}}$
2. $\overline{p} = const$

where $\overline{\rho}(r)$ is the radial distribution of the mean density. Show that Lilley's equation reduces to (Eqn 3):

$$\frac{d^2P}{dr^2} + \left\{\frac{1}{r} - \frac{1}{\rho}\frac{d\overline{\rho}}{dr} + \frac{2k}{(\omega - kW)}\frac{dW}{dr}\right\}\frac{dP}{dr} + \left\{\frac{(\omega - kW)^2}{\overline{a^2}} - k^2 - \frac{n^2}{r^2}\right\}P = RHS \tag{3}$$

To aid in this derivaiton, we find the results for the first-order partial derivatives of parameters relevant to the solution. First, we compute the derivative of the perturbation pressure p' with respect to (WRT) time t:

$$\begin{aligned}
\frac{\partial p'}{\partial t} &= \frac{\partial}{\partial t}\left\{P(r)\exp\left[i(kz + n\theta - \omega t)\right]\right\} \\
&= P\frac{\partial}{\partial t}[i(kz + n\theta - \omega t)]\exp[i(kz + n\theta - \omega t)] \\
&= iP\left[\cancelto{0}{\frac{\partial}{\partial t}(kz)} + \cancelto{0}{\frac{\partial}{\partial t}(n\theta)} - \frac{\partial}{\partial t}(\omega t)\right]\exp[i(kz + n\theta - \omega t)]
\end{aligned}$$

$$\boxed{\frac{\partial p'}{\partial t} = -iP\omega\exp[i(kz + n\theta - \omega t)] = -iP\omega e^{iE}} \tag{4}$$

where $E = kz + n\theta - \omega t$, $\dfrac{\partial E}{\partial t} = -\omega$, and $\dfrac{\partial E}{\partial z} = k$. Next, we compute the derivative of p' WRT the angular direction θ:

$$\begin{aligned}
\frac{\partial p'}{\partial \theta} &= \frac{\partial}{\partial \theta}\left\{P(r)\exp\left[i(kz + n\theta - \omega t)\right]\right\} \\
&= P\frac{\partial}{\partial \theta}[i(kz + n\theta - \omega t)]\exp[i(kz + n\theta - \omega t)] \\
&= iP\left[\cancelto{0}{\frac{\partial}{\partial \theta}(kz)} + \frac{\partial}{\partial \theta}(n\theta) - \cancelto{0}{\frac{\partial}{\partial \theta}(\omega t)}\right]\exp[i(kz + n\theta - \omega t)]
\end{aligned}$$

$$\boxed{\frac{\partial p'}{\partial \theta} = iPn\exp[i(kz + n\theta - \omega t)] = iPne^{iE}} \tag{5}$$

Next, we compute the derivative of p' WRT the axial flow direction z:

$$\begin{aligned}
\frac{\partial p'}{\partial z} &= \frac{\partial}{\partial z}\left\{P(r)\exp\left[i(kz + n\theta - \omega t)\right]\right\} \\
&= P\frac{\partial}{\partial z}[i(kz + n\theta - \omega t)]\exp[i(kz + n\theta - \omega t)] \\
&= iP\left[\frac{\partial}{\partial z}(kz) + \cancelto{0}{\frac{\partial}{\partial z}(n\theta)} - \cancelto{0}{\frac{\partial}{\partial z}(\omega t)}\right]\exp[i(kz + n\theta - \omega t)]
\end{aligned}$$

$$\boxed{\frac{\partial p'}{\partial z} = iPk\exp[i(kz + n\theta - \omega t)] = iPke^{iE}} \tag{6}$$

Next, we compute the derivative of p' WRT the radial direction r:

$$\frac{\partial p'}{\partial r} = \frac{\partial}{\partial r}\left\{P(r)\exp\left[i(kz + n\theta - \omega t)\right]\right\}$$

$$\boxed{\frac{\partial p'}{\partial r} = \frac{dP}{dr}\exp[i(kz + n\theta - \omega t)] = \frac{dP}{dr}e^{iE}} \tag{7}$$

II.A. Term 1

To simplify the derivation, we will apply the solution form individually to each term in Lilley's equation. For the first term, we must apply the multi-derivative operator a total of three times:

$$\begin{aligned}
\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)^3 p' &= \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)^2 \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right) p' \\
&= \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)^2 \left(\frac{\partial p'}{\partial t}+W\frac{\partial p'}{\partial z}\right) \\
&= \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)^2 \left(-P\omega i e^{iE}+PkWie^{iE}\right) \\
&= \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)(-\omega+kW)Pi(e^{iE}) \\
&= \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)(-\omega+kW)Pi(-i\omega e^{iE}+ikWe^{iE}) \\
&= \left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)(-\omega+kW)^2Pi^2(e^{iE}) \\
&= (-\omega+kW)^3Pi^3(e^{iE}) = (-\omega+kW)^3P(-i)(e^{iE}) \\
&= (\omega-kW)^3iP(e^{iE})
\end{aligned}$$

The cubed imaginary number i^3 simplyfies to $-i$ and the -1 is distributed into the cubed factor. This results in the final expression for Term 1:

$$\boxed{\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)^3 p' = i\exp[i(kz+n\theta-\omega t)](\omega-kW)^3P} \tag{8}$$

II.B. Term 2

Application of the solution form to second term (Eqn 9) of Lilley's equation is slightly more involved.

$$\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)(\overline{a^2}\nabla^2 p') \tag{9}$$

Term 2 requires computing the double divergence of perturbation pressure $\nabla^2 p'$:

$$\begin{aligned}
\nabla^2 p' &= \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial p'}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 p'}{\partial \theta^2} + \frac{\partial^2 p'}{\partial z^2} \\
&= \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{dP}{dr}e^{iE}\right) + \frac{1}{r^2}\frac{\partial}{\partial \theta}\left(Pnie^{iE}\right) + \frac{\partial}{\partial z}\left(Pkie^{iE}\right) \\
&= \frac{1}{r}e^{iE}\left(\frac{dP}{dr} + r\frac{d^2P}{dr^2}\right) + \frac{1}{r^2}Pn^2i^2e^{iE} + Pk^2i^2e^{iE} \\
&= e^{iE}\frac{d^2P}{dr^2} + e^{iE}\frac{1}{r}\frac{dP}{dr} + i^2e^{iE}\left(\frac{n^2}{r^2} + k^2\right)P
\end{aligned}$$

Thus, the double divergence of p' can be expressed in the following expression, which is seperated into like differential terms of P:

$$\boxed{\nabla^2 p' = \exp[i(kz + n\theta - \omega t)]\left[\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P\right]} \tag{10}$$

Substituting Eqn 10 into Term 2 (Eqn 9), we can perform the multi-derivative expression to derive the final term. All terms grouped with P as well as a are constant WRT t and z and can be carried outside of the derivative expression.

$$\begin{aligned}
\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)(\overline{a^2}\nabla^2 p') &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)\overline{a^2}e^{iE}\left[\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P\right] \\
&= \overline{a^2}\left[\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P\right]\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)e^{iE} \\
&= \overline{a^2}\left[\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P\right]\left(-\omega ie^{iE} + kWie^{iE}\right) \\
&= -\overline{a^2}ie^{iE}(\omega - kW)\left[\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P\right]
\end{aligned}$$

This results in the final expression for Term 2:

$$\boxed{\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)(\overline{a^2}\nabla^2 p') = -\overline{a^2}i\exp[i(kz + n\theta - \omega t)](\omega - kW)\left[\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P\right]} \tag{11}$$

II.C. Term 3

Applying the solution form to third term of Lilley's equation:

$$\begin{aligned}\frac{d\overline{a^2}}{dr}\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)\frac{dp'}{dr} &= \frac{d\overline{a^2}}{dr}\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)\frac{dP}{dr}e^{iE} \\ &= \frac{d\overline{a^2}}{dr}\frac{dP}{dr}\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)e^{iE} \\ &= \frac{d\overline{a^2}}{dr}\frac{dP}{dr}i(-\omega+kW)e^{iE}\end{aligned}$$

Applying the isentropic relationship assumption for speed of sound and taking the derivative WRT r:

$$\begin{aligned}\frac{d\overline{a^2}}{dr}\frac{dP}{dr}i(-\omega+kW)e^{iE} &= \frac{d}{dr}\left(\frac{\gamma\overline{p}}{\overline{\rho}}\right)\frac{dP}{dr}i(-\omega+kW)e^{iE} \\ &= -\left(\frac{\gamma\overline{p}}{\overline{\rho}^2}\right)\frac{d\overline{\rho}}{dr}\frac{dP}{dr}i(-\omega+kW)e^{iE} \\ &= \left(\frac{\overline{a^2}}{\overline{\rho}}\right)\frac{d\overline{\rho}}{dr}\frac{dP}{dr}i(\omega-kW)e^{iE}\end{aligned}$$

This results in the final expression for Term 3:

$$\boxed{\frac{d\overline{a^2}}{dr}\left(\frac{\partial}{\partial t}+W\frac{\partial}{\partial z}\right)\frac{dp'}{dr}=\overline{a^2}i\exp[i(kz+n\theta-\omega t)](\omega-kW)\frac{1}{\overline{\rho}}\frac{d\overline{\rho}}{dr}\frac{dP}{dr}} \tag{12}$$

II.D. Term 4

Apply solution form to fourth term of Lilley's equation:

$$\begin{aligned}2\overline{a^2}\frac{dW}{dr}\frac{\partial^2 p'}{\partial z\partial r} &= 2\overline{a^2}\frac{dW}{dr}\frac{\partial}{\partial z}\frac{\partial p'}{\partial r} = 2\overline{a^2}\frac{dW}{dr}\frac{\partial}{\partial z}\left(\frac{dP}{dr}e^{iE}\right) \\ &= 2\overline{a^2}\frac{dW}{dr}\frac{dP}{dr}\frac{\partial}{\partial z}\left(e^{iE}\right) = 2\overline{a^2}\frac{dW}{dr}\frac{dP}{dr}ike^{iE}\end{aligned}$$

This results in the final expression for Term 4:

$$\boxed{2\overline{a^2}\frac{dW}{dr}\frac{\partial^2 p'}{\partial z\partial r}=\overline{a^2}i\exp[i(kz+n\theta-\omega t)]2k\frac{dW}{dr}\frac{dP}{dr}} \tag{13}$$

II.E. Lilley's Equation Solution

To derive the final form of Lilley's equation, we combine Terms 1 through 4:

$$Eqn\ 8 - Eqn\ 11 - Eqn\ 12 + Eqn\ 13 = S(\vec{x},t)$$

Which becomes the following in expanded form:

$$\begin{aligned}&\left\{ie^{iE}(\omega-kW)^3P\right\}-\left\{-\overline{a^2}ie^{iE}(\omega-kW)\left[\frac{d^2P}{dr^2}+\frac{1}{r}\frac{dP}{dr}-\left(\frac{n^2}{r^2}+k^2\right)P\right]\right\} \\ &\quad -\left\{\overline{a^2}ie^{iE}(\omega-kW)\frac{1}{\overline{\rho}}\frac{d\overline{\rho}}{dr}\frac{dP}{dr}\right\}+\left\{\overline{a^2}ie^{iE}2k\frac{dW}{dr}\frac{dP}{dr}\right\}=S(\vec{x},t)\end{aligned}$$

Dividing both sides of the equation by the term $\overline{a^2} i e^{iE}(\omega - kW)$:

$$\frac{(\omega - kW)^2}{\overline{a^2}}P + \frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \left(\frac{n^2}{r^2} + k^2\right)P - \frac{1}{\overline{\rho}}\frac{d\overline{\rho}}{dr}\frac{dP}{dr} + \frac{1}{\omega - kW}2k\frac{dW}{dr}\frac{dP}{dr} = \frac{S(\vec{x},t)}{\overline{a^2} i e^{iE}(\omega - kW)}$$
$$\frac{d^2P}{dr^2} + \frac{1}{r}\frac{dP}{dr} - \frac{1}{\overline{\rho}}\frac{d\overline{\rho}}{dr}\frac{dP}{dr} + \frac{1}{\omega - kW}2k\frac{dW}{dr}\frac{dP}{dr} + \frac{(\omega - kW)^2}{\overline{a^2}}P - \left(\frac{n^2}{r^2} + k^2\right)P = \frac{S(\vec{x},t)}{\overline{a^2} i e^{iE}(\omega - kW)}$$

Grouping like terms of P, we obtain the final form of the general solution of Lilley's Equation:

$$\boxed{\frac{d^2P}{dr^2} + \left\{\frac{1}{r} - \frac{1}{\overline{\rho}}\frac{d\overline{\rho}}{dr} + \frac{1}{\omega - kW}2k\frac{dW}{dr}\right\}\frac{dP}{dr} + \left\{\frac{(\omega - kW)^2}{\overline{a^2}} - \left(\frac{n^2}{r^2} + k^2\right)\right\}P = \frac{S(\vec{x},t)}{\overline{a^2} i e^{iE}(\omega - kW)}} \tag{14}$$

where $E = kz + n\theta - \omega t$.

III. Problem 2 – General Solution for Jet Flow Far-Field

Next, we will solve the homogeneous form of Lilley's equation (Eqn 14) to determine the general solution for the pressure fluctuation outside of the jet in the ambient medium where the sources vanish. This region can be considered to be the far-field of the jet, so the solution must be chosen to ensure decaying solutions (outgoing waves).

Since this is far-field, ambient flow, we will assume that the mean velocity and density in this region are equal to the ambient mean values:

$$\begin{aligned} W(r) &= W_\infty = const \\ \overline{\rho} &= \overline{\rho_0} = const \end{aligned} \tag{15}$$

which allows us to simplify the homogeneous equation:

$$\frac{d^2P}{dr^2} + \left\{\frac{1}{r} - \frac{1}{\overline{\rho}}\frac{d\overline{\rho}}{dr} + \frac{1}{\omega - kW}2k\frac{dW}{dr}\right\}\frac{dP}{dr} + \left\{\frac{(\omega - kW)^2}{\overline{a^2}} - \left(\frac{n^2}{r^2} + k^2\right)\right\}P = 0$$
$$\frac{d^2P}{dr^2} + \left\{\frac{1}{r} - \frac{1}{\overline{\rho_0}}\cancelto{0}{\frac{d\overline{\rho_0}}{dr}} + \frac{1}{\omega - kW_\infty}2k\cancelto{0}{\frac{dW}{dr}}\right\}\frac{dP}{dr} + \left\{\underbrace{\frac{(\omega - kW_\infty)^2}{\overline{a^2}} - k^2}_{\lambda^2} - \frac{n^2}{r^2}\right\}P = 0$$

Resulting in the homogeneous Bessel equation:

$$\boxed{\frac{d^2P}{dr^2} + \left\{\frac{1}{r}\right\}\frac{dP}{dr} + \left\{\lambda_\infty^2 - \frac{n^2}{r^2}\right\}P = 0} \tag{16}$$

where the constant term is collected in λ_∞:

$$\lambda_\infty = \sqrt{\frac{(\omega - kW_\infty)^2}{\overline{a^2}} - k^2}$$

Substituting the definition of wavenumber $k = \dfrac{\omega}{\overline{a}}$, the expression can be reduced:

$$
\begin{aligned}
\lambda_\infty &= \sqrt{\frac{\left(\omega - \frac{\omega}{\bar{a}} W_\infty\right)^2}{\bar{a}^2} - \frac{\omega^2}{\bar{a}^2}} \\
&= \sqrt{\frac{\omega^2 \left(1 - \frac{W_\infty}{\bar{a}}\right)^2}{\bar{a}^2} - \frac{\omega^2}{\bar{a}^2}} \\
&= \sqrt{\frac{\omega^2}{\bar{a}^2} \left[\left(1 - \frac{W_\infty}{\bar{a}}\right)^2 - 1\right]} \\
&= \sqrt{\frac{\omega^2}{\bar{a}^2} \left(1 - 2\frac{W_\infty}{\bar{a}} + \frac{W_\infty^2}{\bar{a}^2} - 1\right)} \\
&= \sqrt{\frac{\omega^2}{\bar{a}^2} \left(\frac{W_\infty^2}{\bar{a}^2} - 2\frac{W_\infty}{\bar{a}}\right)} \\
&= \sqrt{\frac{\omega^2}{\bar{a}^2} \frac{W_\infty}{\bar{a}} \left(\frac{W_\infty}{\bar{a}} - 2\right)}
\end{aligned}
$$

Substituting the definition of Mach number $M_\infty = \dfrac{W_\infty}{\bar{a}}$, we obtain:

$$
\lambda_\infty = \sqrt{\frac{\omega^2}{\bar{a}^2} M_\infty (M_\infty - 2)} \tag{17}
$$

where M_∞ is the freestream Mach number and λ_∞ is an imaginary quantity for $M_\infty < 2$.

Eqn 16 is of the same form as Bessel's ordinary differential equation and can be solved using the first $H_n^{(1)}(\lambda_\infty r)$ and second $H_n^{(2)}(\lambda_\infty r)$ order Hankel functions:

$$
P = C H_n^{(1)}(\lambda_\infty r) + D H_n^{(2)}(\lambda_\infty r) \tag{18}
$$

where C and D are arbitrary constants and the first and second order Hankel functions are equal to the outgoing $H_n^+(x)$ and incoming $H_n^+(x)$ wave solutions, respectively, and $x = \lambda_\infty r$:

$$
H_n^{(1)}(\lambda_\infty r) = H_n^+(x) = J_n(x) + iY_n(x) \tag{19}
$$

$$
H_n^{(2)}(\lambda_\infty r) = H_n^-(x) = J_n(x) - iY_n(x) \tag{20}
$$

where the Hankel functions are composed of the first and second order Bessel functions J_n and Y_n, respectively.

Taking the limit of $H_n^+(x)$ and $H_n^-(x)$ as x and r approach infinity (far-field condition):

$$
\begin{aligned}
\lim_{x \to \infty} H_n^+(x) &= (-i)^{n+1} \frac{e^{ix}}{x} \\
\lim_{x \to \infty} H_n^+(x) &= i^{n+1} \frac{e^{-ix}}{x}
\end{aligned}
$$

In the far-field, flow velocity is subsonic ($M_\infty < 1$), making λ_∞ imaginary according to Eqn 17. This, in turn, makes the limit of $H_n^+(x)$ a diminishing exponential and the limit of $H_n^-(x)$ an increasing exponential, which is impossible. Thus, Eqn 16 reduces to Eqn 21 for final solution of the far-field outside of the potential core of the jet:

$$
\boxed{P = C H_n^{(1)}(\lambda_\infty r)} \tag{21}
$$

IV. Problem 3 – General Solution for Jet Potential Core

In this section, we will solve the homogeneous form of Lilley's equation (Eqn 14) to determine the general solution for the potential core region of the jet. We will assume that the mean velocity and density in this region are constant and equal to the jet exit values:

$$W(r) = W_j = const \tag{22}$$
$$\overline{\rho} = \overline{\rho_j} = const \tag{23}$$

These assumptions are similar to those made in Section III, which allows the reduction of Eqn 14 to the Bessel equation:

$$\boxed{\frac{d^2P}{dr^2} + \left\{\frac{1}{r}\right\}\frac{dP}{dr} + \left\{\lambda_j^2 - \frac{n^2}{r^2}\right\}P = 0} \tag{24}$$

with

$$\lambda_j = \sqrt{\frac{\omega^2}{\overline{a}^2}M_j(M_j - 2)} \tag{25}$$

where $M_j = \dfrac{W_j}{\overline{a}}$ is the jet exit Mach number and λ_j is an imaginary quantity for $M_j < 2$.

For this potential core solution, we assume the Bessel equation is equivalent to the first J_n and second Y_n order Bessel functions:

$$P = AJ_n(\lambda_j r) + BY_n(\lambda_j r) \tag{26}$$

where A and B are arbitrary constants.

The second order Bessel function Y_n is defined as having a singularity where $\lambda_j r = 0$, so inside the potential core region where $0 < r < R_j$, we can leave the singular Y_n out of Eqn 24, resulting in the final expression for the solution inside of the potential core:

$$\boxed{P = AJ_n(\lambda_j r)} \tag{27}$$

V. Problem 4 – Matching of Far-Field and Potential Core Solutions

In Section III we derived a solution for the pressure fluctuations (P^+) outside of the jet potential core in the far-field (Eqn 21) and in Section IV we derived a solution for the pressure fluctuations (P^-) inside the jet potential core (Eqn 27). Ideally, we would like to combine these near and far solutions by joining them at the edge of the potential core.

One method of accomplishing this solution is to replace the potential core boundary with a vortex sheet located at $r = R_j$. To successfully join the two solutions, we must match them at this vortex sheet according the two following conditions:

1. Pressure fluctuations on either edge of the vortex sheet must match:

$$P^+ = P^- \tag{28}$$

2. Radial velocity fluctuations accross the vortex sheet must match:

$$V_r^+ = V_r^- \tag{29}$$

To achieve the two constraints, we require two independent equations. These equations will be related through the variable V_r, which is the radial component of velocity at the vortex sheet:

$$V_r = \nu \exp[i(-\omega t + kz)] \tag{30}$$

where ν is a constant value of velocity in the radial direction that is modulated by the expression $\exp[i(-\omega t + kz)]$, which is a function of time and the jet flow direction z.

For the first constraint, we will derive a relationship between V_r and the pressure fluctuation using the momentum equation:

$$\overline{\rho_0}\frac{D\vec{V}}{Dt} = -\nabla P \tag{31}$$

Because the matching conditions apply only across the vortex sheet in the radial direction, we look at the radial component of Eqn 31 by taking the dot product of both sides WRT the radial unit vector $\hat{r}$:

$$\begin{aligned}
\overline{\rho_0}\frac{D\vec{V}}{Dt}\cdot\hat{r} &= -\nabla P\cdot\hat{r} \\
\overline{\rho_0}\frac{DV_r}{Dt} &= -\frac{\partial P}{\partial r}
\end{aligned}$$

The total time derivative of V_r is (see Section II):

$$\begin{aligned}
\frac{DV_r}{Dt} &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)V_r \\
&= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)\nu\exp[i(-\omega t + kz)] \\
&= -i(\omega - kW)\nu\exp[i(-\omega t + kz)] \\
&= -i(\omega - kW)V_r
\end{aligned}$$

Substituting $\frac{DV_r}{Dt}$ into Eqn 31, we obtain:

$$\overline{\rho_0}i(\omega - kW)V_r = \frac{dP}{dr} \tag{32}$$

which can be rearranged in terms of V_r, allowing us to satisfy our second constraint:

$$\boxed{V_r = \frac{\frac{dP}{dr}}{\overline{\rho_0}i(\omega - kW)}} \tag{33}$$

For the second equation, we will define the radial velocity fluctuation V_r instead as the time derivative of the radial displacement ζ of fluid particles located at the vortex sheet.

$$V_r = \frac{D\zeta}{Dt} \tag{34}$$

where

$$\zeta(z,t) = \chi\exp[i(-\omega t + kz)] \tag{35}$$

and where χ is a constant value of radial position that is modulated by the same expression $\exp[i(-\omega t + kz)]$, which governs the oscillation of the vortex sheet.

Taking the time derivative of ζ (see Section II):

$$\begin{aligned}
V_r = \frac{D\zeta}{Dt} &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)\zeta \\
&= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)\chi\exp[i(-\omega t + kz)] \\
&= -\chi\exp[i(-\omega t + kz)]i(\omega - kW) \\
&= -\zeta i(\omega - kW)
\end{aligned}$$

Applying the solution for V_r to the outer and inner vortex sheet velocity fluctuations:

$$V_r^+ = -\zeta i(\omega - kW_\infty) \tag{36}$$
$$V_r^- = -\zeta i(\omega - kW_j) \tag{37}$$

Rearranging the solution of V_r for $-\zeta$, and setting Eqns 36 and 37 equal to each other:

$$-\zeta = \frac{V_r^+}{i(\omega - kW_\infty)} = \frac{V_r^-}{i(\omega - kW_j)} \tag{38}$$

Consider a case in which the real jet is replaced by a vortex sheet at $r = R_j$. If the solutions are to be matched at the vortex sheet, describe what matching conditions should be applied. Give both the physical description and the mathematical expressions.