

# MAE 298 Aeroacoustics – Homework #2

## Lilley's Equation Solution and Application to Jet Noise

Logan D. Halstrom

*Graduate Student*

*Department of Mechanical and Aerospace Engineering*

*University of California, Davis, CA 95616*

**Abstract about homework**

### Nomenclature

$R_j$	Axisymmetric jet radius	$E$	Exponential term: $kz + n\theta - \omega t$
$W_j$	Jet exit mean velocity	$P$	
$\bar{\rho}_j$	Jet exit mean density	$k$	
$W_0$	Ambient mean velocity	$\omega$	
$\bar{\rho}_0$	Ambient mean density	$n$	
$W(r)$	Radial distribution of axial velocity	$r$	
$\bar{a}^2(r)$	Radial distribution of speed of sound squared	$\theta$	
$p'$	Perturbation pressure	$z$	
$i$	Imaginary number $\sqrt{-1}$		

### I. Background

Lilley's equation for a parallel axisymmetric flow (Eqn 1):

$$\left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right)^3 p' - \left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) \left(\bar{a}^2 \nabla^2 p'\right) - \frac{d\bar{a}^2}{dr} \left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) \frac{dp'}{dr} + 2\bar{a}^2 \frac{dW}{dr} \frac{\partial^2 p'}{\partial z \partial r} = S(\vec{x}, t) \quad (1)$$

$$\text{where } \nabla^2 \equiv \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2}$$

$W(r)$  is the radial distribution of axial velocity and  $\bar{a}^2(r)$  is the radial distribution of speed of sound squared.

### II. Problem 1 – Solution to Lilley's Equation

Seek solutions of Lilley's equation in the form:

$$p'(r, \theta, z, t) \sim P(r) \exp[i(kz + n\theta - \omega t)] \quad (2)$$

Assume:

1.  $\bar{a}^2 = \frac{\gamma \bar{p}}{\bar{\rho}}$
2.  $\bar{p} = \text{const}$

where  $\bar{\rho}(r)$  is the radial distribution of the mean density. Show that Lilley's equation reduces to (Eqn 3):

$$\frac{d^2 P}{dr^2} + \left\{ \frac{1}{r} - \frac{1}{\bar{\rho}} \frac{d\bar{\rho}}{dr} + \frac{2k}{(\omega - kW)} \frac{dW}{dr} \right\} \frac{dP}{dr} + \left\{ \frac{(\omega - kW)^2}{a^2} - k^2 - \frac{n^2}{r^2} \right\} P = RHS \quad (3)$$

To aid in this derivaiton, we find the results for the first-order partial derivatives of parameters relevant to the solution. First, we compute the derivative of the perturbation pressure  $p'$  with respect to (WRT) time  $t$ :

$$\begin{aligned} \frac{\partial p'}{\partial t} &= \frac{\partial}{\partial t} \{P(r) \exp[i(kz + n\theta - \omega t)]\} \\ &= P \frac{\partial}{\partial t} [i(kz + n\theta - \omega t)] \exp[i(kz + n\theta - \omega t)] \\ &= iP \left[ \frac{\partial}{\partial t} (kz) + \frac{\partial}{\partial t} (n\theta) - \frac{\partial}{\partial t} (\omega t) \right] \exp[i(kz + n\theta - \omega t)] \end{aligned}$$

$$\boxed{\frac{\partial p'}{\partial t} = -iP\omega \exp[i(kz + n\theta - \omega t)] = -iP\omega e^{iE}} \quad (4)$$

where  $E = kz + n\theta - \omega t$ ,  $\frac{\partial E}{\partial t} = -\omega$ , and  $\frac{\partial E}{\partial z} = k$ . Next, we compute the derivative of  $p'$  WRT the angular direction  $\theta$ :

$$\begin{aligned} \frac{\partial p'}{\partial \theta} &= \frac{\partial}{\partial \theta} \{P(r) \exp[i(kz + n\theta - \omega t)]\} \\ &= P \frac{\partial}{\partial \theta} [i(kz + n\theta - \omega t)] \exp[i(kz + n\theta - \omega t)] \\ &= iP \left[ \frac{\partial}{\partial \theta} (kz) + \frac{\partial}{\partial \theta} (n\theta) - \frac{\partial}{\partial \theta} (\omega t) \right] \exp[i(kz + n\theta - \omega t)] \end{aligned}$$

$$\boxed{\frac{\partial p'}{\partial \theta} = iPn \exp[i(kz + n\theta - \omega t)] = iPn e^{iE}} \quad (5)$$

Next, we compute the derivative of  $p'$  WRT the axial flow direction  $z$ :

$$\begin{aligned} \frac{\partial p'}{\partial z} &= \frac{\partial}{\partial z} \{P(r) \exp[i(kz + n\theta - \omega t)]\} \\ &= P \frac{\partial}{\partial z} [i(kz + n\theta - \omega t)] \exp[i(kz + n\theta - \omega t)] \\ &= iP \left[ \frac{\partial}{\partial z} (kz) + \frac{\partial}{\partial z} (n\theta) - \frac{\partial}{\partial z} (\omega t) \right] \exp[i(kz + n\theta - \omega t)] \end{aligned}$$

$$\boxed{\frac{\partial p'}{\partial z} = iPk \exp[i(kz + n\theta - \omega t)] = iPk e^{iE}} \quad (6)$$

Next, we compute the derivative of  $p'$  WRT the radial direction  $r$ :

$$\frac{\partial p'}{\partial r} = \frac{\partial}{\partial r} \{P(r) \exp[i(kz + n\theta - \omega t)]\}$$

$$\boxed{\frac{\partial p'}{\partial r} = \frac{dP}{dr} \exp[i(kz + n\theta - \omega t)] = \frac{dP}{dr} e^{iE}} \quad (7)$$

### II.A. Term 1

To simplify the derivation, we will apply the solution form individually to each term in Lilley's equation. For the first term, we must apply the multi-derivative operator a total of three times:

$$\begin{aligned} \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)^3 p' &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)^2 \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right) p' \\ &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)^2 \left(\frac{\partial p'}{\partial t} + W\frac{\partial p'}{\partial z}\right) \\ &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)^2 (-P\omega e^{iE} + PkW e^{iE}) \\ &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right) \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right) (-\omega + kW)Pi(e^{iE}) \\ &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right) (-\omega + kW)Pi(-i\omega e^{iE} + ikW e^{iE}) \\ &= \left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right) (-\omega + kW)^2 Pi^2(e^{iE}) \\ &= (-\omega + kW)^3 Pi^3(e^{iE}) = (-\omega + kW)^3 P(-i)(e^{iE}) \\ &= (\omega - kW)^3 iP(e^{iE}) \end{aligned}$$

The cubed imaginary number  $i^3$  simplifies to  $-i$  and the  $-1$  is distributed into the cubed factor. This results in the final expression for Term 1:

$$\boxed{\left(\frac{\partial}{\partial t} + W\frac{\partial}{\partial z}\right)^3 p' = i \exp[i(kz + n\theta - \omega t)](\omega - kW)^3 P} \quad (8)$$

## II.B. Term 2

Application of the solution form to second term (Eqn 9) of Lilley's equation is slightly more involved.

$$\left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) (\bar{a}^2 \nabla^2 p') \quad (9)$$

Term 2 requires computing the double divergence of perturbation pressure  $\nabla^2 p'$ :

$$\begin{aligned} \nabla^2 p' &= \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial p'}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 p'}{\partial \theta^2} + \frac{\partial^2 p'}{\partial z^2} \\ &= \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{dP}{dr} e^{iE} \right) + \frac{1}{r^2} \frac{\partial}{\partial \theta} (P n i e^{iE}) + \frac{\partial}{\partial z} (P k i e^{iE}) \\ &= \frac{1}{r} e^{iE} \left( \frac{dP}{dr} + r \frac{d^2 P}{dr^2} \right) + \frac{1}{r^2} P n^2 i^2 e^{iE} + P k^2 i^2 e^{iE} \\ &= e^{iE} \frac{d^2 P}{dr^2} + e^{iE} \frac{1}{r} \frac{dP}{dr} + i^2 e^{iE} \left( \frac{n^2}{r^2} + k^2 \right) P \end{aligned}$$

Thus, the double divergence of  $p'$  can be expressed in the following expression, which is separated into like differential terms of  $P$ :

$$\boxed{\nabla^2 p' = \exp[i(kz + n\theta - \omega t)] \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right]} \quad (10)$$

Substituting Eqn 10 into Term 2 (Eqn 9), we can perform the multi-derivative expression to derive the final term. All terms grouped with  $P$  as well as  $a$  are constant WRT  $t$  and  $z$  and can be carried outside of the derivative expression.

$$\begin{aligned} \left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) (\bar{a}^2 \nabla^2 p') &= \left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) \bar{a}^2 e^{iE} \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right] \\ &= \bar{a}^2 \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right] \left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) e^{iE} \\ &= \bar{a}^2 \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right] (-\omega i e^{iE} + k W i e^{iE}) \\ &= -\bar{a}^2 i e^{iE} (\omega - kW) \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right] \end{aligned}$$

This results in the final expression for Term 2:

$$\boxed{\left(\frac{\partial}{\partial t} + W \frac{\partial}{\partial z}\right) (\bar{a}^2 \nabla^2 p') = -\bar{a}^2 i \exp[i(kz + n\theta - \omega t)] (\omega - kW) \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right]} \quad (11)$$

### II.C. Term 3

Applying the solution form to third term of Lilley's equation:

$$\begin{aligned}\frac{\overline{da^2}}{dr} \left( \frac{\partial}{\partial t} + W \frac{\partial}{\partial z} \right) \frac{dp'}{dr} &= \frac{\overline{da^2}}{dr} \left( \frac{\partial}{\partial t} + W \frac{\partial}{\partial z} \right) \frac{dP}{dr} e^{iE} \\ &= \frac{\overline{da^2}}{dr} \frac{dP}{dr} \left( \frac{\partial}{\partial t} + W \frac{\partial}{\partial z} \right) e^{iE} \\ &= \frac{\overline{da^2}}{dr} \frac{dP}{dr} i(-\omega + kW) e^{iE}\end{aligned}$$

Applying the isentropic relationship assumption for speed of sound and taking the derivative WRT  $r$ :

$$\begin{aligned}\frac{\overline{da^2}}{dr} \frac{dP}{dr} i(-\omega + kW) e^{iE} &= \frac{d}{dr} \left( \frac{\gamma \overline{p}}{\overline{\rho}} \right) \frac{dP}{dr} i(-\omega + kW) e^{iE} \\ &= - \left( \frac{\gamma \overline{p}}{\overline{\rho}^2} \right) \frac{d\overline{p}}{dr} \frac{dP}{dr} i(-\omega + kW) e^{iE} \\ &= \left( \frac{\overline{a^2}}{\overline{\rho}} \right) \frac{d\overline{p}}{dr} \frac{dP}{dr} i(\omega - kW) e^{iE}\end{aligned}$$

This results in the final expression for Term 3:

$$\boxed{\frac{\overline{da^2}}{dr} \left( \frac{\partial}{\partial t} + W \frac{\partial}{\partial z} \right) \frac{dp'}{dr} = \overline{a^2} i \exp[i(kz + n\theta - \omega t)] (\omega - kW) \frac{1}{\overline{\rho}} \frac{d\overline{p}}{dr} \frac{dP}{dr}} \quad (12)$$

### II.D. Term 4

Apply solution form to fourth term of Lilley's equation:

$$\begin{aligned}2\overline{a^2} \frac{dW}{dr} \frac{\partial^2 p'}{\partial z \partial r} &= 2\overline{a^2} \frac{dW}{dr} \frac{\partial}{\partial z} \frac{\partial p'}{\partial r} = 2\overline{a^2} \frac{dW}{dr} \frac{\partial}{\partial z} \left( \frac{dP}{dr} e^{iE} \right) \\ &= 2\overline{a^2} \frac{dW}{dr} \frac{dP}{dr} \frac{\partial}{\partial z} (e^{iE}) = 2\overline{a^2} \frac{dW}{dr} \frac{dP}{dr} i k e^{iE}\end{aligned}$$

This results in the final expression for Term 4:

$$\boxed{2\overline{a^2} \frac{dW}{dr} \frac{\partial^2 p'}{\partial z \partial r} = \overline{a^2} i \exp[i(kz + n\theta - \omega t)] 2k \frac{dW}{dr} \frac{dP}{dr}} \quad (13)$$

### II.E. Lilley's Equation Solution

To derive the final form of Lilley's equation, we combine Terms 1 through 4:

$$Eqn\ 8 - Eqn\ 11 - Eqn\ 12 + Eqn\ 13 = RHS$$

Which becomes the following in expanded form:

$$\begin{aligned} & \left\{ i e^{iE} (\omega - kW)^3 P \right\} - \left\{ -\overline{a^2} i e^{iE} (\omega - kW) \left[ \frac{d^2 P}{dr^2} + \frac{1}{r} \frac{dP}{dr} - \left( \frac{n^2}{r^2} + k^2 \right) P \right] \right\} \\ & - \left\{ \overline{a^2} i e^{iE} (\omega - kW) \frac{1}{\bar{\rho}} \frac{d\bar{\rho}}{dr} \frac{dP}{dr} \right\} + \left\{ \overline{a^2} i e^{iE} 2k \frac{dW}{dr} \frac{dP}{dr} \right\} = RHS \end{aligned}$$

## II.F. Read Case Data

Parameters pertaining to the run conditions are acquired by processing the output of the “qinfo” script, which provides Mach number  $M$ , Reynold’s number  $Re$ , freestream specific heat ratio  $\gamma$ , freestream static temperature  $T_\infty$ , freestream density  $\rho_\infty$ , freestream dynamic pressure  $q_\infty$ , OVERFLOW angle of attack  $\alpha_{OVR}$ , and OVERFLOW sideslip angle  $\beta_{OVR}$ , among other parameters. Freestream parameters are converted into the grid reference unit of inches.

Parameters for simple pendulum oscillations including maximum amplitude  $\theta_{max}$  and oscillation frequency  $f$  are read from the case’s “Scenario.xml” file.  $\alpha_{OVR}$  is converted from the capsule reference frame to the parachute reference frame according to Eqn. 14.

$$\alpha_G = 180^\circ - \alpha_{OVR} \quad (14)$$

## II.G. Condition Raw Data

Raw OVERFLOW data must also be conditioned before aerodynamic calculations are performed. All force data is read from “aero.chute” which is produced by the “overlst -p” function call. Raw data is first trimmed of any duplicates. Force and moment coefficients are then corrected for actual reference parameters as show in Eqn. 16 that may be different than those contained in “mixsur.i”, which is used as input for surface integration during OVERFLOW solver iterations.

$$\begin{aligned} C_F &= C_{F,OVR} \frac{A_{ref,old}}{A_{ref,new}} \\ C_M &= C_{M,OVR} \frac{A_{ref,old}}{A_{ref,new}} \frac{l_{ref,old}}{l_{ref,new}} \end{aligned} \quad (15)$$

Cases that are run with a half-body geometry symmetric about the y-axis are then corrected for symmetry effects. Forces and moments that are expected to be non-zero must be doubled, to account for half of the geometry missing, and forces that are expected to be zero due to geometry symmetry must be forced to zero.

$$\begin{aligned} C_X &= 2 \cdot C_{X,OVR} \\ C_Y &= 0 \\ C_Z &= 2 \cdot C_{Z,OVR} \\ C_l &= 0 \\ C_m &= 2 \cdot C_{m,OVR} \\ C_n &= 0 \end{aligned} \quad (16)$$

Because these simulations are time accurate and have relative mesh motion ( $DYNMCS = .T.$ ), it is necessary to calculate the simulation time from the iteration history. Non-dimensional simulation time can be calculated by multiplying the iteration number  $Iter$  subtracted by the iteration where grid motion began  $i_{start}$  by the non-dimensional time step  $DTPHYS$ , as shown in Eqn. 17. Dimensional time in seconds can then be calculated by multiplying by the reference length of one inch and dividing by the reference velocity, which is equal to the freestream velocity  $V_\infty$ .

$$\begin{aligned} t_{non} &= (Iter - i_{start}) \cdot DTPHYS \\ t &= \left( \frac{L_{ref}}{V_{ref}} \right) t_{non} \end{aligned} \quad (17)$$

## III. Process Dynamic Parameters

With OVERFLOW outputs corrected and brought into the time domain, actual aerodynamic properties can now be calculated.

### III.A. Prescribed Motion Dynamics

First, attitude angles and rates are calculated according to the prescribed motion equation contained in “Scenario.xml”. For these simulations, where oscillations are contained within the y-plane, the only attitude parameters to calculate are those pertaining to pitching motion  $\theta$  and aerodynamic clock angle  $\phi$ .  $\theta$  and the associated rates can be calculated based on the simulation time for the simple, harmonic pendulum according to Eqn. 18 and for the Wind Tunnel Test (WTT) Track case according to Eqn. 19. The simple pendulum is constrained to simple, harmonic motion defined by a frequency and an amplitude, while the WTT Track case is designed to approximate the motion of the parachute wind tunnel test and is derived using a Fast Fourier Transform (FFT).

$$\begin{aligned}\omega &= 2\pi f \\ \theta(t) &= \theta_{max} \cdot \sin(\omega t) \\ \dot{\theta}(t) &= \theta_{max} \cdot \omega \cdot \cos(\omega t) \\ \ddot{\theta}(t) &= -\theta_{max} \cdot \omega^2 \cdot \sin(\omega t)\end{aligned}\tag{18}$$

where  $\dot{\theta}$  is the pitch rate of the parachute,  $\theta_{max}$  is the maximum angular amplitude of the oscillation,  $f$  is the frequency of the oscillation in Hertz, and  $t$  is the simulation time.

$$\begin{aligned}\omega_1 &= 2\pi(0.186246), \omega_2 = 2\pi(0.200573) \\ A_1 &= 0.011695, A_2 = 4.268271, A_3 = 1.001533, A_4 = 4.983758 \\ \theta(t) &= -A_1 \sin(\omega_1 t) - A_2 \cos(\omega_1 t) + A_3 \sin(\omega_2 t) + A_4 \cos(\omega_2 t) \\ \dot{\theta}(t) &= -A_1 \omega_1 \cos(\omega_1 t) + A_2 \omega_1 \sin(\omega_1 t) + A_3 \omega_2 \cos(\omega_2 t) - A_4 \omega_2 \sin(\omega_2 t) \\ \ddot{\theta}(t) &= A_1 \omega_1^2 \sin(\omega_1 t) + A_2 \omega_1^2 \cos(\omega_1 t) - A_3 \omega_2^2 \sin(\omega_2 t) - A_4 \omega_2^2 \cos(\omega_2 t)\end{aligned}\tag{19}$$

where  $\omega$  are the natural frequencies of the oscillation and  $A$  are the associated Fourier coefficients.

The rotational (tangential) velocity of the parachute  $V_T$  is calculated from the rate of pitching motion  $\dot{\theta}$  and the fixed rotational radius  $R = 2714.61in$ , as depicted in Eqn. 20.  $V_T$  is in the normal body direction ( $Z_b$ -axis) and will always have the same sign as the pitching motion. Eqn. 20 produces the tangential velocity of the actual parachute, so this value will be negated to produce the relative wind velocity in the tangential direction, which is demonstrated in Fig. ?? (Derivation of the tangential velocity can be found later on in Eqn. 38).

$$V_T = R\dot{\theta}\tag{20}$$

### III.B. Aerodynamic Parameters

With all of the velocities and angular rates calculated, actual aerodynamic parameters such as effective velocity  $V_{eff}$  and effective angle of attack  $\alpha_{eff}$  can be determined. Referencing Fig. ??, it can be seen that the angle  $V_\infty$  makes with the body axis  $X_b$  of the parachute at any given time is equivalent to  $\alpha_G - \theta$ . Thus, the resulting relative air velocity components in the body frame axes are computed according to Eqn. 21.

$$\begin{aligned}V_{X_b,rel} &= -V_\infty \cos(\alpha_G - \theta) \\ V_{Z_b,rel} &= V_\infty \sin(\alpha_G - \theta) - V_T\end{aligned}\tag{21}$$

where velocity components are of the air surrounding the parachute and are in the opposite respective directions of the actual motion of the parachute.

From the component velocities, the effective relative velocity of the parachute  $V_{eff}$  can be calculated as the magnitude as in Eqn. 22.

$$V_{eff} = \sqrt{[V_{X_b,rel}]^2 + [V_{Z_b,rel}]^2}\tag{22}$$

The effective angle of attack is calculated from the ratio of the body velocity components.

$$\alpha_{eff} = \tan^{-1} \left[ \frac{V_{Z_b,rel}}{-V_{X_b,rel}} \right]\tag{23}$$



where  $V_{X_b,rel}$  is negated due to the direction of the  $X_b$ -axis to maintain consistency with the definition of angle of attack.

For a better comparison with the wind tunnel test, angle of attack is also calculated in the total frame, as seen in Eqn. 24.

$$\alpha_T = \cos^{-1} \left[ \frac{V_{X_b,rel}}{V_{eff}} \right] \quad (24)$$

Because there is no side-slip in the CFD model,  $\alpha_T$  is effectively the absolute value of  $\alpha$  in these simulations.

Aerodynamic rates ( $\dot{\alpha}$ ,  $\dot{\alpha}_T$ ) are calculated from the  $\alpha$ ,  $\alpha_T$  data series using numerical differencing, and is shown in Eqn. 25.

$$\dot{\alpha}_{T,i} = \frac{\alpha_{T,i+1} - \alpha_{T,i-1}}{2\Delta t} \quad (25)$$

Central differencing is used to calculate all middle derivatives, while forward and backward differencing are used to calculate the first and last derivatives in the time series, respectively.

To more closely compare to the wind tunnel test,  $\dot{\alpha}_T$  is non-dimensionalized according to Eqn. 26 by a reference length equal to the parachute skirt radius and the effective velocity  $V_{eff}$  of the parachute at the given instant in time.

$$\dot{\alpha}_{T,non} = \frac{|\dot{\alpha}_T| l_{ref}}{2V_{eff}} \quad (26)$$

Finally, the aerodynamic clock angle  $\phi$ , which is defined in Eqn. 27, must be calculated to allow coordinate rotations into the missile frame.

$$\phi = \tan^{-1} \left( \frac{-V_{Y_b,rel}}{V_{Z_b,rel}} \right) \quad (27)$$

Though  $V_{Y_b,rel}$  is always zero for the in-plane oscillations of the CFD simulations (as will  $C_Y$ ), the changing sign of  $V_{Z_b,rel}$  makes it such that  $\phi$  will be either 0 or  $2\pi$ , making the missile rotation matrix in Eqn. 31 in Section III.C either the identity or negative identity matrix.

### III.C. Force and Moment Coefficient Rotations

Because the body and missile coordinate systems will be used to compare CFD results to the WTT, a series of coordinate frame rotations are required to transform OVERFLOW  $C_X$ ,  $C_Z$  to  $C_A$ ,  $C_{N,m}$ . However, all coefficients must first be corrected for inaccurate non-dimensionalization by the freestream dynamic pressure. Because OVERFLOW does not bookkeep the relative velocity of the grids during the simulation, all coefficients are non-dimensionalized simply by the freestream dynamic pressure  $q_\infty$  though they should actually be non-dimensionalized by the dynamic pressure of the flow in the frame of the actual parachute, defined in Eqn. 28.

$$q = \frac{1}{2} \rho V_{eff}^2 \quad (28)$$

Assuming incompressible flow ( $\rho = \rho_\infty$ ), this correction can be expressed simply as the ratio of the squares of the freestream and effective velocities, as derived in Eqn. 29.

$$\begin{aligned} C_F &= \frac{F}{q} = \frac{F}{q_\infty} \cdot \frac{q_\infty}{q} = C_{F,OV} \cdot \frac{q_\infty}{q} \\ \frac{q_\infty}{q} &= \frac{\frac{1}{2} \rho_\infty V_\infty^2}{\frac{1}{2} \rho_\infty V_{eff}^2} = \left( \frac{V_\infty}{V_{eff}} \right)^2 \end{aligned} \quad (29)$$

After the non-dimensionalization correction, forces can then be rotated first from the inertial frame to the body frame, and then from the body frame into the missile frame

The rotation into the body frame shown in Fig. ?? and derived in Eqn. 30 is a simple, 2D rotation about the pitch axis for the CFD cases, as the yaw angle is always zero.

$$\begin{bmatrix} C_A \\ C_N \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} C_X \\ C_Z \end{bmatrix} \quad (30)$$

The missile frame rotation in Eqn 31, as previously explained, manifests as either the identity or negative identity matrix, effectively causing a discontinuity in the normal force coefficient  $C_N$ .

$$\begin{bmatrix} C_{Y,m} \\ C_{N,m} \end{bmatrix} = \begin{bmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{bmatrix} \begin{bmatrix} C_Y \\ C_N \end{bmatrix} \quad (31)$$

## IV. Apparent Mass Adjustment

To maintain consistency with the wind tunnel test, an apparent mass adjustment is applied the body forces to account for the displacement of the air trapped within the parachute as it pitches.

### IV.A. Apparent Force

The apparent mass of the trapped air  $m_a$  is calculated as the mass of the air within a hemisphere of diameter  $L_{ref}$  multiplied by a scale factor, as seen in Eqn. 32.

**apparent force diagram**

$$m_a = \frac{1}{2} \left[ \frac{4}{3} \pi \left( 0.7 \frac{L_{ref}}{2} \right)^3 \right] \cdot \rho_\infty \quad (32)$$

The force the parachute exerts on this mass is then calculated according to Newton's second law, using the axial and normal accelerations in Eqn. 33 derived from the prescribed equations of motion.

$$\bar{a}_b = R\dot{\theta}^2 \hat{X}_b + R\ddot{\theta} \hat{Z}_b \quad (33)$$

By Newton's third law, this force must be subtracted from the aerodynamic forces acting on the parachute calculated by surface integration in OVERFLOW, resulting in Eqn. 34 for the final expression for the aerodynamic forces on the parachute.

$$\Sigma \bar{F} = \bar{F}_{aero} - m_a \bar{a}_b \quad (34)$$

### IV.B. Accelerations

To calculate the forces required to accelerate the apparent mass, it is necessary to derive the equations of motion pertaining to the parachute's oscillation. Representing the parachute as a single point rotating at a distance  $R$  behind its attachment point in the negative  $\hat{X}_b$  direction, the expression for its position is given in Eqn. 35.

$$\bar{r}_b = -R \hat{X}_b \quad (35)$$

The velocity of the parachute due to oscillations in  $\theta$  can be determined by differentiating Eqn. 35 with respect to time. There is no change in the radius in these simulations, so the relative acceleration term in Eqn. 36 goes to zero, leaving the derivative of the  $\hat{X}_b$  unit vector.

$$\bar{v}_b = \frac{B d}{dt} \bar{r}_b = \cancel{\frac{B d}{dt}} - R \frac{d\hat{X}_b}{dt} \quad (36)$$

From Eqn. 30,  $\hat{X}_b$  can be substituted into terms of inertial frame unit vectors and simplified as in Eqn. 37.

$$\begin{aligned}
&= -R \cdot \frac{d}{dt} (\cos \theta \hat{X}_i - \sin \theta \hat{Z}_i) \\
&= -R \left( -\dot{\theta} \sin \theta \hat{X}_i - \dot{\theta} \cos \theta \hat{Z}_i \right) \\
&= R \dot{\theta} \underbrace{(\sin \theta \hat{X}_i + \cos \theta \hat{Z}_i)}_{\hat{Z}_b}
\end{aligned} \tag{37}$$

$$\boxed{\bar{v}_b = R \dot{\theta} \hat{Z}_b}$$

Finally, the velocity can be differentiated to produce the acceleration components in the body axis system.

$$\begin{aligned}
\bar{a}_b &= \frac{{}^B d}{dt} \bar{v}_b = \frac{{}^B d}{dt} (R \dot{\theta} \hat{Z}_b) = R \left[ \ddot{\theta} \hat{Z}_b + \dot{\theta} \frac{d\hat{Z}_b}{dt} \right] \\
&= R \left[ \ddot{\theta} \hat{Z}_b + \dot{\theta} \frac{d}{dt} (\sin \theta \hat{X}_i + \cos \theta \hat{Z}_i) \right] = R \left[ \ddot{\theta} \hat{Z}_b + \dot{\theta} (\dot{\theta} \cos \theta \hat{X}_i - \dot{\theta} \sin \theta \hat{Z}_i) \right] \\
&= R [\ddot{\theta} \hat{Z}_b + \dot{\theta}^2 \underbrace{(\cos \theta \hat{X}_i - \sin \theta \hat{Z}_i)}_{\hat{X}_b}]
\end{aligned} \tag{38}$$

$$\boxed{\bar{a}_b = R \dot{\theta}^2 \hat{X}_b + R \ddot{\theta} \hat{Z}_b}$$