MAE 298 Aeroacoustics – Final Project Aeroacoustics of Aerospace Vehicles During Launch and Ascent

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Abstract about lit study

Nomenclature

0	Subscript for quiescent parameters	V_n	Velocity component in surface normal direction
e	Subscript for emission parameters	M	Mach number
L	Subscript for loading parameters	c	Speed of Sound
t	Time	$\overline{ ho}$	Mean density
au	Retarded time	p	Pressure
ret	Evaluated at retarded time	\widetilde{p}	Pressure (Discontinuous across data surface)
ĥ	Unit surface normal vector	\overline{p}	Mean pressure
\vec{r}	Radial direction vector	p'	Perturbation pressure
î	Radial unit vector	ΔP	Pressure difference from CFD solution
r	Magnitude of radial vector $ \vec{r} $	L	Pressure loading
θ	Angle between \hat{n} and \hat{r}	$\overline{\partial}$	Generalized derivative
$ec{V}$	Function of surfaces within a fluid space	δ	Dirac delta function
$ec{V}$	General velocity vector	FW-H	Ffowcs Williams-Hawkings
V_r	Velocity component in radial direction		

I. Introduction

In this assignment, we will first derive the wave equation for sound generated by a moving body using generalized differentiation. This process is similar to that used by Ffowcs Williams and Hawkings in the derivation of their equation describing aeroacoustic noise. Equations like these are useful for computing aeroacoustic effects once source terms such as body loading are calculated.

Secondly, we will derive Farassat's Formulation 1A of the Ffowcs Williams-Hawkings (FW-H) equation for pressure due to loading noise beginning with Formulation 1. This result will

demonstrate the process for solving the FW-H equation and may be applied to the other terms to obtain Farassat's complete Formulation 1A of the FW-H equation.

II. Main Discussion

In this section, we will derive the wave equation for sound generated by a moving body, which is known as the Kirchhoff formula for moving surfaces. The general acoustic wave equation with no source term (homogeneous) is expressed as follows:

$$\boxed{\frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} - \nabla^2 p = 0} \tag{1}$$

We will used generalized differentiation to show that the wave equation whose sound is generated by an arbitrary moving body f = 0 can be expressed as follows:

$$\left[\frac{1}{c^2} \frac{\overline{\partial}^2 \widetilde{p}}{\partial t^2} - \overline{\nabla}^2 \widetilde{p} = -\left[\frac{M_n}{c} \frac{\partial p}{\partial t} + p_n \right] \delta(f) - \frac{1}{c} \frac{\partial}{\partial t} \left[M_n p \delta(f) \right] - \nabla \cdot \left[p \hat{n} \delta(f) \right] \right]$$
(2)

where \hat{n} is the unit normal vector on the surface and $p_n = \nabla p \cdot \hat{n}$.

We will use Greens function of the wave equation in the unbounded space to find the unknown function $p(\vec{x},t)$, which exists everywhere in space. This result is called the Kirchhoff formula for moving surfaces.

II.A. Sub topic

The moving surface generating acoustic effects will be defined as $f(\vec{x},t) = 0$ and will be referred to as a "data surface". The definition of f allows the fluid to be defined at all points in space, which can be divided into three regions:

Data Surface: f = 0Exterior Region: f > 0Interior Region: f > 0

We will also assume that the gradient of f is the surface outward unit normal vector \hat{n} :

$$\nabla f = \hat{n} \tag{3}$$

Under this definition of f, the fluid will exist in all regions of unbounded space and thus enable the usage of Green's function to solve for the sound field.

We will solve for the sound field in the surface exterior region, but we will assume that the fluid that extends into the interior region has the same conditions as an undisturbed, quiescent medium. We can define a pressure variable with these properties as \tilde{p} using the embedding technique as follows:

$$\widetilde{p} = \begin{cases} p, & f > 0 \\ 0, & f < 0 \end{cases}$$
 (4)

Thus, pressure will vary outside of the data surface and will be undisturbed inside. Though not continuous, this function of pressure will be defined at all points in an unbounded space and thus applicable to Green's function.

III. Conclusions

We have now completed the derivation of an aeroacoustic equation describing the pressure fluctuations around a moving body and then demonstrated a method of solving equations such as these. In a real life application, the equation being solved would be more complicated and include more source terms, like the complete FW-H equation does. This equation can be solved by applying Farassat's Formulation 1A and implementing the solution numerically. Source inputs for loading noise (ΔP) would be obtained from an unsteady CFD simulation and used to inform the aeroacoustic solution.

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