

(Green Rep. Formula) Let  $u$  be a Helmholtz soln (of radiating type for exterior case), then

$$\text{for } x \in \mathbb{R}^d \setminus \bar{\Omega} \text{ (exterior)} \quad u(x) = \pm \int_{\partial\Omega} [u_n(y)\Phi(x,y) - u(y) \frac{\partial\Phi(x,y)}{\partial n}] ds_y$$

Pf: • Interior pf:  
same as GGF for Laplace op, since singularity same.  
• exterior uses rad. cond. (CK Thm 2.4).

• exterior:  $\uparrow$   
BIE soln. for  $u^+$ :  $u^+ = D\tau$  satis. Helmholtz in  $\mathbb{R}^d \setminus \bar{\Omega}$ , radiative.

solves exterior BVP if  $(I + 2D)\tau = 2f = -2u^+|_{\partial\Omega}$  inc. field.  
 $2u^+$  by JR3

However, there's a problem:  $\text{Nul}(I + 2D)$  not injective for certain  $K$ , so nonunique soln, numerically cond # -> 0  
Interpretation. Thm: say  $K^2 = E_j^{(n)}$  for which  $\begin{cases} \Delta u = k^2 u \text{ in } \Omega \\ u_n = 0 \text{ on } \partial\Omega \end{cases}$  has discrete set of Neumann eigenvalues of  $\Delta$ .  
then.  $\dim \text{Nul}(I + 2D) > 0$ .

Pf: SLP  $u := S\delta$ , JR2 says  $0 = u^- = (D^T + \frac{1}{2})\delta$   
so nontriv.  $u \Rightarrow$  nontriv.  $\delta$ .  
 $\Rightarrow \dim \text{Nul}(I + 2D^T) > 0$

By Fredholm alternative for cpt op  $D$ ,  $\dim \text{Nul}(I + 2D) > 0$ .

Show evolving signals of  $2D$  as  $kc$  increases. -- Matlab. where hit -1 get Neumann eigenvalues, +1 " Dirichlet".  
got to here.

Stop interior Neumann resonances from plaguing scattering soln: use rep.  $u^+ = (D - iyS)\tau$ ,  $y > 0$   
This solves ext. BVP if  $(I + 2D - 2iyS)\tau = -2u^+|_{\partial\Omega}$ . Brakhage-Werner, Leis, Panich (1966)  
Thm:  $\frac{1}{2}I + D - iyS$  injective  $\forall k > 0$   $2u^+$  by JR3, 1.

Pf: suppose  $(\frac{1}{2} + D - iyS)\tau = 0$ , wish to show  $\tau \equiv 0$ .

say  $V := (D - iyS)\tau$ , then  $V^+ = 0$  by construction of BIE.

$\Rightarrow V = 0$  in  $\mathbb{R}^d \setminus \bar{\Omega}$  by uniqueness of ext. Dir. BVP. for radiative solns  
 $\Rightarrow V_n^+ = 0$  on  $\partial\Omega$ .

$$\begin{aligned} JR1, 3 \Rightarrow V^+ &= -\tau \\ JR2, 4 \Rightarrow V_n^+ &= -iy\tau \end{aligned} \quad \left. \begin{aligned} &\text{(a)} \\ &\text{(a)} \end{aligned} \right\}$$

$$\text{GT1 in } \Omega \text{ gives } \int_{\partial\Omega} \bar{V}^- V_n^- ds = \int_{-\Lambda} \bar{V} \Delta v + \bar{\nabla} \bar{V} \cdot \bar{\nabla} v dx$$

$$\underbrace{iy \int |\tau|^2 ds}_{\text{from (a)}}$$

since  $y \neq 0$ , take Im part shows  $\tau \equiv 0$ , QED.

$$\int_{\Omega} -k|v|^2 + |\nabla v|^2 dx \text{ pure real.}$$

In practice, choose  $y \approx k$ .

Note:  $D - iyS$  has a singular kernel since  $S(s, t) = \frac{1}{2\pi} \ln |y(s) - y(t)| + O(1)$ . Need quadrature rule for singular integrand.

Several ways to handle this:

i) Martensen-Kussmann (Kress M91) (CK books)

give quadr weights, on uniform periodic nodes

for integration of  $\ln(\sin^2 \frac{t-t_j}{2}) f(t)$ ,  $f$  anal,

where  $t_j$  is one of the nodes. A or  $2\pi$  periodic log sing.

Then  $D(t, t_j) - iy S(t, t_j) = \ln(4 \sin^2 \frac{t-t_j}{2}) K_1(t, t_j)$

+  $k_2(t, t_j)$

quad weights for unit periodic nodes,

which can handle smooth +  $(\ln(t-s))$  smooth.

→ spectral convergence, excellent, less work needed to derive weights. L split.

- diagonal values  $D(t_i, t_i)$  or  $S(t_i, t_i) (\in \infty)$  not needed

- algebraic conv. up to  $\sim 10^{th}$  order, but prefactor bad.

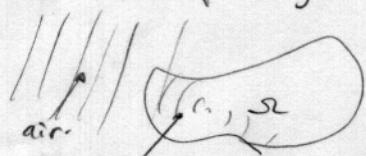
iii) Alpert (1999) hybrid Gauss-trapezoidal: use non-equispaced nodes near singularity.

In vs. large systems ( $N > 10^4$ ) it becomes better to use Fast Multipole Method to compute action of  $D$  or  $S$  on density  $\tau$ , a fast matrix-vector product in  $O(N \ln N)$  flops.

The lin. sys. is then solved via iterative meth. which require only  $x \rightarrow Ax$  products (a few tens of them).

FMM key idea: field due to arb. # sources inside ball can be rep. outside ball by fixed # L.I. forces.

• Transmission prob: eg acoustic scatt off medium of diff. wave speed.



medium, refractive index  $n = \frac{k}{k_0}$  (speed n times slower).

Say  $u^s = S_0 \sigma_0$  outside to be equal to densities, radii,  $\sigma_0, \sigma$  unknowns

$$\left. \begin{aligned} (\Delta + k_0^2) u &= 0 && \text{in } \mathbb{R}^2 \setminus \bar{\Omega} \\ (\Delta + k^2) u &= 0 && \text{in } \Omega \end{aligned} \right\} \text{we say } u = u^i + u^s$$

$u^i = \begin{cases} e^{ik_0 r} & \text{outside} \\ 0 & \text{inside} \end{cases}$

$u$  cont. on  $\partial\Omega$

$u^i$  cont. " "

$$\left. \begin{aligned} u^s &\text{ obeys PDE} \\ u^i &= u^s && \text{on } \partial\Omega \\ u^i &= u^s && \text{inside} \end{aligned} \right\} \text{then } S_0 \sigma_0 - S \sigma = f$$

$$K_{00} \sigma_0 - K_{00} \sigma = u^i - u^s$$

$$U_n^s - U_n^i = u_n^i - u_n^s$$

$$(D_0^T \nabla \frac{1}{2}) \sigma_0 - (D^T + \frac{1}{2}) \sigma = g$$

$$\text{ie } \begin{bmatrix} S_0 & -S \\ D_0^T + \frac{1}{2} & -D^T + \frac{1}{2} \end{bmatrix} \begin{bmatrix} \sigma_0 \\ \sigma \end{bmatrix} = \begin{bmatrix} f \\ g \end{bmatrix}$$

coupled BIEs., but not 2nd kind since not of form

$$\left( \begin{bmatrix} I & 0 \\ 0 & I \end{bmatrix} + \begin{bmatrix} K_{00} & K_{00} \\ K_{00} & K_{00} \end{bmatrix} \right) \begin{bmatrix} \sigma_0 \\ \sigma \end{bmatrix} = \begin{bmatrix} f \\ g \end{bmatrix}$$

Correct way is use  $u^s = D_0 \tau + S_0 \sigma$  outside  
 $u^s = D \tau + S \sigma$  inside.

which gives 2nd kind & cancels singularities. (trick of Rokhlin 1983)

$$\begin{bmatrix} D_0^T + \frac{1}{2} - (D - \frac{1}{2}) & S_0 - S \\ T_0 - T & D_0^T + \frac{1}{2} - (D^T - \frac{1}{2}) \end{bmatrix} \begin{bmatrix} \tau \\ \sigma \end{bmatrix} = \begin{bmatrix} f \\ g \end{bmatrix}$$

$\Rightarrow I + \text{compact.}$

use Nyström to get  $\tau, \sigma$  then recon. us field.

## Debugging

i) bug  $\approx$  needle in haystack

$\rightarrow$  break haystack into <sup>smallest poss</sup> pieces, test them as you go.

Eg. defining domain w/  $y_i$  bdy pts,  $N_j$  norms, plot it before you move on.

ii) examine visually everything you can.

- debugging
- test seq. for BWLP.
- MPS
  - complete
  - closed sets
  - in  $\mathbb{R}^d$
  - least-sq. BVP
- EV-

Recall BIE:  $(I+2D)\mathbf{r} = \mathbf{f}$  for exterior radiation BVP  $U^S$ , but  $2D$  has eigen -1 wherever  $k$  is interior Neumann eigenvalue causes cond  $(I+2D) \rightarrow \infty$  at such k, even though exterior  $U^S$  doesn't blow up  $\rightarrow$  numerical problem.

GWLP. last time (4)

(skip 5)

MPS

$$\text{approx } U(\mathbf{x}) = \sum_{j=1}^N \alpha_j \phi_j(\mathbf{x}).$$

basis funcs, particular solns.

to PDE, but not the BCs.

Then adjust  $\{\alpha_j\}_{j=1}^N := \vec{\alpha}$  coeff. vector to get BCs close to correct.

eg. replace  $\Delta u = 0$  in  $\Omega$ ,  $u = f$  on  $\partial\Omega$ .

Recall complex anal:  $f(z)$  anal. at 0, has radius of conv.

recalling Taylor  $f(z) = f(0) + z f'(0) + \frac{z^2}{2!} f''(0) \dots$  convex in disc,  $|z| < r$ . outside  $|z| > r$ .



coeffs fixed by beh. at 0.

unfinished.

2) Then (Faber, Szegő, Walsh):

if  $f$  anal. in  $D \setminus \overline{\Omega}$ ,  $\exists$  seg. polys.  $p_n(z)$   $n=0, 1, 2, \dots$

st.  $|f(z) - p_n(z)| \leq c K^{-n}$

for some  $K > 1$ .

In fact  $K$  is 'conformal dist' to nearest sing.



The pol polys will not in general be Taylor series;  
coeffs change each time

is every harm in the  
Re of analytic func?

Skipped

IMP5 for interior BVP:  $\begin{cases} (\Delta + k^2)u = 0 & \text{in } \Omega \\ u = f & \text{or } \partial\Omega \end{cases}$

$k=0 : \phi_n(x) = e^{inx} r^{int}$ ,  $n \in \mathbb{Z}$  (Laplace)

$k>0 : \phi_n(x) = e^{inx} J_n(kr)$

Choose truncations  $-N \leq n \leq N$ , i.e.  $2N+1$  basis funcs.

Given  $\vec{x} \in \mathbb{C}^{2N+1}$ ,  $u^{(n)}(\vec{x}) = \sum_{n=-N}^N \alpha_n \phi_n(x)$  is interior PDE soln.

BC error func is  $U^{(n)}|_{\partial\Omega} - f$ , want as small as poss.

$$\Rightarrow \min_{\vec{x}} \int_{\Omega} \left| \sum_n \alpha_n \phi_n(x) - f(x) \right|^2 dx \quad \text{square of } L^2(\Omega) \text{ norm.}$$

quadrature nodes  $y_j \in \Omega$ , weights  $w_j$   $\int_{\Omega} g(y) dy \approx \sum_{j=1}^M w_j g(y_j)$

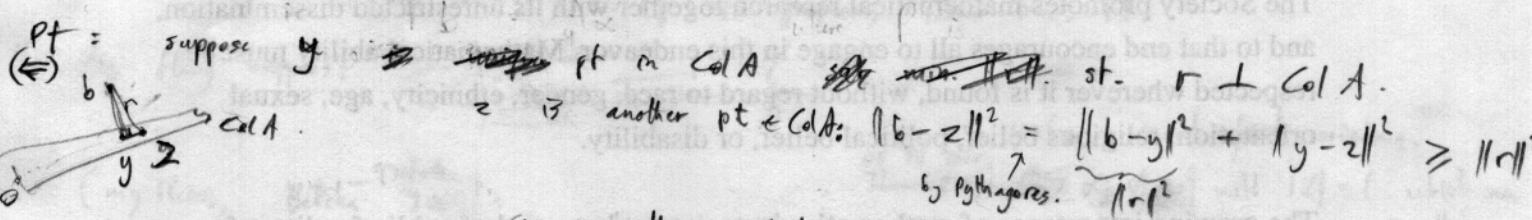
$$\sum_j w_j \left| \sum_n \alpha_n \phi_n(y_j) - f(y_j) \right|^2 = \left\| \sum_n A_{jn} \alpha_n - b_j \right\|_2^2 \quad \text{where } A_{jn} := \sqrt{w_j} \phi_n(y_j) \text{ rect. matrix.}$$

For what is  $\|A\vec{x} - \vec{b}\|^2$  minimized? linear least squares problem. (linalg.), usually overdetermined you may solve by Matlab's \ command.

What is actually going on here?  $\vec{r} = A\vec{x} - \vec{b}$  resid. vec.

Then: a vector  $\vec{x}$  minimizes norm  $\|\vec{r}\|_2$  iff  $\vec{r} \perp \text{Col } A$ , i.e.  $A^* \vec{r} = 0$

(mult. by  $A^*$ ) i.e.  $A^* A \vec{x} = A^* \vec{b}$  (Normal Eqns.)

$\Leftrightarrow$  suppose  $y \neq \vec{r}$  pt in Col  $A$  st.  $\vec{r} \perp \text{Col } A$ .  
  
 $\Rightarrow$  another pt  $\in \text{Col } A$ :  $\|\vec{b} - z\|^2 = \|\vec{b} - y\|^2 + \|\vec{y} - z\|^2 \geq \|\vec{r}\|^2$  by Pythagoras.

Solve Normal eqns. by SVD  $A = U \Sigma V^*$   $A^* = V \Sigma^* U^*$

$$A^* A = V \Sigma^* \Sigma V^*$$

$$\therefore V \Sigma^* \Sigma^* \vec{x} = V \Sigma^* \vec{b}$$

$V$  invertible,  $\Sigma^*$  has no zero entries if  $A$  full rank  $\Rightarrow \Sigma^* \vec{x} = U^* \vec{b}$

If  $A$  not full rank,  $\Sigma^{-1}$  is used with  $\vec{x}$

Lastalg algorithms are backwards stable.  $\therefore$  solves exactly for some  $A + \delta A$  with  $\frac{\|E\|}{\|A\|} = O(\epsilon_{mach})$ .

$A^+$ , pseudo inverse

summary: set up matrix  $A$  whose cols. are the basis funcns eval. at bdry nodes, (3)  
(rhs. vector)

solve  $A\vec{x} = \vec{b}$  in least-sq. sense.

Evaluate soln.  $\phi_{\vec{x}} = \sum \alpha_n \phi_n(x)$  at whatever  $x \in \Omega$  you need.

Eigen prob:

$$\begin{cases} \text{find } E \text{ s.t. } -\Delta u = Eu \text{ in } \Omega. \\ \text{3 nontriv } u \text{ s.t. } u = 0 \text{ on } \partial\Omega \end{cases} \quad \leftarrow \text{Helm } (\Delta + k^2)u = 0 \text{ with eigen } E = k^2$$

MPS: namely: want non-triv  $\vec{x}$  s.t.  $u^{(n)} = \sum \alpha_n \phi_n(x)$  vanishes on  $\partial\Omega$

if such  $\vec{x}$  exists,  $k^2$  is an eigenvalue &  $u^{(n)}$  is corr. efunc  $\phi_n$ .

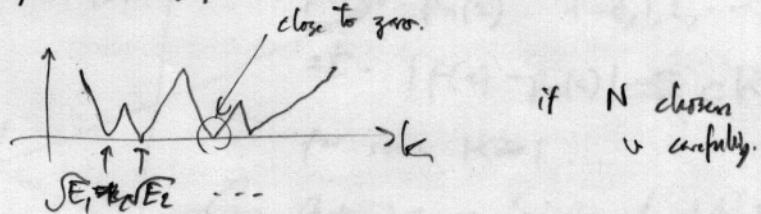
$$\text{mass: vanishing by } \|u^{(n)}\|_{L^2(\partial\Omega)}^2 \stackrel{\text{1st}}{=} \sum w_i |\phi_n(y_i) \alpha_n|^2 = \|A\vec{x}\|_2^2$$

$$\int_{\Omega} |\sum \alpha_n \phi_n(y)|^2 dy$$

we minimize  $\|Ax\|$  while holding  $\|\vec{x}\| = 1$  to prevent trivial solns.

$$\text{but } \min_{\vec{x}} \frac{\|Ax\|}{\|\vec{x}\|} = \sigma_1 \text{ min. sing. val of } A.$$

Now sweep  $k$ , plotting  $\sigma_1 = \sigma_1(k)$ :

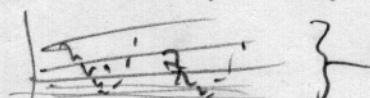


if  $N$  chosen  
v carefully.

get to here

However to increase accuracy of rep. of eigenfunc, want more  $N$ .

Interesting thing happens:



Solu. (my firs, Bachel Trefethen 2009).

$$\min_{\vec{x}} \frac{\|u\|_{L^2(\partial\Omega)}}{\|u\|_{L^2(\Omega)}}$$

intuition: norm:  $\|u\|_{L^2(\Omega)}^2 = \int_{\Omega} |u|^2 dx \approx \frac{1}{I} \sum_{j=1}^I |u(z_j)|^2$ .  $\rightarrow 0$  as  $I \rightarrow \infty$   
can approx. quite crudely via I int. pts.  $z_j$

useless: cannot find minimum.  
as  $N$  incr,  
there exist  $\sum \alpha_n \phi_n$  with  $\|\vec{x}\| = 1$  which are  
appr. small in  $\Omega$ , i.e.  $\|u\|_{L^2(\Omega)} \rightarrow 0$

Eg: Bessel funcns behave like

$$J_n(z) = \frac{1}{n!} \left(\frac{z}{2}\right)^n$$

$\in O(z^n)$