

Nystrom beautifully efficient, & inherits the convergence of the quadr. scheme applied to  $K(s, \cdot)u(\cdot)$ ,  
 for certain well-behaved ops  $K$ , called 'compact'.  
 If drop  $I$  in (16) can apply to 45°-kind (HW4), solving for  $u_j^{(n)}$ , but there is no interpolation formula (N).

Compact operators: may have  $\infty$ -dim range but behave 'like' finite-dim. ops. length

consider func  $f$  in  $\text{eg } C[a, b]$  as point in topological space,  $X = C[a, b]$ , w/ metric (norm)  $\|f\|$ ,  
 $\text{eg. } \|f\|_\infty = \sup_{s \in [a, b]} |f(s)|$

- seq  $(f_n)_{n=1,2,\dots}$  bounded if  $\|f_n\| \leq C (< \infty) \quad \forall n = 1, 2, \dots$  note: seq. goes on forever. (long time)
- seq.  $(f_n)$  converges to  $f \in X$  if  $\forall \varepsilon > 0$  no matter how small,  $\exists N$  st.  $\|f_n - f\| < \varepsilon \quad \forall n \geq N$

Say  $f_n$  were points in finite-dim space, eg  $\mathbb{R}^m$ , we'd apply Bolzano-Weierstrass.

Thm: in finite-dim normed lin. space, every bounded seq. contains a convergent subseq. also goes on forever

Eg. try it in  $\mathbb{R}$ : the only way to avoid having some limit pt. escape to  $\pm \infty$ .

B-W proof uses this fact successively on each coord 1, 2, ..., in  $\mathbb{R}^m$

However  $\infty$ -dim space such as  $C[a, b]$ ,  $L^\infty[a, b]$ , etc. are fundamentally different.

Eg. seq.  $\{\sin(nx)\}_{n=1}^\infty$  is bounded in  $C[0, 2\pi]$  or  $L^\infty[0, 2\pi]$ , but has no convergent subseq.

Eg. any other o.n.b.

E.g.  $\ell_2$ : space of sequences  $\{a_1, a_2, \dots\}$  with  $\sum_{j=1}^\infty |a_j|^2 < \infty$  has o.n.b.

Convergence: Seq.  $(x_n)$  in  $\infty$ -dim space  $X$  converges to  $x$  if  $\forall \varepsilon > 0 \exists N$  st.  $\|x_n - x\| < \varepsilon \quad \forall n \geq N$

Defns: Lin. op.  $K: X \rightarrow Y$  between normed lin. spaces  $X$  &  $Y$  is  
 compact if given any bounded seq.  $(x_n)$  in  $X$ , the seq.  $(Kx_n)$  contains a convergent subseq.  
 (I.e.  $K$  maps bounded sets to 'precompact sets' (which must contain convergent subseq's))

Eg. if  $K$  has finite-dim range, B-W  $\Rightarrow (Kx_n)$  has conv. subseq., so  $K$  cpt. as if, finite-dim.  
 bounded

$$\text{eg. } (Ku)(s) = \int_a^b \sum_{i=1}^m \alpha_i(s) \beta_i(t) u(t) dt$$

finite sum of rank-1 kernels.

Conversely  $K = \text{Id}$  with  $X=Y$   $\infty$ -dim spaces maps  $(x_n) \rightarrow (x_n)$  so need not contain conv. subseq.,  
 ie  $\text{Id}$  is compact iff  $Y$  finite-dim. not cpt.

Thm: Compact ops are bounded. pf suppose not, then  $\exists$  seq.  $(x_n)$  with  $\|x_n\|=1 \quad \forall n=1, 2, \dots$   
 $\text{s.t. } \|Kx_n\| > n$   
 But this cannot contain convergent subseq.  $\Rightarrow K$  not compact.

But, bounded  $\not\Rightarrow$  compact, viz.  $\text{Id}$ .

$K$  Compact: finite-dim, or  $\infty$ -dim in a way that components made arbitrarily small as dims  $\rightarrow \infty$ .

- Coding tips.
- $\|K\|_2$  upper bound via Hilbert-Schmidt (② 10/21)  $\Rightarrow$  then  $k(s,t) = \frac{1}{|s-t|^r}$  is H-S for  $0 < r < 1/2$
- Compactness. (③ 10/21)

Why compactness useful?

- 1) Thm (Fredholm Alternative). Let  $K: X \rightarrow X$  be compact on normed lin. space  $X$ . Then either i) for each  $f \in X$ ,  $(I - K)f = f$  has unique soln.  $u \in X$   
 or ii) hom. eqn.  $(I - K)u = 0$  has nontrivial soln.
- If case i) holds,  $(I - K)^{-1}$ , whose existence is asserted there, is bounded.

- Ie,  $(I - K)$  surjective (onto) iff injective (triv. nullspace). Cf. invertible matrix thm:  $A\vec{x} = \vec{f}$  soln exists, unique, iff  $A\vec{x} = \vec{0}$  has only triv. soln.

- Means  $(I + \text{compact})$  behaves like finite-dim square matrix (either invertible or not).

Useful: if  $\exists$   $\epsilon$  s.t.  $(I - K)u = 0$  has only triv. soln, ie  $1$  is not an eigenval of  $K$ , then 2nd kind IE solns always exist, unique.

NB: 2) we can

- 2) we can prove convergence rate of Nyström method for compact  $K$  is  $\|u^{(n)} - u\|_\infty \leq C \|Ku - K_n u\|$

Defn:  $(x_n)$  Cauchy convergent if  $\forall \epsilon > 0$ ,  $\exists N$  st.  $\|x_n - x_m\| \leq \epsilon$  for all  $n, m > N$ . "all pairs get arbitrarily close". applied to  $K(s, \cdot)u(\cdot)$

Tests for compactness:

- 1) Thm: Let  $K_1, K_2, \dots$  be seq. of compact ops  $X \rightarrow Y$  st.  $\lim_{n \rightarrow \infty} \|K_n - K\| = 0$ , then  $K$  compact.
- Rmk: eg. cpt if seq. of finite-dim ops which conv. in norm to it, eg  $U_j \rightarrow \int_0^1 u_j$  in  $L_2$ :  $K_n$  is  $U_j \rightarrow \int_0^1 U_j$   $j \in \mathbb{N}$ , so  $K = K_n$
- Pf: let  $(x_n)$  bounded seq.  $\exists$  subseq  $(x_{n_k})$  st.  $(K_{n_k} x_{n_k})$  converges.  $\exists$  subseq of this,  $(x_{n_{k_j}})$  st.  $(K_{n_{k_j}} x_{n_{k_j}})$  conv. ... etc.

$$\begin{matrix} x_n \\ x_{1n} & \xrightarrow{\text{def}} & \dots & \xrightarrow{\text{def}} & x_n \\ \vdots & & & & \end{matrix}$$

$T_k x_{n_k}$

$\vdots$

choose 'diagonal subseq'  $x_{n_{k_j}}$ : for each  $k \in \mathbb{N}$ ,  $(T_k x_{n_{k_j}})$  conv. since subseq of  $x_{n_{k_j}}$ .

If we can show  $(T_{n_{k_j}})$  conv. to  $t$ , we've proved  $K$  cpt since for any  $(x_n)$  we found  $(T_{n_{k_j}} x_n)$  has conv. subseq.

Let  $\epsilon > 0$ :  $\exists k$  st.  $\|T_k - T_{k+1}\| < \frac{\epsilon}{3M}$  where  $M = \sup\|x_n\|$

let  $k_1$  be st.  $\|T_{k_1} x_{n_{k_j}} - T_{k_1} x_{n_{k_{j+1}}}\| \leq \frac{\epsilon}{3}$  ( $n_{k_j}, n_{k_{j+1}} > k_1$ ) (since  $T_{k_1} x_{n_{k_j}}$  (it's Cauchy))

then  $\|T_{n_{k_j}} x_{n_{k_{j+1}}} - T_{n_{k_{j+1}}} x_{n_{k_{j+1}}}\| \leq \underbrace{\|T_{n_{k_j}} - T_{k_1}\|}_{\epsilon/3} + \underbrace{\|T_{k_1} x_{n_{k_j}} - T_{k_1} x_{n_{k_{j+1}}}\|}_{\epsilon/3} + \underbrace{\|T_{k_1} x_{n_{k_{j+1}}} - T_{n_{k_{j+1}}} x_{n_{k_{j+1}}}\|}_{\epsilon/3} \leq \epsilon$

so  $(T_{n_{k_j}})$  Cauchy  $\Rightarrow$  conv. to element in  $Y$  by completeness.

Ask who looked at each other's work? Who's looked at debriefings? 10/23. (2)

## Coding tips:

- if have param. such as max n ( $n=1 \dots 20$ ), set it once at start of code & have all ~~loop~~ reference this.

- switches in code to avoid duplication (don't make many nearly-identical codes; more errors, harder to maintain, read) → times. at start: use handles  $f = @()$  ... Eval. later but never need to quadr. schemes. type

modular :

$$\begin{bmatrix} x \\ w \end{bmatrix} = \text{gauss}(n)$$

↑ nodes      ↑ wei      ↑ order. n

main code eg Nystöm  
loop over n  
use  $x, w$ . → quadr.

defines a uniform interface.

have other quadr. schemes return in same way:

$$\begin{bmatrix} x \\ w \end{bmatrix} = \text{compttrap}(n)$$

equal(newtoncotes( $\frac{n}{2}$ ))

(Finish cpt)

[Lec. 10]

10/28/08.

- Cpt ops  $K$ :
- Fredholm Alt: unique solvability of  $\lambda(I-K)u=f$  if  $\text{Nul}(I-K)=\{0\}, \dots$
  - Nystöm conv. at same rate as quadr. scheme.
- What means for spectrum of  $K$ ? 1 is not eigenvalue. ( $\lambda=1: Ku=u$  for some  $u$ ).

Pf: i) show that  $K$  maps bounded seq. ( $u$ ) to equicontinuous' backed seq.

e.g.  $u_n(t) = x^n$  continuous on  $[0,1]$  for each  $n$ ,  $\lim_{n \rightarrow \infty} u_n(t) = x$  equicontinuous on  $[0,1]$  ... why not?

ii) Then by Arzela-Ascoli theorem such sets are compact. (contain compact subseq.) ... proof by diagonalization &  $\epsilon$ - $\delta$  argument.

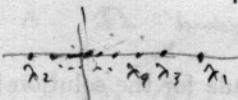
Show i):  $|(Ku)(s)| \leq M|b-a| \max_{s,t} |k(s,t)|$

where  $M := \sup_{s,t} \|u(s)\|$ .

$k$  cont on square:  $\forall \epsilon > 0 \exists \delta$  st.  $|k(s,t) - k(r,t)| \leq \frac{\epsilon}{M(b-a)}$  whenever  $|s-r| < \delta$

Then  $|(Ku)(s) - (Ku)(r)| \leq \left| \int_a^b (k(s,t) - k(r,t)) u(t) dt \right| \leq |b-a| \max_{|s-r|<\delta} |k(s,t) - k(r,t)| \|u\| \leq \epsilon$

Spectral Then for compact ops:  $K$  compact: has spectrum of discrete, finite-multiplicity eigenvalues, with 0 the only limit pt.



no nasty 'continuous spectrum, etc.'

since  $\lambda_n \rightarrow 0$ . think of <sup>ball</sup> ellipsoid w/ shrinking semi-axes.

( $\lambda_n = \sigma_n$  SVD, if Hermitian self-adjoint)



## Lec 10 -

• finish cpt ops.

Tools for  
Laplace's eqn.

$$\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2}$$

$$\Delta u = 0$$

$\Leftrightarrow u$  a harmonic func.

steady-state heat flow, diffusion  
electrostatics

div. op. kills

GT1, 2-

plan: finish cpt.  
(20 mins).

Final soln. in 2d,  $\Delta$   
Div. Thm & Adm.

GRF -

write as layer pots,  
define LPs, as int. op.

formula for  $\frac{\partial \Phi}{\partial y}(x, y)$ .

change-of-variables  
for  $y(t)$  parametr.  
curve.

1) Fundamental soln.  $\Phi(x, y) = -\frac{1}{2\pi} \ln |x-y|$

$$x = (x_1, x_2) \quad y = (y_1, y_2)$$

aka 'free-space Greens func': will give integral kernel for solution op.

Fact: for each  $y \in \mathbb{R}^2$ ,  $\Phi(\cdot, y)$  is harmonic in  $\mathbb{R}^2 \setminus \{y\}$

$$\text{ie } \Delta_x \Phi(x, y) = 0 \quad \forall x \neq y$$

Check: set  $y=0$  without loss of generality, then  $\frac{\partial}{\partial x_1} \ln|x| = \frac{1}{2} \frac{\partial}{\partial x_1} \ln(x_1^2 + x_2^2)$

$$= \frac{1}{2(x_1^2 + x_2^2)} \cdot 2x_1 = \frac{x_1}{|x|^2} \quad x \neq 0. \text{ etc.}$$

since  $\Phi$  symm,  $\Delta_y \Phi(x, y) = 0 \quad \forall y \neq x$ .

2) Divergence Thm:  $\Omega \subset \mathbb{R}^2$  bounded open domain w/ boundary  $\partial\Omega$ ,  $\vec{a} = (a_1(x), a_2(x))$  vector field  
1 cont. on  $\partial\Omega$  &  $C^1$  cont. inside,

$$\int_{\Omega} \vec{\nabla} \cdot \vec{a} \, dx = \int_{\partial\Omega} \hat{n} \cdot \vec{a} \, ds$$

$$\text{div } \vec{a} := \frac{\partial a_1}{\partial x_1} + \frac{\partial a_2}{\partial x_2}$$

$$\text{arclength on } \partial\Omega$$

$\Omega \rightarrow \hat{n}$  outward unit normal.

outward flux through  $\partial\Omega$ .  $\int \hat{n}(y) \cdot \vec{a}(y) \, dy$  explicitly.

Greens Thms:  $u, v$  be sufficiently smooth funcs defined in  $\bar{\Omega}$  & closure  $\bar{\Omega} + \partial\Omega$

$$\int_{\Omega} (u \Delta v + \nabla u \cdot \nabla v) \, dx = \int_{\partial\Omega} u \left[ \frac{\partial v}{\partial n} \right] \, ds \quad (\text{GT1})$$

$$:= v_n := \hat{n} \cdot \nabla v$$

$$\int_{\Omega} u \Delta v - v \Delta u \, dx = \int_{\partial\Omega} u v_n - v u_n \, ds \quad (\text{GT2})$$

pf: do as WS?

$$\vec{\nabla} \cdot (u \vec{v}) = \text{prod rule } \nabla u \Delta v + \nabla u \cdot \nabla v.$$

apply Div. Thm to  $\vec{a} = \vec{v}$  to get GT1.

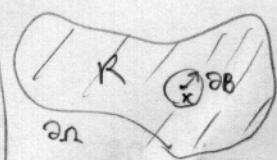
GT2 by subtracting GT1 with  $u \leftrightarrow v$  from itself.

sufficient smoothness is  $\bar{\Omega}$  having  $C^1$  smooth bdry,  $u, v \in C^2(\bar{\Omega})$ , however can be loosened to  $\bar{\Omega}$  w/ corners.

Corollary: setting  $u=1$  in GT1; If  $v$  harmonic in  $\Omega$  then  $\int_{\partial\Omega} v_n \, ds = 0$  (zero-flux)

Green's Representation formula: Let  $u$  be harmonic in  $\Omega$ , then for each  $x \in \Omega$ ,  
 $u(x) = \int_{\partial\Omega} u_n(y) \Phi(x, y) - u(y) \frac{\partial \Phi}{\partial n}(x, y) ds_y$   
 interior values from  
 i.e. integral kernels acting on bdy funcns.  $u, u_n$ .

Pf:  $\partial B(x; r) = \text{circle rad. } r \text{ about } x$  with <sup>inwards</sup>



in  $R = \Omega \setminus \overbrace{B(x; r)}^{\text{closed ball}}$   $\Phi(x, y)$  harmonic as func. of  $y$ .

GT2 in  $R$  for  $u$ ,  $v = \Phi(x, \cdot)$ :

$$\int_R u \Delta_y \Phi(x, y) - \Phi(x, y) \Delta u dy = \int_{\partial R} u(y) \frac{\partial \Phi}{\partial n_y} - \Phi(x, y) \frac{\partial u}{\partial n}(y) ds_y$$

<sup>so</sup>  $= 2\pi + \partial B$   
 inwards-pointing  $n$ .

$$\begin{aligned} \int_{\partial R} \Phi(x, y) u(y) - \frac{\partial \Phi}{\partial n_y}(x, y) u(y) ds_y &= \int_{\partial B(x; r)} \underbrace{\frac{\partial \Phi(x, y)}{\partial n_y} u(y)}_{y \text{ on } \partial B} - \underbrace{\Phi(x, y) u(y)}_{\frac{1}{2\pi} \ln r} ds_y \\ &= \underbrace{\frac{1}{2\pi} \int_{\partial B} u(y) ds_y}_{+ \frac{1}{2\pi} \ln r} + \underbrace{\frac{1}{2\pi} \ln r \int_{\partial B} u(y) ds_y}_{= 2\pi \cdot u(y) \text{ for some } y \in \partial B \text{ by MVT.}} \end{aligned}$$

<sup>by zero-flux</sup>

Finally take  $\lim_{r \rightarrow 0}$ , get  $\lim_{r \rightarrow 0} u(y) \Big|_{y \in \partial B} = u(x)$ .

Cor: since  $\Phi \in \frac{\partial \Phi}{\partial n_y}$  analytic func of  $1^{st}$  variable,  $u$  analytic in  $\Omega$  regardless how nasty the bdy data  $u, u_n$  is.

Given func  $\sigma$  on  $\partial\Omega$ ,  
 Define single-layer integral op  $u$

double-

$$(S\sigma)(x) := \int_{\partial\Omega} \Phi(x, y) \sigma(y) dy$$

$$(D\tau)(x) = \int_{\partial\Omega} \frac{\partial \Phi(x, y)}{\partial n_y} \tau(y) dy$$

then GRF says.  $u = S\sigma + D\tau$  if choose  $\sigma = u_n|_{\partial\Omega}$

Other consequences of GRF: let  $u$  be harmonic,

Choose  $\Omega = B(x; r)$  then  $u(x) = \ln r \int_{\partial\Omega} u(y) ds_y - \frac{1}{2\pi r} \int_{\partial\Omega} u(y) ds_y$  : value at center = mean value of  $u$  over  $\partial B(x; r)$

(2) Maximum principle).

1) Computation:

$$\text{given } x, y, \vec{n}_y : \frac{\partial \vec{\phi}(x, y)}{\partial \vec{n}_y} = -\frac{1}{2\pi} \vec{n}_y \cdot \vec{\nabla}_y \ln |x-y| \quad \frac{(x-y)}{|x-y|^2}$$

2)  $\int_{\partial D}$  param by  $y(t)$ ,  $t \in (0, 2\pi)$ :

for some bdry func  $g$

$$\int_{\partial D} g(y) ds_y = \int_0^{2\pi} g(y(t)) \underbrace{|y'(t)|}_{\text{speed}} dt.$$

