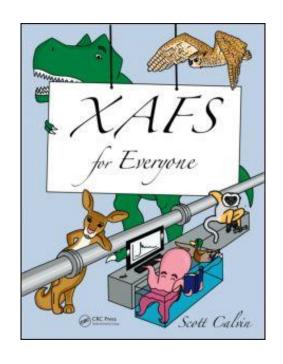
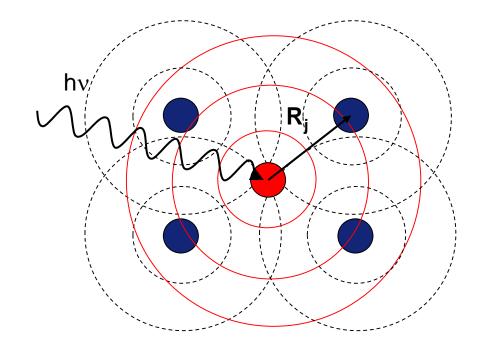


ESRF | The European Synchrotron

XAFS FUNDAMENTALS, DATA REDUCTION AND ANALYSIS

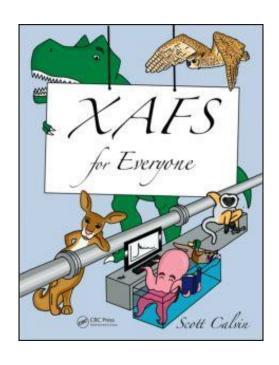
Vera Cuartero BM23 (ESRF) cuartero@esrf.fr

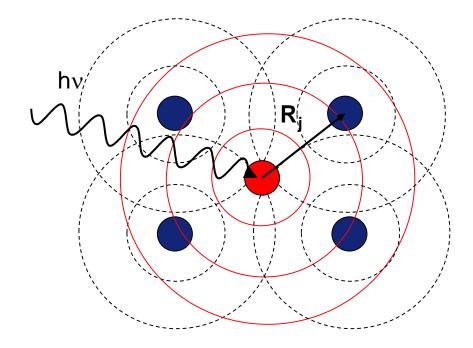




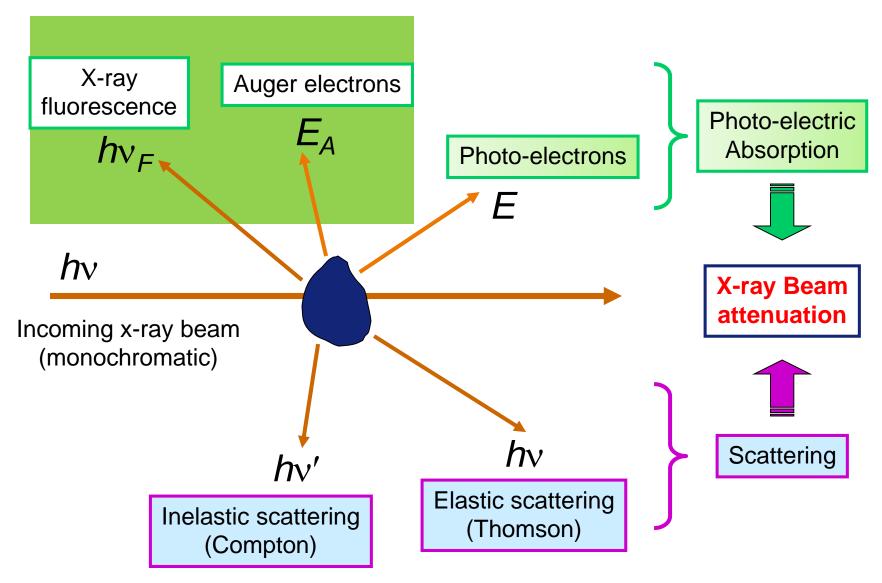
XAFS FUNDAMENTALS, DATA REDUCTION AND ANALYSIS

- 1. Fundamentals of X-ray Absorption Fine Structure (XAFS)
- 2. Measurement of Cu foil at BM23
- 3. Data analysis Cu foil
- 4. Real sample preparation, CuO
- 5. Data analysis, comparison



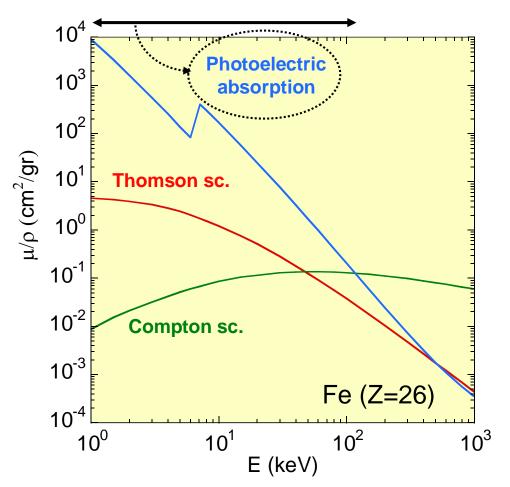


INTERACTION OF X-RAYS WITH MATTER (E< 1 MeV)



ATOMIC CROSS SECTION (E< 1 MeV)

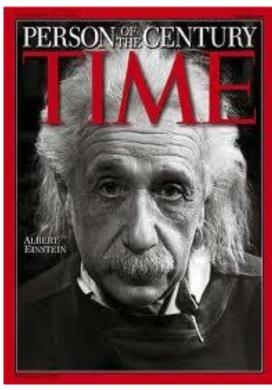
- Quantifies the strength of the X-ray matter interaction
- For photon energies between 1 and 30 keV the photoelectric absorption is the leading interaction with matter



$$\frac{\mu/\rho = (N_A/A)\sigma_a}{\sigma_a = \sigma_{pe} + \sigma_{coh} + \sigma_{incoh} + \sigma_{pp}}$$

X-RAY ABSORPTION PROCESS: PHOTOELECTRIC EFFECT

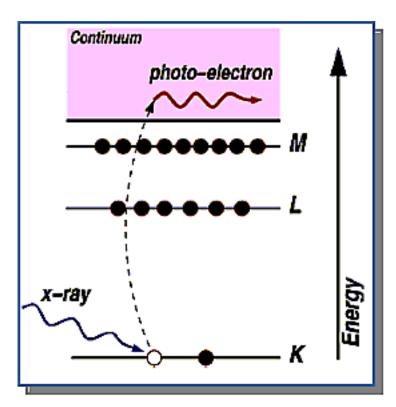
Hard X-rays (light with energy $1 \le E \le 100$ keV or wavelength $0.12 \le \lambda \le 12$ Å) are absorbed by all matter through the photoelectric-effect:



ay has an energy e binding energy of (**K=1**s; **L**_I=**2**s; L_{II}, photon is absorbed s ejected from the

xcited state with an core hole) and any en to the ejected

Nobel Prize 1921



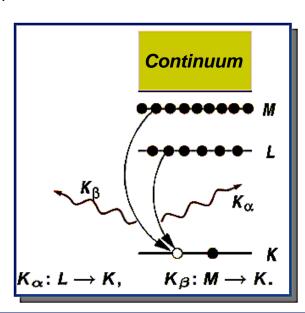
The excited core-hole will relax back to the "ground state" of the atom by two main mechanisms: X-ray Fluorescence, for hard x-ray regime (> 2 keV) and Auger Effect for lower energy x-ray absorption.

DE-EXCITACTION PROCESSES

When x-rays are absorbed by the photoelectric effect, the excited core-hole will relax back to a "ground state" of the atom. A higher level core electron drops into the core hole, and a *fluorescent x-ray* or *Auger electron* is emitted.

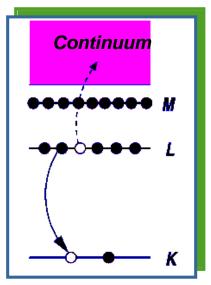
X-ray Fluorescence:

A higher energy core-level electron fills the deeper core-hole, ejecting an x-ray of energy equal to the difference of the core-levels.



Auger Effect:

A higher energy core-level electron drops into the core-hole and a second electron is emitted into the continuum from another corelevel.



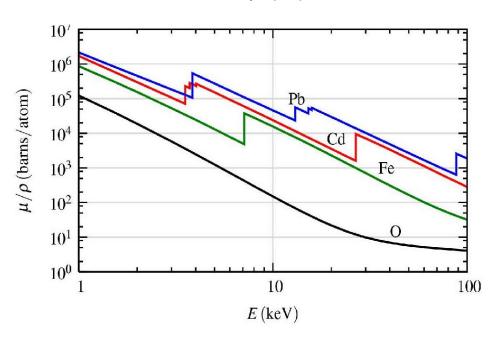
X-ray fluorescence and Auger emission occur at discrete energies that are characteristic of the absorbing atom.

THE ABSORPTION COEFFICIENT μ

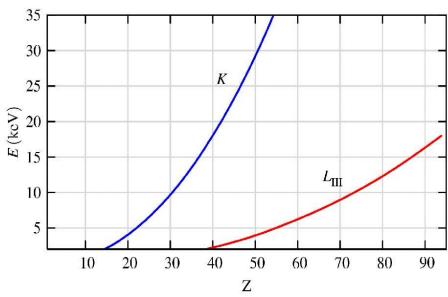
The probability that an x-ray beam passing through a material of thickness t will be absorbed is given by the *linear absorption coefficient*, μ:

$$\mupproxrac{
ho {f Z^4}}{{f A}{f E}^3}$$

Good contrast $\mu(E) \propto Z^4 / E^3$

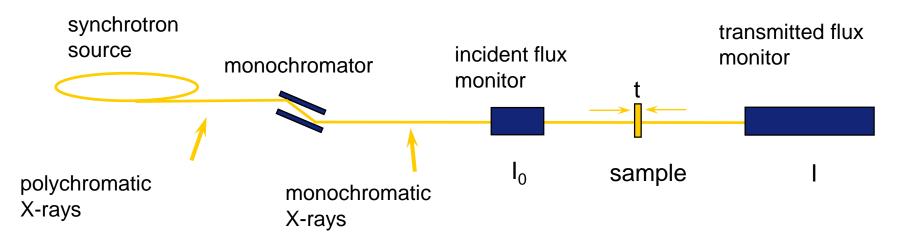


Element specific $E_K \propto Z^2$



X-RAY ABSORPTION MEASUREMENT

XAFS measures the energy dependence of $\mu(E)$ near and above the absorption edge of a selected element.



Transmission:

The absorption is measured directly by measuring what is transmitted through the sample:

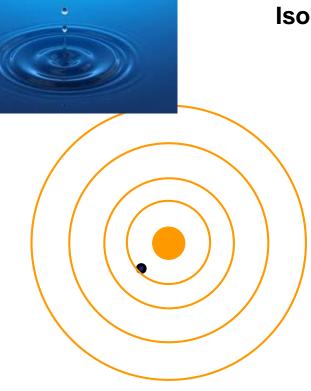
$$I = I_0 e^{-\mu (E)t}$$

$$\mu(E) t = -\ln (I/I_0)$$
 Beer-Lambert Law

Fluorescence:

The re-filling the deep core hole is detected. Typically the fluorescent x-ray is measured.

XAFS: SIMPLE PHYSICAL DESCRIPTION



Isolated "bare" atom

 X-ray photon with enough energy ejects one core (photo)electron (photoelectric effect)

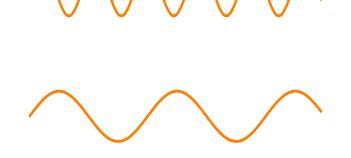
$$E_k=E_{h\nu}-E_0=rac{p^2}{2m}=rac{(\hbar k)^2}{2m}$$
 Kinetic energy of the p.e.

$$k = \sqrt{[2m(E_{h\nu} - E_0)/\hbar^2]}$$
 wavevector of the p.e.

$$\lambda = 2\pi/k$$
 wavelength of the p.e.

 The photoelectron can be described by a wave function approximated by a spherical wave

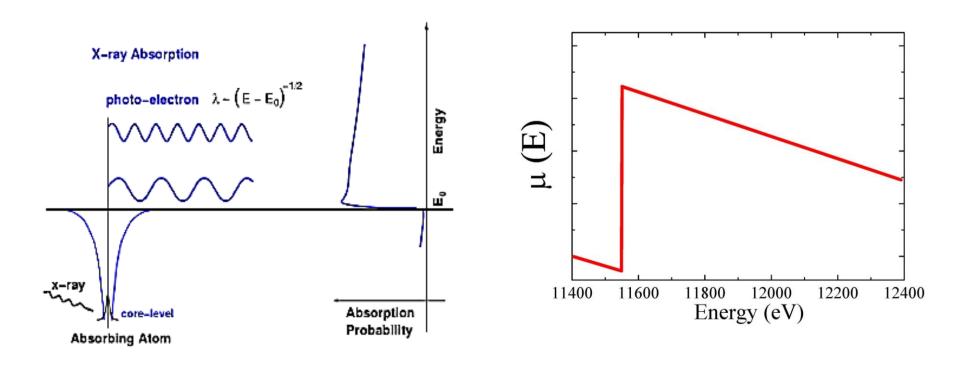






XAFS: SIMPLE PHYSICAL DESCRIPTION

Isolated "bare" atom

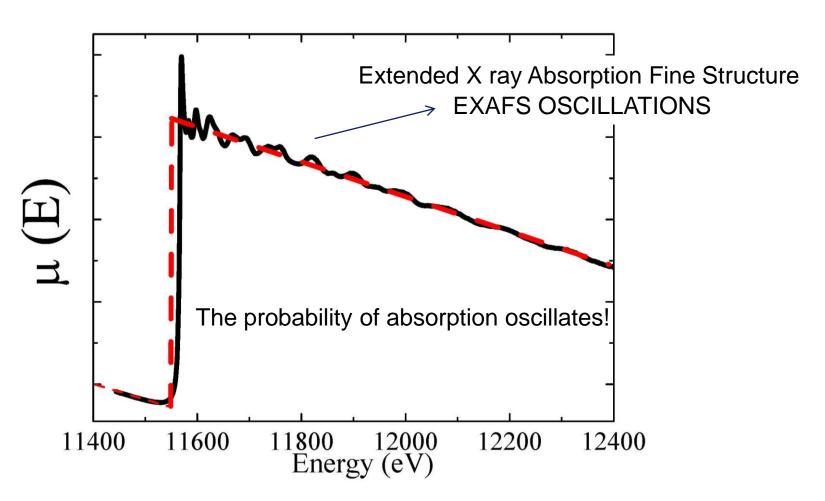


The atom absorbs an x-ray of energy E_0 destroying a core-electron of energy E_0 and creating a photo-electron with energy $(E-E_0)$.

 μ (E) has a sharp step at the core-level binding energy E_0 and is a smooth function of energy above this absorption edge.

XAFS: SIMPLE PHYSICAL DESCRIPTION

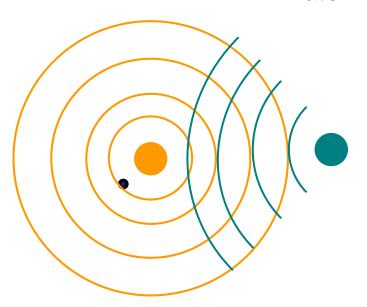
Condensed Matter



Proximity of neighboring atoms strongly modulates the absorption coefficient

EXAFS QUALITATIVELY: CONDENSED MATTER

Diatomic molecule



- The photoelectron can scatter from a neighbouring atom giving rise to an incoming spherical wave coming back to the absorbing atom
- The outgoing and ingoing waves may interfere.



The ejected photo-electron propagates, can scatter from a neighbouring atom and returns to the absorbing atom.

Interference between out-going and back-scattered photo-electron waves at the core site give rises to **modulation** of the absorption coefficient.

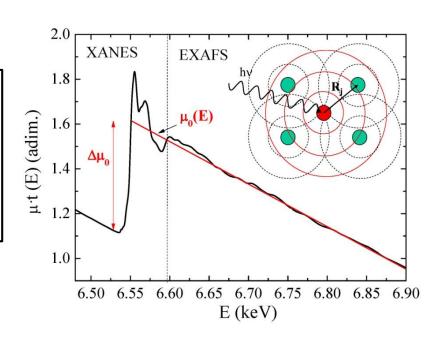
EXAFS DEFINITION

What? Oscillatory behaviour of the of the X-ray absorption as a function of photon energy beyond an absorption edge

When? Non isolated atoms

Why? Proximity of neighboring atoms strongly modulates the absorption coefficient

The EXAFS oscillations are an interference effect of the photo-electron with itself, due to the presence of neighboring atoms.



EXAFS oscillations:

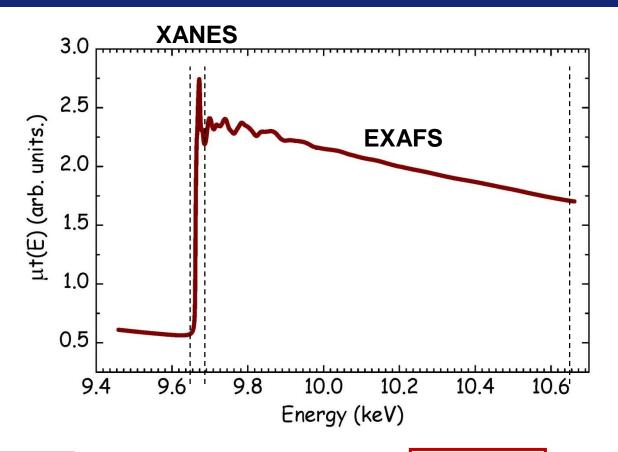
$$\chi(E) = \frac{\mu(E) - \mu_0(E)}{\Delta \mu_0(E_0)} - \frac{\mu(E) - \mu_0(E)}{\Delta \mu_0(E_0)}$$

 $\mu_0(\mathsf{E})$ smooth "bare atom" background

 $\Delta\mu_0(\mathsf{E}_0)$ edge step or jump



XAFS NOMENCLATURE: XANES AND EXAFS



X-ray
Absorption
Near
Edge
Structure

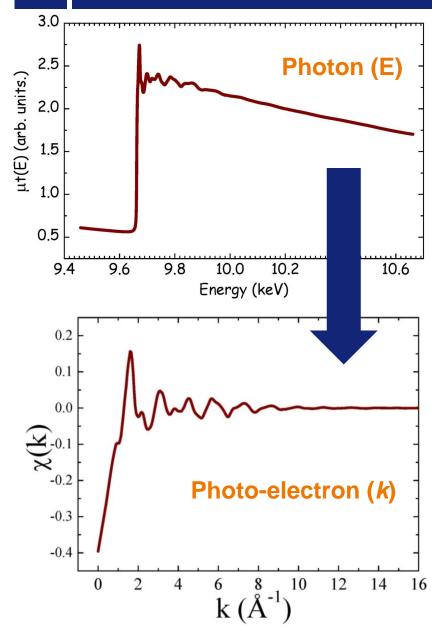
up to ~ 60 eV above the edge

Extended X-ray Absorption Fine Structure

from ~ 80 eV to 1200 eV above the edge



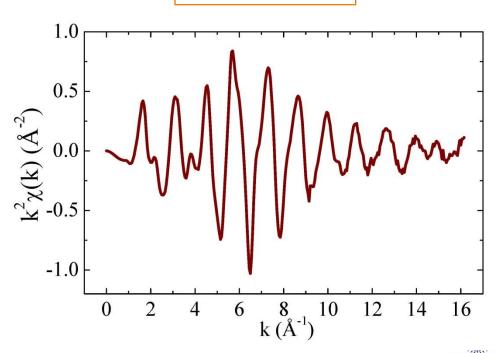
EXAFS OSCILLATIONS, $\chi(k)$



EXAFS depends on the wave nature of the photo-electron

$$\sqrt{\frac{2m(E - E_0)}{\hbar^2}} = k = \frac{2\pi}{\lambda}$$

 k^2 or k^3 weighted Oscillations at high-k

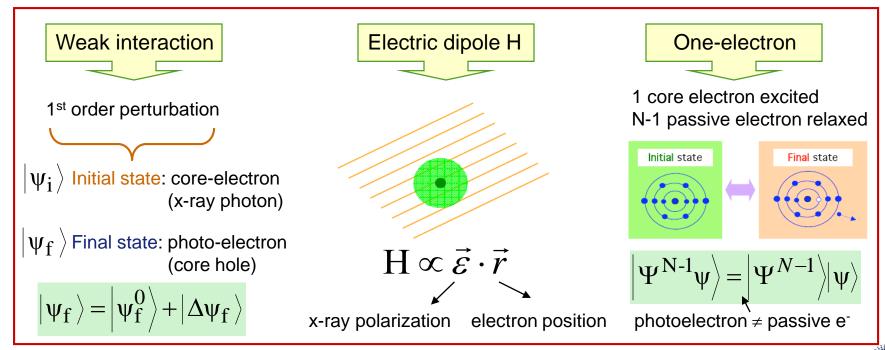


EXAFS: PHYSICAL DESCRIPTION

Fermi's Golden Rule describes $\mu(\omega)$ as a transition between two quantum states :

$$\mu(\omega) \propto \sum_{\rm f} W_{\rm if} \underbrace{\delta(E_{\rm i} - E_{\rm f} + \hbar \omega)}_{\rm Energy \ conservation}$$
 Transition probability $W_{\rm if} \propto \left| \left\langle \psi_{\rm i} \left| H \right| \psi_{\rm f} \right\rangle \right|^2$

Approximations!



EXAFS: PHYSICAL DESCRIPTION

$$\begin{aligned} W_{if} & \propto \left| \left\langle \psi_i \left| H \right| \psi_f \right\rangle \right|^2 & \text{Fermi Golden Rule} \\ & \text{1 active electron} \\ \mu(\omega) & \propto \sum_{f} \left| \left\langle \psi_i \left| \vec{\mathcal{E}} \cdot \vec{r} \right| \psi_f \right\rangle \right|^2 \delta(E_i - E_f + \hbar \omega) \left| \left\langle \Psi_i^{N-1} \left| \left| \Psi_f^{N-1} \right| \right\rangle \right|^2 \\ \text{Dipole interaction} & \psi_f \right\rangle & = \left| \psi_f^0 \right\rangle + \left| \Delta \psi_f \right\rangle \\ & S_0^2 \approx 0.6 - 1 & \text{N-1 electrons relaxation} \end{aligned}$$

$$\chi(k) = \frac{\mu - \mu_0}{\mu_0} = \frac{2\operatorname{Re} \int d\vec{r} (\psi_i \vec{\varepsilon} \cdot \vec{r} \psi_f^{0^*})(\psi_i^* \vec{\varepsilon} \cdot \vec{r} \Delta \psi_f)}{\int d\vec{r} |\psi_i^* \vec{\varepsilon} \cdot \vec{r} \psi_f^{0}|^2} \propto 2\operatorname{Re}(\psi_{\text{scatt}}(0))$$

Wave-function of the scattered photo-electron at the absorbing atom

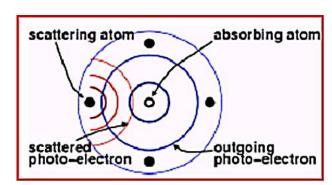
The region where $\psi_i \neq 0$ represents simultaneously the source and the detector for the photoelectron that probes the local structure around the absorber atom

EXAFS: PHYSICAL DESCRIPTION

We can build a simple model for χ from the photo-electron for a two-atomic system:

- 1. Leaving from the absorbing atom
- 2. Scattering from the neighbour atom
- 3. Returning to the absorbing atom

The outgoing photo-electron wave-function →



a neighbour atom at a distance R



 $\psi(k,r) = \frac{e^{ikr}}{kr} + \frac{1}{kr}$

 $\chi(k) \propto 2 \operatorname{Re}(\psi(k, r = 0)) = 2 \operatorname{Re}(i \frac{e^{ikR}}{2kR} [f(k, \pi)e^{i\delta(k)}] \frac{e^{ikR}}{R})$

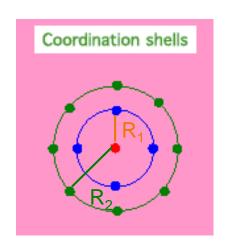
 $\chi(k) = \frac{f(k, \pi)}{kR^2} sin(2kR + \delta(k))$ Spherical wave attenuation

Back-scattering amplitude phase-shift \rightarrow

Back & forth path

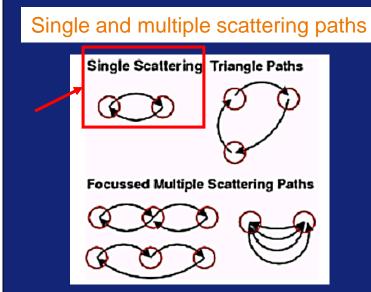
"Neighbour atom"

THE EXAFS EQUATION: MANY ATOMIC SYSTEM

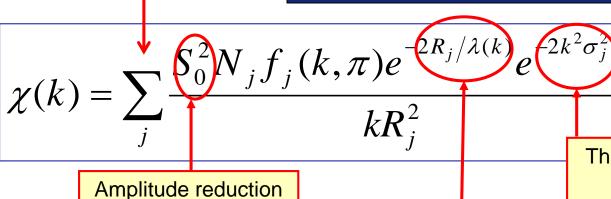


term

Intrinsic losses on the EXAFS spectrum



Multiple-scattering paths are most important when the scattering angle is > 150°



 $\sin(2kR_i + \delta_i(k))$

Thermal and static mean-square disorder in R

Photo-electron mean-free path (including core-hole lifetime)

THE EXAFS EQUATION PARAMETERS

$$\chi(k) = \sum_{j} \frac{S_0^2 N_j f_j(k, \pi) e^{-2R_j/\lambda(k)} e^{-2k^2 \sigma_j^2}}{kR_j^2} \sin(2kR_j + \delta_j(k))$$

Parameters we need to calculate

$$f(k)$$
 scattering amplitude $\delta_j(k)$ phase-shif

Parameters to fit

RDistance to the neighbouring atom

Coordination number of the neighbouring atom

Disorder in the neighbour distance

THE EXAFS EQUATION PARAMETERS

$$\chi(k) = \sum_{j} \frac{S_0^2 N_j f_j(k, \pi) e^{-2R_j/\lambda(k)} e^{-2k^2 \sigma_j^2}}{kR_j^2} \sin(2kR_j + \delta_j(k))$$

R Distance to the neighbouring atom.

Or **half path length**, is the average distance between the absorbing and scattering atoms (single scattering path).

 ${\cal N}$ Coordination number of the neighbouring atom.

Or **degeneracy** of the path, is the identical number of distinct ways, per absorbing atom, that the scattering defined by the path can take place.

 $\sigma^{ extstyle 2}$ Disorder in the neighbour distance.

Or **Mean Square Relative Displacement** is the variance in the half path length, that is, the square of the standard deviation of the half path length: $\sigma^2 = <(r-r_m)^2>$.

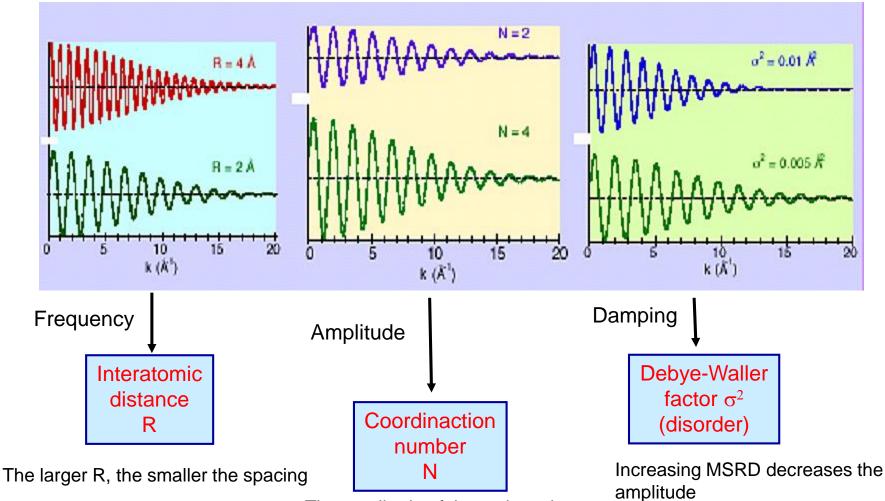
Static and/or thermal disorder.



EXAFS takes place on a time scale much shorter than that of atomic motion, so the measurement serves as an instantaneous snapshot of the atomic configuration

EXAFS: A STRUCTURAL PROBE

If we know the scattering properties of the neighbouring atoms: f(k) and $\delta(k)$, and the mean-free path $\lambda(k)$, we can determine:

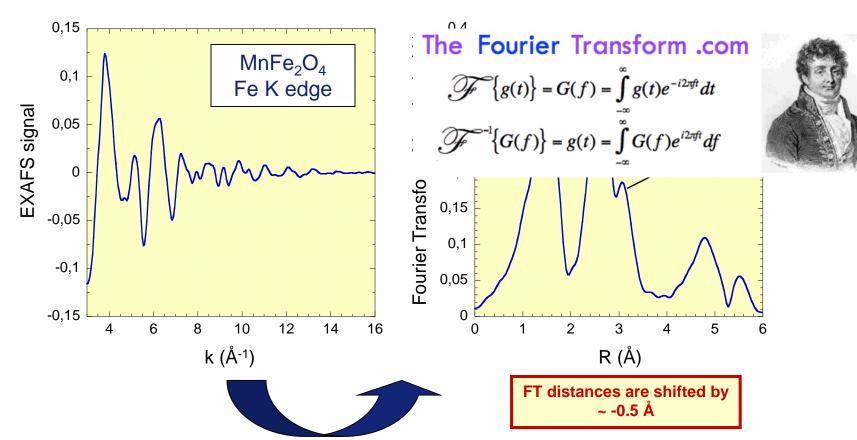


The amplitude of the path scales with the degeneracy.



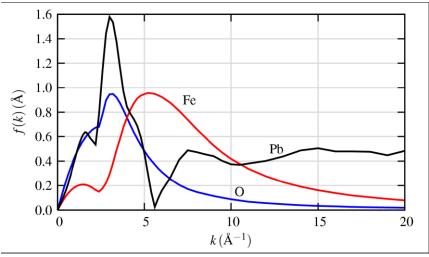
EXAFS: A STRUCTURAL PROBE

- Multiple terms in the EXAFS equation may be significant, each of them being a modulated sinusoid → need to decompose a function into constituent sine waves
- Fourier transformation can be used to descompose a frequency-space signal into its different constituent frequencies. The Fourier transform of an EXAFS spectrum gives a pseudo-radial distribution function.

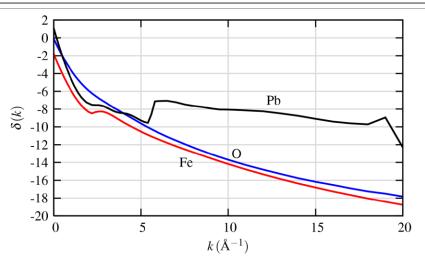


EXAFS: SELECTIVITY OF ATOMIC SPECIES

The scattering amplitude f(k) and phase shift $\delta(k)$ depend on the atomic number Z



The scattering amplitude f (k) peaks at different k values and extends to higher-k for heavier elements. For very heavy elements, there is structure in f(k).



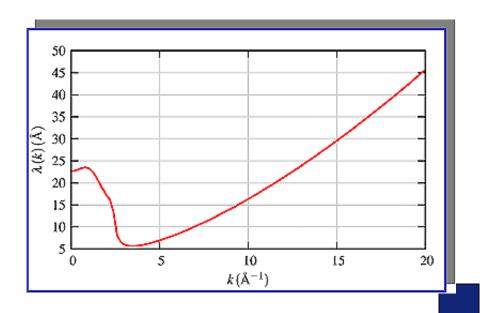
The phase shift $\delta(k)$ shows sharp changes for very heavy elements.

These scattering functions can be accurately calculated (say with the programs FEFF, GNXAS, etc.), and used in the EXAFS modeling.

EXAFS: A SHORT RANGE ORDER PROBE

The photo-electron mean-free path λ (k) limites how far the photo-electron can go:

- Inelastic scattering
- Core-hole finite life-time



The mean-free path λ depends on k

 λ < 25 Å for the EXAFS k-range

The term
$$\frac{e^{-2K_j/\lambda(k)}}{R_j^2}$$
 makes EXAFS a local atomic probe, typically within 10 Å from the absorber

EXAFS ACCURACY

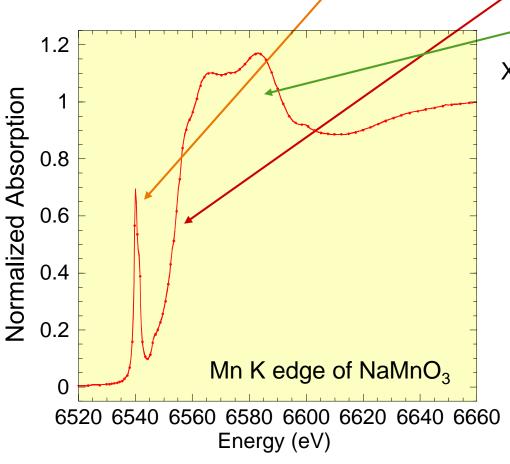
- Distances ±0.01Å The absorber–scatter distance can be measured quite accurately from the frequency of the sinusoidal oscillations.
- <u>Coordination numbers</u>
 <u>Debye-Waller factors</u> ± 20-25% High correlation between them ↔ less accurate
- <u>Scattering Atom</u> Z ± 1 (Z=6-17)
 Z ± 3 (Z=20-35)

Elements that are next to each other in the periodic table has barely distinguishable photoelectron scattering characteristics. Thus, C,N and O are impossible to distinguish and metal scattered can only be placed in the proper row of the periodic table

XANES REGION

What is XANES?

XANES = Pre-edge + Edge + XANES



XANES is extremely sensitive to the chemistry of the absorbing atom:

- > Formal oxidation state
- Coordination environment

XANES probes the angular momentum of the unoccupied electronic states.

$$\Delta \ell = \pm 1$$

$$\Delta s = 0$$

$$\Delta j = \pm 1, 0$$

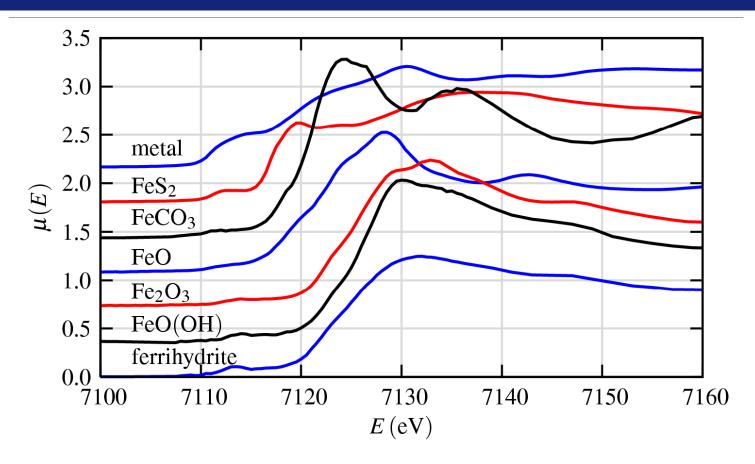
$$\Delta m = 0$$



WHY ARE WE INTERESTED IN XANES REGION?

Region	Transition	Information Content
Pre-edge	Electronic transitions to empty bound states (transition probability controlled by dipolar selection rules). e.g. 1s→3d, 1s→4p (1st transition series metals)	Local coordination environment around the absorbing atom. Dependence on the oxidation state and bonding.
Edge	Defines ionization threshold to continuum states.	Dependence on the oxidation state "Chemical shift": main edge (binding energy) shifts to higher energy with increasing oxidation state.
XANES	Feautures dominated by multiple-scattering resonances of the photoelectrons ejected at low E _{kinetic} .	Atomic positions of neighbours: Interatomic distances and bond angles.

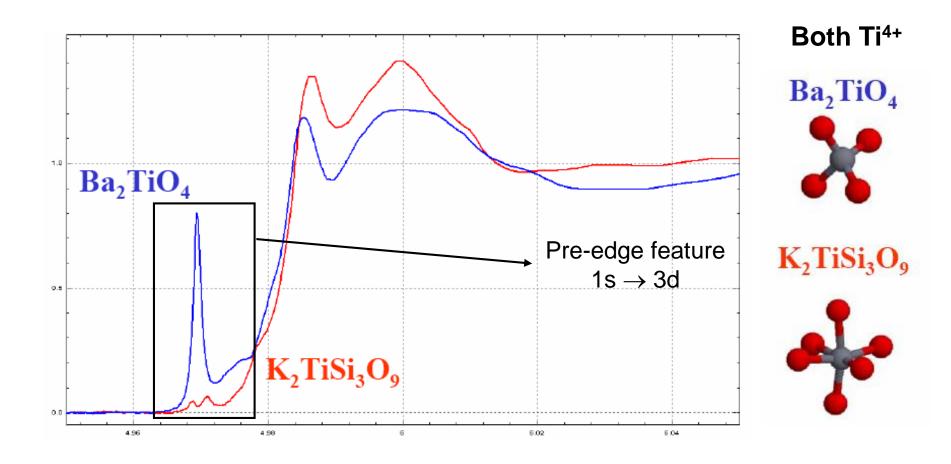
EDGE POSITION: OXIDATION STATE



XANES can be used simply as a fingerprint of phases and oxidation state.

XANES Analysis can be as simple as making linear combinations of "known" spectra to get compositional fraction of these components.

XANES ANALYSIS: LOCAL COORDINATION ENVIROMENT



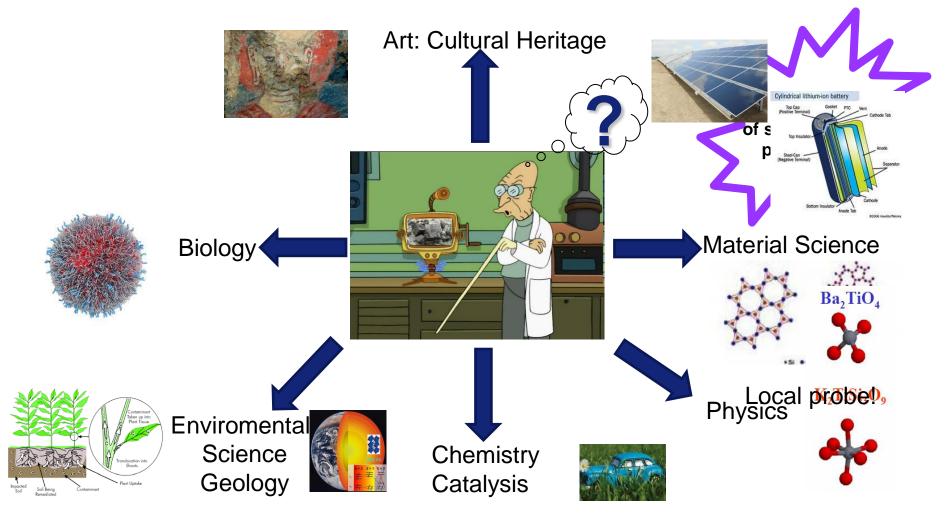
Ti K-edge XANES is highly dependent on the local coordination: (1) Tetrahedral vs. Octahedral

(2) Different neighbors shells beyond the first O-shell



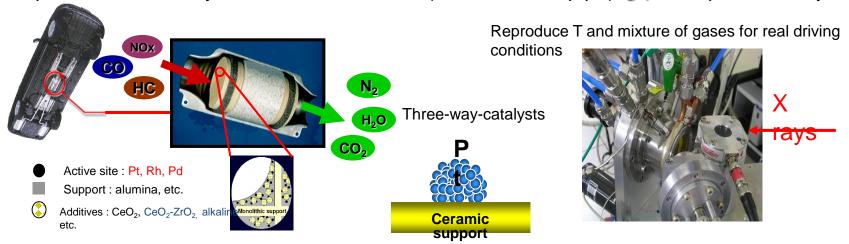
WHAT CAN XAS BE MEASURED ON?

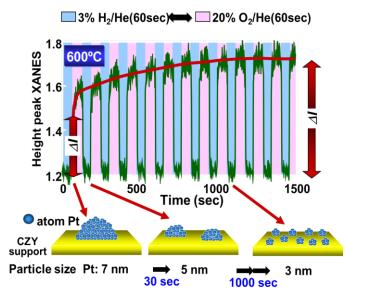
... just about almost anything and with most elements of the periodic table



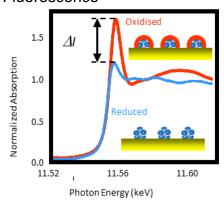
APPLICATIONS: CATALYSIS

Reproduce real controls: 80a % i omorbol cistiche noiso al manauta pt peing percepsos v catalytic





Dynamic and *in situ* study via X ray Absorption in Fluorescence



Experiment performed by TOYOTA.



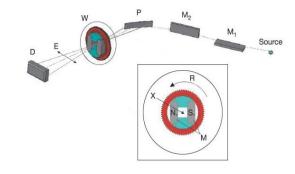
APPLICATIONS: SOLID STATE PHYSICS

Structural changes due to phase transitions

- Temperature
- Pressure
- Magnetic Fields,...

Measurement of femtometre-scale atomic displacements by X-ray absorption spectroscopy

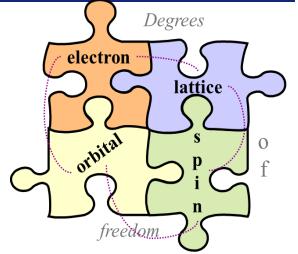
Robert F. Pettifer¹, Olivier Mathon², Sakura Pascarelli², Michael D. Cooke³* & Michael R. J. Gibbs³

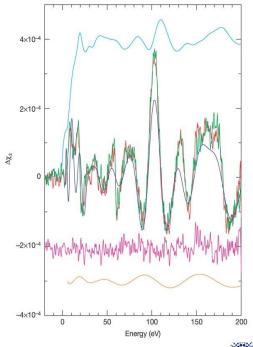


NATURE | VOL 435 | 5 MAY 2005 | www.nature.com/nature

Fe-Co thin films. Periodic strain through magnetostriction.

- Magnetostriction can be detected with EXAFS (diff)!
- Normal exafs accuracy 0.01- 0.001 Å
- Q-exafs → Diff-exafs sensitivity of 1 fm (0.0001 Å)





¹Department of Physics, University of Warwick, Coventry CV4 7AL, UK

²European Synchrotron Radiation Facility (ESRF), BP 220, 38043 Grenoble Cedex 9, France

³Department of Engineering Materials, Sir Robert Hadfield Building, Mappin Street, University of Sheffield, Sheffield S1 3JD, UK

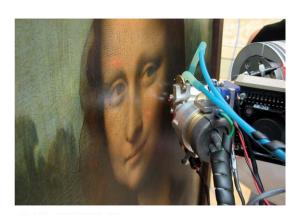
APPLICATIONS: CULTURAL HERITAGE

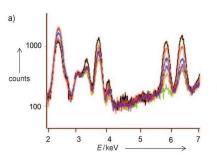
Analyzing Works of Art

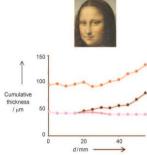
DOI: 10.1002/anie.201001116

Revealing the *sfumato* Technique of Leonardo da Vinci by X-Ray Fluorescence Spectroscopy**

Laurence de Viguerie, Philippe Walter,* Eric Laval, Bruno Mottin, and V. Armando Solé







- Non-invasive
- Binder Concentration
- Pigments Inorganic composition
- Layers Thickness
- Mn, Fe pigments on the glaze
- Pb in every layer
- Glaze layers unfinished

Anal. Chem. 2006, 78, 7484-7492

Blackening of Pompeian Cinnabar Paintings: X-ray Microspectroscopy Analysis

Marine Cotte,*,† Jean Susini,† Nicole Metrich,‡ Alessandra Moscato,[§] Corrado Gratziu,[§] Antonella Bertagnini,[∥] and Mario Pagano[⊥]



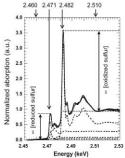


Figure 9. Computation on the "XANES spectra at the S K-edge fit of the experimental data (plain) by combination of cinnabar, sulfur, corderoite, and sypsum spectra (dotted) and estimation of the reduced and oxidized sulfur concentrations by measuring the fluorescence intensity at flow specific energies.

- Why do the red walls of Pompei go black?
- Reaction of chlorine with cinnabar leads to the formation of gray chlorinemercury compounds.
- The sulfation of calcite results in the development of black coatings on the painting surface.



Turning black. A wall painted red in the remains of Pompeii.



TAKE-HOME MESSAGE

Exafs is sensitive to ...



Density of electronic states



- Local structure in non-crystalline matter
- Local environment of an atomic impurity in a matrix of different atomic species
- Study of systems whose local properties differ from the average properties
- Detection of very small distortions of local structure

WHERE TO GO FROM HERE

Books and Review Articles:

Basic Principles and Applications of EXAFS

Handbook of Synchrotron Radiation, Chapter 10, pp. 995-1014, E. A. Stern and S. M. Heald,

E. E. Koch, ed., North-Holland 1983

X-ray Absorption: Principles, Applications, Techniques of EXAFS, SEXAFS and XANES

in Chemical Analysis Vol. 92, D. C. Koningsberger and R. Prins, ed., John Wiley & sons, 1988

X-ray Absorption Fine Structure for Catalysis and Surfaces

World Scientific Series on Synchrotron Radiation Techniques and Applications Vol. 2, ed., Y. Iwasawa, 1996

X-ray absorption spectroscopy in coordination chemistry

J. E. Penner-Hahn, Coordination Chemistry Reviews 190-192, pp. 1101-1123, 1999

Tutorials and other Training Material:

http://xafs.org/Tutorials

http://gbxafs.iit.edu/training/tutorials.html Grant Bunker's tutorials

http://srs.dl.ac.uk/XRS/courses/ tutorial from Daresbury Lab, UK

Software resources:

http://xafs.org/Software

Analysis programs: Atoms + Feff + iFeffit (Theoretical XAFS calculations and fitting) – Univ. Chicago (USA)

Athena, Artemis (EXAFS Data Reduction and fitting

http://gnxas.unicam.it : Gnxas (Theoretical XAFS calculations and fitting) – Univ. Camerino (Italy)

http://www.esrf.fr/computing/scientific/xop/ XOP (DABAX data base, cross-section and lots of general x-ray

Calculations, multipurpose data visualization and analysis)-M. Sanchez del Rio, and R. J. Dejus (ESRF, France)

http://srs.dl.ac.uk/XRS/index.html : Excurv98 (Theoretical XAFS calculations and fitting) – Daresbury Lab (UK)



DATA PROCESSING AND ANALYSIS

Iffefit package:

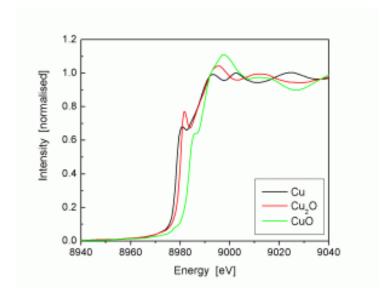
http://cars9.uchicago.edu/ifeffit/

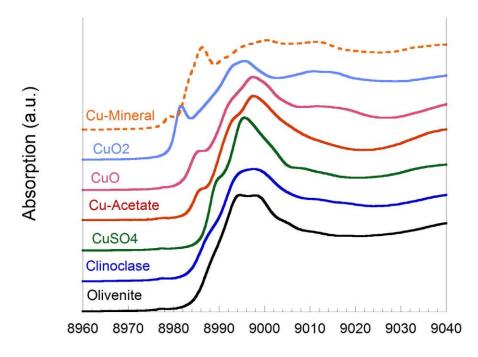
http://bruceravel.github.io/demeter/

Athena: data processing

Artemis: data analysis using Feff

CU FOIL AND CUO ANALYSIS: XANES



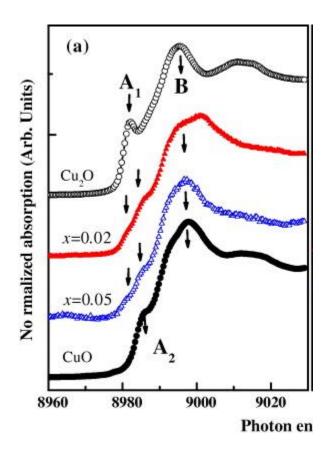


Cu foil Cu⁰, 3d¹⁰ 4s¹ Cu₂O Cu¹⁺, 3d⁹ 4s¹ Cu Tetrah. CuO Cu²⁺, 3d⁸ 4s¹ Cu square planar

CU METAL AND CU OXIDE XAFS

Cu foil Cu⁰, 3d¹⁰ 4s¹ Cu₂O Cu¹⁺, 3d⁹ 4s¹ Cu Tetrah. CuO Cu²⁺, 3d⁸ 4s¹ Cu square planar

The pre-edge features marked A_1 and A_2 , respectively, in the Cu_2O and CuO spectra represent transitions from the 1s \rightarrow axial 4p state, which contain a small portion of the unoccupied 3d states as the quadruple allowed transitions and are close to the 4s state. The feature marked as B is mainly due to the transition from copper 1s to the 4p state



Cu FOIL ANALYSIS

- ✓ Data processed, with Athena.
- ✓ Crystal structure (or reference crystal structure) known.
 - Cu metal, FCC, Fm-3m
 - CuO, monoclinic, C2/c
 - Inorganic Crystal Structure data base https://icsd.fiz-

karlsruhe.de/search/basic.xhtml;jsessionid=5DC14101FB0549F273F08 BC77D85DF9D

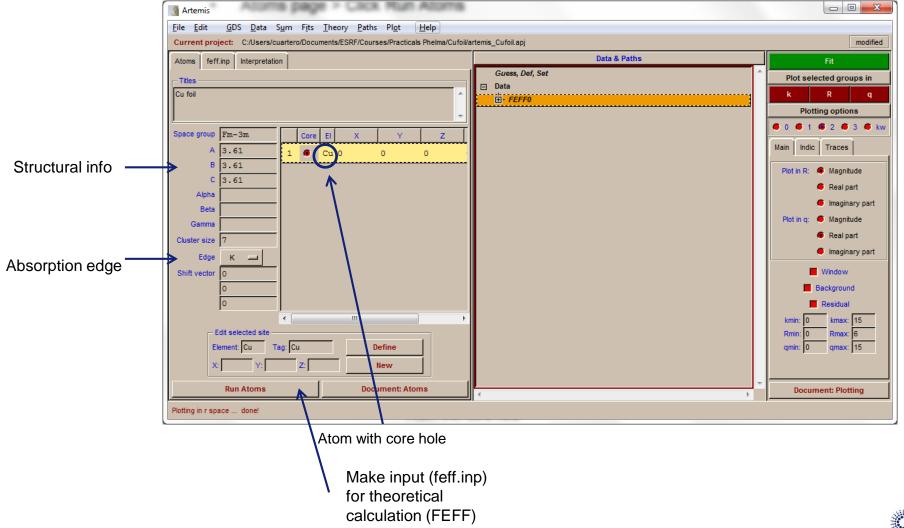
- 1. Create the cluster, with ATOMS.
- 2. Create the scattering paths, with FEFF (integrated in Atoms)
- 3. 1st shell analysis, with ARTEMIS.



Cu FOIL ANALYSIS: ATOMS

1. Create the cluster

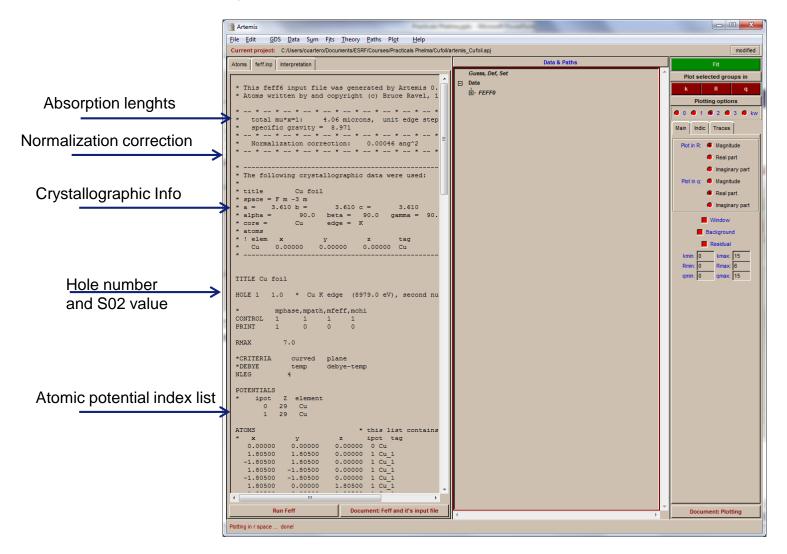
Atoms page > Click Run Atoms



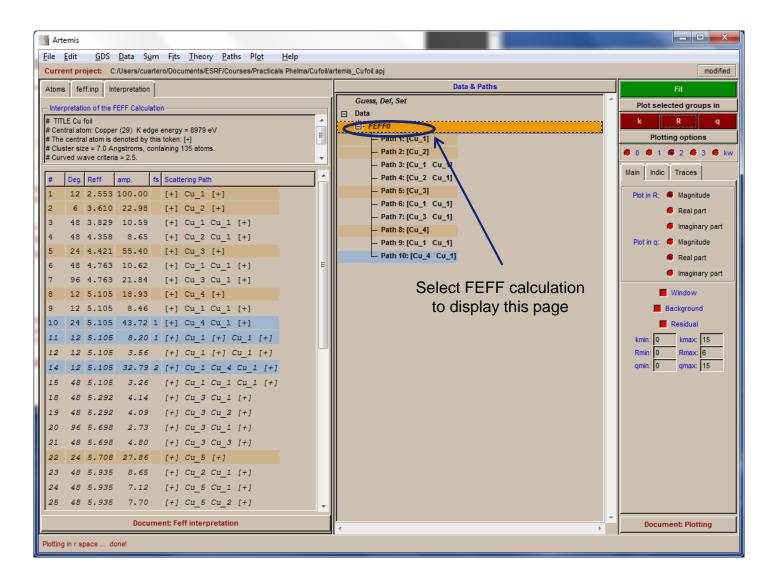
Cu FOIL ANALYSIS: ATOMS

1. Create the cluster

Theory input page feff.inp > Click Run Feff

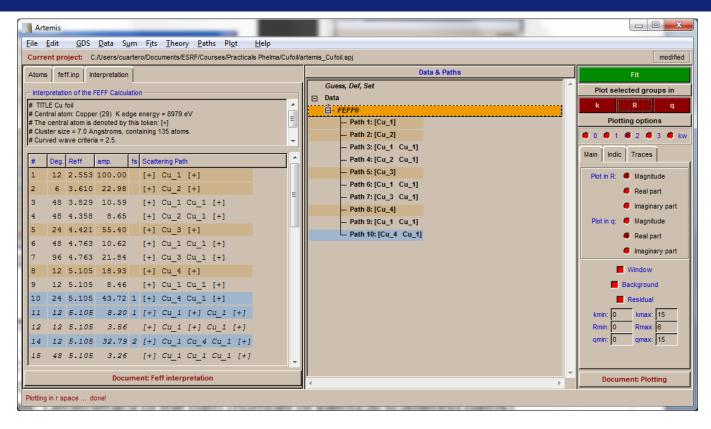


Cu FOIL ANALYSIS: ARTEMIS





Cu FOIL ANALYSIS: ARTEMIS

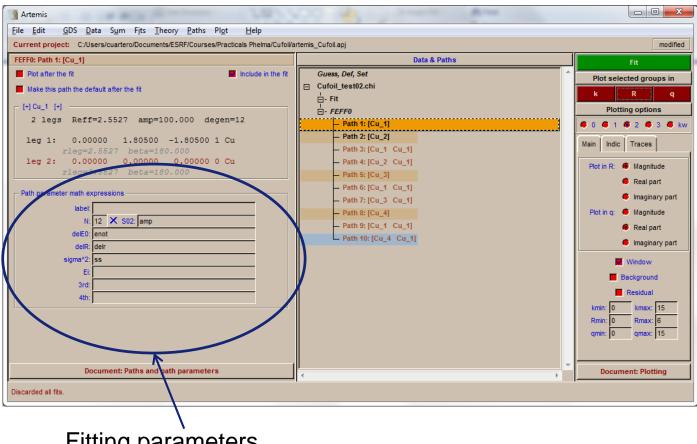


Feff paths

- Deg: Degeneracy of the path (number of identical scattering paths)
- Reff: Inital half path length (bond length for single scattering path)
- amp: Estimate of amplitude of path relative to first path.
- fs: Number of forward scattering events.
- Scattering Path: atoms scattering photoelectron, [+] symbol represents core atom.

Cu FOIL ANALYSIS: ARTEMIS

Path description



Fitting parameters

CUO XTAL STRUCTURE

