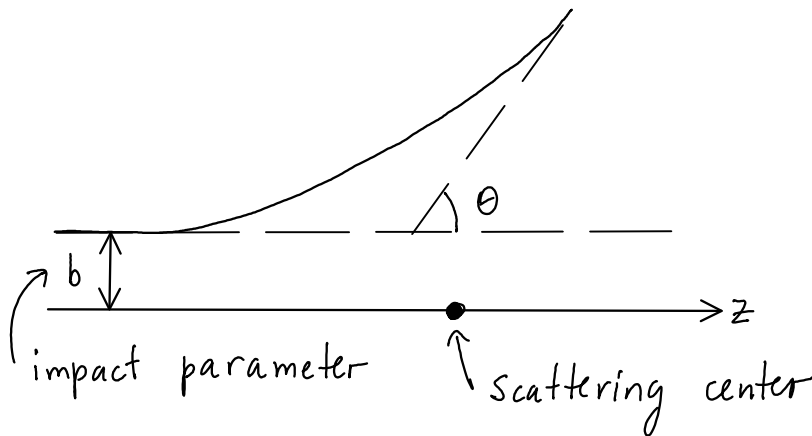
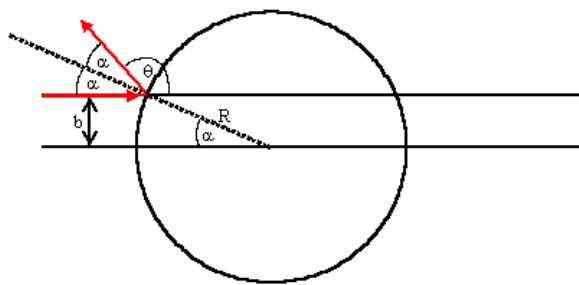


**Scattering**

Classical scattering theory

Problem: given the scattering parameter  $b$ , calculate the scattering angle  $\theta$

**Example: elastic hard-sphere scattering**

Our target is a billiard ball of radius  $R$ , the incident particle is a ball that bounces elastically. The impact parameter is

$$b = R \sin \alpha$$



The scattering angle is

$$\theta = \pi - 2\alpha$$

$$\alpha = \frac{1}{2}(\pi - \theta)$$

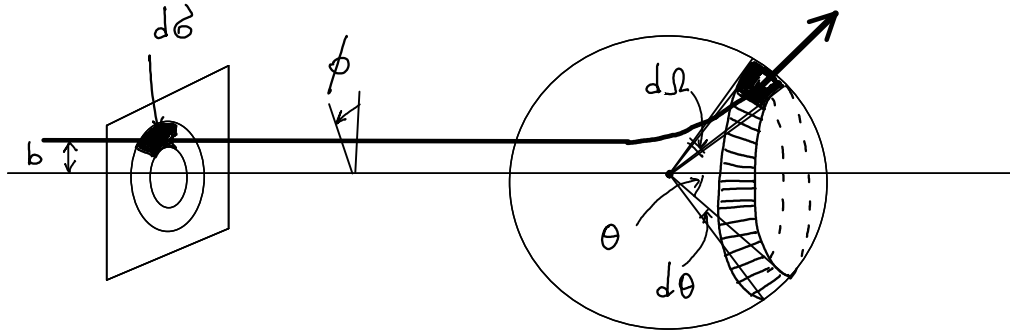
$$b = R \sin\left(\frac{\pi}{2} - \frac{\theta}{2}\right) = R \cos \frac{\theta}{2}$$



Therefore, the scattering angle is:

$$\theta = \begin{cases} 2 \cos^{-1}(b/R) & \text{if } b \leq R \\ 0 & \text{if } b \geq R \end{cases}$$

## Scattering



More general problem: particles incident within an infinitesimal patch of cross-sectional area  $d\sigma$  are scattering into an infinitesimal solid angle  $d\Omega$ .  
The quantity

$$D(\theta) = \frac{d\sigma}{d\Omega}$$

is called the differential (scattering) cross-section.

$$d\sigma = b \, db \, d\phi$$

$$d\Omega = \sin \theta \, d\theta \, d\phi$$

$$D(\theta) = \frac{b}{\sin \theta} \left| \frac{db}{d\theta} \right|$$

Total cross section is defined as the integral of  $D(\theta)$  over all solid angles:

$$\sigma \equiv \int D(\theta) d\Omega$$

**Class exercise: Find the differential and total cross-sections for hard-sphere scattering.**

$$b = R \cos \frac{\theta}{2} \qquad \frac{db}{d\theta} = -\frac{1}{2} R \sin \frac{\theta}{2}$$

$$D(\theta) = \frac{b}{\sin \theta} \left| \frac{db}{d\theta} \right| = \frac{R \cos \frac{\theta}{2}}{\sin \theta} \cdot \frac{1}{2} R \sin \frac{\theta}{2} = \underline{\underline{\frac{R^2}{4}}}$$

$$\sigma = (R^2/4) \int d\Omega = \pi R^2$$

If we have a beam of incident particles, with uniform intensity (luminosity)

$$\mathcal{L} \equiv \text{number of incident particles per} \\ \text{unit area per unit time} \Rightarrow$$

The number of particles entering  $d\sigma$  per unit time is

$$dN = \mathcal{L} d\sigma = \mathcal{L} D(\theta) d\Omega \Rightarrow$$

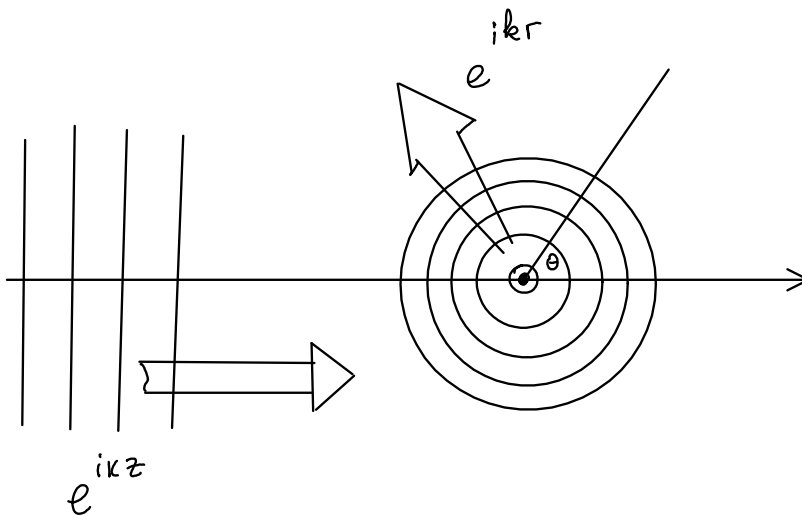
$$D(\theta) = \frac{1}{\mathcal{L}} \frac{dN}{d\Omega}$$

## Quantum scattering theory

Our problem: incident plane wave

$$\psi(z) = A e^{ikz}$$

traveling in Z direction encounters a scattering potential that produces outgoing spherical wave:



Therefore, the solutions of the Schrödinger equation have the general form:

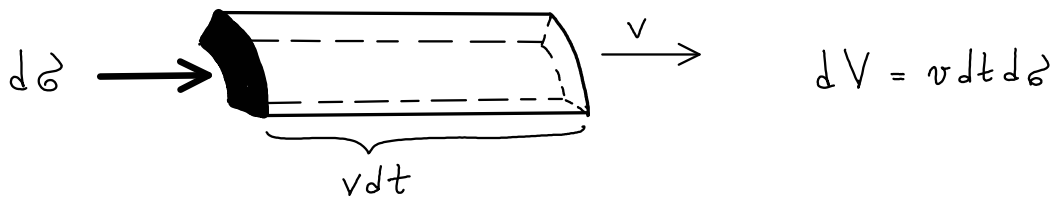
$$\psi(r, \theta) \approx A \left\{ \underset{\substack{\uparrow \\ \text{plane} \\ \text{wave}}}{e^{ikz}} + f(\theta) \underset{\substack{\uparrow \\ \text{spherical} \\ \text{wave}}}{\frac{e^{ikr}}{r}} \right\} \quad \text{for large } r$$

$$k = \frac{\sqrt{2mE}}{\hbar} \quad \leftarrow \begin{array}{l} \text{energy of the incident} \\ \text{particles} \end{array}$$

The quantity  $f(\theta)$  is called scattering amplitude.

It is the probability of scattering in a given direction  $\theta$ .

How is it connected to the differential cross-section?



Volume  $dV$  of incident beam (see above) passes through area  $d\sigma$  in time  $dt$ . The probability that the particle with speed  $v$  passes through this area  $d\sigma$  is

$$dP = |\psi_{\text{incident}}|^2 dV = |A|^2 v dt d\sigma$$

$\nwarrow = A e^{ikz}$

This must be equal to the probability that the particle scatters into the solid angle  $d\Omega$ :

$$dP = |\psi_{\text{scattered}}|^2 dV = \frac{|A|^2 |f|^2}{r^2} (v dt) r^2 d\Omega$$

$\nearrow A \frac{e^{ikz}}{r} f(\theta)$

Therefore  $|A|^2 v dt d\sigma = \frac{|A|^2 |f(\theta)|^2}{r^2} (v dt) r^2 d\Omega$

$$d\sigma = |f(\theta)|^2 d\Omega \Rightarrow$$

$$D(\theta) = \frac{d\sigma}{d\Omega} = |f(\theta)|^2$$

Therefore, to solve the scattering problem, we need to calculate the scattering amplitude  $f(\theta)$ .

## Partial wave analysis

Our potential is spherically symmetric  $\Rightarrow$

The solutions of the Schrödinger equation are

$$\psi(r, \theta, \phi) = R(r) Y_{\ell}^m(\theta, \phi)$$

$u(r) = rR(r)$   
satisfies

spherical harmonic

$$-\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + \left[ V(r) + \frac{\hbar^2}{2m} \frac{\ell(\ell+1)}{r^2} \right] u = Eu$$

(radial Schrödinger equation)

"centrifugal contribution"

Very large  $r$

$\longrightarrow V \rightarrow 0$

and "centrifugal contribution" is negligible

Radial equation becomes

$$-\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} \approx Eu$$

$$\frac{d^2 u}{dr^2} \approx -k^2 u$$

The general solution is

$$u = C e^{ikr} + D e^{-ikr}$$

outgoing spherical wave

incoming spherical wave  $\Rightarrow D = 0$

$$R = \frac{C e^{ikr}}{r}$$