

Computational Physics CheatSheet

Samyak Rai

November 09, 2025

1 Root Finding

Bisection Method

```
def bisection(f, a, b, tol=1e-8, maxiter=100):
    fa, fb = f(a), f(b)
    if fa == 0: return a
    if fb == 0: return b
    if fa * fb > 0:
        raise ValueError("f(a) and f(b) must have opposite signs")
    for _ in range(maxiter):
        c = 0.5*(a+b)
        fc = f(c)
        if abs(fc) < tol or abs(fc) < tol:
            return c
        if fa * fc < 0:
            b, fb = c, fc
        else:
            a, fa = c, fc
    return 0.5*(a+b)
```

Newton Raphson Method

```
def newton(f, df, x0, tol=1e-10, maxiter=50):
    x = x0
    for i in range(maxiter):
        fx = f(x)
        dfx = df(x)
        if abs(dfx) < tol:
            raise ZeroDivisionError("Derivative too small")
        dx = fx / dfx
        x = x - dx
        if abs(dx) < tol:
            return x
    return newton(f, df, x0, tol, maxiter)

def newton_fd(f, x0, tol=1e-10, maxiter=50, h=1e-6):
    def df(x):
        return (f(x+h) - f(x-h)) / (2*h)
    return newton(f, df, x0, tol, maxiter)
```

2 Numerical Integration Techniques

Trapezoidal Rule

```
def trapezoid(f, a, b, n=1000):
    x = np.linspace(a, b, n+1)
    y = f(x)
    h = (b-a)/n
    return (h/2) * (y[0] + 2*np.sum(y[1:-1]) + y[-1])
```

Simpson Rule

```
def simpson(f, n, a, b):
    if n % 2 == 1:
        raise ValueError("n must be even for Simpson's rule")
    h = (b - a) / n
    x = np.linspace(a, b, n + 1)
    y = f(x)
    intgr = h/3*(y[0]+y[-1] + 4*np.sum(y[1:-1:2]) + 2*np.sum(y[2:-2:2]))
    return intgr
```

3 Ordinary Differential Equations

Verlet Method

```
def velocity_verlet(x0, v0, a_func, t):
    t = np.asarray(t)
    n = t.size
    x = np.empty((n,) + np.shape(x0))
    v = np.empty((n,) + np.shape(v0))
    x[0] = x0
    v[0] = v0
    for i in range(n-1):
```

```
        dt = t[i+1] - t[i]
        ai = a_func(x[i], t[i])
        x[i+1] = x[i] + v[i]*dt + 0.5*ai*dt*dt
        a_next = a_func(x[i+1], t[i+1])
        v[i+1] = v[i] + 0.5*(ai + a_next)*dt
    return x, v
```

Runge Kutta 2 (Midpoint)

```
def rk2(f, y0, t):
    y0 = np.asarray(y0, dtype=float)
    t = np.asarray(t)
    n = t.size
    y = np.empty((n,) + y0.shape)
    y[0] = y0
    for i in range(n-1):
        dt = t[i+1] - t[i]
        k1 = f(t[i], y[i])
        k2 = (f(t[i] + 0.5*dt, y[i] + 0.5*dt*k1)
        y[i+1] = y[i] + dt*k2
    return y
```

RK4

```
def force(function, state):
    dxdt = state[1]
    dvxdt = function.x
    dydt = state[3]
    dvydt = function.y
    return np.array([dxdt, dvxdt, dydt, dvydt])
```

```
def rk4_step(f, state, t, dt):
    k1 = f(state, t)
    k2 = f(state + 0.5 * dt * k1, t + 0.5 * dt)
    k3 = f(state + 0.5 * dt * k2, t + 0.5 * dt)
    k4 = f(state + dt * k3, t + dt)
    return state + (dt / 6.0) * (k1 + 2*k2 + 2*k3 + k4)

def Solver(f, state_initial, t0, t_final, dt):
    n_steps = int(np.ceil((t_final - t0) / dt))
    t_eval = np.linspace(t0, t_final, n_steps + 1)
    history = np.empty((n_steps + 1, len(state_initial)))
    history[0] = state_initial

    state = state_initial.copy()
    for n in range(n_steps):
        t = t_eval[n]
        state = rk4_step(f, state, t, dt)
        history[n + 1] = state
    return history, t_eval
```

Predictor Corrector

Predictor-corrector methods combine an explicit *predictor* (cheap, possibly less accurate) with an implicit or more accurate *corrector* that uses the predicted value. A typical workflow is PECE (Predict, Evaluate, Correct, Evaluate) where the final Evaluate updates the derivative used in the next step.

Simple Heun (PECE) — one-step example.

```
Predictor:  $y_{n+1}^{(0)} = y_n + h f(t_n, y_n)$ ,  

Evaluate:  $f_{n+1}^{(0)} = f(t_{n+1}, y_{n+1}^{(0)})$ ,  

Corrector:  $y_{n+1} = y_n + \frac{h}{2} (f(t_n, y_n) + f_{n+1}^{(0)})$ ,  

Re-evaluate:  $f_{n+1}^{(1)} = f(t_{n+1}, y_{n+1})$  (Optional)
```

This is PECE: Predictor (Euler) → Evaluate → Correct (trapezoid) → Evaluate.

One may iterate the corrector (use the newly computed y_{n+1} to recompute f_{n+1} and reapply the Corrector formula) until the change is below a tolerance.

4 1D Boundary Value Problems

Newton Shooting

```
def newton_shooting(func, x_start: float, x_end: float, y_start: float,
                    y_end: float, num_points: int, max_iterations: int):
    slope= 1.0
    dx = (x_end - x_start) / num_points
    for _ in range(max_iterations):
        initial_state = [y_start, slope]
```

```
x_values, y_values = Solver(func, initial_state, x_start, x_end,
                               dx)
error = y_values[-1] - y_end
_, y_values_eps = Solver(func, [y_start, slope+ 1e-6],
                         x_start, x_end, dx)
error_eps = abs(error - y_end)
slope -= error / error_derivative
if abs(error) < 1e-6:
    print(f"Converged with boundary error: {error:.2e}")
break
return x_values, y_values
```

Finite Difference Method

```
def fd_bvp_linear(p, q, r, a, b, alpha, beta, m):
    h = (b-a)/m
    x = np.linspace(a, b, m+1)
    # interior unknowns y1..y_{m-1}
    A = np.zeros((m-1, m-1))
    B = np.zeros(m-1)
    for i in range(1, m):
        xi = x[i]
        ai = 1.0/h**2 - p(xi)/(2*h)
        bi = -2.0/h**2 + q(xi)
        ci = 1.0/h**2 + p(xi)/(2*h)
        idx = i-1
        if idx > 0:
            A[idx, idx-1] = ai
            A[idx, idx] = bi
        if idx < m-2:
            A[idx, idx+1] = ci
            B[idx] = r(xi)
    # incorporate boundary values
    B[0] = -(1.0/h**2 - p(x[1])/(2*h))*alpha
    B[-1] = -(1.0/h**2 + p(x[-2])/(2*h))*beta
    y_interior = np.linalg.solve(A, B)
    y = np.empty(m+1)
    y[0] = alpha
    y[-1] = beta
    y[1:-1] = y_interior
    return x, y
```

```
N = 100
M = 100
Lx, Ly = 1.0, 1.0
dx = Lx/(N+1)
dy = Ly/(M+1)
V = np.zeros((N+2, M+2))
# setting up Boundary Conditions along y
y = np.linspace(0, Ly, M+2)
V[0, :] = np.sin(2*np.pi*y)
# Constant Term
s = (np.zeros((N+2, M+2)))

V_sol = Solver(V, s, dx) # if dx and dy are different change things a bit
```

5.2 Matrix Formulation

```
# Define N, M and delta
N = 100
M = 100
dx = 1.0 / (N + 1)
dy = 1.0 / (M + 1)

# Create 1D finite difference matrices
def create_1d_laplacian(n):
    main_diag = np.full(n, -2)
    side_diag = np.full(n-1, 1)
    return np.diag(main_diag) + np.diag(side_diag, 1) +
    np.diag(side_diag, -1)
```

```
# 1D matrices
A_x = create_1d_laplacian(N)
A_y = create_1d_laplacian(M)
I_N = np.eye(N)
I_M = np.eye(M)

# 2D Laplacian using Kronecker products
D_x = np.kron(I_M, A_x) #  $\partial^2/\partial x^2$ 
D_y = np.kron(A_y, I_N) #  $\partial^2/\partial y^2$ 
L = D_x / (dx ** 2) + D_y / (dy ** 2)

# Boundary condition vectors
V_left = np.sin(2 * np.pi * np.linspace(0, 1, M)) # V(0,:)
V_right = np.zeros(M) # V(N+1,:)
V_bottom = np.zeros(N) # V(:,0)
V_top = np.zeros(N) # V(:,M+1)
```

```
# Pre-scale by dx, dy once
V_left /= dx ** 2
V_right /= dx ** 2
V_bottom /= dy ** 2
V_top /= dy ** 2

# Create source term
x = np.linspace(dx, 1 - dx, N)
y = np.linspace(dy, 1 - dy, M)
X, Y = np.meshgrid(x, y, indexing="ij")

f = 0 * X * Y
f_vec = f.flatten()

# Copy to boundary-corrected RHS
f_bc = f_vec.copy()

for j in range(M):
    for i in range(N):
        k = j * N + i # Flattened index

        if i == 0: # Left boundary
            f_bc[k] = V_left[j]
        if i == N - 1: # Right boundary
            f_bc[k] = V_right[j]
        if j == 0: # Bottom boundary
            f_bc[k] = V_bottom[i]
        if j == M - 1: # Top boundary
            f_bc[k] = V_top[i]

# Solve system
u_vec = np.linalg.solve(L, f_bc)
u_numerical = u_vec.reshape((N, M))

# Plot
```

5 2D Partial Differential Equations

5.1 Iterative Methods

Jacobi Iteration

```
def Jacobi(V, s, delta, max_iter=5000, tol=1e-6):
    for it in tqdm(range(max_iter), desc= "Iteration: "):
        V_new = V.copy()
        for n in range(1, N+1):
            for m in range(1, M+1):
                V_new[n,m] = 0.25*(V[n+1,m] + V[n-1,m] + V[n,m+1] +
V[n,m-1] - s[n,m]*delta*delta)
                if np.linalg.norm(V_new-V)/np.linalg.norm(V) < tol:
                    print(f"Converged at iteration {it}")
                    return V_new
        V = V_new
    return V
```

Gauss-Seidel

```
def GaussSeidel(V, s, delta, max_iter=5000, tol=1e-6):
    for it in tqdm(range(max_iter), desc= "Iteration: "):
        V_prev = V.copy()
        for n in range(1, N+1):
            for m in range(1, M+1):
                V[n,m] = 0.25*(V[n+1,m] + V[n-1,m] + V[n,m+1] + V[n,m-1] -
s[n,m]*delta*delta)
                if np.linalg.norm(V - V_prev)/np.linalg.norm(V) < tol:
                    print(f"Converged at iteration {it}")
                    return V
        V = V_prev
    return V
```

Running the Algorithm for both cases,

```
plt.contourf(Y, X, u_numerical, levels=25, cmap="viridis")
plt.colorbar()
plt.title("Numerical Solution")
plt.show()
```

6 Time Dependent PDEs

Consider the continuity equation,

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0$$

In one dimensional,

$$\frac{\partial \rho(x, t)}{\partial t} + \frac{\partial \rho u(x, t)}{\partial x} = 0$$

Forward Euler (Explicit) Scheme

Taking $u(x, t) = u$ constant we get

$$\rho_j^{n+1} = \rho_j^n - u \frac{\Delta t}{2\Delta x} (\rho_{j+1}^n - \rho_{j-1}^n)$$

This method is unstable as the eigenvalue for the P Matrix is always more than 1, thus norm increases to infinity.

```
def forward_euler(rho0, u, dx, dt, steps):
    rho = rho0.copy()
    N = len(rho)
    for _ in range(steps):
        rho_new = rho.copy()
        rho_new[1:-1] = rho[1:-1] - u*dt/(2*dx)*(rho[2:] - rho[:-2])
        rho = rho_new
    return rho
```

Implicit Euler Method

$$\rho_j^{n+1} = \rho_j^n - u \frac{\Delta t}{2\Delta x} (\rho_{j+1}^{n+1} - \rho_{j-1}^{n+1})$$

Thus we can write it as,

$$T = \begin{bmatrix} 1 & a & 0 & \dots & 0 \\ -a & 1 & a & \dots & 0 \\ 0 & \vdots & \vdots & \ddots & 0 \\ 0 & \dots & 0 & -a & 1 \end{bmatrix}, B = \begin{bmatrix} 0 & 0 & 0 & \dots & -a \\ 0 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & 0 \\ a & 0 & 0 & 0 & 0 \end{bmatrix}$$

where $a = \frac{u\Delta t}{2\Delta x}$

```
def implicit_euler_periodic(rho0, u, dx, dt, steps):
    rho = rho0.copy()
    N = rho.size
    a = u * dt / (2.0 * dx) # as defined in the notes
    # Build A = T + B (dense)
    A = np.eye(N, dtype=float)
    if N > 1:
        off = a * np.ones(N - 1, dtype=float)
        A += np.diag(off, 1) # +a on superdiagonal
        A -= np.diag(off, -1) # -a on subdiagonal
        # periodic wrap contributions (the B matrix)
        A[0, -1] = -a # top-right corner
        A[-1, 0] = a # bottom-left corner
    for _ in range(steps):
        rho = np.linalg.solve(A, rho)
    return rho
```

Leapfrog Method

By taking centre difference in not just space but also time,

$$\rho_j^{n+1} = \rho_j^{n-1} - u \frac{\Delta t}{\Delta x} (\rho_{j+1}^{n+1} - \rho_{j-1}^{n+1})$$

```
def leapfrog(rho_prev, rho_curr, u, dx, dt, steps):
    rho_nm1 = rho_prev.copy()
    rho_n = rho_curr.copy()
    for _ in range(steps):
        rho_np1 = rho_nm1 - u * dt / dx * (rho_n[2:] - rho_n[:-2]) / 2
        rho_nm1 = rho_n, np.concatenate(([0], rho_np1, [0]))
    return rho_n
```

Lax-Friedrich Method

Current step and position is taken as an average of neighbouring positions,

$$\rho^{n+1} = \frac{1}{2} \begin{bmatrix} 0 & 1 - u \frac{\Delta t}{\Delta x} & 0 & \dots & 0 \\ 1 + u \frac{\Delta t}{\Delta x} & 0 & 1 - u \frac{\Delta t}{\Delta x} & \dots & 0 \\ \vdots & & & \ddots & 0 \\ 0 & \dots & 0 & 1 + u \frac{\Delta t}{\Delta x} & 0 \end{bmatrix} \rho^n$$

```
def lax_friedrichs_T(N, u, dx, dt):
    c = u * dt / (2.0 * dx)
```

```
    upper = 0.5 - c # weight for rho_{j+1}
    lower = 0.5 + c # weight for rho_{j-1}
    T = np.zeros((N, N), dtype=float)
    for j in range(N):
        if j == 0:
            T[j, j - 1] = lower
        if j == N - 1:
            T[j, j + 1] = upper
        else:
            T[j, j - 1] = lower
            T[j, j + 1] = upper
    return T
```

```
def evolve_matrix(rho0, T, steps, periodic=False, u=None, dx=None, dt=None):
    rho = rho0.copy()
    N = rho.size
```

```
    if periodic:
        c = u * dt / (2.0 * dx)
        T[0, -1] = 0.5 - c # contribution from rho_{N-1} into rho_0
        T[-1, 0] = 0.5 + c # contribution from rho_0 into rho_{N-1}
        P = T + B
    else:
        P = T
    for _ in range(steps):
        rho = P @ rho
```

```
N = 200
x = np.linspace(0, 1, N, endpoint=False)
dx = x[1] - x[0]
u = 0.1
dt = 0.8 * dx / u
steps = 100

rho0 = np.exp(-(x - 0.5)**2) / 0.02
T = lax_friedrichs(N, u, dx, dt)
rho_p = evolve_matrix(rho0, T, steps, periodic=True, u=u, dx=dx, dt=dt)
```

Lax-Wendroff

```
def lax_wendroff_matrix(N, u, dx, dt):
    c = u * dt / (2 * dx)
    P = np.zeros((N, N))
    for j in range(N):
        P[j, (j - 1) % N] = c + c**2
        P[j, j] = 1 - 2 * c**2
        P[j, (j + 1) % N] = -c + c**2
    return P
```

6.1 Non-Linear Equations**Richtmyer 2-step Lax Wendroff**

```
import numpy as np
import matplotlib.pyplot as plt

N = 1000
a, b = 0, 1
x_full = np.linspace(a, b, N+2, dtype=np.float64)
x = x_full[1:-1]
dx = x[1] - x[0]

u = 0.01 * np.sin(2*np.pi*x)
t_max = 20.0
time = 0.0

fig, axes = plt.subplots(4, 5, sharex=True, sharey=True,
                       figsize=(15,12))

snapshots = 20
snapshot_times = np.linspace(0, t_max, snapshots)
s = 0
for snap_t in snapshot_times:
    while time < snap_t:
        u_max = np.max(np.abs(u)) + 1e-9
    print(f"True Value: {true_value:.5f}")
    print(f"True Value: {true_value:.5f}")
```

Lax-Friedrich Method

$dt = 0.6 * dx / u_{max}$

$dt_dx = dt / dx$

```
# p
f = 0.5 * u**2
u_half = 0.5*(u[1:] + u[:-1]) - 0.5*dt_dx*(f[1:] - f[:-1])
f_half = 0.5 * u_half**2
```

$u[1:-1] = u[1:-1] - dt_dx * (f_half[1:] - f_half[:-1])$

periodic wrap for ghost cells

$u[0] = u[-2]$

$u[-1] = u[1]$

$time += dt$

$r = s // 5$

$c = s \% 5$

$ax = axes[r, c]$

$ax.plot(x, u, color=f'#{8*s:06X}')$

$ax.set_title(f't = {time:.2f}s")$

$ax.grid(True)$

$s += 1$

plt.tight_layout()

plt.show()

Schrodinger Equation

```
def crank_nicolson(psi0, V, dx, dt, hbar=1, m=1, steps=1000):
    N = len(psi0)
```

```
    D = np.diag(-2*np.ones(N)) + np.diag(np.ones(N-1), 1) +
        np.diag(np.ones(N-1), -1)
```

$P = np.diag(V)$

$H = -D + P$

$A = np.eye(N) + 1j * dt / (2) * H$

$B = np.eye(N) - 1j * dt / (2) * H$

$\psi = psi0.copy()$

for _ in range(steps):
 psi = np.linalg.solve(A, B @ psi)

psi /= np.linalg.norm(psi) # normalization

return psi

7 Monte Carlo Methods**7.1 Integration****Area of Circle**

$N = 10000$

$rng = np.random.default_rng(seed=67)$

$x = rng.random(N)$

$y = rng.random(N)$

$inside = (x*x + y*y) < 1$

$cum_inside = np.sum(inside)$

$trials = np.arange(1, N+1)$

$value_circle = 4 * (cum_inside / trials)$

$print(value_circle[-1])$

Volume of N-dim Sphere

$N = 1000000$

$dimension = 10$

N_dim_Sphere(N, dimension):

$rng = np.random.default_rng(seed=70)$

$points = rng.random((N, dimension))$

$points = points*points$

$r = np.sum(points, axis=1)$

$inside = np.where(r <= 1, 1, 0)$

$cum_inside_me = np.sum(inside)$

$trials_ndim = np.arange(1, N+1)$

$value_circle = (2**dimension) * (cum_inside_me / trials_ndim)$

$print(f"Monte Carlo Estimate: {value_circle[-1]:.5f}")$

$true_value = (np.pi**dimension/2) / math.gamma(dimension/2 + 1) # \pi^{d/2}/\Gamma(d/2+1)$

$print(f"True Value: {true_value:.5f}")$

Area under Function

```
def Monte_Carlo_Integration(f, x_lim: np.ndarray, y_lim: np.ndarray, N: int = 10000, seed: int = 42):
    np.random.seed(seed)
    points = np.random.random((N, 2))
```

Making Random Points go into range of x_lim and y_lim

points[:, 0] *= x_lim[1] - x_lim[0]

points[:, 1] *= y_lim[1] - y_lim[0]

points[:, 0] += x_lim[0]

points[:, 1] += y_lim[0]

f_points = f(points[:, 0], 0) # Array of evaluated values on given random points

positive_region = (f_points >= 0) & (points[:, 1] >= 0) & (points[:, 1] <= f_points)

negative_region = (f_points < 0) & (points[:, 1] < 0) & (points[:, 1] >= f_points)

point_contributions = positive_region.astype(int) - negative_region.astype(int)

cumulative_contributions = np.cumsum(point_contributions)

Running estimate at each iteration

box_area = (x_lim[1] - x_lim[0]) * (y_lim[1] - y_lim[0])

iterations = np.arange(1, N + 1)

convergence = (cumulative_contributions / iterations) * box_area

return convergence[-1]

limits_x = np.array([0, 3*np.pi/2])
limits_y = np.array([np.exp(0), np.exp(3*np.pi/2)]) # Min and Max Value of the function in entire range

N = 10000
seed = 67
print(f"Monte Carlo: {Monte_Carlo_Integration(function, limits_x, limits_y, N, seed):.5f}")

7.2 Casting Distributions

To generate a function a distribution $f(y)$, then we need to create a mapping from $x \mapsto y$ such that

$$\frac{dy}{dx} = \frac{1}{f(y)}$$

Then for a probability distribution $\rho(x)$, we can find the integral by,

$$\int f(x) dx = \frac{1}{\rho(x)}$$

Closed Form Integrals

Given $f(y) = \sin(y)$, we can say $\sin(y) dy = dx \Rightarrow y = \arccos(1-x)$.

```
N = 1000000
x = np.random.rand(N)
y = np.arccos(1 - x) # Inverse transform sampling
plt.hist(y, bins=100, density=True, alpha=0.6, label='Sampled')
plt.plot(np.linspace(0, np.pi/2, 100), np.sin(np.linspace(0, np.pi/2, 100))/1, 'r', lw=2, label='sin(y)')
plt.legend(); plt.show()
```

Box Muller Transform

To generate the standard normal distribution,

$$f(y) = \frac{1}{\sigma\sqrt{2\pi}} \exp\left[-\frac{1}{2}\frac{(y-\mu)^2}{\sigma^2}\right] \text{ where } \mu = 0, \sigma = 1$$

There exists no closed form solution to the required integral to obtain the map $x \mapsto y$ so we convert to polar coordinates,

$$f(x)f(y) dx dy \mapsto \frac{1}{2\pi} \exp\left(-\frac{r^2}{2}\right) r dr d\theta$$

Now to generate the probability distribution above,

$$g(r) = \exp\left(-\frac{r^2}{2}\right) r \Rightarrow dx = g(r) dr \Rightarrow r = -\sqrt{2\ln(r)}$$

$$\theta = x \times (2\pi)$$

Thus, we get 2 normal distributions

$$y_1 = r \cos(\theta) \quad y_2 = r \sin(\theta)$$

```
N = 100000

x_1 = np.random.rand(N)
x_2 = np.random.rand(N)

r = np.sqrt(-2*np.log(x_1))
theta = 2*np.pi*x_2

y_1 = r * np.cos(theta)
y_2 = r * np.sin(theta)
normal = lambda x: np.exp(-x**2/2)/(np.sqrt(2*np.pi))
compare = np.linspace(-3, 3, N)

counts1, bins1 = np.histogram(y_1, bins=50)
deltax1 = bins1[1]-bins1[0]
pdf1 = (counts1/N)/deltax1
bin_mid1 = (bins1[:1]+bins1[1:])/2

counts2, bins2 = np.histogram(y_2, bins=50)
deltax2 = bins2[1]-bins2[0]
pdf2 = (counts2/N)/deltax2
bin_mid2 = (bins2[:1]+bins2[1:])/2
```

```
counts2, bins2 = np.histogram(y_2, bins=50)
deltax2 = bins2[1]-bins2[0]
pdf2 = (counts2/N)/deltax2
bin_mid2 = (bins2[:1]+bins2[1:])/2

plt.hist(y_1, bins=50, alpha=0.5, label='cos', color='r')
plt.hist(y_2, bins=50, alpha=0.5, label='sin', color='g')
plt.legend()
plt.plot(bin_mid1, pdf1, label='cos', color='r', alpha = 0.5)
plt.plot(compare, normal(compare), label='normal', color='b')
plt.plot(bin_mid2, pdf2, label='sin', color='g', alpha = 0.5)
plt.legend()
plt.show()
plt.clf()
```

7.3 Random Walks

Metropolis Hastings

```
import numpy as np
import matplotlib.pyplot as plt

# Initialize parameters
n_samples = 100000 # total number of samples to generate
samples = [] # list to store accepted samples
x_current = 0.5 # start point in the middle of [0,1]
sigma = 0.1 # proposal step size (standard deviation)

# Run the Metropolis algorithm
for i in range(n_samples):
    # Propose a new candidate from a normal distribution centered at current x
    x_proposed = x_current + np.random.normal(0, sigma)

    # Reflect if the proposed value goes out of [0,1] (to stay within bounds)
    if x_proposed < 0:
        x_proposed = -x_proposed
    elif x_proposed > 1:
        x_proposed = 2 - x_proposed

    # Compute acceptance ratio (unnormalized)
    p_current = - np.sin(x_current) * np.log(x_current)
    p_proposed = - np.sin(x_proposed) * np.log(x_proposed)

    # Avoid divide-by-zero errors
    if p_current == 0:
        acceptance_ratio = 1
    else:
        acceptance_ratio = p_proposed / p_current

    # Accept or reject based on Metropolis criterion
    if np.random.rand() < min(1, acceptance_ratio):
        x_current = x_proposed # accept move

    # Store the current sample
    samples.append(x_current)

# Convert to NumPy array for analysis
samples = np.array(samples)

# Plot the sampled distribution
plt.figure(figsize=(8, 5))
```

```
plt.hist(samples, bins=100, density=True, alpha=0.6, color='skyblue',
label='Sampled distribution')

# Plot the true (unnormalized) function for comparison
x = np.linspace(0.001, 0.999, 500)
norm_const = np.trapezoid(-np.sin(x)*np.log(x), x)
plt.plot(x, -np.sin(x)*np.log(x) / norm_const, 'r-', lw=2, label='True
(normalized) function')
plt.xlabel('x')
plt.ylabel('Probability density')
plt.legend()
plt.title('Sampling from f(x) = -sin(x) ln(x) using Metropolis
algorithm')
plt.show()
```

8 Neural Networks

Linear model

- We can see that A is

$$A = \begin{bmatrix} 0 & 0 \\ 2 & 1 \end{bmatrix} \quad (4)$$

- This matrix is uniquely determined by the first two conditions
- But then it does not satisfy the third condition
- This means a linear model will not work in this case
- In general, as the number of cases increases, you will get more conditions, while the matrix will have maximum rank of n
- So the system will be overdetermined and we cannot fit the linear model exactly
- One way to overcome this drawback is to go to a nonlinear model

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 8/35

Nonlinear model

- Take the output of a linear model and apply a nonlinear function to it

$$z = g(\mathbf{y}) = \begin{bmatrix} g(y_1) \\ g(y_2) \\ \vdots \\ g(y_m) \end{bmatrix} \quad (5)$$

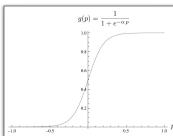
- Defining $g(p)$ gives extra flexibility to fit all the required conditions, ex. the third condition above
- The input and outputs can be standardized for ease of operation
- Set $g(p)$ to lie between 0 and 1

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 9/35

Sigmoid function

- The sigmoid function is a continuous version of the step function

$$g(p) = \frac{1}{1 + e^{-\alpha p}} \quad (6)$$



- α sets the width of the sigmoid function, higher α makes the function sharper

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 10/35

Setting α

- The elements of the nonlinear output \mathbf{z} are

$$z_i = g((\mathbf{Ax})_i) = g\left(\sum_{j=1}^n A_{ij}x_j\right) \quad (7)$$

- Now the argument of function g should span the range of the width of the sigmoid function, so that the output \mathbf{g} takes the full range of 0 to 1
- Assume that the elements of matrix A will be of the order unity
- Then the maximum element of \mathbf{Ax} will be approximately $n\max(x_i)$
- Then take α as

$$\alpha = \frac{10}{n\max(x_i)} \quad (8)$$

- Also, we can shift the input numbers $x_i \rightarrow x_i - \text{med}(x)$ such that the maximum and minimum are equal and opposite in signs
- This way we can handle input with any magnitude \mathbf{x}

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 11/35

Training

- Now the next step is to find the matrix A
- This process is called as training. In this we provide training examples to the model
- These examples are basically pairs of $(\mathbf{x}^k, \mathbf{y}^k)^T$ where \mathbf{x}^k is the k -th input vector and \mathbf{y}^k is its desired output vector
- There are T training examples, labelled by k , and our goal is to set A such that we get \mathbf{y}^k for \mathbf{x}^k
- For this we define an error function f

$$f^k = \|g(\mathbf{Ax}^k) - \mathbf{y}^k\|^2 \text{ for each case } j = 1, 2, \dots, T \quad (9)$$

- The norm of a vector can be defined as the square root of the sum of the squares of its elements, i.e.,

$$f^k = \sqrt{\sum_{i=1}^m \left(\sum_{s=1}^n A_{is}x_s^k - y_i^k \right)^2} \quad (10)$$

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 12/35

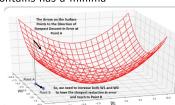
Training

- This function f can be thought of as a function of the matrix A_{ij} because that is the only variable, rest all quantities come from the training set
- Hence, the objective then is to adjust A_{ij} such that we can minimize f for each training example
- f is a positive definite quantity, and this problem is of minimization of f
- A well-known method is known as the steepest descent (gradient descent) method

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 13/35

Gradient descent

- Consider a function of 2 variables, this can be represented as a surface
- This function contains has a minima



- If we have a function $f(\mathbf{x})$ of which we want to find minimum, then starting from some position $\mathbf{x}_0 = (x_0, y_0)$

$$\mathbf{x}_{n+1} = \mathbf{x}_n - \eta \nabla f \quad (11)$$

- η is the step size, also called as "rate of learning"

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 14/35

Steepest descent

- Ex. find the minimum of $f(x, y) = x^2 + 4y^2 - 4x - 24y + 40$
- This actually comes from $(x - 2)^2 + 4(y - 3)^2$, so the minimum is at $(x, y) = (2, 3)$
- Learning rate can be faster with higher η
- However, too high η might lead to method not converging or instability
- Also called steepest descent because change is maximum in direction of gradient
- Convergence is determined when the steps become smaller than some threshold

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 15/35

Training model

- Going back to the case of finding the matrix A in our model
- The function to be minimized is

$$f(A_{ij}) = \sum_{i=1}^m \left[g\left(\sum_{j=1}^n A_{ij}x_j\right) - y_i \right]^2 \quad (12)$$

- given any training pair $(\mathbf{x}^k, \mathbf{y}^k)$ - suppressing the superscript k for simplicity

- Define $b_i = \sum_{j=1}^n A_{ij}x_j$ and $z_i = g(b_i)$, i goes from 1 to m
- Now we want to find the values of matrix elements A_{pq} which minimizes the error f

- For using gradient descent we need

$$\frac{\partial f}{\partial A_{pq}} = \sum_{i=1}^m 2(z_i - y_i) \frac{\partial z_i}{\partial A_{pq}} \quad (13)$$

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 16/35

Training model

- We have

$$\frac{\partial z_i}{\partial A_{pq}} = g'(b_i) \frac{\partial b_i}{\partial A_{pq}} \quad (14)$$

$$\frac{\partial b_i}{\partial A_{pq}} = \frac{\partial}{\partial A_{pq}} \sum_{j=1}^n A_{pj}x_j \quad (15)$$

$$= \sum_{j=1}^m \delta_{pj} \delta_{jq} x_j \quad (16)$$

$$= \delta_{pq} x_q \quad (17)$$

- Also can be shown that

$$g'(p) = \alpha g(p)(1 - g(p)) \quad (18)$$

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 17/35

Training model

- Putting this together we get

$$\frac{\partial f}{\partial A_{pq}} = \sum_{i=1}^m 2(z_i - y_i) \alpha g(b_i)(1 - g(b_i)) \delta_{pq} x_q \quad (19)$$

- which gives

$$\frac{\partial f}{\partial A_{pq}} = 2(z_p - y_p) \alpha g(p)(1 - g(p)) x_q \quad (20)$$

- This makes the calculation of the matrix element easy

Kriti Makarand/Department of Physics, IIT-Hyderabad EP24040 - Computational Physics | 18/35

Algorithm

- Initialize random A of order unity
 - Loop over each training pair
 - Calculate z using the existing A matrix
 - Loop over all elements of A_{pq} and update with
- $$A_{pq} \rightarrow A_{pq} - \eta \left[2(z_p - y_p) \alpha z_p (1 - z_p) x_q \right] \quad (21)$$

- Taking one gradient descent step over all the training examples is called an "epoch"
- Run the training for several epochs till the error is reduced below some threshold or the accuracy of the model reaches desired level
- The accuracy can be defined as the percentage of predictions it gets right
- After going through all training pairs and epochs, the model is trained and can be used for new data

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 19 / 35

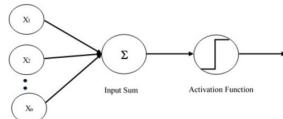
Training model

- This training is not expected to produce exact results
- The outcome of the training will depend on the order in which the test cases are given
- Ideally after giving the training examples multiple times and in random order, the neural network A should converge
- However, this may not be the case always and it may not give good agreement for all the training examples
- After the model is trained on the training set, it can be tested on a testing set to check how well it performs
- If the performance is good, it can be used on new and untested data

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 20 / 35

Hidden layers

- If we are given an input vector with 10 elements, i.e., $n = 10$ and the output if just a single element, i.e., $m = 1$ then
- The matrix A will be 10×1 matrix with 10 parameters to adjust
- This can be represented diagrammatically as



Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 21 / 35

Hidden layers

- If there is a lot of data which is making the problem overconstrained, then we require more freedom in the neural network model
- We can add an intermediate layer of size k
- The input $x \in R^n$ gets multiplied by matrix B which is $k \times n$ giving a $\tilde{z} \in R^k$

$$\tilde{z} = g(Bx) - \frac{1}{2} \quad (22)$$

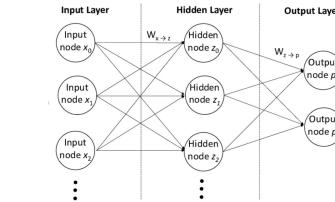
- Since the nonlinear function g is applied here, the $-1/2$ changes the mapping from 0 - 1 to $-1/2$ to $+1/2$
- Now the matrix A of size $m \times k$ is applied on \tilde{z} to get the final output

$$z = \tilde{g}(A\tilde{z}) \quad (23)$$

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 22 / 35

Hidden layers

- The hidden layers can be diagrammatically represented as



Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 23 / 35

Training

- Now both the B and A matrices have to be found by training the model
- The error is again defined as

$$f(A_{ij}, B_{ij}) = \sum_{l=1}^m (z_l - y_l)^2 \quad (24)$$

- Using the steepest gradient technique as before, we can calculate the partial derivative of this w.r.t. the matrix elements A_{ij} and B_{ij}
- It can be shown that

$$\frac{\partial f}{\partial A_{pq}} = 2\tilde{z}(z_p - y_p)z_p(1 - z_p)\tilde{z}_q \quad (25)$$

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 24 / 35

Training

- Similarly, differentiating w.r.t. B_{pq} it can be shown that

$$\frac{\partial f}{\partial B_{pq}} = \sum_{l=1}^m z_l A_{lp}(\frac{1}{2} + \tilde{z}_p)(\frac{1}{2} - \tilde{z}_p)x_q \quad (26)$$

- where

$$z_l = 2\tilde{z}(z_l - y_l)z_l(1 - z_l) \quad (27)$$

- Now we can just modify the update algorithm

$$A_{pq} \rightarrow A_{pq} - \eta \frac{\partial f}{\partial A_{pq}} \quad (28)$$

$$B_{pq} \rightarrow B_{pq} - \eta \frac{\partial f}{\partial B_{pq}} \quad (29)$$

- More generally, the gradient w.r.t. the weights of the neural network nodes are estimated by a method of backward propagation

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 25 / 35

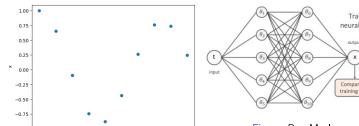
Different models

- A model with a large number of hidden layers is a deep network
- To generate mode data, noise can be added to the data
- The images can also be tranformed, like left-right inversion, or time inversion, etc.
- Convolutional neural network - image data is processed through some convolution in order to account for neighboring/related pixels
- Taking a large hidden layer (high k) versus many hidden layers - mostly based on experience and experimentation
- Steepest descent leads to the local minimum, might not always work or find the best possible network - therefore need to try many different starting points and train multiple times

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 26 / 35

Physics informed neural networks

- Neural networks can also be used to solve differential equations describing physical laws
- For example - suppose we experimentally measure the displacement of a particle as a function of time



- Now we need to design a neural network that can fit a function f(t) that will approximately give us the value of the displacement u(t)

$$f(t) \approx u(t) \quad (37)$$

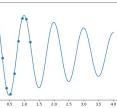
Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 29 / 35

Physics informed neural networks

- With limited data, the neural network will not be able to extrapolate the displacement to later times
- Rather, if we know that the displacement obeys a differential equation then we can get the full solution

$$m \frac{d^2x}{dt^2} + \mu \frac{dx}{dt} + kx = 0 \quad (38)$$

- This is the equation of a damped harmonic oscillator (SHM if $\mu = 0$) and its solution looks as below



Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 30 / 35

PINN

- Consider the damped harmonic oscillator differential equation

$$m \frac{d^2x}{dt^2} + \mu \frac{dx}{dt} + kx = 0 \quad (39)$$

- with the initial condition $x(t = 0) = 1$ and $x'(t = 0) = 0$, i.e. an oscillator starting from the peak displacement
- Setup a neural network NN that takes t as an input and gives x as output

- We have to train the network weight factors θ such that

$$NN(t, \theta) \approx x(t) \quad (40)$$

- We do not need any experimental data to train this network
- The neural network will start generating data itself and the differential equation and boundary conditions will help in training

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 31 / 35

Loss function

- The network is initialized with random weights
- In each epoch a few input collocation times ($t_i, i=1, 2, \dots, N_p$) are chosen randomly over the range of interest (including the boundary point) - it is important to train the model over random time points in order to avoid bias towards early or later times
- The network makes predictions $NN(t_i, \theta)$
- The loss function consists of two parts, the boundary loss term

$$f_b = \lambda_1 \left(NN(t = 0, \theta) - 1 \right)^2 + \lambda_2 \left[\frac{dNN}{dt}(t = 0, \theta) - 0 \right]^2 \quad (41)$$

- and the differential equation loss is (also called physics-loss)

$$f_p = \lambda_3 \sum_{i=1}^{N_p} \left[m \frac{d^2NN}{dt^2}(t_i) + \mu \frac{dNN}{dt}(t_i) + kNN(t_i) \right]^2 \quad (42)$$

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 32 / 35

Grad function

- The weights of the different loss terms $\lambda_1, \lambda_2, \lambda_3$ can be adjusted

```
from torch.autograd import Variable
t_collocation = np.random.uniform(0, 4, size=(1000, 1))
pt_t_collocation = Variable(torch.from_numpy(t_collocation).float(), requires_grad=True)
```

```
# After defining a network NN(time) which takes an input time tensor, its derivative can be calculated
xn=NN(time)
dxndt=torch.autograd.grad(xn,time)
```

```
# Similarly, getting the second derivative will involve taking one more grad on the dxndt variable
```

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 33 / 35

Wave equation

- In order to solve the wave equation for a quantity $u(x, t)$

$$\frac{\partial^2 u}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} = 0 \quad (43)$$

- For this we will define a $NN(x, t)$ and the physics loss will be

$$f_p = \frac{1}{N_p} \sum_i \left[\frac{\partial^2 NN(x_i, t_i)}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 NN(x_i, t_i)}{\partial t^2} \right] \quad (44)$$

- Where the summation is over all collocation points

- For periodic boundary condition the loss is

$$f_b = \sum_j [NN(x_j, t_j) - NN(x_j + L, t_j)]^2 \quad (45)$$

- where x_j is a boundary points and L is the period

Kriti MakwanaDepartment of Physics, IIT-Hyderabad EP24040 - Computational Physics 34 / 35