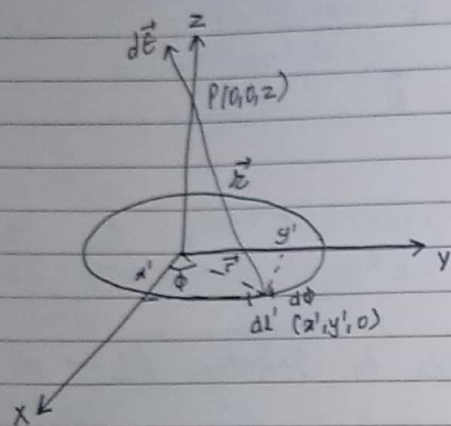


Example: Find the electric field at distance  $z$  above center of a circular loop of radius ' $r$ ' which carries a uniform line charge  $\lambda$ .



Here, the figure illustrates the geometry and coordinates to be used.

From figure,

$$\vec{r} = -x'\hat{i} - y'\hat{j} + z\hat{k}$$

$$= -r\cos\phi\hat{i} - r\sin\phi\hat{j} + z\hat{k}$$

$$\therefore r = (r^2 + z^2)^{1/2}$$

The charge on an elemental length  $dl'$  along a circular loop is,

$$dq = \lambda dl'$$

$$\therefore dq = \lambda r d\phi$$

The electric field at  $P$  due to the charge  $dq$  is,

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^3} \vec{r}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{\lambda r d\phi}{(r^2 + z^2)^{3/2}} (-r\cos\phi\hat{i} - r\sin\phi\hat{j} + z\hat{k})$$

$\therefore$  The net electric field at  $P$  due to the charge on whole circular loop is.

$$\vec{E} = \int d\vec{E}$$

$$= \frac{\lambda r}{4\pi\epsilon_0 (r^2 + z^2)^{3/2}} \int_0^{2\pi} (-r\cos\phi\hat{i} - r\sin\phi\hat{j} + z\hat{k}) d\phi$$

$$= \frac{\lambda r}{4\pi\epsilon_0 (r^2 + z^2)^{3/2}} \left[ \left( -r \int_0^{2\pi} \cos\phi d\phi \right) \hat{i} + \left( -r \int_0^{2\pi} \sin\phi d\phi \right) \hat{j} + \left( z \int_0^{2\pi} d\phi \right) \hat{k} \right]$$

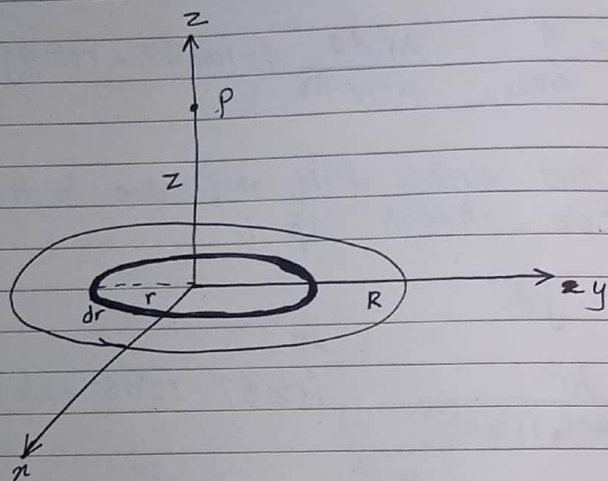
$$= \frac{\lambda r}{4\pi\epsilon_0 (r^2 + z^2)^{3/2}} \times z \times (2\pi) \hat{k} \left[ \int_0^{2\pi} \cos\phi d\phi = 0, \int_0^{2\pi} \sin\phi d\phi = 0 \right]$$

$$= \frac{1}{4\pi\epsilon_0} \lambda (2\pi r) \cdot \frac{z}{(r^2 + z^2)^{3/2}} \hat{k}$$

$$\therefore \vec{E} = \frac{1}{4\pi\epsilon_0} q \frac{z}{(r^2 + z^2)^{3/2}} \hat{k}$$

Example: Find the electric field at a distance  $z$  above the center of a flat circular disc of radius  $R$  which carries uniform surface charge  $\sigma$ .

Soln:



Here, the figure illustrates the geometry and the coordinates to be used.

The disc can be considered as combination of infinite numbers of infinitesimally thin rings.

Consider a ring of radius  $r$  and thickness  $dr$  of this disc.

If  $\sigma$  is the uniform charge density, then the charge on the ring  $dq = \sigma(2\pi r dr)$

The electric field at P due to charge  $dq$  on the ring is,

$$\begin{aligned} d\vec{E} &= \frac{1}{4\pi\epsilon_0} dq \frac{z}{(r^2+z^2)^{3/2}} \hat{k} \\ &= \frac{1}{4\pi\epsilon_0} \sigma(2\pi r dr) \frac{z}{(r^2+z^2)^{3/2}} \hat{k} \\ &= \frac{\sigma z}{2\epsilon_0} \frac{r dr}{(r^2+z^2)^{3/2}} \hat{k} \quad \text{--- (i)} \end{aligned}$$

Hence, the total electric field due to the charge on whole flat circular disc is given by,

$$\begin{aligned} \vec{E}_{\text{disc}} &= \int d\vec{E} \\ &= \frac{\sigma z}{2\epsilon_0} \int_0^R \frac{r dr}{(r^2+z^2)^{3/2}} \hat{k} \end{aligned}$$

$$\text{Put } r^2+z^2 = t^2$$

Then,

$$\begin{aligned} 2r dr &= 2t dt \\ \therefore r dr &= t \cdot dt \end{aligned}$$

$$\text{When } r=0, \quad t=z$$

$$\text{When } r=R, \quad t = \sqrt{R^2+z^2}$$

$$\text{So, } \vec{E}_{\text{disc}} = \int_z^{\sqrt{R^2+z^2}} \frac{t \cdot dt}{t^3}$$



$$= \frac{\sigma z}{2\epsilon_0} \int_z^{\sqrt{R^2+z^2}} \frac{dt}{t^2} \hat{k}$$

$$= \frac{\sigma z}{2\epsilon_0} \left[ -\frac{1}{t} \right]_z^{\sqrt{R^2+z^2}} \hat{k}$$

$$= \frac{\sigma z}{2\epsilon_0} \left[ -\frac{1}{\sqrt{R^2+z^2}} + \frac{1}{z} \right] \hat{k}$$

$$= \frac{\sigma z}{2\epsilon_0} \left[ \frac{1}{z} - \frac{1}{\sqrt{R^2+z^2}} \right] \hat{k}$$

$$\therefore \vec{E}_{disc} = \frac{\sigma}{2\epsilon_0} \left( 1 - \frac{z}{\sqrt{R^2+z^2}} \right) \hat{k}$$

As  $R \rightarrow \infty$ ,

$$\therefore \vec{E}_{disc} = \frac{\sigma}{2\epsilon_0} \hat{k}$$

for points far from the disc,  $z \gg R$ .

$$\vec{E}_{disc} = \frac{\sigma}{2\epsilon_0} \left[ 1 - \frac{1}{(R^2+z^2/z^2)^{1/2}} \right] \hat{k}$$

Using binomial expansion,

$$= \frac{\sigma}{2\epsilon_0} \left[ 1 - \left( 1 + \frac{R^2}{z^2} \right)^{-1/2} \right] \hat{k}$$

Using binomial expansion,

$$= \frac{\sigma}{2\epsilon_0} \left[ 1 - \left( 1 - \frac{1}{2} \frac{R^2}{z^2} + \dots \right) \right] \hat{k}$$

$$\therefore \vec{E}_{disc} = \frac{\sigma}{2\epsilon_0} \times \left( \frac{1}{2} \frac{R^2}{z^2} \right) \hat{k} = \frac{1}{4\pi\epsilon_0} \frac{\sigma \pi R^2}{z^2} \hat{k}$$

$$\therefore \vec{E}_{disc} = \frac{1}{4\pi\epsilon_0} q \frac{1}{z^2} \hat{k}$$

### # Electric Field Lines

Electric field lines describes an electric field in any region of space.

The electric field vector  $\vec{E}$  is tangent to the electric field line at each point.

The number of lines per unit area through a surface perpendicular to the lines is proportional to the magnitude of electric field in that region.

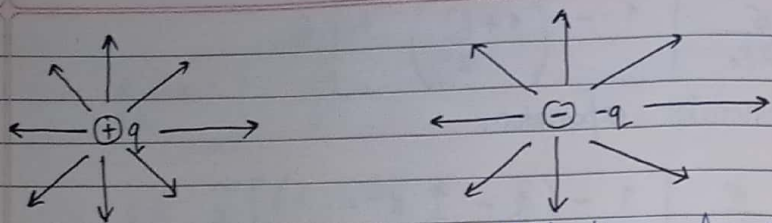
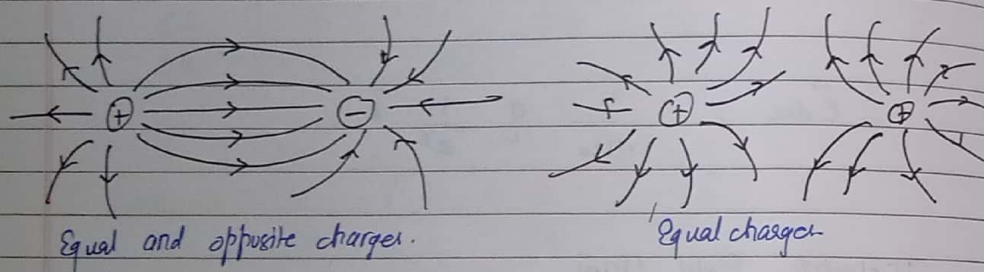


Fig: Representation Electric Field Lines for the field due to single point charge.

For any two point charges, the electric field lines.



### (\*) Properties:

- (i): Electric <sup>field</sup> lines don't intersect.
- (ii): the tangent to electric field line <sup>at a point</sup> gives electric field vector  $\vec{E}$  at that point.
- (iii) The lines begin on a positive charge and end on a negative charge.

### # Electric Flux

The total number of electric field lines passing a given area in unit time is called electric flux.

It is directly proportional to the number of electric field lines that penetrate the surface.

The electric field flux through a surface  $S$  is.

$$\Phi_E = \int_S \vec{E} \cdot d\vec{a}$$

SI unit:  $\text{Nm}^2\text{C}^{-1}$

The electric flux through any closed surface gives the measure of the total charge inside.

### # Gauss Law:

Gauss's law states that, "the total electric flux through any closed surface is equal to  $1/\epsilon_0$  times the total charge enclosed by the surface."

Mathematically

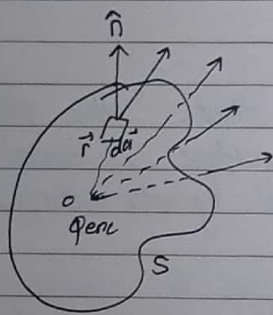
$$\oint_S \vec{E} \cdot d\vec{a} = \frac{1}{\epsilon_0} Q_{\text{enc.}}$$



Q: Two concentric imaginary spherical surfaces of radii  $R$  and  $2R$  respectively surround a positive charge  $q$  located at the centre of the surfaces. When compared to electric flux  $\phi_1$  through the surface of radius  $R$ , the electric flux  $\phi_2$  through the surface of radius  $2R$  is.

- (a)  $\phi_1 = 2\phi_2$   
 (b)  $2\phi_1 = \phi_2$   
 (c)  $\phi_1 = \frac{1}{4}\phi_2$   
 (d)  $\phi_1 = \phi_2$   $\Leftarrow$  Correct answer.  $\phi \propto Q_{enc}$ .

Let us consider an arbitrary shaped which encloses a point charge  $Q_{enc}$ .



The electric field at vector  $\vec{r}$  due to point charge located at origin is.

$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q_{enc}}{r^2} \vec{r}$$

Then, the total electric flux passing through a closed surface  $S$  is.

$$\oint_S \vec{E} \cdot d\vec{a} = \oint_S \left( \frac{1}{4\pi\epsilon_0} \frac{Q_{enc}}{r^2} \vec{r} \right) \cdot (da_r \hat{r} + da_\theta \hat{\theta} + da_\phi \hat{\phi})$$

$$= \frac{Q_{enc}}{4\pi\epsilon_0} \oint_S \frac{1}{r^2} da_r$$

$$= \frac{Q_{enc}}{4\pi\epsilon_0} \oint_S \frac{1}{r^2} (r^2 \sin\theta d\theta d\phi)$$

$$= \frac{Q_{enc}}{4\pi\epsilon_0} \left[ \int_0^\pi \sin\theta d\theta \cdot \int_0^{2\pi} d\phi \right]$$

$$= \frac{Q_{enc}}{4\pi\epsilon_0} \times 2 \times 2\pi \quad \therefore \oint_S \vec{E} \cdot d\vec{a} = \frac{Q_{enc}}{\epsilon_0}$$

This is integrated form of Gauss law.

For multiple charge

$$\oint_S \vec{E} \cdot d\vec{a} = \frac{1}{\epsilon_0} \sum_{i=1}^n Q_i \quad [\text{for discrete distribution}]$$

$$= \frac{1}{\epsilon_0} \int_S d\tau \quad [\text{for continuous distribution}]$$

Applying divergence theorem,

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

This is differential form of Gauss law.