

# Statistical ORBIT DETERMINATION, ASEN5070

Lecture 7

Fundamentals of Orbit Determination Fall 2011, 9/7/2011

Supplemental Reading:

**Chapter 4** 

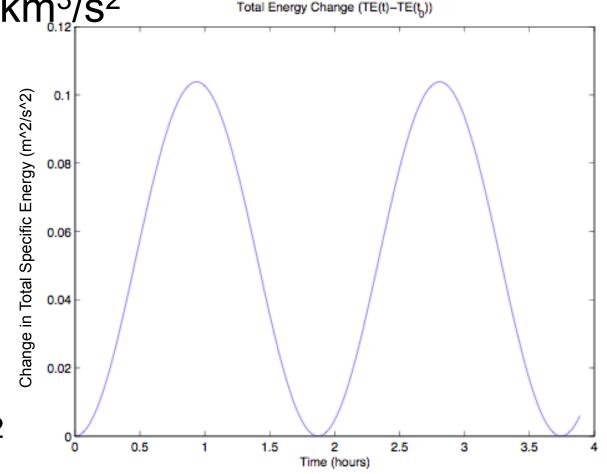


# Accuracy/Consistency of MU

Error in  $\mu$  of 0.5 km<sup>3</sup>/s<sup>2</sup>

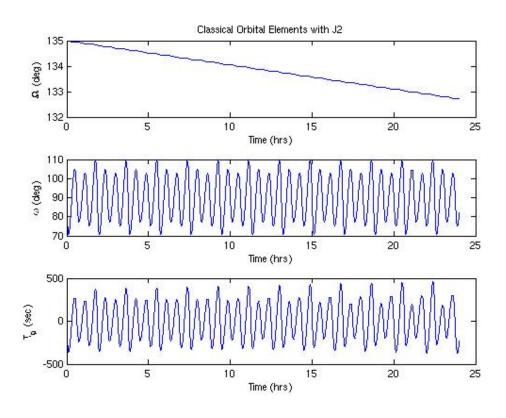
Integrator  $\mu = 398600 \text{ km}^3/\text{s}^2$ 

Energy
Calculation μ = 398600.5 km<sup>3</sup>/s<sup>2</sup>



# Orbit Elements HW #2







# T<sub>p</sub> Calculation

 $T_p$  is determined from the following equations:

$$M = E - e \sin E$$
  $n = \sqrt{\frac{\mu}{a^3}}$   $t - T_p = M/n$ 

However, as time t increases,  $T_p$  is not constrained to an orbital period and thus increases as a step function. To resolve this, MOD  $T_p$  with the orbital period.

$$T_p = MOD(T_p, P)$$

A situation may arise in which the calculation for the mean Anomaly, M, and true anomaly, v, do not agree resulting in the mean anomaly to be past perigee while the true anomaly is behind perigee (this is an artifact of numerical integration).



# T<sub>p</sub> Calculation

To correct this, we will introduce the angle of periapse  $\theta_0$ :

$$\theta_p = nT_p$$

From this, one will notice that the artifacts occur when

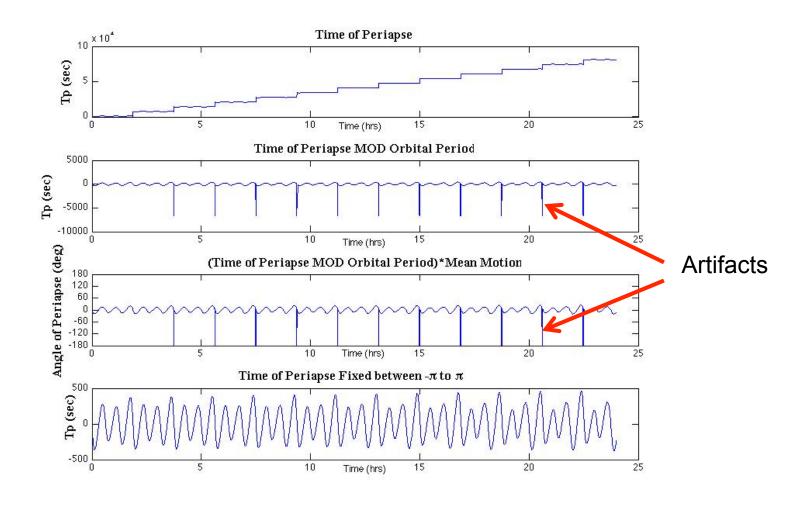
$$-\pi < \theta_p \le \pi$$

Thus, constraining  $\theta_p$  to be between  $-\pi$  to  $\pi$  will remove the artifacts. The angle of periapse  $\theta_p$  can then be converted back to time of periapse  $T_p$  by

$$T_p = \frac{\theta_p}{n}$$



# T<sub>p</sub> Results







```
t(i) = epoch time
  M(i) = E - ecc(i)*sin(E);
  n(i) = sqrt(mu/a(i)^3);
   Determine Tp from T-Tp = M/n
  Tp(i) = t(i) - M(i)/n(i);
% Mod the Epoch Time with the Period
  Tp3(i) = mod(t(i),2*pi/n(i)) - M(i)/n(i);
  Turn that time into an Angle of Periapse
  Tp4(i) = Tp3(i)*n(i);
  Tp5(i) = Tp4(i);
   Brandon's Code
  Tp2(i) = mod(t(i),(2*pi)/n(i))*n(i) - (E - ecc(i)*sin(E));
```

```
% Make sure the angle is between negative pi
to pi
  if Tp2(i) < -pi
     Tp2(i) = Tp2(i) + 2*pi;
     Tp5(i) = Tp5(i) + 2*pi;
  else if Tp2(i) > pi
     Tp2(i) = Tp2(i) - 2*pi;
     Tp5(i) = Tp5(i) - 2*pi;
  end
   Divide angle of Periapse by Mean Motion to
get Time of Periapse
  Tp2(i) = Tp2(i)/n(i);
```

Tp5(i) = Tp5(i)/n(i);





```
mean mot = sqrt(MU./(semi.*semi.*semi));
88 -
89 -
     time peri angle = mod(Tout,(2*pi)./mean mot).*mean mot ...
90
                       - (eccentric - ecc.*sin(eccentric));
91
92 -
     I = find( time peri angle < -pi );</pre>
93 -
     time peri angle(I) = time peri angle(I) + 2*pi;
94
     I = find( time peri angle > pi );
95 -
     time peri angle(I) = time peri angle(I) - 2*pi;
96 -
97
98 –
     time peri = time peri angle./mean mot;
```

#### Linearization of the OD Process

General case, the governing relations involve the non-linear expressions:

$$\dot{\mathbf{X}} = F(\mathbf{X}, t), \qquad \mathbf{X}(t_k) \equiv \mathbf{X}_k$$
 (4.2.1)

$$\mathbf{Y}_i = G(\mathbf{X}_i, t_i) + \boldsymbol{\epsilon}_i; \qquad i = 1, \dots, \ell$$
 (4.2.2)

 $oldsymbol{X_k}$  = the unknown *n*-dimensional state vector at time  $t_k$ 

 $\boldsymbol{Y_i}$  = for i = 1, ..., l is a p-dimensional set of **observations** 

 $\widehat{X}_k$  = best estimate of the unknown value of  $X_k$ 

In general p < n and  $m = p \times l \gg n$ 

#### Linearization of the OD Process

Formulation of:

$$\dot{\mathbf{X}} = F(\mathbf{X}, t), \qquad \mathbf{X}(t_k) \equiv \mathbf{X}_k$$
 (4.2.1)

$$\mathbf{Y}_i = G(\mathbf{X}_i, t_i) + \boldsymbol{\epsilon}_i; \qquad i = 1, \dots, \ell$$
 (4.2.2)

- (1) The inability to observe the state directly
- (2) Non-linear relations between the observations and state
- (3) Fewer observations at any time epoch that there are state vector components p < n
- (4) Errors in the observations represented by  $\epsilon_i$





Replace the **nonlinear** orbit determination problem to estimate the **state vector** with a **linear** orbit determination problem to determine the **deviation** from some reference solution

$$\mathbf{x}(t) = \mathbf{X}(t) - \mathbf{X}^*(t), \quad \mathbf{y}(t) = \mathbf{Y}(t) - \mathbf{Y}^*(t)$$
 (4.2.3)

And thus

$$\dot{\mathbf{x}}(t) = \dot{\mathbf{X}}(t) - \dot{\mathbf{X}}^*(t). \tag{4.2.4}$$



EJ!

Expanding Eqs. (4.2.1) and (4.2.2) in a Taylor series about the reference trajectory yields:

$$\dot{\mathbf{X}}(t) = F(\mathbf{X}, t) = F(\mathbf{X}^*, t) + \left[\frac{\partial F(t)}{\partial \mathbf{X}(t)}\right]^* \left[\mathbf{X}(t) - \mathbf{X}^*(t)\right] + O_F\left[\mathbf{X}(t) - \mathbf{X}^*(t)\right]$$

$$(4.2.5)$$

$$\mathbf{Y}_{i} = G(\mathbf{X}_{i}, t_{i}) + \boldsymbol{\epsilon}_{i} = G(\mathbf{X}_{i}^{*}, t_{i}) + \left[\frac{\partial G}{\partial \mathbf{X}}\right]_{i}^{*} \left[\mathbf{X}(t_{i}) - \mathbf{X}^{*}(t_{i})\right]_{i} + O_{G}\left[\mathbf{X}(t_{i}) - \mathbf{X}^{*}(t_{i})\right] + \boldsymbol{\epsilon}_{i}$$

# Linearization of the OD Process



Neglect higher order terms in Eq. (4.2.5) and apply the following conditions:

$$\dot{X}^* = F(X^*, t) \quad Y_i^* = G(X_i^*, t)$$

Rewriting Eq. (4.2.5):

$$\dot{\mathbf{x}}(t) = A(t)\mathbf{x}(t)$$

$$\mathbf{y}_i = \widetilde{H}_i\mathbf{x}_i + \boldsymbol{\epsilon}_i \qquad (i = 1, \dots, \ell)$$
(4.2.6)

Where:

$$A(t) = \left[\frac{\partial F(t)}{\partial \mathbf{X}(t)}\right]^* \qquad \widetilde{H}_i = \left[\frac{\partial G}{\partial \mathbf{X}}\right]_i^*$$

# Linearization of the OD Process



Hence, the original **non-linear** estimation problem is replaced by the **linear** estimation problem described by Eq. (4.2.6)

$$\mathbf{x}(t) = \mathbf{X}(t) - \mathbf{X}^*(t),$$

$$\mathbf{x}_i = \mathbf{X}(t_i) - \mathbf{X}^*(t_i)$$

$$\mathbf{y}_i = \mathbf{Y}_i - G(\mathbf{X}_i^*, t_i).$$

#### Generally,

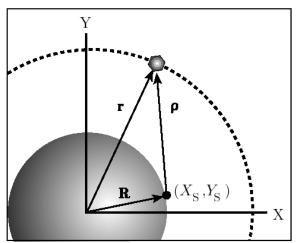
- uppercase X and Y will represent the state and observation vectors
- lowercase **x** and **y** will represent the state and observation deviation vectors



#### Example 4.2.1

Compute the A matrix and the  $\widetilde{H}$  matrix for a satellite in a plane under the influence of only a *central force*. Assume that the satellite is being tracked with range observations,  $\rho$ , from a single ground station. Assume that the station coordinates,  $(X_{\rm S}, Y_{\rm S})$ , and the gravitational parameter are unknown. Then, the state vector,  $\mathbf{X}$ , is given by

$$\mathbf{X} = \left[ egin{array}{c} X \ Y \ U \ V \ \mu \ X_{\mathrm{S}} \ Y_{\mathrm{S}} \end{array} 
ight]$$



where U and V are velocity components and  $X_{\rm S}$  and  $Y_{\rm S}$  are coordinates of the tracking station. From Newton's Second Law and the law of gravitation,

$$\ddot{\mathbf{r}} = -\frac{\mu \mathbf{r}}{r^3}$$



#### Or in component form:

$$\ddot{X} = -\frac{\mu X}{r^3}$$

$$\ddot{Y} = -\frac{\mu Y}{r^3}$$

#### Expressed in first order form:

$$\dot{\mathbf{m}} : \begin{bmatrix} \dot{X} \\ \dot{Y} \\ \dot{Y} \end{bmatrix} = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix} = \begin{bmatrix} U \\ V \\ -\frac{\mu X}{r^3} \\ -\frac{\mu Y}{r^3} \end{bmatrix}$$

$$\dot{\mathbf{X}} = \begin{bmatrix} \dot{V} \\ \dot{V} \\ \dot{V} \end{bmatrix} = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \end{bmatrix} = \begin{bmatrix} U \\ -\frac{\mu X}{r^3} \\ -\frac{\mu Y}{r^3} \\ 0 \end{bmatrix}$$

$$\dot{\mathbf{X}}_{\mathrm{S}} \begin{bmatrix} \dot{\mu} \\ \dot{X}_{\mathrm{S}} \\ \dot{Y}_{\mathrm{S}} \end{bmatrix} = \begin{bmatrix} F_1 \\ F_2 \\ F_4 \end{bmatrix} = \begin{bmatrix} U \\ -\frac{\mu X}{r^3} \\ 0 \\ 0 \end{bmatrix}$$

	$ \frac{\partial F_1}{\partial X} \\ \partial F_2 $	$\frac{\partial F_1}{\partial Y}$	$\frac{\partial F_1}{\partial U}$	$\frac{\partial F_1}{\partial V}$	$\frac{\partial F_1}{\partial \mu}$	$\frac{\partial F_1}{\partial X_{\mathrm{S}}}$	$\frac{\partial F_1}{\partial Y_{\mathrm{S}}}$
	$\frac{\partial F_2}{\partial X}$						$\frac{\partial F_2}{\partial Y_{\rm S}}$
	:	÷	÷	÷	÷	÷	÷
$(t) = \frac{\partial F(\mathbf{X}^*, t)}{\partial \mathbf{X}} =$	:	÷	÷	÷	÷	÷	:
	:	÷	÷	÷	:	÷	÷
	:	÷	÷	÷	÷	÷	÷
	$\frac{\partial F_7}{\partial X}$						$\frac{\partial F_7}{\partial Y_{ m S}}$ .

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	Г						-	*
	0	0	1	0	0	0	0	
	0	0	0	1	0	0	0	
	$-\frac{\mu}{r^3} + \frac{3\mu X^2}{r^5}$	$\frac{3\mu XY}{r^5}$	0	0	$-\frac{X}{r^3}$	0	0	
=	$\frac{3\mu XY}{r^5}$	$-\frac{\mu}{r^3} + \frac{3\mu Y^2}{r^5}$	0	0	$-\frac{Y}{r^3}$	0	0	
	0	0	0	0	0	0	0	
	0	0	0	0	0	0	0	
	0	0	0	0	0	0	0	



The  $\widetilde{H}$  matrix is given by

$$\widetilde{H} = \frac{\partial \rho}{\partial \mathbf{X}} = \begin{bmatrix} \frac{\partial \rho}{\partial X} & \frac{\partial \rho}{\partial Y} & \frac{\partial \rho}{\partial U} & \frac{\partial \rho}{\partial V} & \frac{\partial \rho}{\partial \mu} & \frac{\partial \rho}{\partial X_{\mathrm{S}}} & \frac{\partial \rho}{\partial Y_{\mathrm{S}}} \end{bmatrix}^*$$

where

$$\rho = \left[ (X - X_{\rm S})^2 + (Y - Y_{\rm S})^2 \right]^{1/2}.$$

It follows then that

$$\widetilde{H} = \left[ \begin{array}{cccc} X - X_{\mathrm{S}} & Y - Y_{\mathrm{S}} \\ \hline \rho & \rho \end{array} \right. \quad 0 \quad 0 \quad 0 \quad -\frac{(X - X_{\mathrm{S}})}{\rho} \quad -\frac{(Y - Y_{\mathrm{S}})}{\rho} \quad \right]^*$$



#### **State Transition Matrix**



#### **State Transition Matrix**

$$\dot{\mathbf{x}}(t) = A(t)\mathbf{x}(t)$$

$$\mathbf{y}_{i} = \widetilde{H}_{i}\mathbf{x}_{i} + \boldsymbol{\epsilon}_{i} \qquad (i = 1, \dots, \ell)$$

$$A(t) = \left[\frac{\partial F(t)}{\partial \mathbf{X}(t)}\right]^{*} \qquad \widetilde{H}_{i} = \left[\frac{\partial G}{\partial \mathbf{X}}\right]_{i}^{*}$$

$$(4.2.6)$$

The first of Eq. (4.2.6) represents a system of linear differential equations with time-dependent coefficients. The symbol  $[]^*$  indicates that the values of  $\mathbf{X}$  are derived from a particular solution to the equations  $\dot{\mathbf{X}} = F(\mathbf{X}, t)$  which is generated with the initial conditions  $\mathbf{X}(t_0) = \mathbf{X}_0^*$ . The general solution for this system,  $\dot{\mathbf{x}}(t) = A(t)\mathbf{x}(t)$ , can be expressed as

$$\mathbf{x}(t) = \Phi(t, t_k)\mathbf{x}_k \tag{4.2.7}$$

where  $\mathbf{x}_k$  is the value of  $\mathbf{x}$  at  $t_k$ ; that is,  $\mathbf{x}_k = \mathbf{x}(t_k)$ . The matrix  $\Phi(t_i, t_k)$  is called the state transition matrix and was introduced in Chapter 1, Section 1.2.5.



#### **State Transition Matrix**

1. 
$$\Phi(t_k, t_k) = I$$

2. 
$$\Phi(t_i, t_k) = \Phi(t_i, t_j)\Phi(t_j, t_k)$$
 (4.2.8)

3. 
$$\Phi(t_i, t_k) = \Phi^{-1}(t_k, t_i)$$
.

$$\dot{\Phi}(t, t_k) = A(t)\Phi(t, t_k) \tag{4.2.10}$$

with initial conditions

$$\Phi(t_k, t_k) = I.$$



Given the system: 
$$\dot{x} = ax + by$$
  
 $\dot{y} = ky$ 

Find the STM for:

$$\mathbf{X} = \begin{bmatrix} x \\ y \end{bmatrix}$$

Equations are linear in the dependent variables (x,y) and their derivatives (in space state form)



Equations are linear in the dependent variables (x,y) and their derivatives (in space state form)

$$\dot{X} = AX = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} a & b \\ 0 & k \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

Whose solutions is

$$\mathbf{X}(t) = \Phi(t, t_0) \mathbf{X}_0$$

Where

$$\dot{\Phi}(t,t_0) = A\Phi(t,t_0)$$



Since the differential equations for  $\Phi(t, t_0)$  are linear with constant coefficients we can solve them using Laplace Transforms:

$$\Phi(t,t_0) = \mathcal{L}^{-1}(SI - A)$$

$$SI - A = \begin{bmatrix} s - a & -b \\ 0 & s - b \end{bmatrix} \qquad (SI - A)^{-1} = \begin{bmatrix} \frac{1}{s - a} & \frac{b}{(s - a)(s - b)} \\ 0 & \frac{1}{s - k} \end{bmatrix}$$



If **A** is a constant matrix, there are a number of ways (including Laplace Transforms) to solve the equation

$$\dot{\Phi}(t,t_0) = A\Phi(t,t_0)$$

For example, we may integrate the equations directly:

$$\dot{\Phi}(t,t_0) = \begin{bmatrix} \dot{\phi}_{11} & \dot{\phi}_{12} \\ \dot{\phi}_{21} & \dot{\phi}_{22} \end{bmatrix} = \begin{bmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{bmatrix} \begin{bmatrix} \phi_{11} & \phi_{12} \\ \phi_{21} & \phi_{22} \end{bmatrix}$$

$$= \begin{bmatrix} A_{11}\phi_{11} + A_{12}\phi_{21} & A_{11}\phi_{12} + A_{12}\phi_{22} \\ A_{21}\phi_{11} + A_{22}\phi_{21} & A_{21}\phi_{12} + A_{22}\phi_{22} \end{bmatrix}$$



Note that the columns of  $\dot{\Phi}$  are independent.

Hence, if A = A(t) and we must use numerical integration we could integrate the rows of  $\dot{\Phi}$  independently as n systems of  $n \times 1$  equations as opposed to an  $n \times n$  system of simultaneous equations.



Evaluating the equations for  $\dot{\Phi}$  yields:

$$\dot{\phi}_{11} = a\phi_{11} + b\phi_{21} \tag{1}$$

$$\dot{\phi}_{21} = k\phi_{21} \tag{2}$$

$$\dot{\phi}_{12} = a\phi_{12} + b\phi_{22} \qquad (3)$$

$$\dot{\phi}_{22} = k\phi_{22} \tag{4}$$

with I.C. 
$$\Phi(t, t_0) = I$$

Note that the columns of  $\dot{\Phi}$  are independent i.e., Eqns. (1) and (2) are independent of (3) and (4)



#### Solutions

From Eq. (2) 
$$\frac{d\phi_{21}}{\phi_{21}} = kdt \longrightarrow \ln \phi_{21} = kt + \ln C$$

$$\frac{\phi_{21}}{C} = e^{kt} \longrightarrow \phi_{21} = Ce^{kt}$$

$$\therefore C = 0$$
 and  $\Phi_{21} = 0$ 



Likewise, from Eq. (4)

$$\phi_{22} = Ce^{kt}$$

$$\therefore C = 1$$
 and  $\Phi_{22} = e^{kt}$ 

Thus Eq. (1) becomes

$$\dot{\phi}_{11} = a\phi_{11}$$
 and  $\phi_{11} = e^{at}$ 



Finally, from Eq. (3),

$$\dot{\phi}_{12} = a\phi_{12} + be^{kt}$$
 (5)

The homogeneous equation

$$\dot{\phi}_{12} = a\phi_{12}$$

has the solution

$$\phi_{12} = ce^{at}$$

To get a particular solution note that  $e^{kt}$  has the derivative  $ke^{kt}$ 



So try

$$\phi_{12p} = C_1 e^{kt}$$

substitute the particular equation into Eq. (5)

$$C_1 k e^{kt} = a C_1 e^{kt} + b e^{kt}$$

Then

$$C_1k - aC_1 - b = 0 \longrightarrow C_1 = \frac{b}{k - a}$$



The general solution is the sum of homogeneous and particular solutions

$$\phi_{12} = Ce^{at} + \frac{b}{k-a}e^{kt}$$

$$\therefore C = -\frac{b}{k-a}$$

Hence

$$\phi_{21} = -\frac{b}{k-a}e^{at} + \frac{b}{k-a}e^{kt} = \frac{b}{a-k}(e^{at} - e^{kt})$$

and

$$\Phi(t,t_0) = \begin{bmatrix} e^{at} & \frac{b}{a-k} \left( e^{at} - e^{kt} \right) \\ 0 & e^{kt} \end{bmatrix}$$
Note that @ t=0
$$\Phi(t,t_0) = I$$

$$\Phi(t,t_0)=I$$





#### Using the following definition:

$$\mathbf{y} \equiv \begin{bmatrix} y_1 \\ \vdots \\ y_\ell \end{bmatrix}; \quad H \equiv \begin{bmatrix} \widetilde{H}_1 \Phi(t_1, t_k) \\ \vdots \\ \widetilde{H}_\ell \Phi(t_\ell, t_k) \end{bmatrix}; \quad \boldsymbol{\epsilon} \equiv \begin{bmatrix} \epsilon_1 \\ \vdots \\ \epsilon_\ell \end{bmatrix}$$
(4.2.38)

and if the subscript on  $\mathbf{x}_k$  is dropped for convenience, then Eq. (4.2.37) can be expressed as follows:

$$y = Hx + \epsilon \tag{4.2.39}$$



$$\mathbf{y} = H\mathbf{x} + \boldsymbol{\epsilon} \tag{4.2.39}$$

y is an  $m \times 1$  vector x is an  $n \times 1$  vector

 $\epsilon$  is an  $m \times 1$  vector

H is an  $m \times n$  mapping matrix

 $m = p \times l$  is the total number of observations m > n is an essential condition

Have m unknown observation errors



$$\mathbf{y} = H\mathbf{x} + \boldsymbol{\epsilon} \tag{4.2.39}$$

#### Results in:

m unknown observation errors

m+n total unknowns

*m* equations

The least squares criterion provides us with conditions on the m observation errors that allow a solution for the n state variables,  $X_k$ , at the epoch time  $t_k$ 



#### Review of Variables