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**Lindblad**

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Projekt

nature12016

Sander Stammbach  
Prof. Patrick Plotts

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# 1 Introduction

The most important question in thermodynamics is, how to convert thermal energy into work. For such tasks it exist many classical engines, as example the steam-machine. To quantify heat-engines, its commen to look at the ergotropy. In my master-project I will quantify the a three level maser. The three-level maser is a Quantum heat engine (short: QHE). The work extraction from a classicle heat engine is a moving piston. But in this case it is a driving field.. Experimentally we need two different reservoir. The high-temperature reservoir can be realized by a gas noise lamp and the filter by a wave guide cutting off the lower frequencies.

## 2 Theory

### 2.1 Model

The Basic of our system is a three-level quantum system. This three level system is driven by a hot bath and a cold bath. Those have the temperature  $T_c$  and  $T_h$ . This quantum system lives in a cavity. In my calculation, the thermal bath is constant, so i can use the Lindblad-masterequation. The efficiency is given by the formula:

$$\eta_M = \frac{\nu_s}{\nu_p} \tag{1}$$

The master equation is:

$$\dot{\rho}(t) = \frac{1}{i\hbar}[H, \rho] + \mathcal{L}\rho \quad (2)$$

The first part is the vonNeuman equation. The von Neumanequation is the analog of the Schrödinger-equation but for density matrices. This part of the equation is unitary and therefore the process is reversible. The non-unitary part of the equation  $\mathcal{L}$  have three parts.  $\mathcal{L}_h$  describe the interaction with the hot bath.  $\mathcal{L}_c$  Describe the interaction with the cold bath.  $\mathcal{L}_{cav}$  Describe the Photons which leaves the cavity. so if we have a small  $\kappa$  means less photons will leave and stay in the cavity. we see that in the Fockplott.

The interaction Hamiltonian or Jaynes-Cummings Hamiltonian is:

$$H_{int} = \hbar g(\sigma_{12}a^\dagger + \sigma_{21}a) \quad (3)$$

The Hamiltonian of the photons is, which describes the phonons in the cavity:

$$H_{free} = \sum_{i=1}^3 \hbar\omega_i |i\rangle \langle i| + \hbar\omega_f a^\dagger a \quad (4)$$

The total Hamiltonian

$$H = H_{free} + H_{int} \quad (5)$$

The interaction with the various environmental heat baths is described by the Liouvillian:

$$\begin{aligned} \mathcal{L}\hat{\rho} = & \frac{\gamma_h}{2} \left[ \frac{1}{\exp[\frac{\hbar\omega_h}{k_b T_h}] - 1} + 1 \right] \cdot \left( 2\sigma_{13} \cdot \rho \cdot \sigma_{13}^\dagger - \sigma_{13}^\dagger \sigma_{13} \rho - \rho \sigma_{13}^\dagger \sigma_{13} \right) \\ & + \frac{\gamma_h}{2} \left[ \frac{1}{\exp[\frac{\hbar\omega_h}{k_b T_H}] - 1} \right] \cdot \left( 2\sigma_{31} \cdot \rho \cdot \sigma_{31}^\dagger - \sigma_{31}^\dagger \sigma_{31} \rho - \rho \sigma_{31}^\dagger \sigma_{31} \right) \\ & + \frac{\gamma_c}{2} \left[ \frac{1}{\exp[\frac{\hbar\omega_c}{k_b T_c}] - 1} + 1 \right] \cdot \left( 2\sigma_{23} \cdot \rho \cdot \sigma_{23}^\dagger - \sigma_{23}^\dagger \sigma_{23} \rho - \rho \sigma_{23}^\dagger \sigma_{23} \right) \\ & + \frac{\gamma_c}{2} \left[ \frac{1}{\exp[\frac{\hbar\omega_c}{k_b T_c}] - 1} \right] \cdot \left( 2\sigma_{32} \cdot \rho \cdot \sigma_{32}^\dagger - \sigma_{32}^\dagger \sigma_{32} \rho - \rho \sigma_{32}^\dagger \sigma_{32} \right) \\ & + \kappa \left[ \frac{1}{\exp[\frac{\hbar\omega_f}{k_b T_f}] - 1} + 1 \right] \cdot \left( 2a\rho a^\dagger - a^\dagger a \rho - \rho a^\dagger a \right) \end{aligned}$$

$$\kappa \left[ \frac{1}{\exp[\frac{\hbar \omega_f}{k_b T_f}] - 1} \right] \cdot \left( 2a^\dagger \rho a - a a^\dagger \rho - \rho a a^\dagger \right) \quad (6)$$

When we work with density matrices, its common. Since we are working with density matrices here, it is easy to work with expectation values with  $\langle A \rangle = \text{Tr}[A * \rho]$ .

### 3 Calculation

In our case only  $\omega_f$  interact with the light.

#### 3.1 Software

For the hole implementation of the tree-level-system in a cavity I used qutip. Qutip is library in python, which allows to solve Masterequation pretty fast.