

Divergence of Electrostatic Field (\vec{E})

In order to obtain the value of divergence of electrostatic field (E) that is $\nabla \cdot \vec{E}$, we first calculate the volume integral of the $\nabla \cdot \vec{E}$ which is $\int_V (\nabla \cdot \vec{E}) dV$. Now by applying divergence theorem to the above volume integral we get

$$\boxed{\int_V (\nabla \cdot \vec{E}) dV = \oint_S \vec{E} \cdot d\vec{s}}$$

If we consider a point charge q is inside a sphere of radius r , then the electric flux will cross through the surface of the sphere. Therefore $\oint \vec{E} \cdot d\vec{s}$ represents the flux of electrostatic field \vec{E} through the surface S . In case of spherically symmetric surface,

$$\oint_S \vec{E} \cdot d\vec{s} = \int \frac{1}{4\pi\epsilon_0} \left(\frac{q}{r^2} \hat{r} \right) \cdot (r^2 \sin\theta d\theta d\phi \hat{r})$$

where $d\vec{s} = r^2 \sin\theta d\theta d\phi \hat{r}$

$$\therefore \oint \vec{E} \cdot d\vec{s} = \frac{q}{4\pi\epsilon_0} \int_0^\pi \sin\theta d\theta \int_0^{2\pi} d\phi = \frac{q}{4\pi\epsilon_0} \times 2 \times 2\pi = \frac{q}{\epsilon_0}$$

(12)

For any closed surface we will take the enclosed charge as Q_{enc} . Therefore, the previous relation we can write

as
$$\oint_S \vec{E} \cdot d\vec{s} = \frac{Q_{enc}}{\epsilon_0}$$

The above relation is called Gauss's law in electrostatics in differential form. integral form.

Again
$$\oint_S \vec{E} \cdot d\vec{s} = \int_V (\nabla \cdot \vec{E}) dv = \frac{Q_{enc}}{\epsilon_0}$$

If the volume charge density is ρ then we can write $Q_{enc} = \int_V \rho dv$

Therefore, from the above two relations we get
$$\int_V (\nabla \cdot \vec{E}) dv = \frac{1}{\epsilon_0} \int_V \rho dv$$

From removing volume integral from both sides we get

$$\boxed{\nabla \cdot \vec{E} = \rho / \epsilon_0}$$

This equation is called Gauss's law in electrostatics in differential form.

Application: Find the electric field outside a uniformly charged solid sphere of radius 'R' and total charge 'q'.

Soln: we have to ~~cal~~

calculate the field at point r distance from the sphere.

where $r > R$. Now

draw a Gaussian surface

around the sphere of radius 'r'. By

applying Gauss's law we get

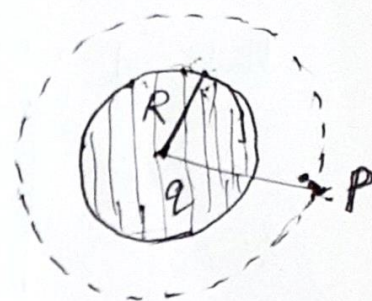
$$\oint_S \vec{E} \cdot d\vec{s} = \frac{Q_{enc}}{\epsilon_0}$$

Here $Q_{enc} = q$

$$\therefore \int_S |\vec{E}| ds = \frac{q}{\epsilon_0} = E \int_S ds = E 4\pi r^2$$

$$\text{Thus, } E 4\pi r^2 = \frac{q}{\epsilon_0}$$

$$\therefore, \vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$



Application: Find the electric field out of a plane carries a uniform surface charge 'σ'.

Soln: In order to obtain the electric field out of the surface, draw a "Gaussian pillbox" extending equal distance above and below

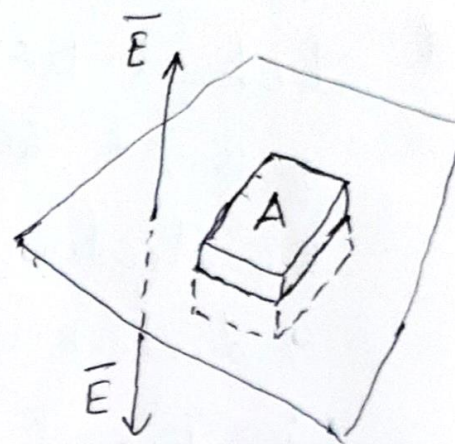
the plane. By applying the (13)
Gauss's law to the surface we get

$$\oint_S \vec{E} \cdot d\vec{s} = \frac{Q_{enc}}{\epsilon_0}$$

$$\therefore \oint |\vec{E}| ds = 2AE$$

In this case $Q_{enc} = \sigma A$,

A is the area of the pillbox at both sides.



\therefore From the above two equations we get

$$2A|\vec{E}| = \frac{\sigma A}{\epsilon_0}$$

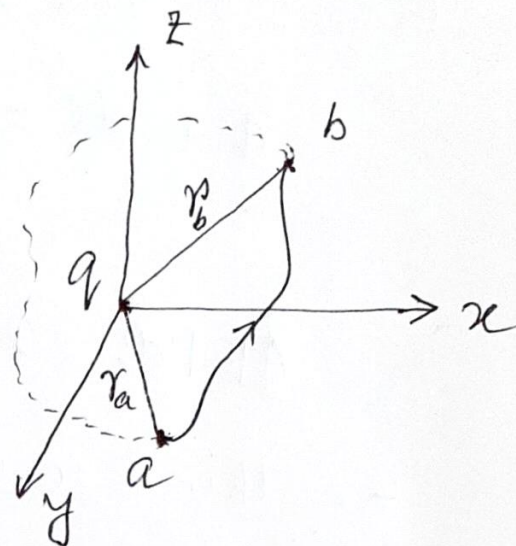
$$\therefore \boxed{|\vec{E}| = \frac{\sigma}{2\epsilon_0}} \text{ or, } \boxed{\vec{E} = \frac{\sigma}{2\epsilon_0} \hat{n}}$$

Where \hat{n} is a unit vector pointing away from the surface.

Curl of \vec{E} ($\nabla \times \vec{E}$)

Let a point charge q is at the origin. Hence, the electric field at a point r distance from origin

$$\text{is } \vec{E} = \frac{1}{4\pi\epsilon_0} \times \frac{q}{r^2} \hat{r}$$



We now calculate line integral of this field from some point 'a' to 'b'

Hence, $\int_a^b \vec{E} \cdot d\vec{l}$ required to obtain

In spherical polar coordinate system

$$d\vec{l} = \hat{r} dr + r d\theta \hat{\theta} + r \sin \theta d\phi \hat{\phi} \quad \left| \begin{array}{l} \hat{r} \cdot \hat{r} = 1 \\ \hat{r} \cdot \hat{\theta} = 0 \\ \hat{r} \cdot \hat{\phi} = 0 \end{array} \right.$$

$$\text{Therefore, } \vec{E} \cdot d\vec{l} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr$$

$$\text{Hence, } \int_a^b \vec{E} \cdot d\vec{l} = \frac{1}{4\pi\epsilon_0} \int_a^b \frac{q}{r^2} dr = \frac{-1}{4\pi\epsilon_0} \frac{q}{r} \Big|_{r_a}^{r_b}$$

$$= \frac{1}{4\pi\epsilon_0} \left(\frac{q}{r_a} - \frac{q}{r_b} \right).$$

where r_a and r_b are distance from the origin to point a and b respectively.

(14)

Here the important thing is that the line integral is plainly independent of path. All it depends on is the two end points, in fact, all that matters is how far a and b from the charge. You may have question why this happened?

Because, \vec{E} points in the radial direction. It costs nothing to move around θ & ϕ directions. Therefore, any contribution from such a displacement is wiped out by the dot product $\vec{E} \cdot d\vec{l}$. The integral around a closed path is obviously zero.

$$r_a = r_b$$

Hence, $\oint \vec{E} \cdot d\vec{l} = 0$ and by applying

Stokes's theorem we get $\boxed{\nabla \times \vec{E} = 0}$ As

$$\text{As } \oint (\nabla \times \vec{E}) \cdot d\vec{s} = 0$$

As $\nabla \times \vec{E} = 0$, the \vec{E} vector is called conservative or irrotational vector.