

# Radiative Corrections

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For further info see:

- M. Vanderhaeghen et al., (2000) Phys. Rev. C 62 (2000) 025501.
- D. Lhuillier, These Univ. Caen, DAPNIA-SPhN-97-01T (1997)
- D. Marchand, These Univ. Blaise Pascal, DAPNIA-SphN-98-04T (1998)
- P.Janssens et al, Nucl. Instr. Meth. A566 (2006) 675-686.

$$\sigma_{exp} = (1 + \delta_{tot})\sigma_{th}. \quad (7)$$

The correction term  $\delta_{tot}$  is negative and depends on the cut in the radiative tail accompanying the scattering process. The internal radiative corrections to VCS are discussed in great detail in [12]. Written in first order, one gets

$$\delta_{tot}^{(1)} = \delta_{vac} + \delta_{ver} + \delta_{rad}, \quad (8)$$

$\delta_{vac}$  accounts for vacuum polarisation diagrams,  $\delta_{ver}$  is the vertex correction and  $\delta_{rad}$  is the correction for radiation in the one additional photon approximation. One can approximately take into account higher order radiative corrections by writing [12]:

$$\sigma_{exp} = \frac{e^{\delta_{ver} + \delta_{rad}}}{(1 - \delta_{vac}/2)^2} \sigma_{th}. \quad (9)$$

For  $Q^2 \gg m^2$ , one can write:

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Functions of :

- $Q^2$
- Beam Energy
- $e^-$  out momentum

$\Delta E_{cm}^c$  depends on MM cut

$$\delta_{rad} \approx \frac{\alpha}{\pi} \left\{ \ln \left( \frac{(\Delta E_{cm}^c)^2}{E_{cm} E'_{cm}} \right) \left[ \ln \left( \frac{Q^2}{m^2} \right) - 1 \right] - \frac{1}{2} \ln^2 \left( \frac{E_{cm}}{E'_{cm}} \right) + \frac{1}{2} \ln^2 \left( \frac{Q^2}{m^2} \right) - \frac{\pi^2}{3} + Sp \left( \cos^2 \frac{\theta_{e',cm}}{2} \right) \right\}, \quad (10)$$

$$\delta_{ver} \approx \frac{\alpha}{\pi} \left\{ \boxed{\frac{3}{2}} \ln \left( \frac{Q^2}{m^2} \right) - 2 - \frac{1}{2} \ln^2 \left( \frac{Q^2}{m^2} \right) + \frac{\pi^2}{6} \right\}, \quad (11)$$

$$\delta_{vac} \approx \frac{2\alpha}{3\pi} \left\{ -\frac{5}{3} + \ln \left( \frac{Q^2}{m^2} \right) \right\}, \quad (12)$$

# Cut dependent variable

been used in the last line ( $E_e'^{el}$  denotes the elastic scattered electron *lab* energy, to distinguish it from  $E_e'$ ). From Eq. (A48), one determines then  $\Delta E_s$  in terms of *lab* quantities from the scattered electron spectrum through

$$\Delta E_s = \eta (E_e'^{el} - E_e') , \quad (\text{A49})$$

where the recoil factor  $\eta$  is given by  $\eta = E_e / E_e'^{el}$ .

$$E_{e'}^{\text{elastic}} = \frac{E_e}{1 + \frac{E_e}{M_{\text{Tg}}(1-\cos \theta)}} .$$

$$\Delta E_{\text{cm}}^c = \frac{\sqrt{M_X^2}}{2} .$$

$\Delta E_s$  determined by resolution of the detector,  
Catches radiative tail corrections,  
Determined by cuts (shown for  $E'$  in elastic scattering,  $MM^2$  for our exp)

# QED radiative corrections

Values of radiative corrections for the first VCS experiment at MAMI.

MAMI VCS kinematics	
Kinematics	$q_{cm} = 0.6 \text{ GeV/c}$ , $\epsilon = 0.61$ , $q'_{cm} = 111 \text{ MeV/c}$ $E = 0.843 \text{ GeV}$ , $E' = 0.528 \text{ GeV}$ , $\theta_e = 53.14^\circ$ $Q^2 = 0.356 \text{ GeV}^2$
$T_{anal}$	+22.0 %
$\delta_{cont}$	-1.2 % (*)
$\delta_V$	-15.6 %
$F_{rad}$	( +94.9 ± 2.0 ) %

(\*) = the value of  $(\delta_1 + \delta_2^{(0)})$  averaged over the two MAMI kinematics of Table 1 in ref. [4].

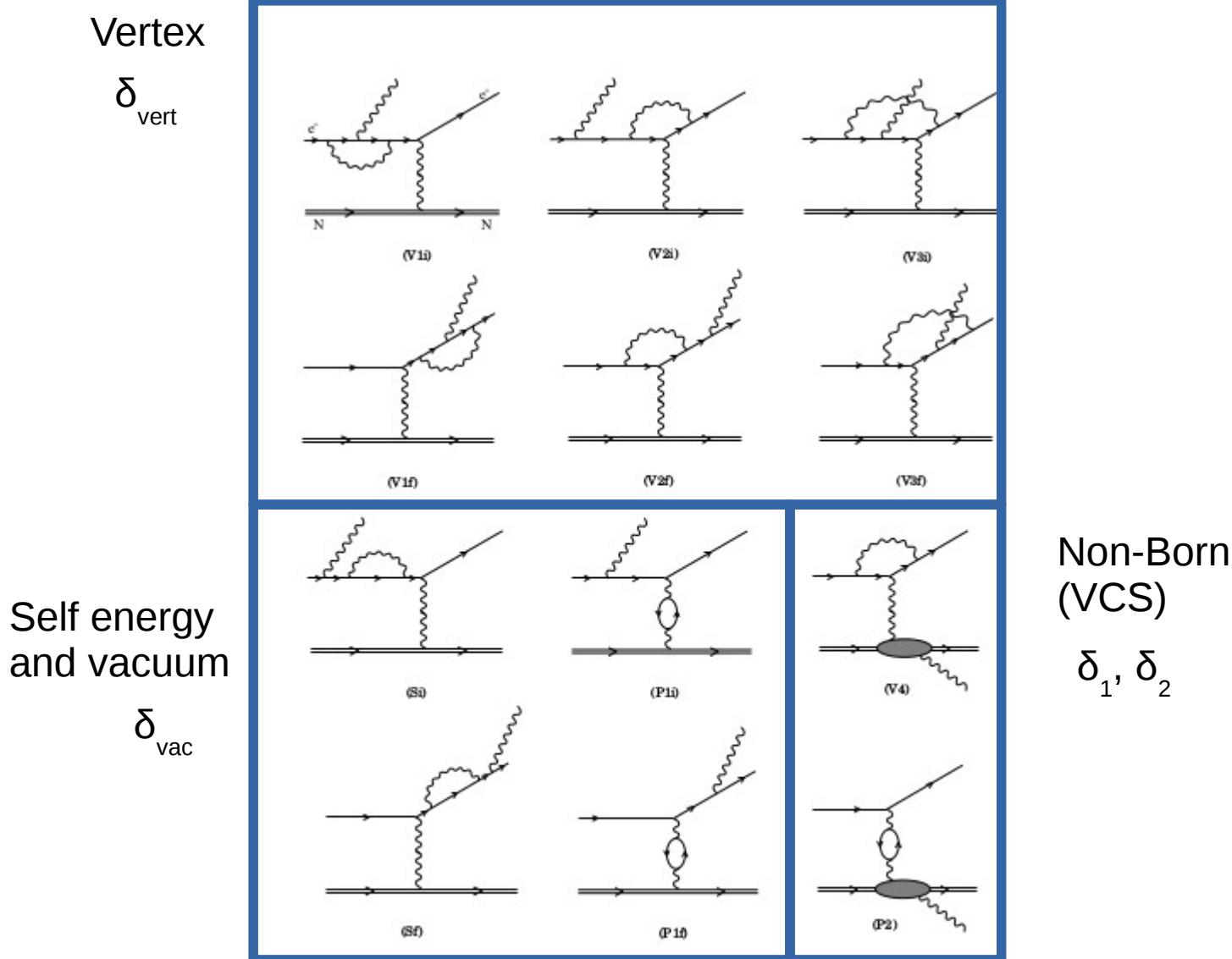
$$\delta_V = \delta_{\text{vertex}} + \delta_{\text{vacuum}} + \delta_{\text{self energy}}$$

A partir de la formule (2.96) du chapitre 2,  $\delta_R$  s'écrit sous la forme:

$$\delta_R = T_{anal} + f(\Delta E_{cut}) \quad (3.55)$$

$T_{anal}$  est un terme analytique fonction uniquement de la cinématique et  $f(\Delta E_{cut})$  est dépendant de la coupure expérimentale  $\Delta E_{cut}$ .

# Internal Corrections



# External Corrections

# Soft photon, Radiative tail

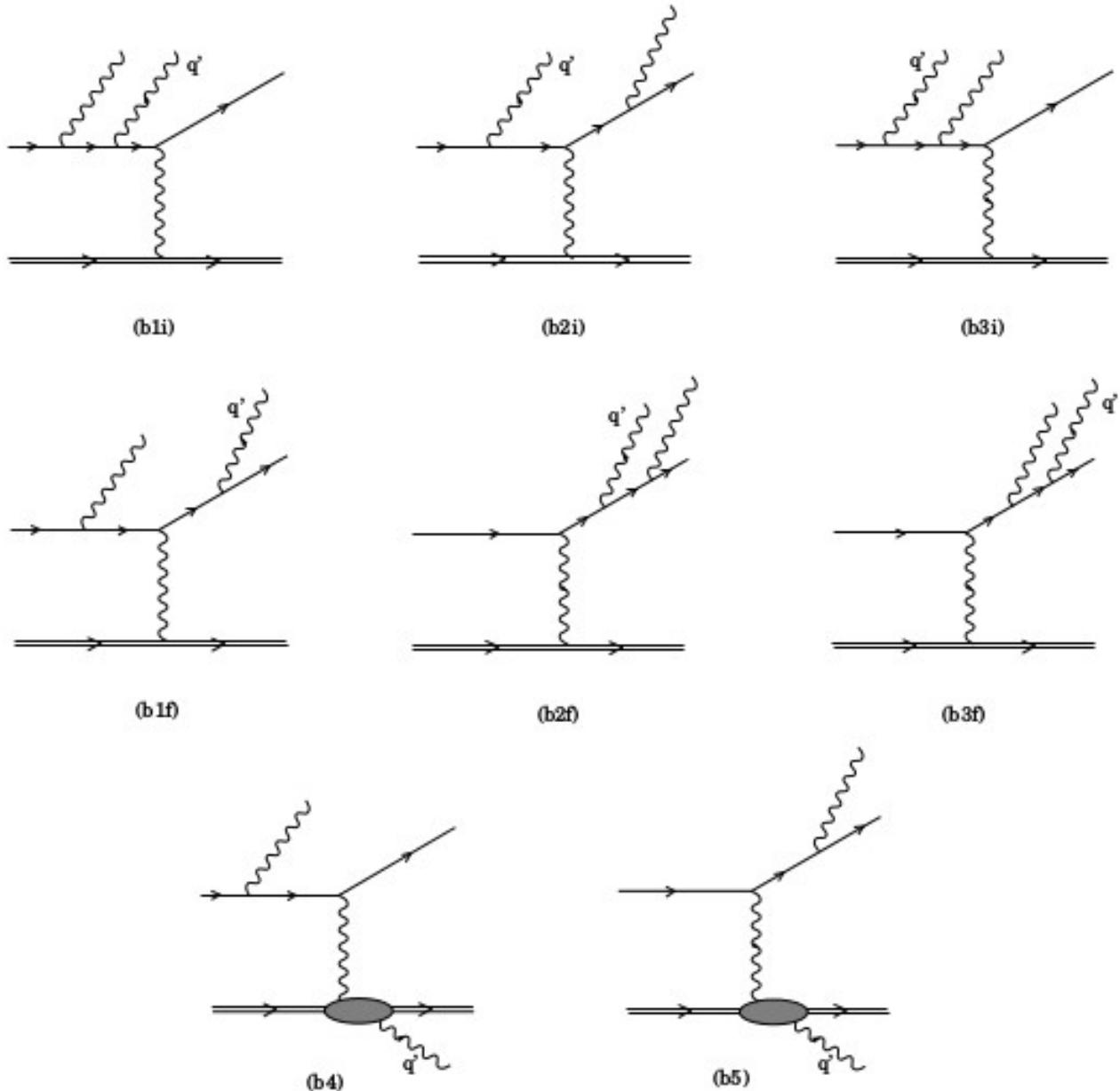


FIG. 3. First order soft photon emission contributions to the  $ep \rightarrow e\gamma\gamma$  reaction.

# Hadron Corrections

$$\left( \frac{d\sigma}{d\Omega'_e} \right)_{TOTAL} = \left( \frac{d\sigma}{d\Omega'_e} \right)_{BORN} \left( 1 + \delta_{vac} + \delta_{vertex} + \delta_R + Z \delta_1 + Z^2 (\delta_2^{(0)} + \delta_2^{(1)}) \right), \quad (A74)$$

where  $\delta_{vac}$ ,  $\delta_{vertex}$  and  $\delta_R$  are given as above (Eq.(A72)). The terms in Eq. (A74) proportional to  $Z$  (hadron charge) and  $Z^2$  contain the corrections from the proton side. The correction  $\delta_1$ , proportional to  $Z$ , was calculated in Ref. [24] as

$$\delta_1 = \frac{2\alpha_{em}}{\pi} \left\{ \ln \left( \frac{4(\Delta E_s)^2}{Q^2 x} \right) \ln \eta + Sp \left( 1 - \frac{\eta}{x} \right) - Sp \left( 1 - \frac{1}{\eta x} \right) \right\}, \quad (A75)$$

where  $\Delta E_s$  and  $\eta$  are given as in Eq. (A49) and where the variable  $x$  is defined by

$$\delta_{cont} = \delta_1 + \delta_2 \quad x = \frac{(Q + \rho)^2}{4M_N^2}, \quad \rho^2 = Q^2 + 4M_N^2, \quad (A76)$$

The correction proportional to  $Z^2$  was split into two parts in Ref. [24]. The contribution  $\delta_2^{(0)}$ , independent of the nucleon form factors was calculated in Ref. [24] as :

$$\begin{aligned} \delta_2^{(0)} &= \frac{\alpha_{em}}{\pi} \left\{ \ln \left( \frac{4(\Delta E_s)^2}{M_N^2} \right) \left( \frac{E'_N}{|\vec{p}'_N|} \ln x - 1 \right) + 1 \right. \\ &\quad \left. + \frac{E'_N}{|\vec{p}'_N|} \left( -\frac{1}{2} \ln^2 x - \ln x \ln \left( \frac{\rho^2}{M_N^2} \right) + \ln x - Sp \left( 1 - \frac{1}{x^2} \right) + 2Sp \left( -\frac{1}{x} \right) + \frac{\pi^2}{6} \right) \right\}, \quad (A77) \end{aligned}$$

where  $\rho$  is defined as in Eq. (A76), and where  $E'_N$  ( $|\vec{p}'_N|$ ) are the *lab* energy (momentum) of the recoiling nucleon. For the lengthier expression of  $\delta_2^{(1)}$ , which depends on the nucleon form factors, we refer to Ref. [24].

# Cross Checks\*

Correction	Helene's Paper	My calculation
$T_{\text{anal}}$	22.0	21.8
$f(\Delta E_{\text{cut}})$ (acct in tail)	-17.3	-16.2
$\Delta_R$ (sum above)	4.7	5.6
$\delta_v$	-15.6	-15.9
$\delta_{\text{cont}}$	-1.2	-1.1

Vanderhaeghen, et al

$E_e$	$\theta_e$	$Q^2$	$\delta_{\text{vertex}}$	$\delta_{\text{vacpol}}$	$\delta_R$	$\delta_1$	$\delta_2^{(0)}$	$\delta_{\text{tot}}$	EXP
0.705	40.66	0.203	-0.1673	0.0208	-0.0453	-0.0067	-0.0018	-0.2003	-0.2025
0.855	52.18	0.418	-0.1881	0.0228	-0.0245	-0.0123	-0.0034	-0.2054	-0.2087
0.705	40.66	0.203	-0.1672	0.0184	-0.0475	-0.0063	-0.0018	-0.1989	-0.1953
0.855	52.18	0.418	-0.1881	0.0195	-0.0273	-0.0118	-0.0035	-0.2111	-0.2059

\* Helene and Marc mention simulations. It is not clear if any values are averages across the phase space

# Simul++ RadCorr

Simul++ has  $T_{\text{anal}}$ , and  $\delta_v$  built in to both PS and Yield\_Sim (cancels?)

Hadron effects not considered ~ 1% effect, can split like  $T_{\text{anal}}$  &  $f(\Delta E)$   
 Exponentiation differences from Marc's paper

$$\left( \frac{d\sigma}{d\Omega'_e} \right)_{VIRTUAL\gamma} + \left( \frac{d\sigma}{d\Omega'_e} \right)_{REALSOFT\gamma} = \left( \frac{d\sigma}{d\Omega'_e} \right)_{BORN} (1 + \delta_{vac} + \delta_{vertex} + \delta_R) , \quad (\text{A71})$$

where  $\delta_{vac}$ ,  $\delta_{vertex}$  and  $\delta_R$  are given by Eqs. (A69), (A68), and (A64)-(A66) respectively.  
 Bringing the three contributions together, leads to the expression (in the  $Q^2 \gg m^2$  limit)

$$\begin{aligned} \delta_{vac} + \delta_{vertex} + \delta_R = & \frac{\alpha_{em}}{\pi} \left\{ \ln \left( \frac{(\Delta E_s)^2}{E_e E'_e} \right) \left[ \ln \left( \frac{Q^2}{m^2} \right) - 1 \right] \right. \\ & \left. + \frac{13}{6} \ln \left( \frac{Q^2}{m^2} \right) - \frac{28}{9} - \frac{1}{2} \ln^2 \left( \frac{E_e}{E'_e} \right) - \frac{\pi^2}{6} + Sp \left( \cos^2 \frac{\theta_e}{2} \right) \right\}, \quad (\text{A72}) \end{aligned}$$

where  $\Delta E_s$ , which is the maximum soft-photon energy in the c.m. system of (recoiling proton + soft-photon), is determined as in Eq. (A49), when applying this formula to the

<sup>3</sup>Note that an incorrect expression is used in [50] for the vacuum polarization contribution due to  $\mu^+ \mu^-$  pairs (Eq. (A5) in their paper).

# Exponentiation

$$\delta_{tot}^{(1)} = \delta_{vac} + \delta_{ver} + \delta_{rad}, \quad (8)$$

$\delta_{vac}$  accounts for vacuum polarisation diagrams,  $\delta_{ver}$  is the vertex correction and  $\delta_{rad}$  is the correction for radiation in the one additional photon approximation. One can approximately take into account higher order radiative corrections by writing [12]:

$$\sigma_{exp} = \frac{e^{\delta_{ver} + \delta_{rad}}}{(1 - \delta_{vac}/2)^2} \sigma_{th}. \quad (9)$$

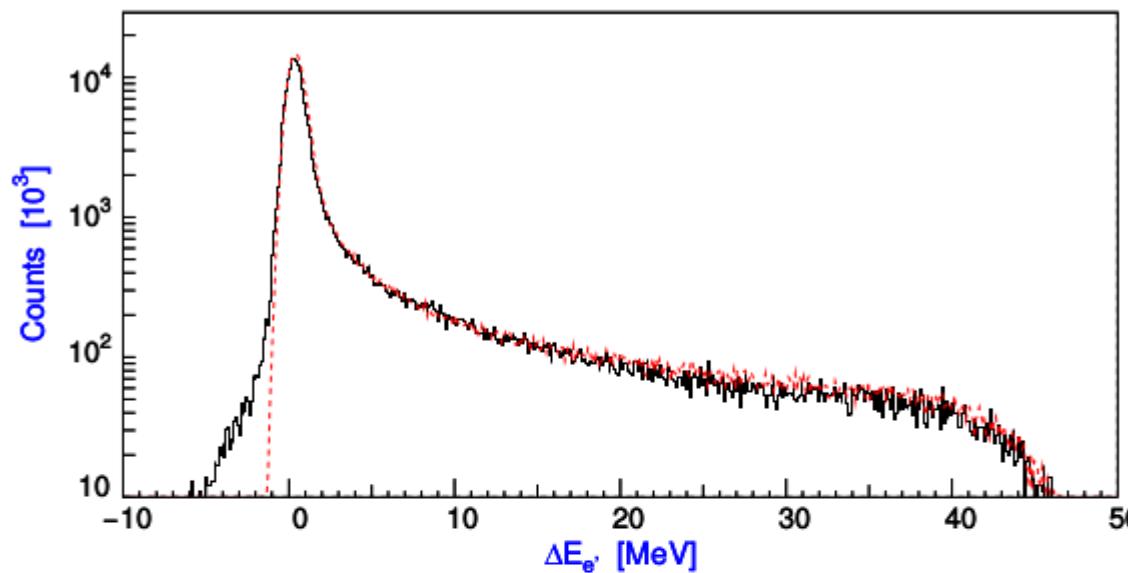
Simul++ applies  $\sigma_{exp} = e^{\delta_{tot}} \sigma_{th}$

$(1-x/2)^{-2}$	$e^x$	Difference
$1 + x + \frac{3x^2}{4} + \frac{x^3}{2} + \frac{5x^4}{16} + \frac{3x^5}{16} + O(x^6)$ (Taylor series) (converges when $ x  < 2$ )	$1 + x + \frac{x^2}{2} + \frac{x^3}{6} + \frac{x^4}{24} + \frac{x^5}{120} + O(x^6)$ (Taylor series) (converges everywhere)	$x^2/4 + O(x^3)$

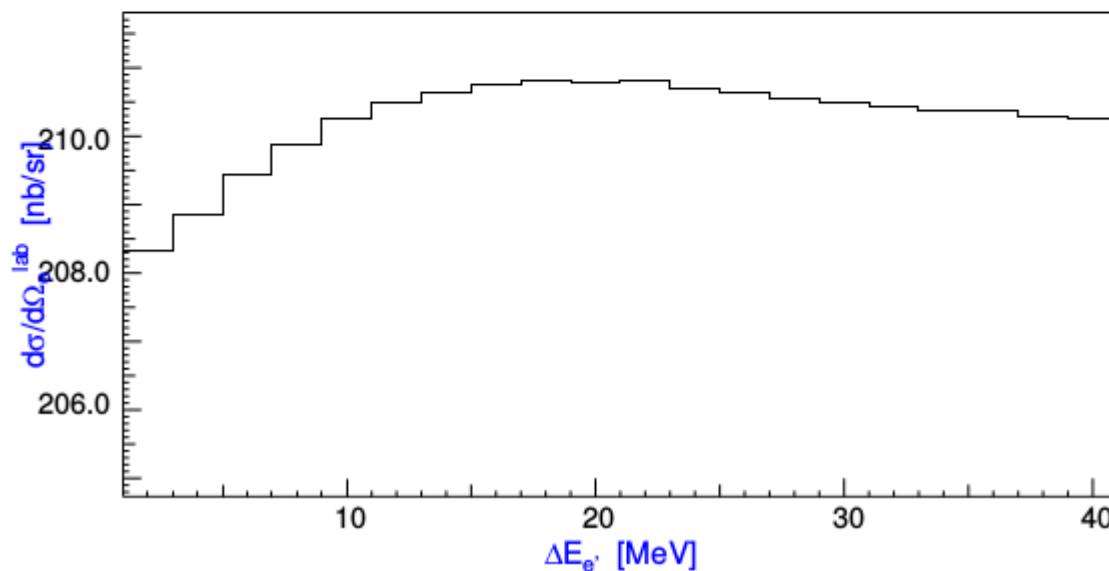
Exponentiating  $\delta_{vac} \sim 2\%$ , difference of 0.04% is negligible

$$f(\Delta E_{\text{cut}})$$

Elastic Peak



Elastic Cross Section



# Corrections for our Kinematics

Correction	KinI	KinII
Tanal *	20.6	20.6
$\Delta_v$ *	-14.8	-14.8
$\delta_{\text{cont}}$	-1.8	-1.8

\* - in Simul++

The correction factors are constant since the electron spectrometer was fixed.  
The change is below uncertainty since the proton spectrometer changed very little.

# Cross Section Results

