

Radiative Corrections

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For further info see:

- M. Vanderhaeghen et al., (2000) Phys. Rev. C 62 (2000) 025501.
- D. Lhuillier, These Univ. Caen, DAPNIA-SPhN-97-01T (1997)
- D. Marchand, These Univ. Blaise Pascal, DAPNIA-SphN-98-04T (1998)
- P.Janssens et al, Nucl. Instr. Meth. A566 (2006) 675-686.

$$\sigma_{exp} = (1 + \delta_{tot})\sigma_{th}. \quad (7)$$

The correction term δ_{tot} is negative and depends on the cut in the radiative tail accompanying the scattering process. The internal radiative corrections to VCS are discussed in great detail in [12]. Written in first order, one gets

$$\delta_{tot}^{(1)} = \delta_{vac} + \delta_{ver} + \delta_{rad}, \quad (8)$$

δ_{vac} accounts for vacuum polarisation diagrams, δ_{ver} is the vertex correction and δ_{rad} is the correction for radiation in the one additional photon approximation. One can approximately take into account higher order radiative corrections by writing [12]:

$$\sigma_{exp} = \frac{e^{\delta_{ver} + \delta_{rad}}}{(1 - \delta_{vac}/2)^2} \sigma_{th}. \quad (9)$$

For $Q^2 \gg m^2$, one can write:

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Functions of :

- Q^2
- Beam Energy
- e^- out momentum

ΔE_{cm}^c depends on MM cut

$$\delta_{rad} \approx \frac{\alpha}{\pi} \left\{ \ln \left(\frac{(\Delta E_{cm}^c)^2}{E_{cm} E'_{cm}} \right) \left[\ln \left(\frac{Q^2}{m^2} \right) - 1 \right] - \frac{1}{2} \ln^2 \left(\frac{E_{cm}}{E'_{cm}} \right) + \frac{1}{2} \ln^2 \left(\frac{Q^2}{m^2} \right) - \frac{\pi^2}{3} + Sp \left(\cos^2 \frac{\theta_{e',cm}}{2} \right) \right\}, \quad (10)$$

$$\delta_{ver} \approx \frac{\alpha}{\pi} \left\{ \frac{3}{2} \ln \left(\frac{Q^2}{m^2} \right) - 2 - \frac{1}{2} \ln^2 \left(\frac{Q^2}{m^2} \right) + \frac{\pi^2}{6} \right\}, \quad (11)$$

$$\delta_{vac} \approx \frac{2\alpha}{3\pi} \left\{ -\frac{5}{3} + \ln \left(\frac{Q^2}{m^2} \right) \right\}, \quad (12)$$

Cut dependent variable

been used in the last line (E_e^{tel} denotes the elastic scattered electron *lab* energy, to distinguish it from E_e'). From Eq. (A48), one determines then ΔE_s in terms of *lab* quantities from the scattered electron spectrum through

$$\Delta E_s = \eta (E_e^{tel} - E_e') , \quad (\text{A49})$$

where the recoil factor η is given by $\eta = E_e / E_e^{tel}$.

$$E_{e'}^{\text{elastic}} = \frac{E_e}{1 + \frac{E_e}{M_{\text{Tg}}(1 - \cos \theta)}} , \quad \Delta E_{cm}^c = \frac{\sqrt{M_X^2}}{2} .$$

ΔE_s determined by resolution of the detector,

Catches radiative tail corrections,

Determined by cuts (shown for E' in elastic scattering, MM^2 for our exp)

QED radiative corrections

Values of radiative corrections for the first VCS experiment at MAMI.

MAMI VCS kinematics	
Kinematics	$q_{cm}=0.6 \text{ GeV}/c$, $\epsilon = 0.61$, $q'_{cm} = 111 \text{ MeV}/c$ $E=0.843 \text{ GeV}$, $E'=0.528 \text{ GeV}$, $\theta_e = 53.14^\circ$ $Q^2 = 0.356 \text{ GeV}^2$
T_{anal}	+22.0 %
δ_{cont}	-1.2 % (*)
δ_V	-15.6 %
F_{rad}	(+94.9 \pm 2.0) %

(*) = the value of $(\delta_1 + \delta_2^{(0)})$ averaged over the two MAMI kinematics of Table 1 in ref. [4].

$$\delta_V = \delta_{\text{vertex}} + \delta_{\text{vacuum}} + \delta_{\text{self energy}}$$

A partir de la formule (2.96) du chapitre 2, δ_R s'écrit sous la forme :

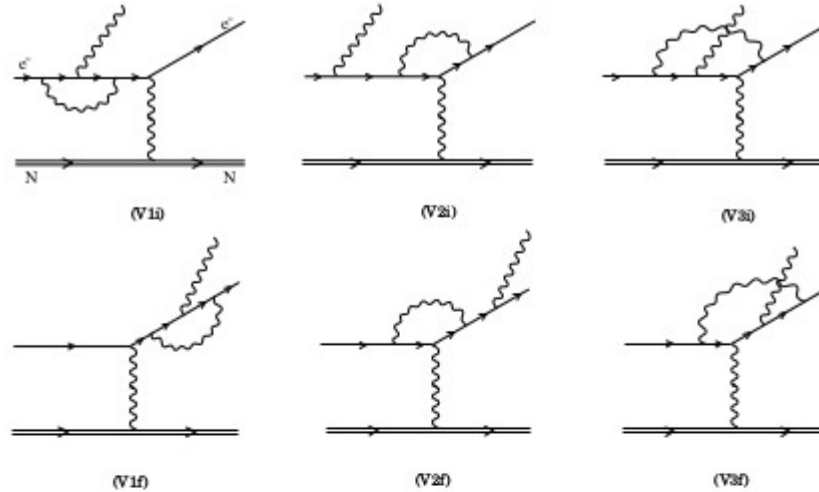
$$\delta_R = T_{anal} + f(\Delta E_{cut}) \quad (3.55)$$

T_{anal} est un terme analytique fonction uniquement de la cinématique et $f(\Delta E_{cut})$ est dépendant de la coupure expérimentale ΔE_{cut} .

Internal Corrections

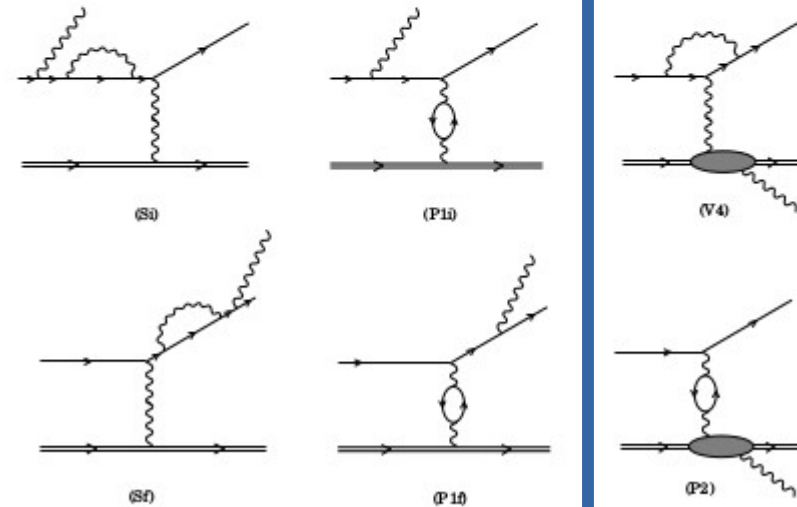
Vertex

δ_{vert}



Self energy
and vacuum

δ_{vac}



Non-Born
(VCS)

δ_1, δ_2

External Corrections

Soft photon,
Radiative tail

$$\delta_R = \delta_{\text{rad}}$$

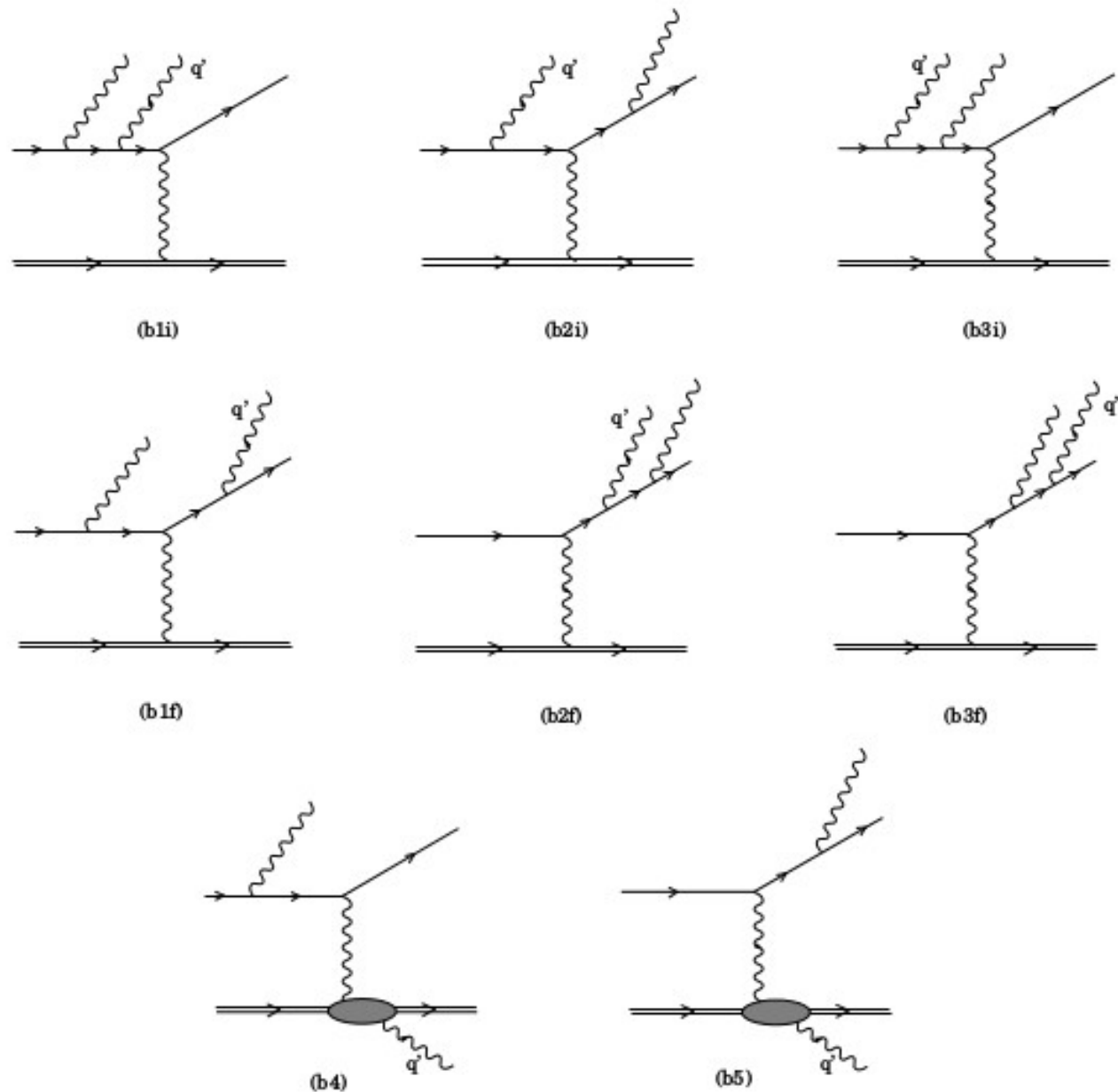


FIG. 3. First order soft photon emission contributions to the $ep \rightarrow e\gamma$ reaction.

Hadron Corrections

$$\left(\frac{d\sigma}{d\Omega'_e}\right)_{TOTAL} = \left(\frac{d\sigma}{d\Omega'_e}\right)_{BORN} \left(1 + \delta_{vac} + \delta_{vertex} + \delta_R + Z \delta_1 + Z^2 (\delta_2^{(0)} + \delta_2^{(1)})\right), \quad (A74)$$

where δ_{vac} , δ_{vertex} and δ_R are given as above (Eq. (A72)). The terms in Eq. (A74) proportional to Z (hadron charge) and Z^2 contain the corrections from the proton side. The correction δ_1 , proportional to Z , was calculated in Ref. [24] as

$$\delta_1 = \frac{2\alpha_{em}}{\pi} \left\{ \ln \left(\frac{4(\Delta E_s)^2}{Q^2 x} \right) \ln \eta + Sp \left(1 - \frac{\eta}{x} \right) - Sp \left(1 - \frac{1}{\eta x} \right) \right\}, \quad (A75)$$

where ΔE_s and η are given as in Eq. (A49) and where the variable x is defined by

$$\delta_{cont} = \delta_1 + \delta_2 \quad x = \frac{(Q + \rho)^2}{4M_N^2}, \quad \rho^2 = Q^2 + 4M_N^2, \quad (A76)$$

The correction proportional to Z^2 was split into two parts in Ref. [24]. The contribution $\delta_2^{(0)}$, independent of the nucleon form factors was calculated in Ref. [24] as :

$$\delta_2^{(0)} = \frac{\alpha_{em}}{\pi} \left\{ \ln \left(\frac{4(\Delta E_s)^2}{M_N^2} \right) \left(\frac{E'_N}{|\vec{p}'_N|} \ln x - 1 \right) + 1 \right. \\ \left. + \frac{E'_N}{|\vec{p}'_N|} \left(-\frac{1}{2} \ln^2 x - \ln x \ln \left(\frac{\rho^2}{M_N^2} \right) + \ln x - Sp \left(1 - \frac{1}{x^2} \right) + 2Sp \left(-\frac{1}{x} \right) + \frac{\pi^2}{6} \right) \right\}, \quad (A77)$$

where ρ is defined as in Eq. (A76), and where E'_N ($|\vec{p}'_N|$) are the *lab* energy (momentum) of the recoiling nucleon. For the lengthier expression of $\delta_2^{(1)}$, which depends on the nucleon form factors, we refer to Ref. [24].

Cross Checks*

Correction	Helene's Paper	My calculation
T_{anal}	22.0	21.8
$f(\Delta E_{\text{cut}})$ (acct in tail)	-17.3	-16.2
Δ_R (sum above)	4.7	5.6
δ_V	-15.6	-15.9
δ_{cont}	-1.2	-1.1

Vanderhaeghen, et al

E_e	θ_e	Q^2	δ_{vertex}	δ_{vacpol}	δ_R	δ_1	$\delta_2^{(0)}$	δ_{tot}	EXP
0.705	40.66	0.203	-0.1673	0.0208	-0.0453	-0.0067	-0.0018	-0.2003	-0.2025
0.855	52.18	0.418	-0.1881	0.0228	-0.0245	-0.0123	-0.0034	-0.2054	-0.2087
0.705	40.66	0.203	-0.1672	0.0184	-0.0475	-0.0063	-0.0018	-0.1989	-0.1953
0.855	52.18	0.418	-0.1881	0.0195	-0.0273	-0.0118	-0.0035	-0.2111	-0.2059

* Helene and Marc mention simulations. It is not clear if any values are averages across the phase space

Simul++ RadCorr

Simul++ has T_{anal} , and δ_v built in to both PS and Yield_Sim (cancels?)

Hadron effects not considered $\sim 1\%$ effect, can split like T_{anal} & $f(\Delta E)$

Exponentiation differences from Marc's paper

$$\left(\frac{d\sigma}{d\Omega'_e}\right)_{\text{VIRTUAL}\gamma} + \left(\frac{d\sigma}{d\Omega'_e}\right)_{\text{REALSOFT}\gamma} = \left(\frac{d\sigma}{d\Omega'_e}\right)_{\text{BORN}} (1 + \delta_{\text{vac}} + \delta_{\text{vertex}} + \delta_R) , \quad (\text{A71})$$

where δ_{vac} , δ_{vertex} and δ_R are given by Eqs. (A69), (A68), and (A64)-(A66) respectively. Bringing the three contributions together, leads to the expression (in the $Q^2 \gg m^2$ limit)

$$\delta_{\text{vac}} + \delta_{\text{vertex}} + \delta_R = \frac{\alpha_{em}}{\pi} \left\{ \ln \left(\frac{(\Delta E_s)^2}{E_e E'_e} \right) \left[\ln \left(\frac{Q^2}{m^2} \right) - 1 \right] + \frac{13}{6} \ln \left(\frac{Q^2}{m^2} \right) - \frac{28}{9} - \frac{1}{2} \ln^2 \left(\frac{E_e}{E'_e} \right) - \frac{\pi^2}{6} + Sp \left(\cos^2 \frac{\theta_e}{2} \right) \right\} , \quad (\text{A72})$$

where ΔE_s , which is the maximum soft-photon energy in the c.m. system of (recoiling proton + soft-photon), is determined as in Eq. (A49), when applying this formula to the

³Note that an incorrect expression is used in [50] for the vacuum polarization contribution due to $\mu^+\mu^-$ pairs (Eq. (A5) in their paper).

Exponentiation

$$\delta_{tot}^{(1)} = \delta_{vac} + \delta_{ver} + \delta_{rad}, \quad (8)$$

δ_{vac} accounts for vacuum polarisation diagrams, δ_{ver} is the vertex correction and δ_{rad} is the correction for radiation in the one additional photon approximation. One can approximately take into account higher order radiative corrections by writing [12]:

$$\sigma_{exp} = \frac{e^{\delta_{ver} + \delta_{rad}}}{(1 - \delta_{vac}/2)^2} \sigma_{th}. \quad (9)$$

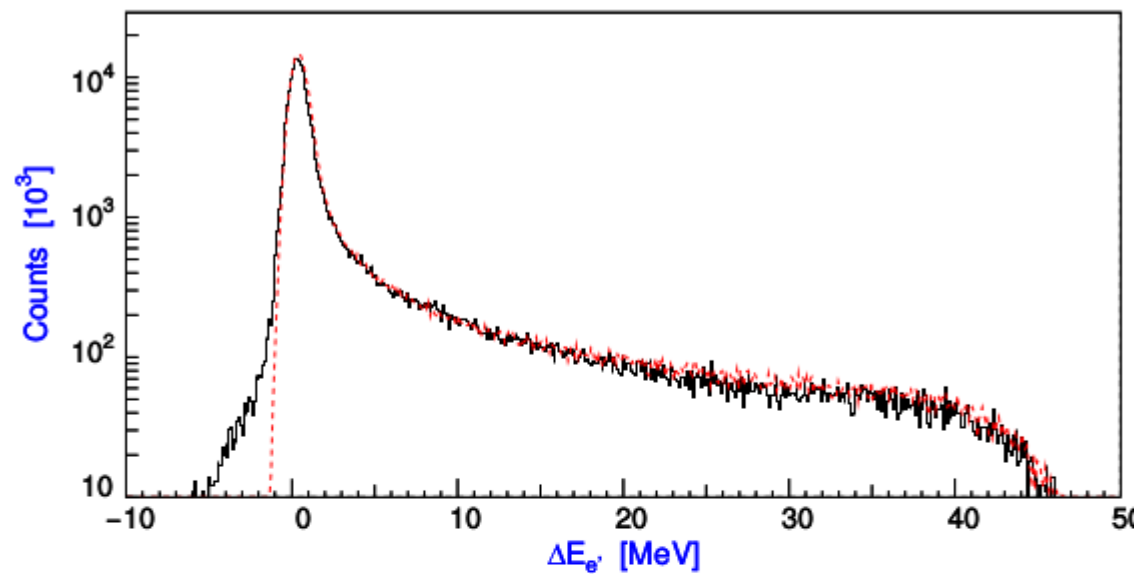
Simul++ applies $\sigma_{exp} = e^{\delta_{tot}} \sigma_{th}$

$(1-x/2)^{-2}$	e^x	Difference
$1 + x + \frac{3x^2}{4} + \frac{x^3}{2} + \frac{5x^4}{16} + \frac{3x^5}{16} + O(x^6)$ <small>(Taylor series)</small> <small>(converges when $x < 2$)</small>	$1 + x + \frac{x^2}{2} + \frac{x^3}{6} + \frac{x^4}{24} + \frac{x^5}{120} + O(x^6)$ <small>(Taylor series)</small> <small>(converges everywhere)</small>	$x^2/4 + O(x^3)$

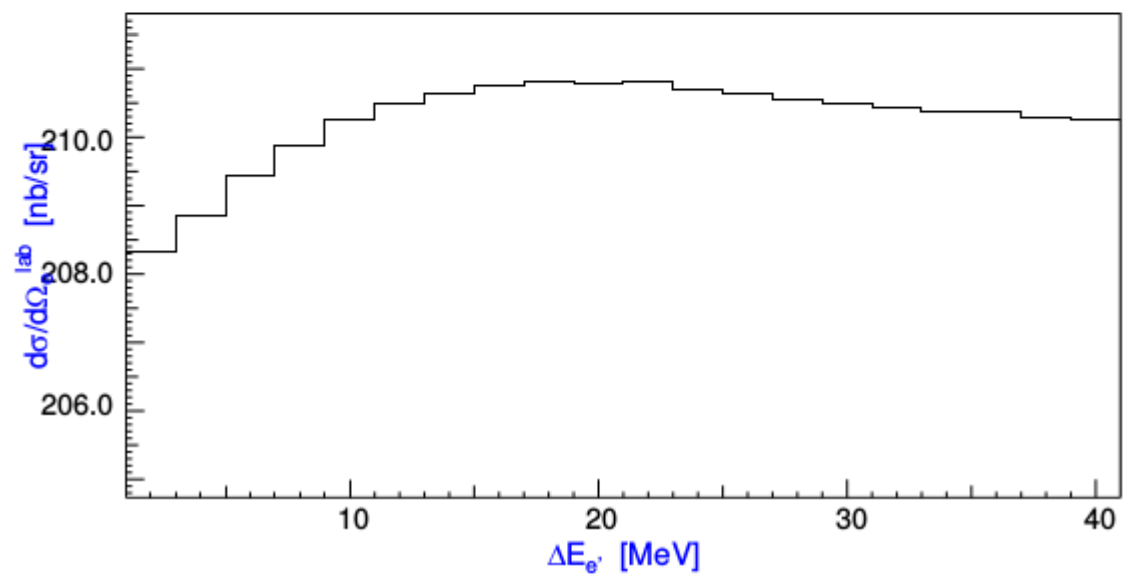
Exponentiating $\delta_{vac} \sim 2\%$, difference of 0.04% is negligible

$$f(\Delta E_{\text{cut}})$$

Elastic Peak



Elastic Cross Section



Corrections for our Kinematics

Correction	KinI	KinII
Tanal *	20.6	20.6
Δ_v *	-14.8	-14.8
δ_{cont}	-1.8	-1.8

* - in Simul++

The correction factors are constant since the electron spectrometer was fixed.
The change is below uncertainty since the proton spectrometer changed very little.

Cross Section Results

