

# 1 Principle of Fermat

Question: Proof that equation 1 and 2 can be reduced to equation 3.

$$L[x, y, z, \dot{x}, \dot{y}, \dot{z}] = n(x, y, z) \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \quad (1)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0, \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} - \frac{\partial L}{\partial y} = 0, \frac{d}{dt} \frac{\partial L}{\partial \dot{z}} - \frac{\partial L}{\partial z} = 0 \quad (2)$$

$$\frac{d}{ds} \left[ n \frac{d\vec{r}}{ds} \right] = \vec{\nabla} n \quad (3)$$

Answer:

First noting that  $ds$  is a small element of distance travelled. Therefore taking into account the  $x$ ,  $y$  and  $z$  direction,  $ds$  is given by:

$$ds = \sqrt{dx^2 + dy^2 + dz^2} \quad (4)$$

A small distance travelled in a trivial direction, lets say  $dx$ , can be approximated by as  $dx = dt \cdot \dot{x}$ . Therefore  $ds$  can be rewritten as:

$$ds = dt \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \quad (5)$$

Rewriting gives:

$$\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} = \frac{ds}{dt} \quad (6)$$

If we combine the equations in equation 2 in vector notation we get:

$$\frac{d}{dt} \left( \frac{\partial}{\partial \dot{x}} \right) L - \nabla L = 0 \quad (7)$$

Rewriting and filling in equation 1 gives:

$$\frac{d}{dt} \left[ \left( \frac{\partial}{\partial \dot{x}} \right) n(x, y, z) \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \right] = \left( \frac{\partial}{\partial x} \right) \left[ n(x, y, z) \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \right] \quad (8)$$

$$\frac{d}{dt} \frac{n \cdot \dot{\vec{r}}}{\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}} = \vec{\nabla} n \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \quad (9)$$

Using equation 6 to replace the  $\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}$  and rewriting the  $\dot{\vec{r}}$  vector gives the following:

$$\frac{d}{dt} \frac{n \cdot \dot{\vec{r}}}{\frac{ds}{dt}} = \vec{\nabla} n \frac{ds}{dt} \quad (10)$$

$$\frac{d}{dt} \frac{n \cdot \frac{d\vec{r}}{dt}}{\frac{ds}{dt}} = \vec{\nabla} n \frac{ds}{dt} \quad (11)$$

Rewriting yields the equation that was to be proved:

$$\frac{d}{ds} \left[ n \frac{d\vec{r}}{ds} \right] = \vec{\nabla} n \quad (12)$$

## 2 Application

### 2.1 Homogeneous medium

Question: Using equation 3, show how light is travelling in a homogeneous medium.

Answer:

Equation 3 can be rewritten using the chain rule:

$$\frac{d\vec{r}}{ds} \frac{d}{ds} n + n \frac{d^2 \vec{r}}{ds^2} = \left( \frac{\partial}{\partial x} \frac{\partial}{\partial y} \frac{\partial}{\partial z} \right) n \quad (13)$$

Note that for a homogeneous medium, the index of refraction,  $n$ , is constant. Therefore  $\frac{dn}{ds} = 0$ ,  $\frac{\partial n}{\partial x} = 0$ ,  $\frac{\partial n}{\partial y} = 0$  and  $\frac{\partial n}{\partial z} = 0$ . Using this in the previous equation yields:

$$n \frac{d^2 \vec{r}}{ds^2} = \vec{0} \quad (14)$$

$$\frac{d^2 \vec{r}}{ds^2} = \vec{0} \quad (15)$$

This implies that the direction and velocity of the light is not changed as the light travels through the medium. Therefore, it travels in a straight line with a constant velocity of  $v = c/n$ .

### 2.2 Snell-Descartes Law

Question: Express first geometrically and then analytically Snell's and Descartes' law of reflection and transmission of the light at the interface between two media of different index of refraction  $n_1$  and  $n_2$ , using the Principle of Fermat and equation 3.

Answer:

#### 2.2.1 Geometrical

The speed of light in a medium is inversely proportional to the refractive index. Therefore, the shortest path (in distance) between two points in materials with different refractive indices is not always the fastest (in time). This phenomenon is nicely described by a 2-dimensional analogy of a beach (see figure 1). The maximum speed on foot on beach is significantly higher than the maximum swimming speed in the water. So if somebody would need to get from a point A on the beach to a point B in the water, the direct route from A to B (dashed line in figure 1) would intuitively be slower than the path with a shorter swimming distance (solid line in figure 1).

It is possible to calculate the fastest route between point A and B. If we add the parameters  $v_1$ ,  $v_2$ ,  $\theta_1$ ,  $\theta_2$ ,  $a$ ,  $b$ ,  $c$  and  $d$  which corresponds respectively to the propagation speed on the beach, the propagation speed in the water, the angle of the path on the beach with the normal, the angle of the path in the water with the normal and distances which can be seen in figure 2.

The time it takes to travel from point A to B,  $t$ , can easily be found dividing the path on the beach and the water, respectively  $l_{beach}$  and  $l_{water}$  by the corresponding speed:

$$t = l_{beach}/v_1 + l_{water}/v_2 \quad (16)$$

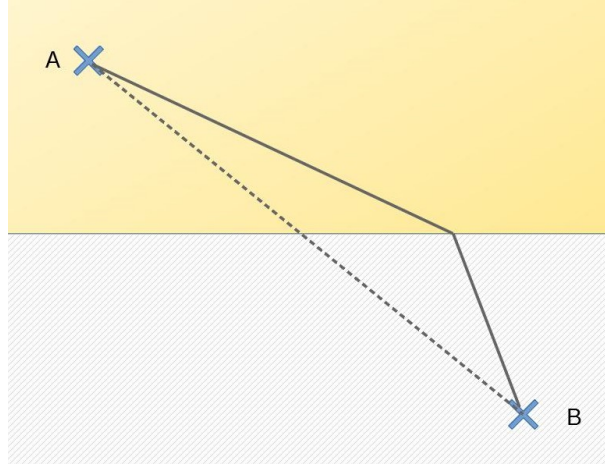


Figure 1: Diagram of a 2-dimensional beach analogy of the interface between two media with different index of refraction. The upper-half corresponds to the beach and the lower-half to the sea. The dashed line corresponds to the direct route between point A and B with the shortest distance. The solid line corresponds to a route that is intuitively faster than the direct route.

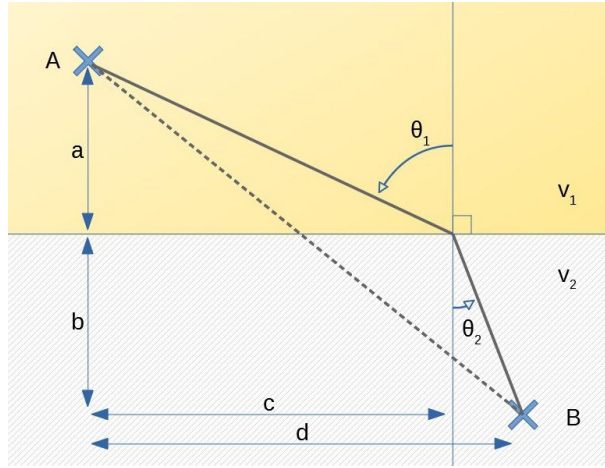


Figure 2: Diagram of beach analogy parameters  $v_1$ ,  $v_2$ ,  $\theta_1$ ,  $\theta_2$ ,  $a$ ,  $b$ ,  $c$  and  $d$ . These correspond respectively to the propagation speed on the beach, the propagation speed in the water, the angle of the path on the beach with the normal, the angle of the path in the water with the normal and distances which can be seen in the diagram.

Using the pythagoras theorem we find:

$$t = \sqrt{a^2 + c^2}/v_1 + \sqrt{b^2 + (d - c)^2}/v_2 \quad (17)$$

If there is a fastest path, there should be an optimum value for  $c$  for which  $dt/dc = 0$ . Therefore, applying the principle of Fermat to equation 17 leads to the following:

$$0 = \frac{c}{v_1 \sqrt{a^2 + c^2}} + \frac{c - d}{v_2 \sqrt{b^2 + (d - c)^2}} \quad (18)$$

Using the trigonometric identity  $\sin(\theta) = (\text{adjacentside})/(\text{diagonalside})$  for right-angled triangle we obtain:

$$0 = \sin(\theta_1)/v_1 - \sin(\theta_2)/v_2 \quad (19)$$

If we rewrite this and use the fact that the speed of light in a medium is given by  $v = c/n$  we obtain Snell-Descartes law:

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2) \quad (20)$$

This equation basically tells us, that for a interface to a higher refractive index, so where the light slows down, the light bends towards the normal.

### 2.2.2 Analytical

For the analytical derivation of Snell-Descartes law we will use a similar diagram as in the geometrical derivation with a coordinate system added as in figure 3a.

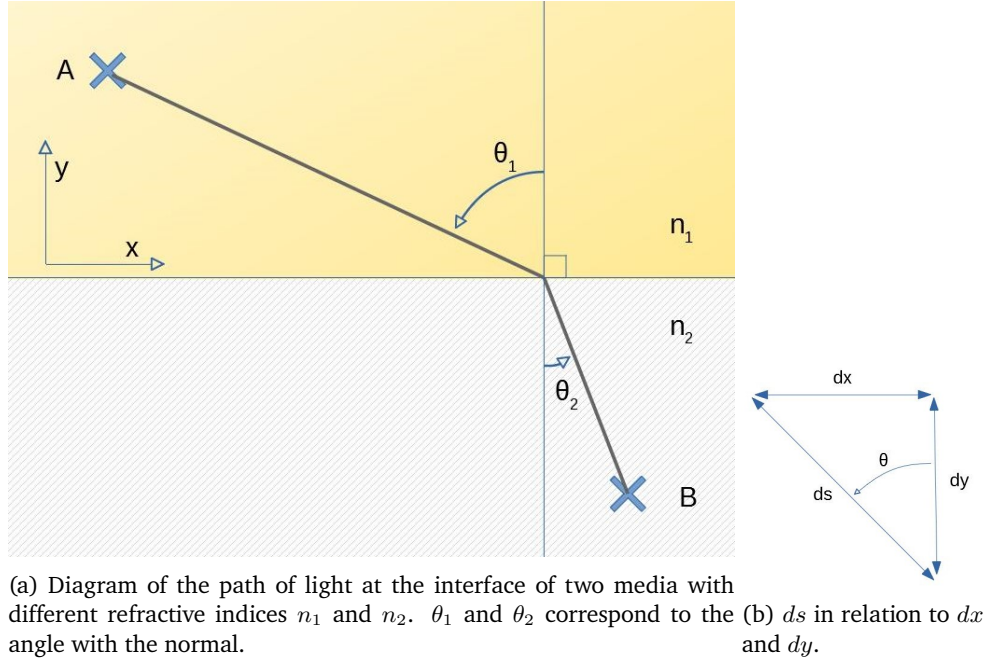


Figure 3

If we write equation 3 for only the x-component and use the fact that  $n$  is independent of  $x$  in our diagram, we get the following:

$$\frac{d}{ds} \left[ n \frac{dx}{ds} \right] = \frac{dn}{dx} \quad (21)$$

$$\frac{d}{ds} \left[ n \frac{dx}{ds} \right] = 0 \quad (22)$$

The  $ds$  in the latter equation is defined as in figure 3b. If we keep a fixed  $ds$  for both media we obtain the following equality:

$$\frac{d}{ds} \left[ n_1 \frac{dx_1}{ds} \right] = \frac{d}{ds} \left[ n_2 \frac{dx_2}{ds} \right] \quad (23)$$

Integrating both sides with respect to  $ds$  and using the trigonometric identity,  $\sin(\theta) = dx/ds$ , yields the Snell-Descartes law:

$$n_1 \frac{dx_1}{ds} = n_2 \frac{dx_2}{ds} \quad (24)$$

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2) \quad (25)$$

## 2.3 Mirage

The mirage is a common phenomenon when the ground is very warm and the temperature of the air decreases with altitude; in this case the density then increases as well as its index of refraction.

*Sketch what is happening.*

In figure 4 a situation in which a mirage might occur has been sketched. In this figure a few parameters were introduced, an  $x$ - and  $z$ -distance respectively denoting the horizontal and vertical distance from the eye of the observer and an angle  $\Theta$  which is the angle between the horizontal and the unbent light path from the observer. A gradient effect is also applied to the image, where the colour is darker the index of refraction is higher.

Since the index of refraction of the air varies with the temperature and the temperature increases as the  $z$ -coordinate increases, the path of light rays will be bent. Thus there will be a light ray coming from the sky bending in such a way that it lands in the eye of an observer. This is the mirage effect.

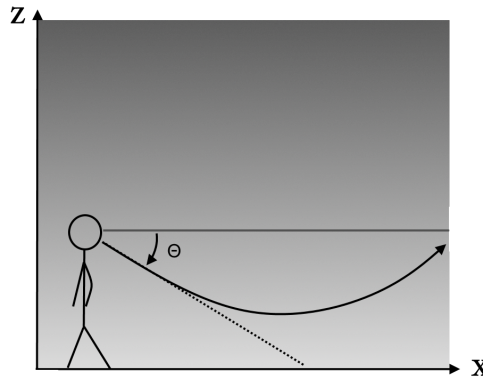


Figure 4: A sketch of situation where a mirage effect occurs.

*Find the relation between the index of refraction with the altitude if we assume that the gradient of the temperature changes linearly with the altitude.*

In equation 26 the temperature at a height  $z$  is defined using a ground temperature  $T_0$  and a lapse rate  $c$  [ $K/m$ ] so that it is linearly decreasing. To relate the index of refraction  $n$  to the temperature at height  $T(z)$ , the Gladstone-Dale relation is used. In equation 27 this relation is shown with a proportionality constant  $K_{air}$ . Finally using the equation of state shown as equation 28 it is possible to relate the index of refraction  $n$  to height  $z$ . Finding the relation for the pressure at height  $P(z)$  is shown below in the derivation and the result is displayed in equation 29. In the derivation the following notation is used,  $R_{sp,air}$  for the specific gas constant of air and  $g_0$  as the gravitational acceleration at  $z = 0$ .

$$T(z) = T_0 - c \cdot z \quad (26)$$

$$n - 1 \propto K_{air} \rho \quad (27)$$

$$\rho = \frac{1}{R_{sp,air}} \frac{P(z)}{T(z)} \quad (28)$$

$$dT/dz \equiv -c$$

$$dP = -g_0 \rho dz$$

$$\begin{aligned}
dP &= g_0 \frac{\rho}{c} dT \\
\frac{dP}{P} &= \frac{g_0}{c \cdot R_{sp,air}} \frac{dT}{T} \\
\int_{P_0}^P \frac{dP}{P} &= \int_{T_0}^T \frac{g_0}{c \cdot R_{sp,air}} \frac{dT}{T} \\
\ln P - \ln P_0 &= [\ln T - \ln T_0] \cdot \frac{g_0}{c \cdot R_{sp,air}} \\
\frac{P}{P_0} &= \left( \frac{T}{T_0} \right)^{\frac{g_0}{c \cdot R_{sp,air}}} \\
P &= P_0 \cdot \left( \frac{T}{T_0} \right)^{\frac{g_0}{c \cdot R_{sp,air}}} \tag{29}
\end{aligned}$$

After deriving equation 29 it is know possible to use both equation 26 and equation 29 to derive the air density  $\rho$  with equation 28. The result is show in equation 30, which will be combined with the Gladstone-Dale relation for a equation that equates the index of refraction  $n$  with the height  $z$ . Shown in equation 31.

$$\begin{aligned}
\rho &= \frac{P_0}{R_{sp,air} \cdot T(z)} \cdot \left( \frac{T(z)}{T_0} \right)^{\frac{g_0}{c \cdot R_{sp,air}}} \\
\rho &= \frac{P_0}{R_{sp,air}} \cdot \frac{T(z)^{\frac{g_0}{c \cdot R_{sp,air}}}}{T(z)} \cdot (T_0)^{-\frac{g_0}{c \cdot R_{sp,air}}} \\
\rho &= [T(z)]^{\frac{g_0}{c \cdot R_{sp,air}} - 1} \frac{P_0}{R_{sp,air}} \cdot T_0^{-\frac{g_0}{c \cdot R_{sp,air}}} \\
\rho &= [T_0 - c \cdot z]^{\frac{g_0}{c \cdot R_{sp,air}} - 1} \frac{P_0}{R_{sp,air}} \cdot T_0^{-\frac{g_0}{c \cdot R_{sp,air}}} \\
\rho &= \frac{P_0}{T_0 R_{sp,air}} \left[ 1 - c \cdot z T_0^{-\frac{g_0}{c \cdot R_{sp,air}}} \right] \\
\rho &= \frac{P_0}{T_0 R_{sp,air}} \left[ 1 - c \cdot z T_0^{-\frac{g_0}{c \cdot R_{sp,air}}} \right] \tag{30}
\end{aligned}$$

$$n(z) = K_{air} \rho + 1 = 1 + \frac{P_0 K_{air}}{T_0 R_{sp,air}} \left[ 1 - c \cdot z T_0^{-\frac{g_0}{c \cdot R_{sp,air}}} \right] \tag{31}$$

Express analytically the trajectory of the light in this situation.

For an analytical solution it is easier to linearise equation 31.

$$\begin{aligned}
n_l(z) &= n_0 + \alpha \cdot z \\
n_l(z) &= n(0) + z \cdot \frac{d}{dz} n(z)|_{z=0}
\end{aligned}$$

To get to an analytical solution; equation 3 needs to be solved. This is done below.

$$\vec{\nabla} \cdot n(z) = \frac{\partial}{\partial z} n(z) \vec{e}_z = \alpha \vec{e}_z$$

$$\begin{aligned}
\frac{d}{ds} \left[ n(z) \frac{d\vec{r}}{ds} \right] &= \vec{\nabla} \cdot n(z) = \alpha \vec{e}_z \\
\frac{d}{ds} \left[ n(z) \frac{d\vec{x} + d\vec{z}}{\sqrt{dx^2 + dz^2}} \right] &= \alpha \vec{e}_z \\
\frac{d}{ds} \left[ n(z) \frac{dx}{\sqrt{dx^2 + dz^2}} \right] &= 0 \\
\frac{d}{ds} \left[ n(z) \frac{dz}{\sqrt{dx^2 + dz^2}} \right] &= \alpha
\end{aligned}$$

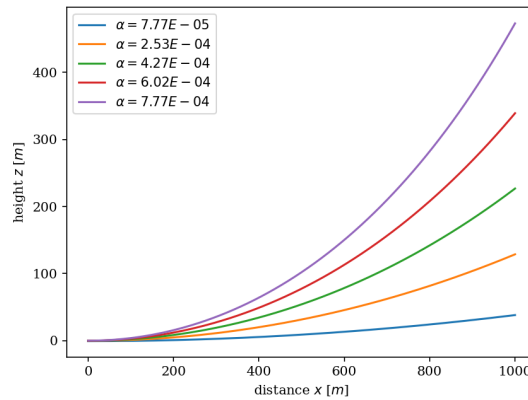
$$\begin{aligned}
n(z) \cdot \frac{dx}{\sqrt{dx^2 + dz^2}} &= c \\
n(z)^2 \cdot \frac{dx^2}{dx^2 + dz^2} &= c^2 \\
n(z)^2 \cdot dx^2 &= c^2(dx^2 + dz^2) \\
\frac{dz}{dx} &= \frac{\sqrt{n(z)^2 - c^2}}{c} = \frac{\sqrt{(n_0 + \alpha \cdot z)^2 - c^2}}{c}
\end{aligned}$$

The above first-order nonlinear ordinary differential equation has been solved using wolfram and is displayed below in 32. In this equation  $k$  is an arbitrary constant and  $c$  is the total path length in meters.  $\alpha$  [ $m^{-1}$ ] is the defined linearisation constant.

$$z(x) = -\frac{c^2 \exp(\alpha x + \alpha k)/c + \exp -(\alpha x + \alpha k)/c - 2n_0}{2\alpha} \quad (32)$$

Plot several of these trajectories for different relevant cases.

Below in figure ?? a few trajectories for different  $\alpha$ -values are plotted. These values range from  $\alpha = 7.77 \cdot 10^{-5}$  to the same constant a magnitude larger. The linearisation  $\alpha$  is constant equal to  $d/dz \cdot n(z)$  evaluated at  $z = 25[m]$ . Our initial guess for  $\alpha$  was  $\alpha = 7.77 \cdot 10^{-7}$ , but this resulted in a distance to the mirage of about half a kilometer. We felt that this does not represent an actual mirage effect and thus decided plot these alfa values. To get these values we had to rethink our lapse rate to be  $c = 7 [Km^{-1}]$ .



Find the closest distance from an observer where a mirage can be visible.

In the figure above an  $\alpha$ -value was used, this value was equal to the derivate of our index of refraction function  $n(z)$  evaluated at  $x = 0$ . This value was dependant on the ground temperature  $T_0$ , the ground pressure  $P_0$  and the lapse rate of the temperature  $c$ . For these the following values were used:  $T_0 = 323.15 K$ ,  $P_0 = 101325 Pa$ , a lapse rate of  $c = 7 K/m$  and an initial angle of looking down  $\theta = 15$  degrees. With these conditions you would see the mirage about 100 to 200 m in front of you.

Forsee what could happen on the north pole when a warm wind is present.

On the North Pole with a warm wind blowing an inverse effect as described before would be seen, since the ground is cold and the air heats up as the height increases. The index of refraction  $n$  would then decrease the higher off the ground. Light would then bend downward instead of up allowing you to see the ground when you look at the sky.

## 2.4 Optical fibre

### 2.4.1

Question: If we assume an incoming beam in the  $\vartheta_{x,y}$  plane at start crossing  $\vartheta_x$  with an angle  $\theta_0$ . Show that the light beam will stay in this plane and that, by solving the Euler-Lagrange equation 3, the equation for the

beam path can be written as:

$$n \sin(i) = a \quad (33)$$

Answer:

For this question it is assumed that the optical fibre is cylindrical with its axis  $\vartheta_x$  (see figure 5). The incident-surface will therefore be in the  $\vartheta_{y,z}$  plane.

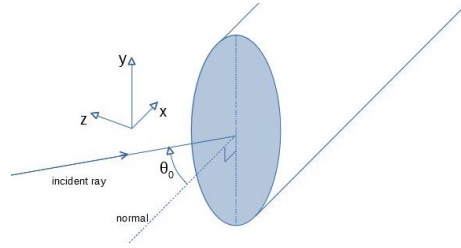


Figure 5: 3-dimensional diagram of fibre with incoming light beam making an angle  $\theta_0$  with the  $\vartheta_x$  axis.

When the incoming beam is in the  $\vartheta_{x,y}$  plane, the angle with the normal is  $\theta_0$  with  $\vartheta_x$  in the  $\vartheta_y$  direction and  $\theta_z = 0$  in the  $\vartheta_z$  direction. If we use Snell-Descartes law (equation 25) for  $\theta_z$  with  $\theta_{z,1} = 0$  we obtain  $\theta_{z,2} = k \cdot \pi [rad]$  with  $k = -1, 0, 1, 2, \dots$ . Therefore, the beam of light will stay in the  $\vartheta_{x,y}$  plane at the interface when entering the optical fibre.

It is the property of a cylinder that the outer surface is always perpendicular with the radius. For this light beam. Since the light beam is travelling in the  $\vartheta_{x,y}$  direction, the outer surface it reaches will be in the  $\vartheta_{x,z}$  plane. The angle with the surface normal in the  $\vartheta_z$  axis is still zero. From the internal reflection it follows that the angle in the  $\vartheta_z$  axis will remain zero. Therefore, for every reflection at the outer surface, the beam will stay in the  $\vartheta_{x,y}$  direction.

Since we know that the beam will stay in the  $\vartheta_{x,y}$  plane, the problem can be simplified to a 2-dimensional problem (see figure 6). The infinitely small segment  $ds$  is related to  $ds$ ,  $dy$ ,  $dx$  and  $i$  as in figure 6.

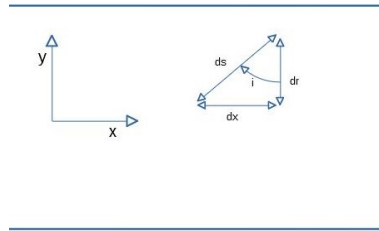


Figure 6: 2-dimensional fibre including the relation between  $ds$ ,  $ds$ ,  $dy$ ,  $dx$  and  $i$ .

We can write equataion 3 for only the x-component and use that  $n$  is independent of  $x$ :

$$\frac{d}{ds} \left[ n \frac{dx}{ds} \right] = \frac{dn}{dx} \quad (34)$$

$$\frac{d}{ds} \left[ n \frac{dx}{ds} \right] = 0 \quad (35)$$

If we now integrate with respect to  $ds$  on both sides we get the following:

$$n \frac{dx}{ds} = a \quad (36)$$



With  $a$  an arbitrary constant. Using the trigonometric identity  $\sin(i) = dx/ds$  we obtain the requested equation:

$$n \sin(i) = a \quad (37)$$

### 2.4.2

Question: Solve equation 33 for  $n(r) = n_0\sqrt{1 - \alpha^2 r^2}$  with  $\alpha < 1/r$  a constant and  $r$  the radius of the fibre. You should end up with an analytical function.

Answer:

From figure 6 and trigonometric identities it follows that  $\sin(i) = dx/ds$ . Using the Pythagoras theorem and filling this in in equation 33 we get:

$$n \frac{dx}{\sqrt{dr^2 + dx^2}} = a \quad (38)$$

Filling in  $n(r)$  and some rewriting results in the following:

$$n_0 \sqrt{1 - \alpha^2 r^2} \frac{dx}{\sqrt{dr^2 + dx^2}} = a \quad (39)$$

$$n_0^2 (1 - \alpha^2 r^2) \frac{dx^2}{dr^2 + dx^2} = a^2 \quad (40)$$

$$n_0^2 (1 - \alpha^2 r^2) dx^2 = a^2 (dr^2 + dx^2) \quad (41)$$

$$(n_0^2 - n_0^2 \alpha^2 r^2 - a^2) dx^2 = a^2 dr^2 \quad (42)$$

$$dx^2 = \frac{a^2 dr^2}{n_0^2 - n_0^2 \alpha^2 r^2 - a^2} \quad (43)$$

$$dx = \frac{a dr}{\sqrt{n_0^2 - n_0^2 \alpha^2 r^2 - a^2}} \quad (44)$$

If we now integrate both sides we will get:

$$\int dx = \int \frac{a}{\sqrt{n_0^2 - n_0^2 \alpha^2 r^2 - a^2}} dr \quad (45)$$

$$x = \frac{a}{\sqrt{n_0^2 - a^2}} \int \frac{1}{1 - \left( \frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \right)^2} dr \quad (46)$$

To solve this integral we need the following substitution:

$$\frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} = \sin(u) \quad (47)$$

$$dr = \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \cos(u) du \quad (48)$$

$$u = \arcsin \left( \frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \right) \quad (49)$$

Using the substitution we get:

$$x = \frac{a}{n_0 \alpha} \int \frac{\cos(u)}{\sqrt{1 - \sin^2(u)}} du \quad (50)$$

If we use that  $1 - \sin^2(x) = \cos^2(x)$ , we obtain:

$$x = \frac{a}{n_0 \alpha} \int \frac{\cos(u)}{\sqrt{\cos^2(u)}} du \quad (51)$$

$$x = \frac{a}{n_0 \alpha} \int du \quad (52)$$

$$x = \frac{a}{n_0 \alpha} u \quad (53)$$

Inverting the substitution and doing some rewriting yields the following equation:

$$x = \frac{a}{n_0 \alpha} \arcsin \left( \frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \right) \quad (54)$$

$$\sin \left( \frac{n_0 \alpha r}{a} \right) = \frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \quad (55)$$

$$r(x) = \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \sin \left( \frac{n_0 \alpha r}{a} \right) \quad (56)$$

The Since there are no discontinuities in  $n(r)$  and  $dn(r)/dr$ , we would expect that there are also no discontinuities in  $r(x)$  and  $dr(x)/dx$ . Since  $r = |y|$  for the light beam in the  $\vartheta_{x,y}$  plane. The only possible solution of  $y(x)$  is the following:

$$y(x) = \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \sin \left( \frac{n_0 \alpha r}{a} \right) \quad (57)$$

*Plot: a numerical solution of the problem.* To plot a numerical solution we use equation 44 and rewrite it to get a fraction  $dr/dx$ . We will then plot this function using the python script below. There is one problem however, since we used equation 44 we can only get the part of the solution where  $(dr/dx)^2$  was positive.

```
#Importin the used libraries
import numpy as np
from numpy.lib import scimath
import math
import matplotlib.pyplot as plt

#Defining a few constants
n0 = 1.7
R = 1
alpha = 1/(R)
a = 1
N = 50 #steps

x = np.linspace(0, 2*R, N)
dx = x[1] - x[0]

def index(r):
    return n0**2 * (1 - (alpha*r)**2)

def dr_squared(r):
    return (n0**2 * (1 - (alpha*r)**2) - a**2)/a**2

def analytical(x):
    return np.sqrt(n0**2 - a**2)/(n0*alpha)*np.sin((n0*alpha*x)/a)

r = np.empty(N) + 0j
increments = np.empty(N) + 0j
```

```

r[0] = 0
for i in range(1,N):
    increment = dr_squared(r[i-1])
    print(np.sqrt(increment))
    increments[i] = increment
    r[i] = r[i-1] + np.real(np.sqrt(increment))*dx - np.imag(np.sqrt(increment))*dx

plt.plot(x, r, marker=".", label="positive_solution")
plt.xlabel(r"distance_along_center_[$R$]")
plt.ylabel(r"distance_from_center_[$R$]")
plt.plot(x, analytical(x), label="analytical_solution")
plt.legend()
plt.show()

```

Using this code the following plot for the numerical solution was made 7. Where the horizontal and vertical axis are normalised to  $R$ , the radius of the fiber.

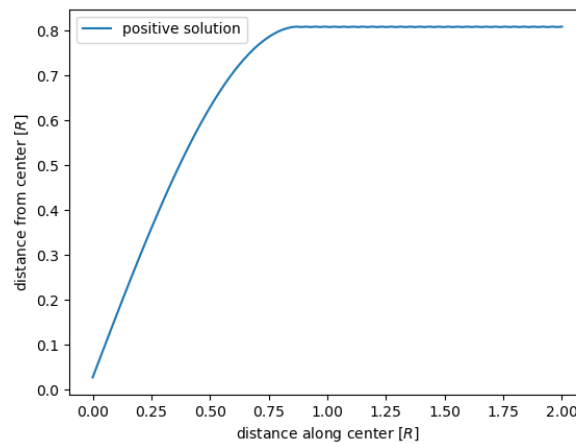


Figure 7: The numerical solution

Compare: the numerical solution of the problem with an analytical one. The above Python code was also used to plot the analytical solution. Both are shown in figure 8. As can be seen in the figure the numerical solution approximates the analytical solution pretty well, up to the point where the derivative becomes negative.

Explain: what is the advantage of such a fiber? When using a alpha-peak profile for the refractive index of the fiber it is possible to use more modes of light more efficiently since for a certain shape factor  $\alpha$  the multimode distortion for this profile is the smallest possible. When more modes of light can be used the cable has a higher bandwidth.

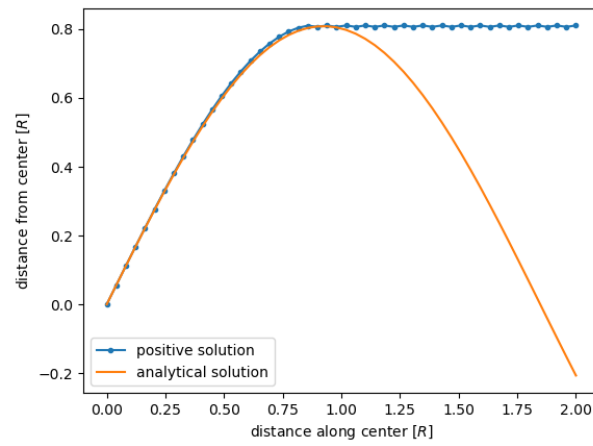


Figure 8: The analytical and numerical solution