

Propagation of Light

R van Loon, J Sangers

Today

1 Principle of Fermat

Question: Proof that equation 1 and 2 can be reduced to equation 3.

$$L[x, y, z, \dot{x}, \dot{y}, \dot{z}] = n(x, y, z) \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \quad (1)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0, \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} - \frac{\partial L}{\partial y} = 0, \frac{d}{dt} \frac{\partial L}{\partial \dot{z}} - \frac{\partial L}{\partial z} = 0 \quad (2)$$

$$\frac{d}{ds} \left[n \frac{d\vec{r}}{ds} \right] = \vec{\nabla} n \quad (3)$$

Answer:

First noting that ds is a small element of distance travelled. Therefore taking into account the x , y and z direction, ds is given by:

$$ds = \sqrt{dx^2 + dy^2 + dz^2}$$

A small distance travelled in a trivial direction, lets say dx , can be approximated by as $dx = dt \cdot \dot{x}$. Therefore ds can be rewritten as:

$$ds = dt \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}$$

Rewriting gives:

$$\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} = \frac{ds}{dt} \quad (4)$$

If we combine the equations in equation 2 in vector notation we get:

$$\frac{d}{dt} \left(\frac{\partial}{\partial \dot{\vec{r}}} \right) L - \nabla L = 0$$

Rewriting and filling in equation 1 gives:

$$\frac{d}{dt} \left[\left(\frac{\partial}{\partial \dot{\vec{r}}} \right) n(x, y, z) \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \right] = \left(\frac{\partial}{\partial \vec{r}} \right) \left[n(x, y, z) \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} \right]$$

$$\frac{d}{dt} \frac{n \cdot \dot{\vec{r}}}{\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}} = \vec{\nabla} n \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}$$

Using equation 4 to replace the $\sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2}$ and rewriting the $\dot{\vec{r}}$ vector gives the following:

$$\frac{d}{dt} n \cdot \frac{\dot{\vec{r}}}{\frac{ds}{dt}} = \vec{\nabla} n \frac{ds}{dt}$$

$$\frac{d}{dt} n \cdot \frac{\frac{d}{dt} \vec{r}}{\frac{ds}{dt}} = \vec{\nabla} n \frac{ds}{dt}$$

Rewriting yields the equation that was to be proved:

$$\frac{d}{ds} \left[n \frac{d\vec{r}}{ds} \right] = \vec{\nabla} n$$

2 Application

2.1 Homogeneous medium

Question: Using equation 3, show how light is travelling in a homogeneous medium.

Answer:

Equation 3 can be rewritten using the chain rule:

$$\frac{d\vec{r}}{ds} \frac{d}{ds} n + n \frac{d^2 \vec{r}}{ds^2} = \left(\frac{\partial}{\partial x} \frac{\partial}{\partial y} \frac{\partial}{\partial z} \right) n$$

Note that for a homogeneous medium, the index of refraction, n , is constant. Therefore $\frac{dn}{ds} = 0$, $\frac{\partial n}{\partial x} = 0$, $\frac{\partial n}{\partial y} = 0$ and $\frac{\partial n}{\partial z} = 0$. Using this in the previous equation yields:

$$n \frac{d^2 \vec{r}}{ds^2} = \vec{0}$$

$$\frac{d^2 \vec{r}}{ds^2} = \vec{0}$$

Using the definition of ds (equation 4) in the latter equation yields:

$$\frac{d^2 \vec{r}}{dt^2 (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)} = \vec{0}$$

$$\frac{d^2 \vec{r}}{dt^2} = \vec{0}$$

Integrating both sides of the equation with respect to t twice gives us the equation for $\vec{r}(t)$ which describes the uniquely the path of the light as a function of time:

$$\frac{d\vec{r}}{dt} = A$$

$$\vec{r} = A t + B$$

We assume that at a time $t = 0$, when we start to study a beam of light in the homogeneous medium, the initial position and velocity of the beam are $\vec{r}(0) = B = \vec{r}_0$ and $\dot{\vec{r}}(0) = A\dot{\vec{r}}_0$ respectively. Therefore we find:

$$\vec{r} = \dot{\vec{r}}_0 t + \vec{r}_0$$

This is a linear function of time and implies that the direction and velocity of the light is not changed as the light travels through the medium. Therefore, it travels in a straight line with a constant velocity.

2.2 Snell-Descartes Law

Question: Express first geometrically and then analytically Snell's and Descartes' law of reflection and transmission of the light at the interface between two media of different index of refraction n_1 and n_2 , using the Principle of Fermat and equation 3.

Answer:

2.2.1 Geometrical

The speed of light in a medium is inversely proportional to the refractive index. Therefore, the shortest path (in distance) between two points in materials with different refractive indices is not always the fastest (in time). This phenomenon is nicely described by a 2-dimensional analogy of a beach (see figure 1). The maximum speed on foot on beach is significantly higher than the maximum swimming speed in the water. So if somebody would need to get from a point A on the beach to a point B in the water, the direct route from A to B (dashed line in figure 1) would intuitively be slower than the path with a shorter swimming distance (solid line in figure 1).

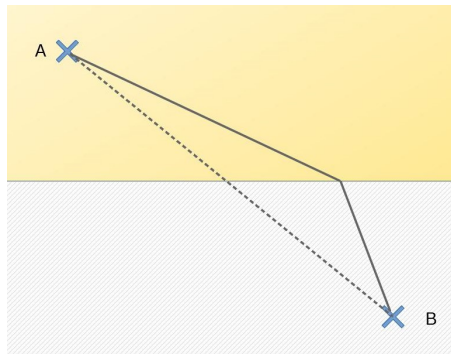


Figure 1: Diagram of a 2-dimensional beach analogy of the interface between two media with different index of refraction. The upper-half corresponds to the beach and the lower-half to the sea. The dashed line corresponds to the direct route between point A and B with the shortest distance. The solid line corresponds to a route that is intuitively faster than the direct route.

To find the angle of refraction at the interface of two media we consider two parallel beams of coherent light that reach an interface, oriented in the x,z -plane, such as in figure 2. The two beams propagate in the x,y -plane and have an incident angle of θ_1 , the distance between the rays is Δx in the x -direction and the refracted angle of the two beams is θ_2 . The latter angle is equal for both beams of light, since the two media have a constant refractive index.

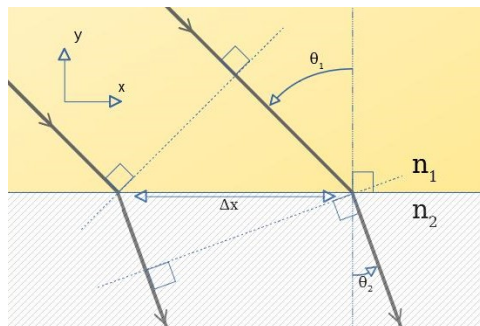


Figure 2: Diagram of two parallel beams of coherent light reaches an interface with an angle between two media with different refractive indices, n_1 and n_2 . The beams have an incident angle θ_1 and refracted angle θ_2 . Δx corresponds to the distance between the beams in the x direction.

The phase change between two beams of light is given by the product of the passed time, Δt , and the angular frequency, ω . Therefore, $\Delta\phi$ is given by:

$$\Delta\phi = \omega\Delta t_2 - \omega\Delta t_1$$

$$\frac{\Delta\phi}{\omega} = \Delta t_2 - \Delta t_1$$

With Δt_1 and Δt_2 , the passed time between the two dotted lines in figure 2. We can calculate Δt_1 and Δt_2 using the speed of the light in the media, $v = c/n$, and the difference in path length in both media, ΔL_1 and ΔL_2 (see figure 3 for clarification of the used parameters). ΔL can be calculated using trigonometric identities with the corresponding angles:

$$\begin{aligned}\frac{\Delta\phi}{\omega} &= \Delta t_2 - \Delta t_1 \\ &= \frac{\Delta L_2}{v_2} - \frac{\Delta L_1}{v_1} \\ &= \frac{n_2\Delta L_2}{c} - \frac{n_1\Delta L_1}{c} \\ &= \frac{n_2[\sin(\theta_2)\Delta x]}{c} - \frac{n_1[\sin(\theta_1)\Delta x]}{c} \\ \frac{c}{\omega}\frac{\Delta\phi}{\Delta x} &= n_2 \sin(\theta_2) - n_1 \sin(\theta_1)\end{aligned}$$

Since we don't expect any discontinuity in phase by the refraction, the light beams in the second medium should also be in phase, just as the coherent light beams in medium 1. This is denoted by the dashed lines in the figures. This implies that $\Delta\phi = 0$ in the latter equation and leads to Snell-Descartes law:

$$\begin{aligned}\frac{c}{\omega}\frac{\Delta\phi}{\Delta x} &= 0 = n \sin(\theta_2) - n \sin(\theta_1) \\ n_2 \sin(\theta_2) &= n_1 \sin(\theta_1)\end{aligned}$$

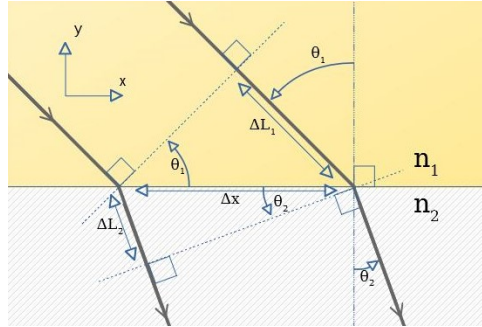


Figure 3: Diagram like in figure 2 with extra parameters ΔL_1 and ΔL_2 corresponding to the path difference in the different media by the two beams of light.

2.2.2 Analytical

For the analytical derivation of Snell-Descartes law we will use a similar diagram as in the geometrical derivation with a coordinate system added as in figure 4a.

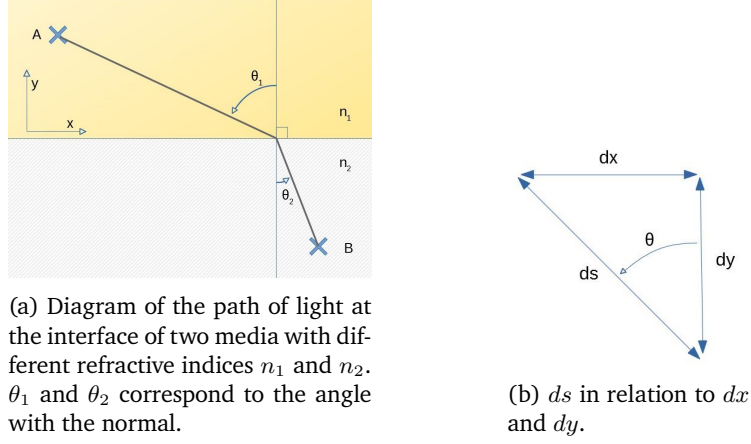


Figure 4

If we write equation 3 for only the x-component and use the fact that n is independent of x in our diagram, we get the following:

$$\frac{d}{ds} \left[n \frac{dx}{ds} \right] = \frac{dn}{dx}$$

$$\frac{d}{ds} \left[n \frac{dx}{ds} \right] = 0$$

The ds in the latter equation is defined as in figure 4b. If we keep a fixed ds for both media we obtain the following equality:

$$\frac{d}{ds} \left[n_1 \frac{dx_1}{ds} \right] = \frac{d}{ds} \left[n_2 \frac{dx_2}{ds} \right]$$

Integrating both sides with respect to ds and using the trigonometric identity from equation ?? yields the Snell-Descartes law:

$$n_1 \frac{dx_1}{ds} = n_2 \frac{dx_2}{ds}$$

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2) \quad (5)$$

2.3 Mirage

Question: The mirage is a common phenomenon when the ground is very warm and the temperature of the air decreases with altitude; in this case the density then increases as well as its index of refraction.
Sketch what is happening.

Answer:

In figure 5 a situation in which a mirage occurs has been sketched. In this figure a few parameters were introduced, an x - and z -distance respectively denoting the horizontal and vertical distance from the feet of the observer and an angle Θ which is the angle between the horizontal and the unbent light path from the observer. A gradient effect is also applied to the image, where the colour is darker the index of refraction is higher.

Since the index of refraction of the air varies with the temperature and the temperature increases as the z -coordinate increases, the path of light rays will be bent. Thus there will be a light ray coming from the sky bending in such a way that it lands in the eye of an observer. This is the mirage effect where light follows a different path than one might expect.

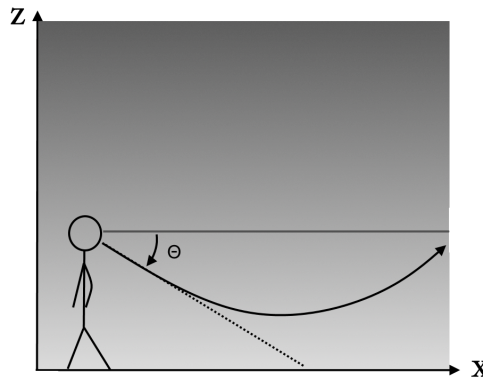


Figure 5: A sketch of situation where a mirage effect occurs.

Question: Find the relation between the index of refraction with the altitude if we assume that the gradient of the temperature changes linearly with the altituder.

Answer:

In equation 6 the temperature at a height z is defined using a ground temperature T_0 and a lapse rate c [K/m] so that it is linearly decreasing. To relate the index of refraction n to the temperature at height $T(z)$ and pressure $P(z)$, the Edlén equation is used this equation is outlined in [1]. In equation 7 this relation is shown. Finding the relation for the pressure at height $P(z)$ is unnecessary as the pressure is assumed constant, the $(n - 1)_S$ term is shown in 8 here σ is the wavelegth of the light. The result of combining all three equations results in a usable relation between the index of refraction n and the height above the ground z . This is shown in equation 9.

$$T(z) = T_0 - c \cdot z \quad (6)$$

$$(n - 1)_{TP} = \frac{P \cdot (n - 1)_S}{96095.43} \cdot \frac{1 + 10^{-8}(0.601 - 00972 \cdot T) \cdot P}{1 + 0.0036610 \cdot T} \quad (7)$$

$$(n - 1)_S \cdot 10^{-8} = 8342.54 + \frac{2406147}{130 - \sigma^{-2}} + \frac{15998}{38.9 - \sigma^{-2}} \quad (8)$$

$$(n - 1)_{TP} = \frac{P \cdot 10^8 \cdot (8342.54 + \frac{2406147}{130 - \sigma^{-2}} + \frac{15998}{38.9 - \sigma^{-2}})}{96095.43} \cdot \frac{1 + 10^{-8}(0.601 - 00972 \cdot (T_0 - c \cdot z)) \cdot P}{1 + 0.0036610 \cdot (T_0 - c \cdot z)} \quad (9)$$

Question: Express analytically the trajectory of the light in this situation.

Answer:

For an analytical solution it is easier to linearise equation 9. Where $\alpha [m^{-1}]$ is a linearisation constant.

$$n_l(z) = n_0 + \alpha \cdot z$$

$$n_l(z) = n(0) + z \cdot \frac{d}{dz}n(z)|_{z=0}$$

To get to an analytical solution; equation 3 needs to be solved. This is done below.

$$\vec{\nabla} \cdot n(z) = \frac{\partial}{\partial z}n(z)\vec{e}_z = \alpha\vec{e}_z$$

$$\frac{d}{ds} \left[n(z) \frac{d\vec{r}}{ds} \right] = \vec{\nabla} \cdot n(z) = \alpha\vec{e}_z$$

$$\frac{d}{ds} \left[n(z) \frac{d\vec{x} + d\vec{z}}{\sqrt{dx^2 + dz^2}} \right] = \alpha\vec{e}_z$$

$$\frac{d}{ds} \left[n(z) \frac{dx}{\sqrt{dx^2 + dz^2}} \right] = 0$$

$$\frac{d}{ds} \left[n(z) \frac{dz}{\sqrt{dx^2 + dz^2}} \right] = \alpha$$

$$n(z) \cdot \frac{dx}{\sqrt{dx^2 + dz^2}} = c$$

$$n(z)^2 \cdot \frac{dx^2}{dx^2 + dz^2} = c^2$$

$$n(z)^2 \cdot dx^2 = c^2(dx^2 + dz^2)$$

$$\frac{dz}{dx} = \frac{\sqrt{n(z)^2 - c^2}}{c} = \frac{\sqrt{(n_0 + \alpha \cdot z)^2 - c^2}}{c}$$

The above first-order nonlinear ordinary differential equation has been solved using wolfram and is displayed below in 10. In this equation k is an arbitrary constant and c is the total path length in meters. $\alpha [m^{-1}]$ is the defined linearisation constant.

$$z(x) = -\frac{c^2 e^{(\alpha x + \alpha k)/c} + e^{-(\alpha x + \alpha k)/c} - 2n_0}{2\alpha} \quad (10)$$

Question: Plot several of these trajectories for different relevant cases.

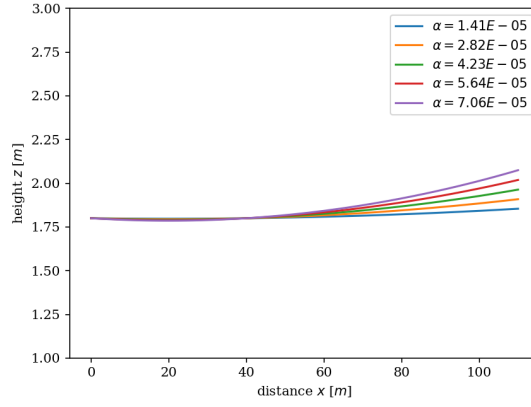
Answer:

Below in figure ?? a few trajectories for different α -values are plotted. These values range from $\alpha = 1.41 \cdot 10^{-5} [m^{-1}]$ to the same constant five times larger. The linearisation α is constant equal to $d/dz \cdot n(z)$ evaluated as the slope of the line between n_0 the index of refraction at $z = 0 [m]$ and n_1 for $z = 3 [m]$. Our initial guess for α was $\alpha = 1.41 \cdot 10^{-5}$, but this resulted in a distance to the mirage of about 50 $[m]$. We felt that this might an actual mirage effect and thus decided plot these alfa values. To get these values we had a ground temperature of $T_0 = 344 K$ and a temperature of $T_1 = 293 K$, we chose these values since the might represent a hot day somewhere near the equator. This is still not an actual representation for a mirage since the temperature and the index of refraction doe not change linearly with the height above the ground.

Question: Find the closest distance from an observer where a mirage can be visible.

Answer:

From the plot above we can see that the closest a mirage can occur is 50 meters from the observer. This



scenario is still not an actual representation for a mirage since the temperature and the index of refraction do not change linearly with the height above the ground.

Question: Forsee what could happen on the north pole when a warm wind is present.

Answer:

On the North Pole with a warm wind blowing an inverse effect as described before would be seen, since the ground is cold and the air heats up as the height increases. The index of refraction n would then decrease the higher off the ground. Light would then bend downward instead of up. If we assume the temperature gradient, and therefore the refractive index gradient, is constant in the direction parallel to the ground, there are two things likely to happen. (1) Since the path of the light has its curvature in the same direction as the curvature of the earth allowing, the light bends around the curvature of the earth and objects that are beyond the horizon can be seen. (2) When looking up, the downward bending of the light shows objects that are actually on the ground. This results in objects that seem to fly.

2.4 Optical fibre

2.4.1

Question: If we assume an incoming beam in the $\vartheta_{x,y}$ plane at start crossing ϑ_x with an angle θ_0 . Show that the light beam will stay in this plane and that, by solving the Euler-Lagrange equation 3, the equation for the beam path can be written as:

$$n \sin(i) = a \quad (11)$$

Answer:

For this question it is assumed that the optical fibre is cylindrical with its axis ϑ_x (see figure 6). The incident-surface will therefore be in the $\vartheta_{y,z}$ plane.

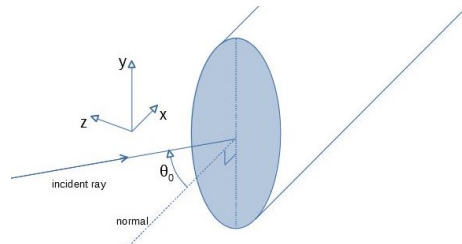


Figure 6: 3-dimensional diagram of fibre with incoming light beam making an angle θ_0 with the ϑ_x axis.

When the incoming beam is in the $\vartheta_{x,y}$ plane, the angle with the normal is θ_0 with ϑ_x in the ϑ_y direction and $\theta_z = 0$ in the ϑ_z direction. If we use Snell-Descartes law (equation 5) for θ_z with $\theta_{z,1} = 0$ we obtain $\theta_{z,2} = k \cdot \pi [\text{rad}]$ with $k = -1, 0, 1, 2, \dots$. Therefore, the beam of light will stay in the $\vartheta_{x,y}$ plane at the interface when entering the optical fibre.

It is the property of a cylinder that the outer surface is always perpendicular with the radius. For this light beam. Since the light beam is travelling in the $\vartheta_{x,y}$ direction, the outer surface it reaches will be in the $\vartheta_{x,z}$ plane. The angle with the surface normal in the ϑ_z axis is still zero. From the internal reflection it follows that the angle in the ϑ_z axis will remain zero. Therefore, for every reflection at the outer surface, the beam will stay in the $\vartheta_{x,y}$ direction.

Since we know that the beam will stay in the $\vartheta_{x,y}$ plane, the problem can be simplified to a 2-dimensional problem (see figure 7). The infinitely small segment ds is related to ds , dy , dx and i as in figure 7.

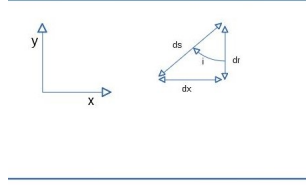


Figure 7: 2-dimensional fibre including the relation between ds , ds , dy , dx and i .

We can write equation 3 for only the x-component and use that n is independent of x , and therefore $\frac{dn}{dx} = 0$:

$$\frac{d}{ds} \left[n \frac{dx}{ds} \right] = \frac{dn}{dx} = 0 \quad (12)$$

If we now integrate with respect to ds on both sides we get the following:

$$n \frac{dx}{ds} = a \quad (13)$$

With a an arbitrary constant. From equation ?? it follows that $\sin(i) = dx/ds$. Using this, we obtain the requested equation:

$$n \sin(i) = a \quad (14)$$

2.4.2

Question: Solve equation 11 for $n(r) = n_0 \sqrt{1 - \alpha^2 r^2}$ with $\alpha < 1/R$, a constant, R the radius of the fibre and r the distance to the ϑ_x -axis. You should end up with an analytical function.

Answer:

From figure 7 and equation ?? it follows that $\sin(i) = dx/ds$. Using this and the Pythagoras theorem in equation 11 yields:

$$n \frac{dx}{\sqrt{dr^2 + dx^2}} = a$$

Filling in $n(r)$ and some rewriting results in the following:

$$n_0 \sqrt{1 - \alpha^2 r^2} \frac{dx}{\sqrt{dr^2 + dx^2}} = a$$

$$n_0^2 (1 - \alpha^2 r^2) \frac{dx^2}{dr^2 + dx^2} = a^2$$

$$n_0^2 (1 - \alpha^2 r^2) dx^2 = a^2 (dr^2 + dx^2)$$

$$\begin{aligned}
(n_0^2 - n_0^2 \alpha^2 r^2 - a^2) dx^2 &= a^2 dr^2 \\
dx^2 &= \frac{a^2 dr^2}{n_0^2 - n_0^2 \alpha^2 r^2 - a^2} \\
dx &= \frac{a dr}{\sqrt{n_0^2 - n_0^2 \alpha^2 r^2 - a^2}}
\end{aligned} \tag{15}$$

If we now integrate both sides we will get:

$$\begin{aligned}
\int dx &= \int \frac{a}{\sqrt{n_0^2 - n_0^2 \alpha^2 r^2 - a^2}} dr \\
x &= \frac{a}{\sqrt{n_0^2 - a^2}} \int \frac{1}{1 - \left(\frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \right)^2} dr
\end{aligned}$$

To solve this integral we need the following substitution:

$$\begin{aligned}
\frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} &= \sin(u) \\
dr &= \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \cos(u) du \\
u &= \arcsin \left(\frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \right)
\end{aligned}$$

Using the substitution we get:

$$x = \frac{a}{n_0 \alpha} \int \frac{\cos(u)}{\sqrt{1 - \sin^2(u)}} du$$

If we use that $1 - \sin^2(x) = \cos^2(x)$, we obtain:

$$\begin{aligned}
x &= \frac{a}{n_0 \alpha} \int \frac{\cos(u)}{\sqrt{\cos^2(u)}} du \\
x &= \frac{a}{n_0 \alpha} \int du \\
x &= \frac{a}{n_0 \alpha} u
\end{aligned}$$

Inverting the substitution and doing some rewriting yields the following equation:

$$\begin{aligned}
x &= \frac{a}{n_0 \alpha} \arcsin \left(\frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \right) \\
\sin \left(\frac{n_0 \alpha x}{a} \right) &= \frac{r n_0 \alpha}{\sqrt{n_0^2 - a^2}} \\
r(x) &= \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \sin \left(\frac{n_0 \alpha x}{a} \right)
\end{aligned}$$

Since there are no discontinuities in $n(r)$ and $dn(r)/dr$, we would expect that there are also no discontinuities in $r(x)$ and $dr(x)/dx$. Since $r = |y|$ for the light beam in the $\vartheta_{x,y}$ plane. The only possible solution of $y(x)$ is the following:

$$y(x) = \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \sin \left(\frac{n_0 \alpha x}{a} \right) \tag{16}$$

If we now consider that the angle of the light beam at $x = 0$, θ_i , is determined. We must consider the derivative of the latter equation:

$$\frac{d y(x)}{d x} = \frac{n_0 \alpha}{a} \frac{\sqrt{n_0^2 - a^2}}{n_0 \alpha} \cos\left(\frac{n_0 \alpha x}{a}\right)$$

For $x = 0$:

$$\frac{d y(x)}{d x} = \sqrt{n_0^2 - a^2}$$

Given the trigonometric identity $\tan(\theta) = (\text{opposite} - \text{side})/(\text{adjacent} - \text{side})$. It follows that:

$$\frac{d y}{d x} = \sqrt{n_0^2 - a^2} = \tan(\theta_i)$$

If we now consider n_0 to be constant for the given fibre and a dependent on θ_i , the initial angle can be integrated in equation 16.

$$a^2 = n_0^2 - \tan^2(\theta_i) \quad (17)$$

Question: Plot a numerical solution of the problem.

Answer:

To plot a numerical solution we use equation 15 and rewrite it to get a fraction dr/dx . We will then plot this function using the python script below. There is one problem however, since we used equation 15 we can only get the part of the solution where $(dr/dx)^2$ was positive.

```
# -*- coding: utf-8 -*-
"""
Created on Wed May 6 17:53:40 2020

@author: vanlo
"""

import numpy as np
import matplotlib.pyplot as plt

plt.close() # Close previous plots

"""Define function to plot the analytical solution
to the problem using only the x """
def analytical(x):
    return np.sqrt(n0**2 - a**2)/(n0*alpha)*np.sin((n0*alpha*x)/a)

""" Define function to find the slope of the function at every step in x-step.
Input: previous value of r """
def drdx(r):
    # Added an incline variable in order to let the function change direction.
    # If incline = 1, the function is going up until it reaches it's peak,
    # else, the function needs to go down until it's peak
    if incline == 1:
        drdx_1 = abs(np.sqrt((n0**2 - n0**2*alpha**2*r**2 - a**2)*a**(-2)))
    else:
        drdx_1 = -abs(np.sqrt((n0**2 - n0**2*alpha**2*r**2 - a**2)*a**(-2)))
    # If the slope has a value close to 0,
    # we know the function has reached a peak (or is near a peak)
    # --> We change value of peak to 1 if value for slope is near 0.
    # If not, peak = 0 --> not at a peak
    if abs(drdx_1) < 0.01:
```

```

        peak = 1
    else:
        peak = 0
    return [drdx_1, peak]

# Define constants also using complex part in order to perform complex roots:
n0 = 1.7 + 0*1j      # Refractive index at r = 0
R = 1                # Outer radius of fiber
alpha = 1/R + 0*1j   # inverse of R
a = 1 + 0*1j         # Arbitrary constant
N = int(2*1e3)       # Number of steps for numeric calculation

incline = 1          # Variable to define if function is going down or going up

x = np.linspace(0, 10*R, N) # Array for which to find r values
dx = x[1] - x[0] + 0*1j     # Step size in x-direction
drdx_array = np.array([])   # Empty array to fill with correct
                             # slopes for every x-step.

r = np.zeros(N) + np.zeros(N)*1j
# Empty array to fill up with correct r-values.
# Also complex part for computation
r[1] = dx # We let the beam come in at an angle of 1/2*Pi.

peak_array = np.empty(N)    # Array for peak values to keep track
                             # previous peak values

# Numeric calculation of r:
for i in range(1, N):
    # Calculate slope and add to array
    drdx_i = drdx(r[i-1])[0]
    r[i] = r[i-1] + drdx_i*dx
    # See if function is at a peak or near a peak, and add to array
    peak_i = drdx(r[i-1])[1]
    peak_array[i-1] = peak_i
    # If functions is at a peak and hasn't just passed a peak,
    # change incline so it changes direction
    # in order to let function oscillate
    if peak_i == 1 and (1 not in peak_array[i-6:i-1]):
        incline = -1* incline

# plots of numeric and analytical solution
plt.plot(x, analytical(x), label='analytical', linewidth=3.)
plt.plot(x, r, label='Numeric', linestyle="--", marker=".", markevery=10, color="black")
plt.legend()
plt.xlabel('x/R')
plt.ylabel('y/R')
plt.show()

```

Figure 8 shows a plot of the numerical solution using the python code above. The horizontal and vertical axis are normalised to R , the radius of the fibre.

Question: Compare the numerical solution of the problem with an analytical one.

Answer:

The above Python code was also used to plot the analytical solution. Both are shown in figure 9. As can be

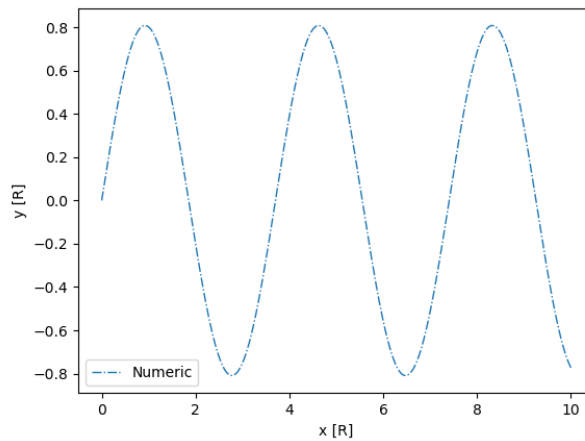


Figure 8: The numerical solution

seen in the figure the numerical solution approximates the analytical solution pretty well, up to the point where the derivative becomes negative.

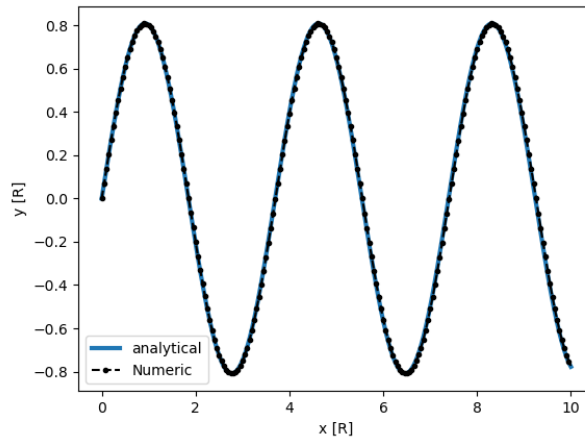


Figure 9: The analytical and numerical solution

Question: Explain: What is the advantage of such a fibre?

Answer:

Consider a straight piece of cylindrical fibre with a constant refractive index and a length l (see figure 10 for a schematic diagram). A light beam that gets send in the fibre parallel to the axis of the fibre will not reflect and follow the shortest route with length l to the other end of the fibre. If a beam of light is enters the fibre with an angle θ_i to the axis, the beam will reflect at the outer surface and the path the light will travel to the end of the fibre will be of length $l/\cos(\theta_0)$.

Since the speed of light is finite, this difference in path length causes a delay for light rays that enter the fibre in different angles. Since it is difficult or even impossible to get a fully coherent light signal, this causes a signal that is transmitted in the fibre to spread over a larger time interval. This time-spread of signals is disadvantageous if someone would want to send as much information through the fibre as possible.

The given refractive index will give a light path as in equation 16. Considering the shape of this sine in comparison to the situation with constant refractive index, this sine-shaped path is much shorter. Therefore the time-spread of signals is shorter for the situation with the refractive index gradient. This allows

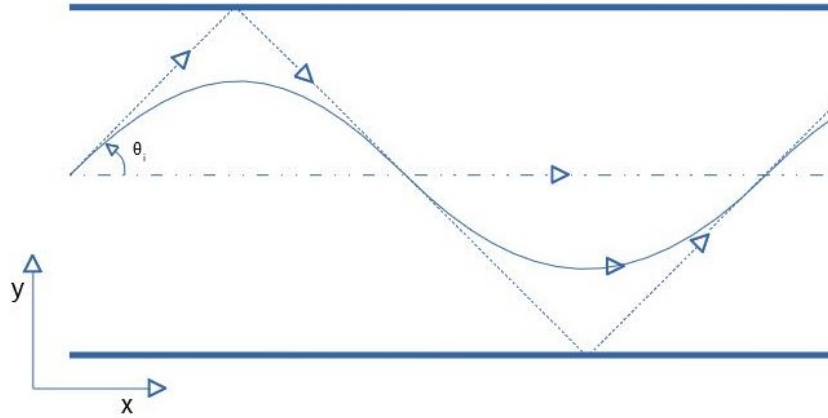


Figure 10: 2-dimensional diagram of an optical fibre with possible light paths. The straight line with the dots and dashes corresponds to an optical path parallel to the axis of a fibre with a constant refractive index. The dashed line corresponds to the optical path for a beam entering with an angle θ_i for a fibre with constant refractive index. The solid line corresponds to the optical path for a beam entering with an angle θ_i for a fibre with a refractive index gradient given by $n(r) = n_0\sqrt{1 - \alpha^2 r^2}$.

you to send more information through the fibre in a time unit.

References

- [1] K P Birch and M J Downs. An updated edlén equation for the refractive index of air. *Metrologia*, 30(3):155–162, jan 1993.