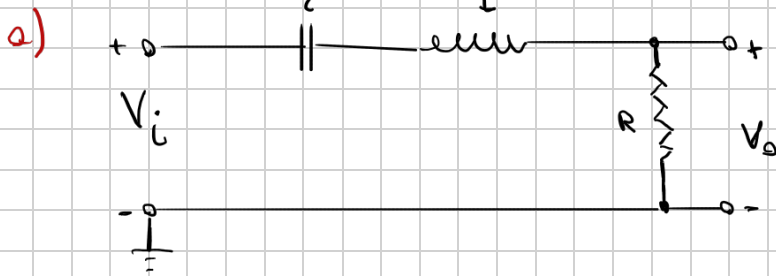


$$[s = \sigma + j\omega]$$



$$H(s) = \frac{V_o}{V_i}$$

$$Z_c = \frac{1}{sC} \quad Z_L = sL \quad Z_R = R$$

$$V_o = Z_R$$

$$V_i = Z_c + Z_L + Z_R$$

$$H(s) = \frac{Z_R}{Z_c + Z_L + Z_R} = \frac{R}{\frac{1}{sC} + sL + R} = \frac{R \cdot sC}{s^2 CL + RSC + 1} = \frac{R sC}{s^2 + \frac{R}{L}s + \frac{1}{CL}}$$

$$T(s) = \frac{\omega_0/Q \cdot s}{s^2 + \frac{\omega_0}{Q}s + \omega_0^2} \rightarrow \text{Poles}$$

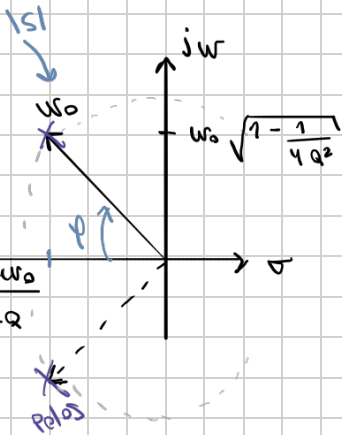
$$s^2 + \frac{\omega_0}{Q}s + \omega_0^2 = 0$$

$$s = \omega_0 \left(\frac{-1}{2Q} \pm j \sqrt{1 - \frac{1}{4Q^2}} \right)$$

$$\omega_0 = \frac{1}{\sqrt{CL}}$$

$$Q = \frac{\sqrt{CL}}{R}$$

Normalizado $\omega_0 = 1$



Veamos que:

$\frac{-\omega_0}{2Q}$ Indica el amortiguamiento (parte real)

$\omega_0 \sqrt{1 - \frac{1}{4Q^2}}$ Freq. oscilatoria (pte. im.)

$$\cos \varphi = \frac{\sigma}{|s|}$$

$$|s| = \omega_0 \rightarrow \cos \varphi = \frac{\omega_0}{\frac{\omega_0}{2Q}} = \frac{1}{2Q}$$

con $Q = 1/2$ se juntan los polos

P/valores altos de Q $\varphi \rightarrow 90^\circ$, p/valores bajos $\varphi \rightarrow 0^\circ$

$$Q = \frac{1}{2 \cos \varphi}$$

Dejamos fijo $\omega_0 \rightarrow \omega_0^2 = 1 = \frac{1}{CL} \rightarrow C = \frac{1}{L}$

$$\varphi = \cos^{-1}(2Q)$$

→ Estado estacionario: $s = j\omega$? Este sirve p/ analisis en frecuencia

$$T(j\omega) = \frac{\omega_0/Q - j\omega}{(j\omega)^2 + j\omega \frac{\omega_0}{Q} + \omega_0^2} = \frac{\omega_0/Q - j\omega}{\omega_0^2 - \omega^2 + j\omega \frac{\omega_0}{Q}}$$

$$|T| = \frac{\omega \cdot \omega_0/Q}{\sqrt{(\omega_0^2 - \omega^2)^2 + \left(\frac{\omega \omega_0}{Q}\right)^2}}$$

• A bajas frecuencias $\omega \ll \omega_0$

$$\rightarrow |T(j\omega)| = \frac{\omega}{Q \omega_0}$$

• A altas frecuencias $\omega \gg \omega_0$

$$\rightarrow |T(j\omega)| = \frac{\omega_0}{Q \omega}$$

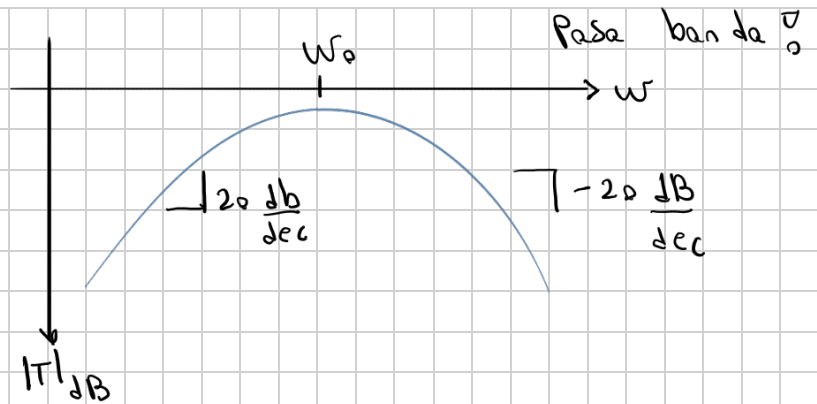
$$|T| = \frac{\omega \cdot \omega_0/Q}{\sqrt{(\omega_0^2 - \omega^2)^2 + \left(\frac{\omega \omega_0}{Q}\right)^2}}$$

$$|T| = \frac{\omega \cdot \omega_0/Q}{\sqrt{(\omega_0^2 - \omega^2)^2 + \left(\frac{\omega \omega_0}{Q}\right)^2}}$$

- En resonancia $\omega = \omega_0$

$$|T(j\omega)| = 1$$

Gratificamos el patron de dB.



→ Analisis de fase

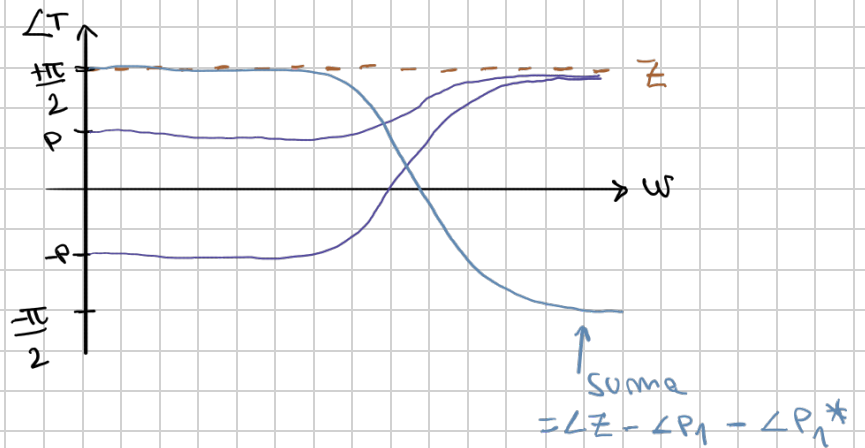
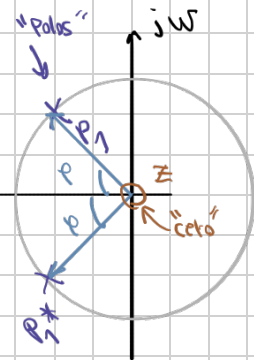
$$T(s) = \frac{s \cdot \omega_0 / Q}{s^2 + s \frac{\omega_0}{Q} + \omega_0^2}$$

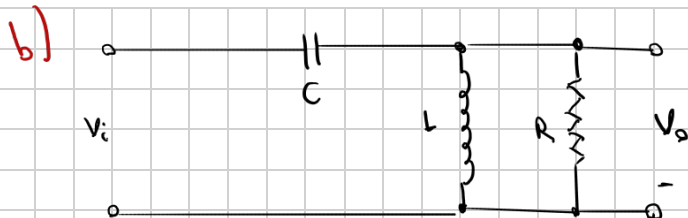
$$T(s) = \frac{s \omega_0 / Q}{(s - p_1)(s - p_1^*)}$$

↖ ↗
 Poles

⇒ graficamente:

Podemos ver que cuando $\omega \rightarrow \infty$ cada polo aporta $-\pi/2$ a la fase y el cero aporta $\pi/2$





↓ Pq?

$$V_o = Z_{LR} = \frac{Z_L \cdot Z_R}{Z_L + Z_R}$$

$$V_i = Z_C + Z_{LR} = Z_C + \frac{Z_L \cdot Z_R}{Z_L + Z_R}$$

$$= \frac{1}{sC} + \frac{sLR}{sL + R}$$

• SC (SL+R)

$$H(s) = \frac{Z_{LR}}{Z_C + Z_{LR}} = \frac{\frac{sLR}{sL + R}}{\frac{1}{sC} + \frac{sLR}{sL + R}} = \frac{s^2 cLR}{s^2 + \frac{s}{RC} + \frac{1}{LC}}$$

$$= \frac{s^2}{s^2 + \frac{s}{RC} + \frac{1}{LC}}$$

$\omega_0 = \frac{1}{\sqrt{LC}}$ $\frac{\omega_0}{Q} = \frac{1}{RC}$ $Q = \frac{RC}{\sqrt{LC}}$

Estado estacionario $\rightarrow s = j\omega$

$$\Rightarrow T(j\omega) = - \frac{\omega^2}{\omega_0^2 - \omega^2 + j\omega \frac{\omega_0}{Q}}$$

$$|T| = \frac{\omega^2}{\sqrt{(\omega_0^2 - \omega^2)^2 + \left(j\omega \frac{\omega_0}{Q}\right)^2}}$$

• A bajas f $\omega \ll \omega_0$

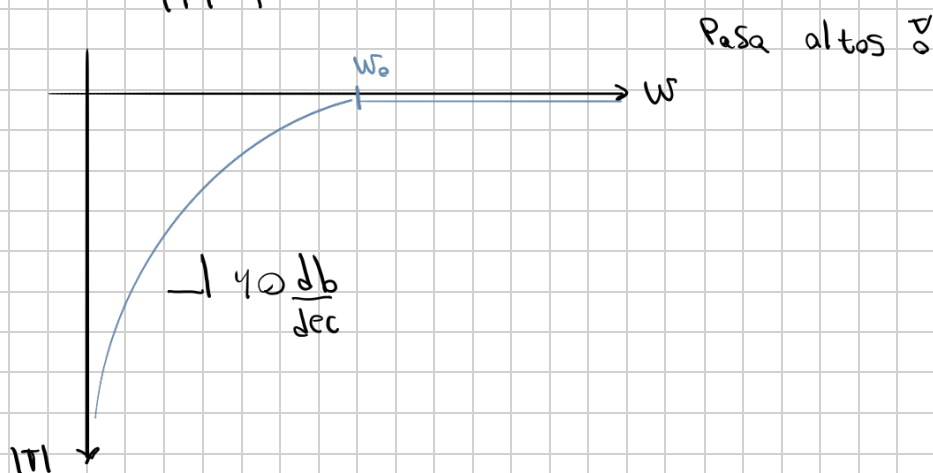
$$|T| = \frac{\omega^2}{\omega_0^2}$$

• A altas f $\omega \gg \omega_0$

• Resonancia ($\omega = \omega_0$) $|T| = 1$

$$|T| = 1$$

Gráficamente el
patron de dB \rightarrow

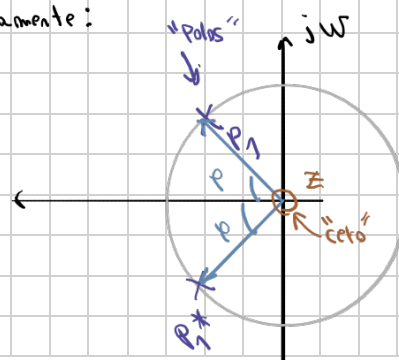


→ Analisis de fase

$$T(s) = \frac{s^2}{s^2 + s \frac{\omega_0}{Q} + \omega_0^2}$$

$$\Rightarrow T(s) = \frac{s^2}{(s - p_1)(s - p_1^*)}$$

⇒ Gráficamente:



Mismo analisis que
parte a)

Podemos ver que
cuando $\omega \rightarrow \infty$ cada polo
aporta $-\pi/2$ a la fase y el
cero aporta $\pi/2$

