

# Scientific Machine Learning - Theoretical aspects

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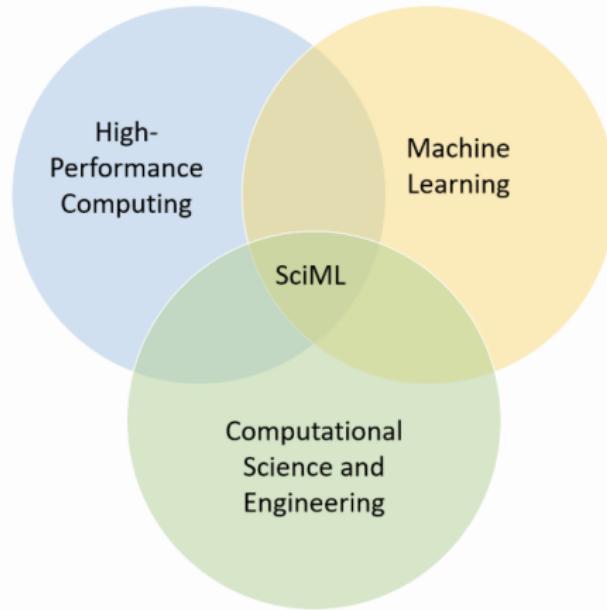
# Outline

- 1 Introduction to Scientific Machine Learning (SciML)
- 2 Several partial differential equations
- 3 Approximation of functions by neural networks

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# Introduction to SciML



In recent years, the combination of numerical methods and machine learning has gained an ever increasing interest as a research field within Numerical Mathematics and Scientific Computing

# Introduction to SciML

## What is Scientific Machine Learning ?

- SciML is an emerging discipline within the data science community
- SciML seeks to address **domain-specific** data challenges and extract insights from scientific data sets through innovative methodological solutions
- SciML draws on tools from both Machine Learning (ML) and scientific computing to develop new methods for **robust** learning and data analysis
- SciML will be critical in driving the next wave of data-driven scientific discovery in the physical and engineering sciences
- SciML is multidisciplinary and leverages expertise from applied and computational mathematics, computer science, and the physical sciences

# Introduction to SciML

## Why Scientific Machine Learning ?

- New innovations in ML and Big Data are beginning to drive advances in scientific disciplines
- But the full potential of these techniques for data-driven discovery has yet to be fully realized
- One barrier to data-driven discovery is that existing methods often do not meet the needs of scientific users
- Application-agnostic algorithms, or those designed for more traditional ML applications such as image or natural language processing, can not typically be directly applied to scientific data sets and require non-trivial, task-specific modifications

# Introduction to SciML

## Why Scientific Machine Learning ?

- In many applications only limited or low-quality labels are available, while massive unlabeled (often class imbalanced) data sets are common
- Scientific data are often high-dimensional, noisy, heterogeneous, low-signal-to-noise, and multiscale
- Models should [respect or incorporate physical laws](#), constraints, and other scientific domain knowledge
- Robust methods and an ability to quantify uncertainty are required for scientific rigor

### Data-driven ML

- Machine Learning (ML) has become increasingly popular across science
- Traditionally, scientific research has revolved around theory and experiment
- One designs a well-defined theory and then continuously refines it using experimental data
- With rapid advances in the field of ML and the availability of increasing amounts of scientific data, data-driven approaches have become increasingly popular
- With ML, an existing theory is not required, and instead a ML algorithm can be used to analyse a scientific problem using data alone

### ML for scientific problems

- But do ML algorithms actually *understand* the scientific problems they are trying to solve?
- PINNs are a powerful way of incorporating physical principles into ML

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### ML for scientific problems

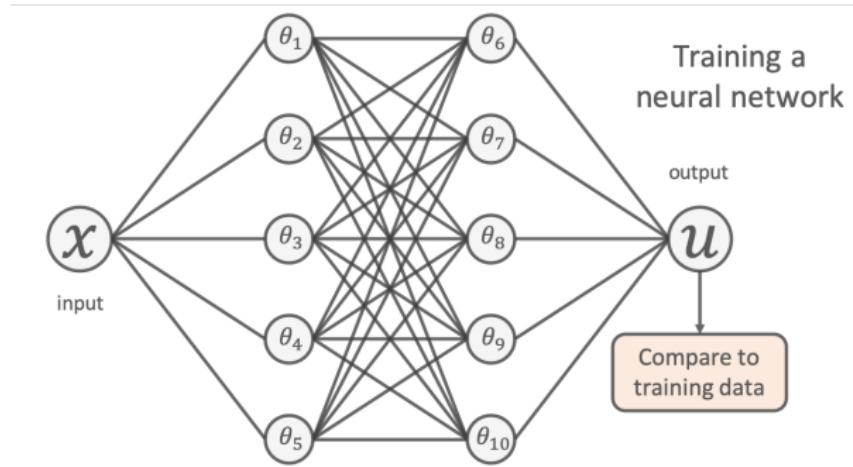
- But do ML algorithms actually *understand* the scientific problems they are trying to solve?
- PINNs are a powerful way of incorporating physical principles into ML

### Learning to model experimental data

- Imagine we are given some experimental data points that come from some unknown physical phenomenon, e.g. the orange points
- A common scientific task is to find a model which is able to accurately predict new experimental measurements given this data

## Learning to model experimental data

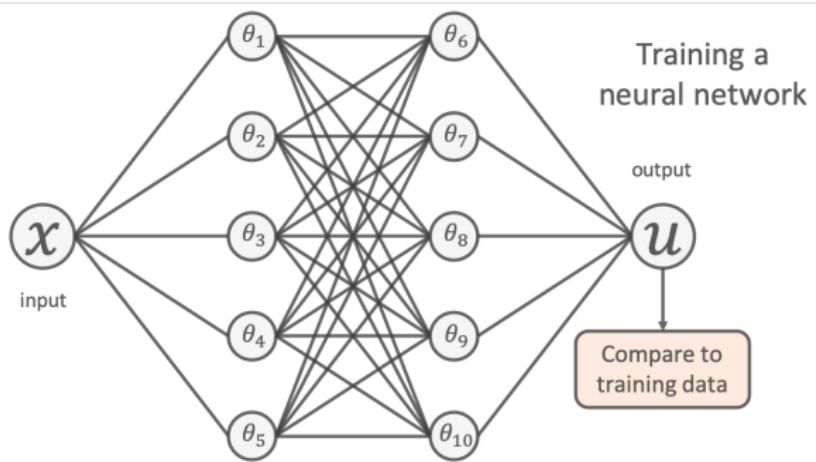
- One popular way of doing this using ML is to use a neural network (NN)
- Given the location of a data point as input (denoted  $x$ ), a NN can be used to output a prediction of its value (denoted  $u$ )



## Learning to model experimental data

- To learn a model, the network's free parameters (denoted by the  $\theta$ s) are tuned
- The goal is to have network's predictions closely match the available experimental data

$$\min \frac{1}{N} \sum_{i=1}^N (u_{\text{NN}}(x_i; \theta) - u_{\text{true}}(x_i))^2$$



### Learning to model experimental data

- The problem is, using a purely data-driven approach like this can have significant downsides
- Whilst the NN accurately models the physical process within the vicinity of the experimental data, it fails to generalise away from this training data

### The rise of scientific machine learning (SciML)

- By only relying on the data, one could argue the NN has not truly *understood* the scientific problem
- What if we inform the NN with some knowledge about the physics of this process?
- For example, that the data points are actually measurements of the position of a damped harmonic oscillator

## The rise of scientific machine learning (SciML)

- There is a well-known differential model of this system

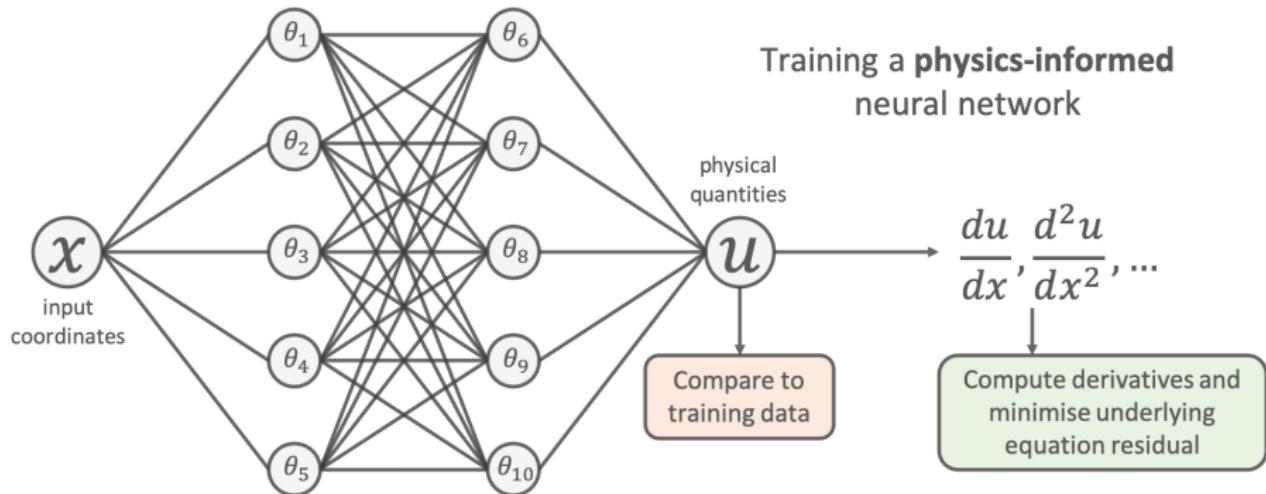
$$m \frac{d^2 u}{dx^2} + \mu \frac{du}{dx} + ku = 0$$

where  $m$  is the mass of the oscillator,  $\mu$  is the coefficient of friction and  $k$  is the spring constant.

- The goal is to look for ways to include this type of prior scientific knowledge into our ML workflow
- The formalism of PINNs is one way to achieve such an adapted ML workflow

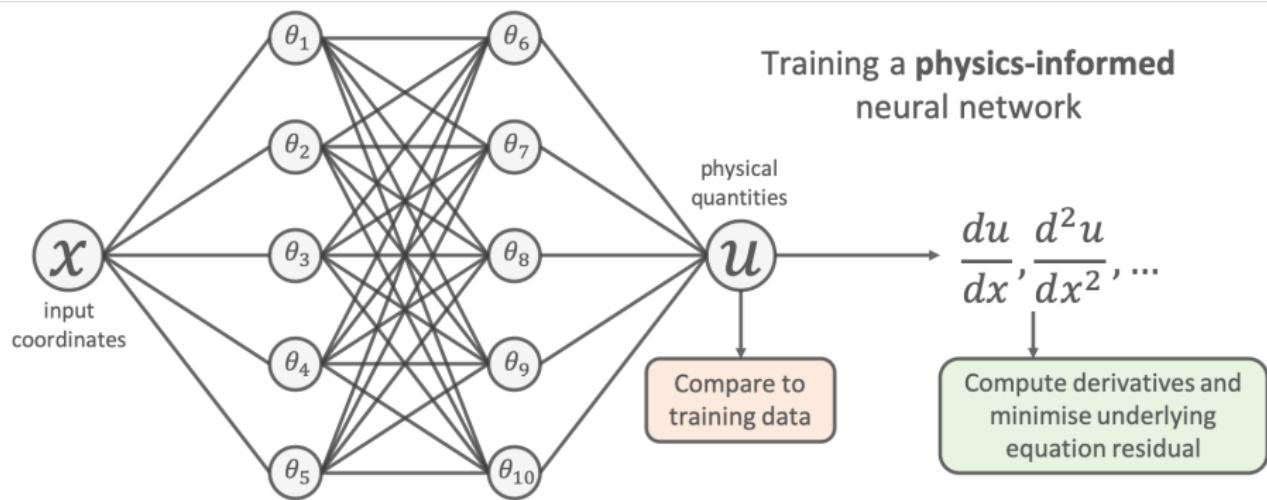
## Principles

- Add the known differential equations directly into the loss function when training the neural network



## Principles

- ① Sample a set of input training locations  $\{x_j\}$  and pass them through the network
- ② Compute gradients of the network's output with respect to its input at the  $\{x_j\}$ s
- ③ Compute residual of the underlying differential equation using these gradients
- ④ Add residual as an extra term in the loss function



## Principles

- Definition of the loss function

$$\min \frac{1}{N} \sum_{i=1}^N (u_{\text{NN}}(x_i; \theta) - u_{\text{true}}(x_i))^2 + \frac{1}{M} \sum_{j=1}^M \left( [m \frac{d^2}{dx^2} + \mu \frac{d}{dx} + k] u_{\text{NN}}(x_j; \theta) \right)^2$$

- The additional *physics loss* tries to ensure that the solution learned by the network is consistent with the known physics

- M. Raissi, P. Perdikaris and G.E. Karniadakis

Physics-informed neural networks: a deep learning framework for solving forward and inverse problems involving nonlinear partial differential equations  
Journal of Computational Physics, Vol. 378, pp. 686-707 (2019)  
<https://doi.org/10.1016/j.jcp.2018.10.045>

- Authors introduce the concept of PINNs
- Exploit Deep Neural Networks (DNN) and leverage their well known capability as universal function approximators
- PINNs are NN that are trained to solve supervised learning tasks while respecting any given laws of physics described by general nonlinear ODEs or PDEs
- For this, differentiate NN with respect to their input coordinates and model parameters
- PINNs can be used for solving two main classes of problems: data-driven solution and data-driven discovery of PDEs
- They can deal with time-dependent problems with continuous time and discrete time PINNs

- PINNs are a type of **universal function approximator** that can embed the knowledge of any physical laws
- They overcome the low data availability of some biological and engineering systems that makes most state-of-the-art ML techniques lack robustness, rendering them ineffective in these scenarios
- The prior knowledge of general physical laws acts in the training of NNs as a **regularization agent** that limits the space of admissible solutions, increasing the correctness of the function approximation
- This way, embedding this prior information into a NN results in enhancing the information content of the available data
- This in turn facilitates the learning algorithm to capture the right solution and to generalize well even with a low amount of training examples

- Most of the physical laws that govern the dynamics of a system can be described by systems of ODEs or PDEs
- In general, these governing equations cannot be solved exactly
- Numerical methods must be used (such as finite differences, finite elements and finite volumes)
- In this setting, these governing equations must be solved while accounting for prior assumptions, linearization, and adequate time and space discretization
- Solving the governing equations of physical phenomena using DL has emerged as a new field of SciML, leveraging the universal approximation and high expressivity of neural networks
- In general, DNNs could approximate any high-dimensional function given that sufficient training data are supplied
- However, such networks do not consider the physical characteristics underlying the problem
- The level of approximation accuracy provided by these networks is still heavily dependent on careful specifications of the problem geometry as well as the initial and boundary conditions
- Without this preliminary information, the solution is not unique and may lose physical correctness

- PINNs are designed to be trained to satisfy the given training data as well as the imposed governing equations
- In this fashion, a NN can be guided with training data that do not necessarily need to be large and complete
- PINNs allow for addressing a wide range of problems in computational science and represent a pioneering technology leading to the development of new classes of numerical solvers for PDEs
- PINNs can be thought of as a meshfree alternative to traditional approaches and new data-driven approaches for model inversion and system identification
- Notably, the trained PINN network can be used for predicting the values on simulation grids of different resolutions without the need to be retrained

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# Poisson equation

## ■ Dirichlet boundary conditions

$$\begin{cases} -u''(x) + c(x)u(x) = f(x), & x \in ]-1, 1[, \\ u(-1) = 0, \quad u(1) = 0. \end{cases}$$

Exact solution for  $c = 0$  and  $f(x) = \pi^2 \sin(\pi x)$ ,  $-1 \leq x \leq 1$ ,

$$u(x) = \sin(\pi x), \quad -1 \leq x \leq 1.$$

## ■ Dirichlet/Neumann boundary conditions

$$\begin{cases} -u''(x) + c(x)u(x) = f(x), & x \in ]-1, 1[, \\ u(-1) = 0, \quad u'(1) = 4. \end{cases}$$

Exact solution for  $c = 0$  and  $f(x) = -2$ ,  $-1 \leq x \leq 1$ ,

$$u(x) = (x + 1)^2, \quad -1 \leq x \leq 1.$$

## ■ Dirichlet/Robin boundary conditions

$$\begin{cases} -u''(x) = f(x), & x \in ]0, 1[, \\ u(-1) = 0, \quad u'(1) = u(1). \end{cases}$$

Same exact solution for  $c = 0$  and  $f(x) = -2$ ,  $-1 \leq x \leq 1$ .

## Heat equation

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) - \frac{\partial^2 u}{\partial x^2}(t, x) = 0, & x \in ]-1, 1[, t > 0, \\ u(t, -1) = u(t, 1) = 0, & t > 0, \\ u(0, x) = u_0(x), & x \in ]-1, 1[. \end{cases}$$

$u_0$  is the initial condition

# Advection equation

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) + a \frac{\partial u}{\partial x}(t, x) = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R} \\ u(0, x) = u_0(x), & x \in \mathbb{R}, \end{cases}$$

$u_0$  is the initial condition,  $a$  is a constant,  $a > 0$ .

- If  $u_0$  belongs to  $C^1(\mathbb{R})$ , existence and uniqueness of a classical solution  $u$  belonging to  $C^1(\mathbb{R}^+ \times \mathbb{R})$

$$u(t, x) = u_0(x - at).$$

# Inviscid Burgers equation / Viscid Burgers equation

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) + u(t, x) \frac{\partial u}{\partial x}(t, x) = 0, & (t, x) \in \mathbb{R}^+ \times \mathbb{R} \\ u(0, x) = u_0(x), & x \in \mathbb{R}, \end{cases}$$

$u_0$  is the initial condition.

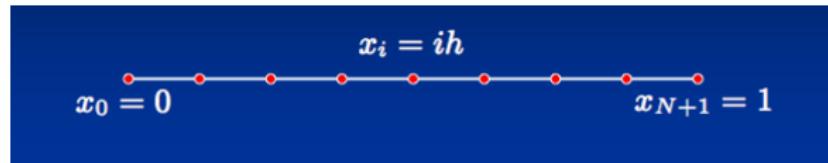
- If  $u_0$  belongs to  $C^1(\mathbb{R})$ , bounded over  $\mathbb{R}$  with a derivative bounded over  $\mathbb{R}$ ,  
if  $u'_0 \geq 0$  : existence and uniqueness of a classical solution  $u$  belonging to  $C^1(\mathbb{R}^+ \times \mathbb{R})$ .

$$\begin{cases} \frac{\partial u}{\partial t}(t, x) + u(t, x) \frac{\partial u}{\partial x}(t, x) = \nu \frac{\partial^2 u}{\partial x^2}(t, x), & (t, x) \in \mathbb{R}^+ \times \mathbb{R} \\ u(0, x) = u_0(x), & x \in \mathbb{R}, \end{cases}$$

$u_0$  is the initial condition,  $\nu$  is a constant,  $\nu > 0$ .

**Hopf-Cole transformation**, transform the viscous Burgers equation into a linear heat equation.

## Forward Euler scheme

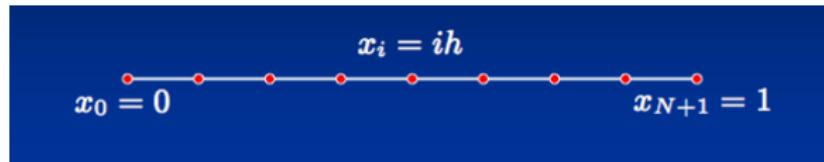


### Forward Euler scheme

1

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} - \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{h^2} = 0, \quad i = 1, \dots, N, \quad n = 0, \dots, M-1,$$

# Forward Euler scheme



## Forward Euler scheme

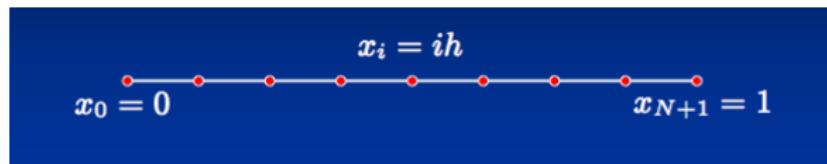
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2

$$u_i^0 = u_0(x_i), \quad i = 1, \dots, N,$$

# Forward Euler scheme



## Forward Euler scheme

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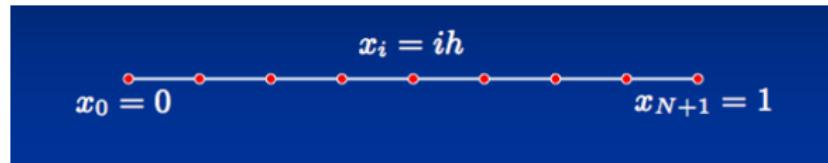
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$$u_i^0 = u_0(x_i), \quad i = 1, \dots, N,$$

3

$$u_0^n = u_{N+1}^n = 0, \quad n = 0, \dots, M.$$

# Forward Euler scheme



## Forward Euler scheme

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$$\frac{u_i^{n+1} - u_i^n}{\Delta t} - \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{h^2} = 0, \quad i = 1, \dots, N, \quad n = 0, \dots, M-1,$$

2

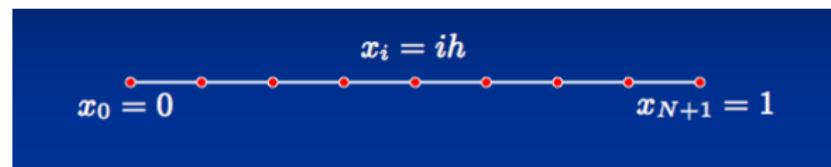
$$u_i^0 = u_0(x_i), \quad i = 1, \dots, N,$$

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$$u_0^n = u_{N+1}^n = 0, \quad n = 0, \dots, M.$$

$$\lambda = \frac{\Delta t}{h^2}$$

# Backward Euler scheme

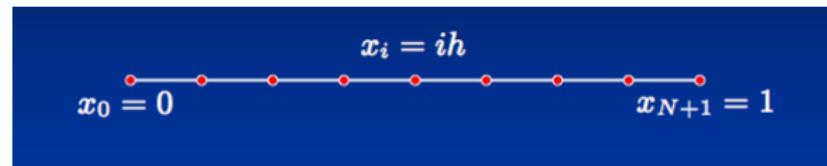


Backward Euler scheme

1

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} - \frac{u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}}{h^2} = 0, \quad i = 1, \dots, N, \quad n = 1, \dots, M-1,$$

# Backward Euler scheme



Backward Euler scheme

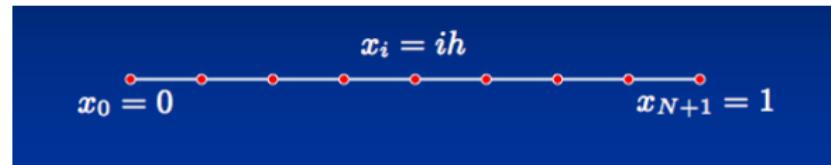
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$$u_i^0 = u_0(x_i), \quad i = 1, \dots, N,$$

# Backward Euler scheme



Backward Euler scheme

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$$\frac{u_i^{n+1} - u_i^n}{\Delta t} - \frac{u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}}{h^2} = 0, \quad i = 1, \dots, N, \quad n = 1, \dots, M-1,$$

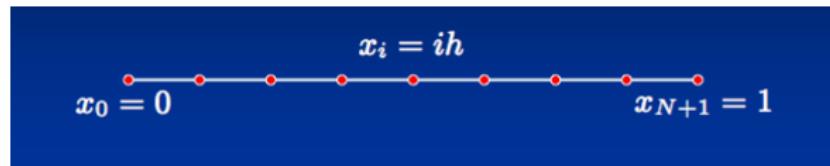
2

$$u_i^0 = u_0(x_i), \quad i = 1, \dots, N,$$

3

$$u_0^n = u_{N+1}^n = 0, \quad n = 0, \dots, M.$$

# Upwind scheme / Lax-Wendroff scheme



1

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + a \frac{u_j^n - u_{j-1}^n}{h} = 0,$$

2

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + a \frac{u_{j+1}^{n+1} - u_{j-1}^{n+1}}{2h} - \frac{a^2 \Delta t}{2} \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{h^2} = 0.$$

# $L^2$ stability analysis

- 1 Let  $u$  be a 1-periodic function, the Fourier coefficients  $\hat{u}(k)$ ,  $k \in \mathbb{Z}$  are complex numbers defined by the integrals

$$\hat{u}(k) = \int_0^1 u(x) e^{-2\pi i k x} dx.$$

- 2 Parseval formula

$$\|u\|_2^2 = \sum_{k \in \mathbb{Z}} |\hat{u}(k)|^2.$$

- 3 Von Neumann analysis, there exists a constant  $K > 0$  independent of  $\Delta t$  and  $h$  such that for all  $n$ ,  $1 \leq n \leq M$ ,

$$\|u^n\|_2 = \left( \sum_{j=1}^N h |u_j^n|^2 \right)^{1/2} \leq K \|u^0\|_2 = \left( \sum_{j=1}^N h |u_j^0|^2 \right)^{1/2}.$$

- 4 Stability, consistency, convergence - Lax-Richtmyer theorem.

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## Back to the future

Let  $n$  in  $\mathbb{N}^*$ ,  $\sigma$  a function from  $\mathbb{R}$  to  $\mathbb{R}$ ,  $M([0, 1]^n)$  the space of finite, signed regular Borel measures and  $C([0, 1]^n)$  the space of continuous functions on  $[0, 1]^n$ .

Definition.

We say that  $\sigma$  is **sigmoidal** if

$$\lim_{t \rightarrow -\infty} \sigma(t) = 0, \quad \lim_{t \rightarrow +\infty} \sigma(t) = 1,$$

we say that  $\sigma$  is **discriminatory** if for a measure  $\mu$  in  $M([0, 1]^n)$

$$\int_{[0,1]^n} \sigma(y^t x + \theta) d\mu(x) = 0$$

for all  $y$  in  $\mathbb{R}^n$  and  $\theta$  in  $\mathbb{R}$  implies that  $\mu = 0$ . □

A sigmoidal function **can not be** a polynomial function.

## Back to the future

Let  $n$  in  $\mathbb{N}^*$ ,  $\sigma$  a function from  $\mathbb{R}$  to  $\mathbb{R}$ ,  $M([0, 1]^n)$  the space of finite, signed regular Borel measures and  $C([0, 1]^n)$  the space of continuous functions on  $[0, 1]^n$ .

Theorem 1, Cybenko (1989).

Let  $\sigma$  be any continuous discriminatory function. The finite sums of the form

$$G(x) = \sum_{j=1}^N \alpha_j \sigma(y_j^t x + \theta_j)$$

are dense in  $C([0, 1]^n)$ . □

Lemma 1, Cybenko (1989).

Any bounded, measurable sigmoidal function,  $\sigma$ , is discriminatory. In particular, any continuous sigmoidal function is discriminatory. □

# Approximation by ReLU Neural Networks

$$g(x) = \begin{cases} 2x, & 0 \leq x \leq 1/2, \\ 2(1-x), & 1/2 < x < 1. \end{cases}$$

